

m(B*) - m(B) = 45.2 MeV

 B^{0}/B^{0} from a Y(4s) decay are supposed to be in an entangled state _

$$
\Psi = \frac{1}{\sqrt{2}} \left[|B^0(p)\rangle |\overline{B}^0(-p)\rangle - |\overline{B}^0(p)\rangle |B^0(-p)\rangle \right]
$$

If one B decays, the common wave function collapses and the B^0/B^0 are in a defined state. γβcτ/r(B⁰) ~ 5x10¹⁰ => well separated spatially

_

Our measurements of Δm_d and TDCPV are based on the entanglement (B-tag)

- 1) Can we demonstrate the entanglement (e.g. checking Bell's inequality ?)
- 2) How certain are we that the entanglement is always 100% ?

 $Y(4s) \rightarrow B^0 \overline{B^0} \gamma$ decoherence due to interaction with (BSM) background fields _

Such effects could lead to systematic errors of our TDCPV measurements

Accelerator

From KEKB to SuperKEKB

Belle (II)

B⁰ B⁰ mixing _

Due to weak interaction a $B⁰$ can transform into its antiparticle Formally this is described by a new (weak) base of $\mathsf{B^0}_\mathsf{L}$ and $\mathsf{B^0}_\mathsf{H}$

$$
|B_{\text{L,H}}\rangle = p|B_q^0\rangle \pm q|\overline{B}_q^0\rangle \qquad \text{ |p/q| = 1 (CP cons.)}
$$

These two states interfere and resulting in time dependent oscillations

$$
P(B^0 \rightarrow \overset{-}{B^0}) = \frac{1}{2} \Gamma \exp(-\Gamma t) (1 - \cos(\Delta m t))
$$

$$
P(B^0 \rightarrow B^0) = \frac{1}{2} \Gamma \exp(-\Gamma t) (1 + \cos(\Delta m t))
$$

 Δ m: mass difference of $\mathsf{B^0_H}$ and $\mathsf{B^0_L}$

CP violation

Weak interaction is also a source for CP violation

CP violation is a consequence of the complex phase of the CKM matrix

It happens if two (or more) weak amplitudes interfere

 $A_f(\underline{B}^0 \rightarrow f) = A_1 \exp(i\phi_w) + A_2 \exp(i\phi_s)$ $A_f(B^0 \to I) = A_1 \exp(i\phi_w) + A_2 \exp(i\phi_s)$
 $A_b(\overline{B}^0 \to \overline{I}) = A_1 \exp(-i\phi_w) + A_2 \exp(i\phi_s)$ => $|A_f|^2 \neq |A_b|^2$ \Rightarrow $|A_f|^2 \neq |A_b$

In decays this happens either by interference of different decay amplitudes (tree and penguin) or (more important) by interference of mixing and decay

Precise measurements of CP violation are the primary goals of the Belle and Belle II experiments

 $(1,0)$

$$
V_{\rm CKM} \equiv V_L^u V_L^{d\dagger} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \qquad \begin{matrix} \bar{\varphi}, \bar{\eta} \\ \bar{V}_{cd} V_{ub}^* \\ \bar{V}_{cd} V_{cb}^* \end{matrix} \qquad \qquad \begin{matrix} \bar{\varphi}, \bar{\eta} \\ \bar{V}_{cd} \\ \bar{V}_{cb} \end{matrix} \qquad \qquad \begin{matrix} \bar{\varphi}, \bar{\eta} \\ \bar{V}_{cd} \\ \bar{V}_{cb} \end{matrix}
$$

Phenomenology

 $Y(4s)$ created at t=0 and decays in B⁰/B⁰ _

Tag: one Bº decays at t₁ in a way that we know whether it is Bº or Bº (e.g. $\mu^+\nu$ D⁻) (in practice: use a BDT or NN to determine the decay state) _

The otherB⁰ (always in the opposite state because of entanglement) starts oscillating

Sig: the other B⁰ decays at t_2 in the signal mode we are interested in

From the spatial separation and the known boost we determine $\Delta t = t_2 - t_1$

Phenomenology

Oscillations:

$$
P(B^0 \rightarrow B^0) = \frac{1}{2} \Gamma \exp(-\Gamma \Delta t) (1 - \cos(\Delta m \Delta t))
$$

$$
A_{osc} = \frac{N(B^0 \to B^0) - N(B^0 \to \overline{B^0})}{N(B^0 \to B^0) + N(B^0 \to \overline{B^0})} = \cos(\Delta m \Delta t)
$$

In reality: take into account mistag, resolution, background

Mixing induced (or time dependent) CP violation (TDCPV): _

'Golden mode' $(q=+1); B^0(q=-1) \rightarrow J/\psi K_s$ _

 $P(\Delta t, q) = \frac{1}{4} \Gamma \exp(-\Gamma \Delta t) \{1 + q \ [\text{S} \sin(\Delta m \Delta t) - \text{C} \cos(\Delta m \Delta t)]\}$

For J/
$$
\psi
$$
 K_s $S = \sin(2\beta)$ $C \sim 0$

$$
a_{CP} = \frac{N(B^0 \to f_{CP}) - N(B^0 \to f_{CP})}{N(B^0 \to f_{CP}) + N(B^0 \to f_{CP})} = S \sin(\Delta m \Delta t)
$$

Classical Aspect type correlation experiment A. Aspect et al, Phys. Rev. Lett 49, 1804 (1982)

 $\Delta p \cdot \Delta q \geq \frac{1}{2}t$

spin-singlet state of photons or particles: $\frac{1}{\sqrt{2}}[\ket{\uparrow}_1\ket{\Downarrow}_2 - \ket{\Downarrow}_1\ket{\uparrow}_2]$

correlation coeffs in data vs QM optimum relative angles 22.5° and 67.5°

FIG. 3. Correlation of polarizations as a function of the relative angle of the polarimeters. The indicated errors are ± 2 standard deviations. The dotted curve is not a fit to the data, but quantum mechanical predictions for the actual experiment. For ideal polarizers, the curve would reach the values ± 1 .

- . Bell's Theorem (via Clauser, Horne, Shimony, and Holt):
	- correlation coeff: $E(\vec{a}, \vec{b}) = \frac{R_{++}(\vec{a}, \vec{b}) + R_{--}(\vec{a}, \vec{b}) R_{+-}(\vec{a}, \vec{b}) R_{-+}(\vec{a}, \vec{b})}{R_{++}(\vec{a}, \vec{b}) + R_{--}(\vec{a}, \vec{b}) R_{+-}(\vec{a}, \vec{b}) + R_{-+}(\vec{a}, \vec{b})}$
	- $S = E(\vec{a}, \vec{b}) E(\vec{a}, \vec{b}') + E(\vec{a}', \vec{b}) + E(\vec{a}', \vec{b}')$
	- $|S| \le 2$ for any local realistic model; $S_{QM} = \pm 2\sqrt{2}$ for optimal settings

 $S = 2.697 \pm 0.015$; cf. $S_{OM} = 2.70 \pm 0.05$

Replace \uparrow by B⁰ and \downarrow by B⁰

Bell's inequality with B⁰B⁰ _

Replace the angle between the polarizers by $\Delta m \Delta t$

S = 2.70 (for 22.5 and 67.5°)

 $\Delta p \cdot \Delta q \geq \frac{1}{2}k$

(A. Go & Chung Li, quant-ph/0310192v1 (2003): $S = 2.725 \pm 0.167 \pm 0.092$ (Belle data))

Non QM test: Assume complete decoherence (both B oscillate independently from t=0)

S = 2.61 S(decoherent) < 2 for $\Delta m/\Gamma > 2$

What's wrong? => high correlation due to short lifetime

Bertlmann, Bramon, Garbarino, Hiesmayr, Phys. Lett. A 332, 355-360 (2004)

crucial parameter $x_d = \Delta m_d/\Gamma_d$: rate of oscillation relative to decay

Bell test impossible if $x < 2.0$:

Furthermore:

we rely on the random decays of both B

no (reasonable) way to **actively** determine \leq a,b -> Δ m Δ t

A kind of Maxwell's demon could tune hidden parameters ($\mathsf{t}_1, \mathsf{t}_2$, decay type) so that QM and Bell's inequality is emulated, despite it's not QM and local. No practical way to close this loophole

We can still calculate effects of special decoherent or non local models Fit (modified) time dependence to data

• Spontaneous decoherence (SD):

entanglement is lost at t=0 for a certain fraction of events decoherence fraction ζ

• Pompili Selleri model (PS) (Eur. Phys. J. C14, 469 (2000)):

local realism, which reproduces oscillation phenomenology

• Lindblad type decoherence

coherence is gradually lost within a certain decoherence time

Time dependence if decoherent

 $\Delta p \cdot \Delta q \geq \frac{1}{2} t$

B⁰/B⁰ Oscillations: Probability to measure a same sign (SS) event: _

a) Full coherence:

$$
P(B^0B^0, \overline{B^0B^0}) = \frac{1}{2}\Gamma \exp(-\Gamma t_2)(1 - \cos \Delta m t_2)
$$

b) Full decoherence (Spontaneos decoherence, SD)

$$
P(B^0B^0, \overline{B^0B^0}) = \frac{1}{2}\Gamma \exp(-\Gamma t_2)(1 - \frac{1}{2}\frac{\Delta m + 2\Gamma^2}{\Gamma^2 + \Delta m^2}\cos\Delta m t_2 + \frac{1}{2}\frac{\Gamma\Delta m}{\Gamma^2 + \Delta m^2}\sin\Delta m t_2)
$$

Damping by
$$
\frac{1}{2}\frac{\Delta m + 2\Gamma^2}{\Gamma^2 + \Delta m^2} \sim 0.81
$$

Additional SIN-term:
$$
\frac{1}{2} \frac{\Gamma \Delta m}{\Gamma^2 + \Delta m^2} \sim 0.24
$$

(using PDG averages: $\Delta m = 0.505 + 0.002 \text{ ps}^{-1}$, $\Gamma = 0.658 + 0.002 \text{ ps}^{-1}$)

Time Dependence (SD)

Like & unlike

The damping is probably difficult to measure, as it could be interpreted as mistag

The SIN term (or phase shift) should be measurable

Similar damping and phase shifts occur in measurements of TDCPV

In principle the damping (if ignored) would lead to a wrong measurement of $sin(2\beta)$, This might be compenated out if the mistag is taken from the oscillation measurement (same damping)

The phase shift would lead to a cross talk between S and C terms

coherent - incoherent

Time integrated oscillations: measure $\chi = \frac{1}{2} \Delta m^2 / (\Delta m^2 + \Gamma^2)$ using dilepton events

Use Δ m from LHCb as reference (no coherence at LHCb!)

Coherence: $R = (N^{++} + N^{-})/(N^{++} + N^{-} + N^{+-}) = \chi$

SD: $R = (N^{++} + N^{-})/(N^{++} + N^{-} + N^{+-}) = 2(\chi - \chi^2)$

Indirect Indicators: TDCPV

 $\cos \Delta m t_2$

Compare S_{CP} and C_{CP} measurements at LHCb (always decoherent) and at Y(4s)

H.-G. Moser, Quantum Observables for Collider Physics, Firence, 6.11.-10.11. 2023

 $\Delta p \cdot \Delta q \geq \frac{1}{2}k$

A.Go & Belle, Phys. Rev. Lett 99 (2007) 131802

Belle's most current sin $2\phi_1$, $|\lambda|$, τ_B , Δm_d measurement at the time:

- 152×10^6 $B\overline{B}$ pairs
	- $5\times$ the discovery dataset
	- \bullet $\frac{1}{5}$ x the eventual dataset
- \bullet 5417 CP- and 177368 flavoureigenstate B -decay candidates
- sample purities vary 63-98% depending on the decay mode
- multivariate flavour-tagging of the other B decay; $\epsilon_{\text{eff}} = 28.7\%$
- $\Delta m_d = (0.511 \pm 0.005 \pm 0.006) \text{ ps}^{-1}$ cf. (0.5065 ± 0.0019) ps⁻¹ PDG23
- We then adapted this in various ways ...

- restrict $177368 \rightarrow 84823$ flavour eigenstates, choosing only $B^0 \to D^{*-}\ell^+\nu$ where the lepton explicitly determines the B-flavour
- restrict 84823 \rightarrow 8565 by choosing only the best flavour tags of the other B : highest of 7 purity categories; leptons only
- \bullet signal relies on $D^{*-} \to \overline{D}{}^0 \pi^-$ tag: energy release $Q \ll m_\pi \ll m_D$

• estimate background under peak using sideband region:

H.-G. Moser, Quantum Observables for Collider Physics, Firence, 6.11.-10.11. 2023

- Background subtraction, deconvolution of Δt resolution, mistag.....
	- ...and fitting for the B^0 lifetime:

Fit functions

QM: standard quantum mechanical entanglement SD: spontaneous decoherence PM: A. Pompili & F. Selleri, Eur. Phys. J. C14, 469 (2000)

Results

fit: float Δm_d subject to WA-sans-(Belle+BaBar): (0.496 \pm 0.014) ps⁻¹

• "SD fraction": $(1 - \zeta_{B0\overline{B}0})A_{QM} + \zeta_{B0\overline{B}0}A_{SD}$, $\zeta_{B0\overline{B}0} = 0.029 \pm 0.057$

• Pompili-Selleri class: QM-like states, stable mass, flavor correlations; QM predictions for *single B-mesons* preserved

 $\Delta p \cdot \Delta q \geq \frac{1}{2}k$

What do we learn?

SD excluded by 13 σ , but more relevant is the fraction of decoherent events

 $f = (1-\zeta) A_{OM} + \zeta A_{SD}$ $\zeta = 0.029 \pm 0.057$

A fraction of ~10% is still possible!

This could lead to shift of our Scp measurements by

 $\Delta S_{\text{cp}} \sim 0.012$ (@ Y(4s))

The total systematic errors of the Belle II J/ ψ K_s analysis is 0.014 !

largest single systematic error?

Repeat Belle analysis with higher statistics, more channels, better resolution

$$
B^0\to D^-\pi^+, D^{*-}\pi^+, D^{*-}\rho^+
$$

Make use of better vertex resolution and smaller interaction region:

 $\gamma\beta\tau c = 0.125$ mm Not perfect yet, but some chance to limit t_1

Transverse separation ~50 µm Vertex resolution $\sigma r_e s \sim 20 \mu m$

Discrimination Power

Access to t_1 adds a new dimensions and should result higher sensitivity

Entanglement: depends only on Δt Decoherence: depends on t_1 and Δt

Setting a lower limit on t_1 could also make a EPR type measurements possible (randomize)

Conclusions

- 'Ascent' style experiments to check Bell's inequality are not possible with $\rm Y(4s) \rightarrow B^0$ $\rm \overline{B^0}$ _
	- $-$ no active measurement (random decay of the B⁰): conspiracy loophole!
	- $-$ short B⁰ lifetime induces correlations which violate Bell's inequality even for a local realistic scenario
- QM and alternative models can be tested fitting the time dependence of B^0 oscillations. Belle analysis: alternative scenarios excluded by 13σ (SD) and 5.1σ (PS)
- A fraction of ~10% of decoherent events is still compatible with the data
- Possible systematic error to out TDCPV measurements (so far not taken into account!)
- Belle II has the potential to improve on this
- Questions to theory: what mechanisms (SM or BSM) could lead to a loss of coherence?

