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## Searches for invisible new particles at Belle II

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Belle II has unique sensitivity for a broad class of models postulating the existence of dark matter particles with masses in the MeV–GeV range. We present the recent world-leading physics results from Belle II searches for several non-SM particles: Z' bosons, axion-like particles and dark scalars, through their decays to dark or SM particles; as well as long-lived (pseudo) scalars in *B*-meson decays; and invisible  $\alpha$ -particle in  $\tau$ -decays.

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## 1. INTRODUCTION

Several astrophysical observations suggest the existence of Dark Matter (DM), a component of matter
that does not interact through strong or electromagnetic forces. Although DM constitutes approximately 42
85% of the total matter in our Universe, its nature 43
remains unknown. To date DM is one of the most 44
compelling phenomena in support for physics beyond 45
the Standard Model (SM).

If DM consists of particles feebly coupled with SM 47 14 particles, it could be produced in SM particle annihi-48 15 lations at high-energy colliders. However, the lack of 49 16 evidence of non-SM physics at the electroweak scale, 50 17 leads to hypothesize sub-GeV DM particles feebly in-51 18 teracting with SM particles through non-SM media-52 19 tors. Sub-GeV DM as well as the non-SM mediators 53 20 constitute the dark sector (DS), and efforts to detect 54 21 them have been actively pursued at beam dump and 55 22 high-intensity frontier experiments. 56 23

Belle II [2] is a high-intensity frontier experiment 57 24 that operates at the SuperKEKB  $e^+e^-$  asymmetric-en-58 25 ergy collider [3]. SuperKEKB is a second generation 59 26 B-factory designed to reach the instantaneous lumi-60 27 nosity of  $6 \times 10^{35}$  cm<sup>-2</sup>s<sup>-1</sup> thanks to the usage of the  $\bullet$ 1 28 large crossing-angle nano-beam scheme. The Belle II 62 29 detector [1] is a magnetic spectometer, composed of 63 30 several sub-detectors installed around the interaction 64 31 point, that allow to reconstruct all the products of the 65 32  $e^+e^-$  interaction. During the first data taking run 66 33 (2019-2022), 424 fb<sup>-1</sup> of data were collected, of which 34  $363 \text{ fb}^{-1}$  at the at the energy in the center-of-mass 35  $\sqrt{s} = 10.58$  GeV. 36

Thanks to the excellent reconstruction capabilities for low multiplicity and missing energy signatures, and dedicated triggers, Belle II has a unique sensitivity to the DS and will progressively lead its exploration at the luminosity frontier [17].

## 2. RECENT DARK SECTOR RESULTS AT BELLE II

#### **2.1.** Search for an invisible Z'.

The Z' is a light gauge boson introduced by the  $L_{\mu} - L_{\tau}$  model [4–6]. We search for the invisible decay of Z' through the process  $e^+e^- \rightarrow \mu^+\mu^- Z'(\rightarrow \text{inv.})$ , where Z' is radiated off one of the muons. The Z'could decay invisibly to SM neutrinos, with a branching fraction of  $\mathcal{B}(Z' \to \text{inv.}) \sim 33\%$ , or to kinematically accessible DM candidates with  $\mathcal{B}(Z' \to \text{inv.}) =$ 100%. Signal appears as a narrow enhancement in the recoil mass against the two final-state muons, in events where nothing else is detected. The main backgrounds are QED radiative di-lepton and four-lepton final states, which are suppressed using a neural-network trained simultaneously for all Z' masses [9], and fed with kinematic variables sensitive to the origin of the missing energy: in the signal, the Z' is produced as final-state radiation (FSR); in the background, the missing energy is due to neutrinos or undetected particles. From 2D template fits to the recoil mass squared, in bins of recoil polar angle, we did not observe any significant excess in  $79.7 \text{ fb}^{-1}$  of data, and we set 90% C.L. upper limits on the coupling of the  $L_{\mu} - L_{\tau}$ model, g', as a functions of the Z' mass,  $M_{Z'}$ . We exclude the region favored for the  $(g-2)_{\mu}$  anomaly in the mass range  $0.8 < M_{Z'} < 5$  GeV/ $c^2$  for the fully invisible  $L_{\mu} - L_{\tau}$  model (Fig. 1) [7, 8].

# **2.2.** Search for $e^+e^- \rightarrow \mu^+\mu^- X(\rightarrow \tau^+\tau^-)$

We search for a  $X \to \tau^+ \tau^-$  resonance, where X could be a Z', a leptophilic dark spin-0 particle (scalar) S or an axion-like particle (ALP), in  $e^+e^- \to \mu^+\mu^-\tau^+\tau^-$  events, with  $\tau$  decaying to one charged particle. The S is an hypothetical particle that couples preferentially to charged leptons through Yukawa-like couplings [11]. Axion-like particles are pseudo-scalars that appear in many SM extensions [12, 13].

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Figure 1. Observed 90% C.L. upper limits and corresponding expected limits on the g' coupling as a function of the Z' mass, assuming  $\mathcal{B}(Z' \to \text{inv.}) = 100\%$ .

Similarly to the  $Z' \to inv.$  analysis, we search for 75 a narrow enhancement in the recoil mass against two 76 oppositely charged muons, in four-track events with 77 zero net charge. Standard Model backgrounds are sup-78 pressed with eight neural-networks fed with kinematic 79 variables sensitive to the X-production mechanism as 80 FSR off one of the two muons, and trained in different 81 X-mass regions. From extended maximum likelihood 82 fits to the recoil mass distribution, we did not observe 83 any significant excess in  $62.8 \text{ fb}^{-1}$  of data. We derive 84 world-leading 90% C.L. upper limits on the S-coupling 85  $\xi$  for  $m_S > 6.5 \text{ GeV}/c^2$ , and on the ALP-lepton cou-86 pling  $|C_{\ell\ell}|/\Lambda$ , assuming equal ALP-couplings to the 87 three lepton families and zero couplings to all other 88 particles (Fig. 2) [10]. 80

# **2.3.** Search for $e^+e^- \rightarrow \mu^+\mu^- X(\rightarrow \mu^+\mu^-)$

We search for a  $X \to \mu^+ \mu^-$  resonance in  $e^+ e^- \to \mu^{13}$ 90  $\mu^+\mu^-\mu^+\mu^-$  events as a narrow enhancement in the 14 91 dimuon mass distribution in four-track events with15 92 zero net charge and no extra-energy. The dominant<sup>116</sup> 93 background are SM four-muon final-state processes.<sup>117</sup> 94 Background is suppressed applying five neural-net118 95 works fed with kinematic variables sensitive to the19 96 X-production mechanism as FSR off one of the  $two_{120}$ 97 muons, and on the presence of a resonance in both the21 98 candidate and the recoil muon pairs, and trained in dif122 99 ferent X-mass ranges. From extended maximum likeli123 100 hood fits to the dimuon mass distribution, we did not124 101 observe any significant excess in 178 fb<sup>-1</sup> of data, and 125 102 we set 90% C.L. upper limits on the cross section of 126 103 the process. We interpret the results obtained on the27 104 cross section as 90% C.L. limits on the g' coupling of 128 105 the  $L_{\mu}-L_{\tau}$  model, and on the coupling of a muonphiliq 29 106 scalar S with muons [14]. Despite the small data-set<sub>130</sub> 107 used, we obtain competitive results with the existing131 108 limits on q' from BABAR and Belle, which performed 32 109 the analysis with 514  $\text{fb}^{-1}$  and 643  $\text{fb}^{-1}$  respectively<sub>133</sub> 110 We set the first limits for the muonphilic scalar modeh34 111 from a dedicated search (Fig. 3). 112 135



Figure 2. Observed 90% C.L. upper limits (UL) and corresponding expected limits as a function of the mass on (top) the leptophilic scalar coupling  $\xi$ , and on (bottom) the ALP-lepton coupling  $|C_{\ell\ell}|/\Lambda$ .

# 2.4. Search for a long-lived (pseudo) scalar in $b \rightarrow s$ transitions

Extensions of the SM introduce a new light scalar S that may give mass to DM particles. The scalar S would mix with the SM Higgs boson through a mixing angle  $\theta$ , and would be naturally long-lived for small values of  $\theta$ .

We search for  $B^0 \to K^{*0} (\to K^+ \pi^-) S$  and  $B^+ \to$  $K^+S$  events, with  $S \to x^+x^-$  and  $x = e, \mu, \pi, K$ . The signal signature is a prompt decay of the kaon and two opposite-charged particles from a displaced vertex. The signal yield is extracted through extended maximum likelihood fits to the reduced invariant mass of S,  $m_{S \to xx}^{reduced} = \sqrt{M_{S \to xx}^2 - 4m_x^2}$ , in order to improve the modeling of the signal width close to the kinematic thresholds. Main background components are the combinatorial  $e^+e^- \rightarrow q\bar{q}$ , suppressed by requiring a kinematics similar to B-meson expectations, peaking  $K_S^0$ , vetoed in the invariant mass  $m_{S \to \pi\pi}$ , and further peaking backgrounds, suppressed by tightening the displacement selections. We did not observed any significant excess in  $189 \text{ fb}^{-1}$  of data, and we set the first model-independent limits at 95% C.L. on  $\mathcal{B}(B \to KS) \times \mathcal{B}(S \to x^- x^+)$  as a function of the scalar mass  $m_S$  for different S-lifetimes (Fig. 4) [15].



Figure 3. Observed 90% C.L. upper limits and corresponding expected limits as a function of the mass on (top) the g' coupling of the  $L_{\mu} - L_{\tau}$  model, and on (bottom) the muonphilic dark scalar model.

#### **2.5.** Search for the $\tau \to \ell \alpha$ decay

Charged-lepton flavour violation (LFV) is allowed 136 in various extensions of the SM, however it has never 137 been observed. In these extensions, the LFV processes 163 138 could be mediated by a new hypothetical  $\alpha$ -boson. We<sub>164</sub> 139 search for an invisible  $\alpha$  produced in the  $\tau \to \ell \alpha$  de<sub>ies</sub> 140 cay, with  $\ell = e, \mu$ , in  $e^+e^- \to \tau^+\tau^-$  events. In the cen-141 ter-of-mass frame,  $\tau$  pairs are produced back-to-back 142 so that each  $\tau$ -decay products are contained in two 143 separate hemispheres. The tag hemisphere contains 144 three charged hadrons from  $\tau_{\text{tag}} \to h^- h^+ h^- \nu_{\tau}$ , with  $h = \pi$ , K, while the *signal* hemisphere contains only 145 146 one charged lepton from the  $\tau_{sig}^- \to \ell^- \alpha$  decay. 147

For this analysis,  $\tau \to \ell \nu_{\tau} \bar{\nu_l}$  is an irreducible background. However, the lepton momentum has a broad distribution for the background, while it depends only on the  $\alpha$ -mass for the signal and it appears as a bump over the irreducible background.

In particular, we search for an excess over the normalized lepton energy spectrum  $x_{\ell}$  of  $\tau \rightarrow \ell \nu_{\tau} \bar{\nu}_{l}$ , where  $x_{\ell} = 2E_{\ell}^{*}/m_{\tau}$ , performing template fits. The energy  $E_{\ell}^{*}$  is defined in the approximate rest frame of  $\tau_{sig}$ , i.e. where  $E_{\tau} \approx \sqrt{s}/2$  and the  $\tau_{sig}$  direction is opposite to the  $\tau_{\text{tag}}$  direction. We did not find any significant excess in 62.8 fb<sup>-1</sup> of data, and we set world-leading 95% C.L. upper limits to  $\mathcal{B}(\tau \to \ell \alpha)/\mathcal{B}(\tau \to \ell \bar{\nu}_{\ell} \nu_{\tau})$ , as a function of the  $M_{\alpha}$  mass (Fig. 5) [16].



Figure 4. Observed 95% C.L. on  $\mathcal{B}(B \to KS) \times \mathcal{B}(S \to l^+l^-/h^+h^-)$  as a function of the scalar mass  $m_S$  for different lifetimes  $c\tau$ .

#### 3. SUMMARY

We presented the latest Belle II world-leading results from DS searches, published with partial data-sets of 424 fb<sup>-1</sup> collected to date. New results with improved analyses and more data are coming.



Figure 5. Upper limits at 95% C.L. on the ratio  $\mathcal{B}(\tau \rightarrow e\alpha)/\mathcal{B}(\tau \rightarrow e\bar{\nu}_e \nu_{\tau})$ .

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- 166
   [1] T. Abe et al. (Belle II Collaboration), arXiv:1011.035263

   167
   [physics.ins-det] (2010).
   184
- 168 [2] E. Kou et al. (Belle II Collaboration), Prog. Theor.185
   Exp. Phys. 2019, 123C01 (2019).
   186
- 170
   [3] K. Akai, K. Furukawa, and H. Koiso (SuperKEKB187

   171
   Accelerator Team), Nucl. Instrum. Meth. A 907, 188188

   172
   (2018).
- 173 [4] X. G. He, G. C. Joshi, H. Lew, and R. R. Volkas, Phys190 174 Rev. D 43, R22 (1991).
- 175 [5] B. Shuve and I. Yavin, Phys. Rev. D 89, 113004192 176 (2014).
- [6] W. Altmannshofer, S. Gori, S. Profumo, and F. S194
   Queiroz, J. High Energy Phys. 12, 106 (2016).
- [7] I. Adachi et al. (Belle II Collaboration), Phys. Rev196
   Lett. 124, 141801 (2020).
- [8] I. Adachi et al. (Belle II Collaboration), Phys. Rev198
   Lett. 130, 231801 (2023).

- [9] F. Abudinén et al., Eur. Phys. J. C 82, 121 (2022).
- [10] I. Adachi et al. (Belle II Collaboration) Phys. Rev. Lett. 131, 121802 (2023).
- [11] B. Batell, N. Lange, D. McKeen, M. Pospelov, and A. Ritz, Phys. Rev. D 95, 075003 (2017).
- [12] M. Bauer, M. Neubert, S. Renner, M. Schnubel, and A. Thamm, J. High Energy Phys. 09 (2022) 56.
- [13] M. Bauer, M. Neubert, and A. Thamm, J. High Energy Phys. 12 (2017) 044.
- [14] D. Forbes, C. Herwig, Y. Kahn, G Krnjaic, C. M. Suarez, N. Tran, and A. Whitbeck Phys. Rev. D 107, 116026 (2023).
- [15] I. Adachi et al. (Belle II Collaboration) Phys. Rev. D 108, L111104 (2023).
- [16] I. Adachi et al. (Belle II Collaboration), Phys. Rev. Lett. 130, 181803 (2023).
- [17] L. Aggarwal et al. (Belle II Collaboration) arXiv:2207.06307 [hep-ex] (2022)