

Introduction to the Physics of Particle Flavors



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Belle II Explorer
June 17, 2024



What is flavor?

- Flavor**: indicates the different types of Standard Model (SM) fermions, both quarks and leptons.



WIKIPEDIA
The Free Encyclopedia

mass	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$
charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$
spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$
	u up	c charm	t top
	up type quarks		
	$\approx 4.7 \text{ MeV}/c^2$	$\approx 96 \text{ MeV}/c^2$	$\approx 4.18 \text{ GeV}/c^2$
	$-\frac{1}{3}$	$-\frac{1}{3}$	$-\frac{1}{3}$
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$
	d down	s strange	b bottom
	down type quarks		
	$\approx 0.511 \text{ MeV}/c^2$	$\approx 105.66 \text{ MeV}/c^2$	$\approx 1.7768 \text{ GeV}/c^2$
	-1	-1	-1
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$
	e electron	μ muon	τ tau
	charged leptons		
	$< 1.0 \text{ eV}/c^2$	$< 0.17 \text{ MeV}/c^2$	$< 18.2 \text{ MeV}/c^2$
	0	0	0
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino
	neutrinos		

Flavour in particle physics

Flavour quantum numbers

- Isospin: I or I_3
- Charm: C
- Strangeness: S
- Topness: T
- Bottomness: B'

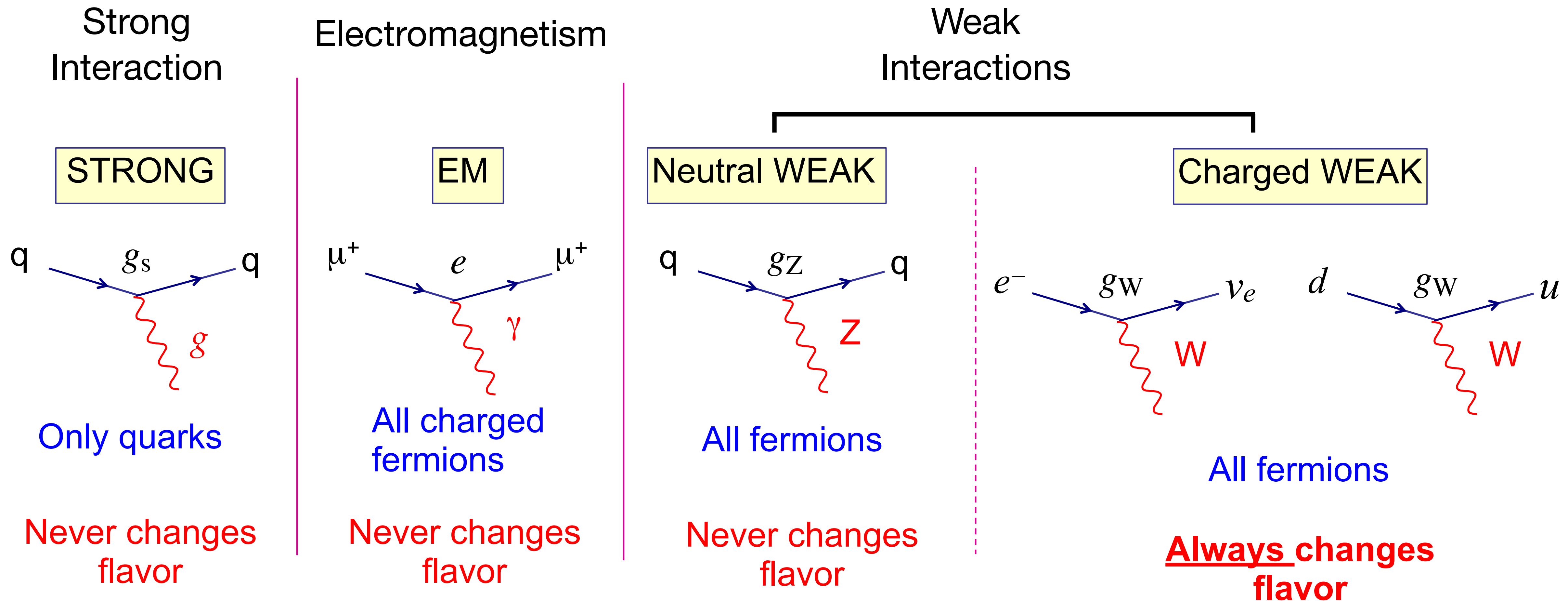
Related quantum numbers

- Baryon number: B
- Lepton number: L
- Weak isospin: T or T_3
- Electric charge: Q
- X-charge: X

Combinations

Standard Model Vertices

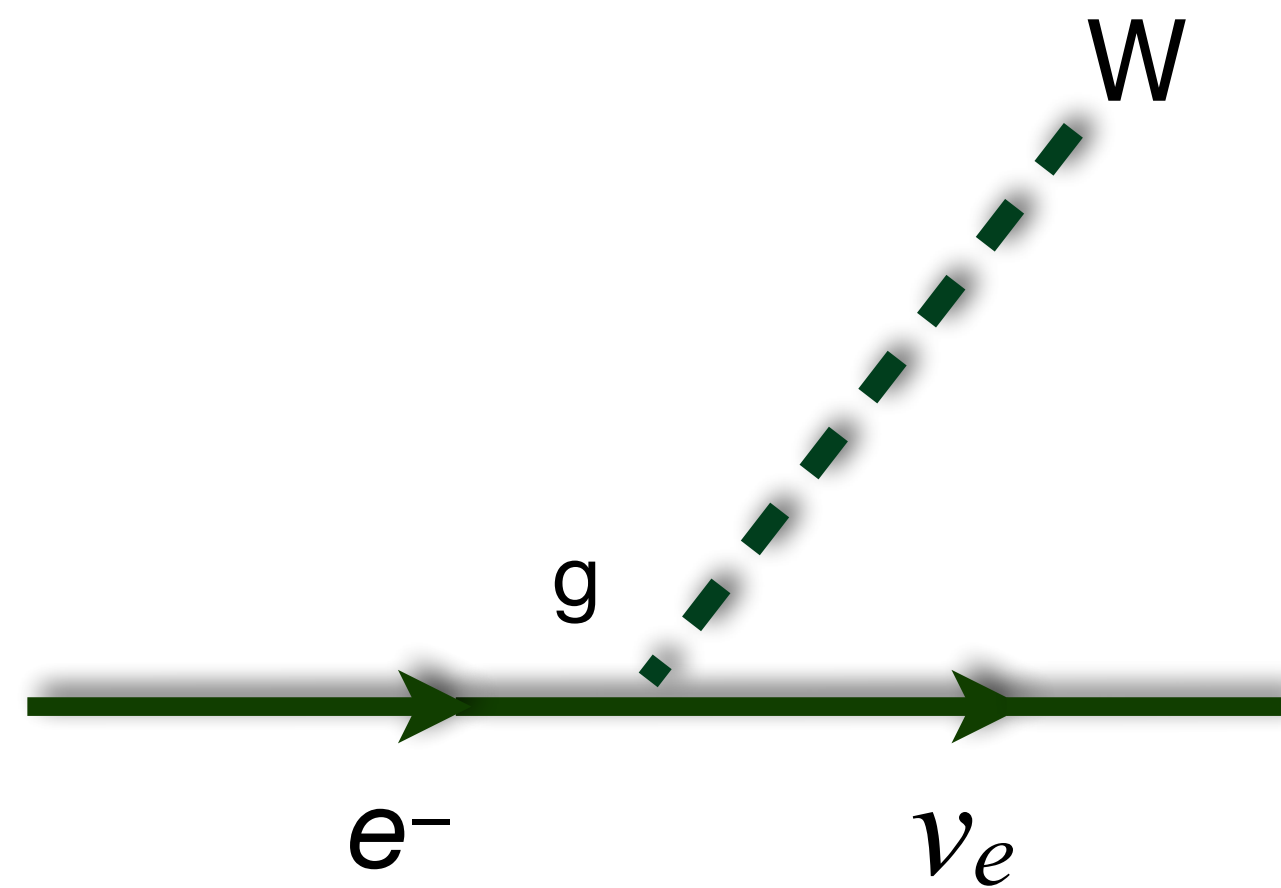
- Interaction of **gauge bosons** with **fermions** described by SM vertices
- Properties of the **gauge bosons** and **nature of the interaction** between the bosons and fermions determine the properties of the interaction



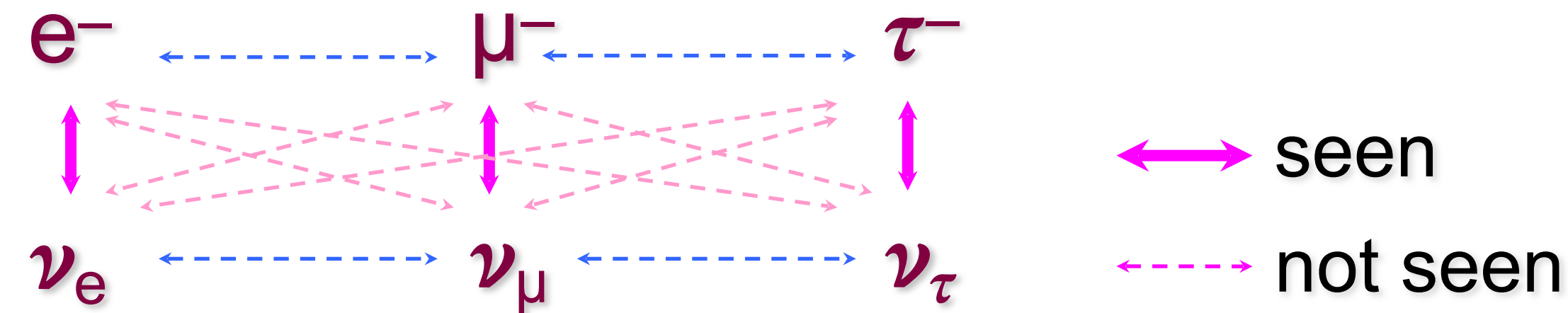
Flavor Change

Only the weak interaction allows flavor change

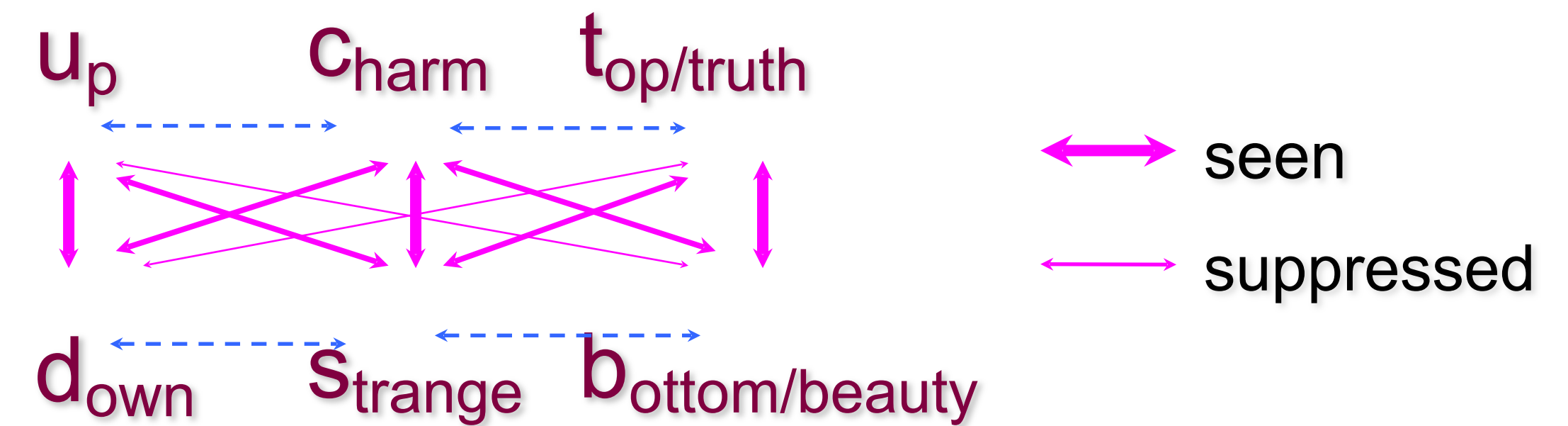
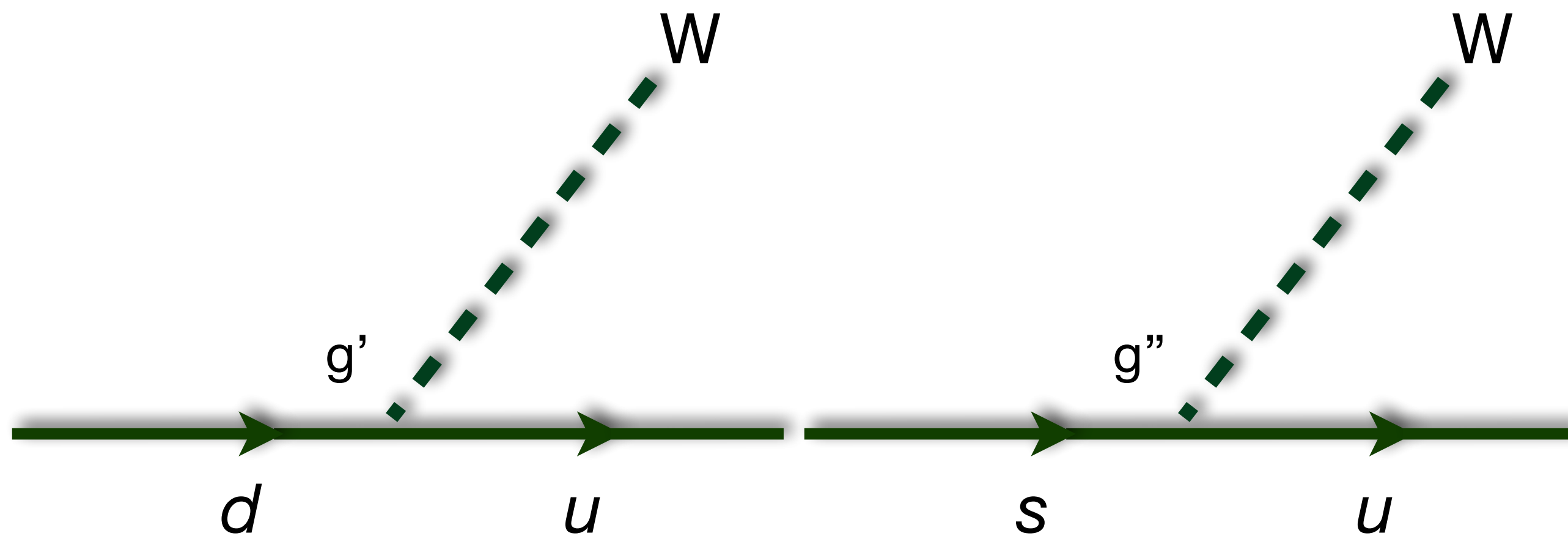
Leptons: no generation crossing



W^\pm "charged current"



Quarks: generation crossing

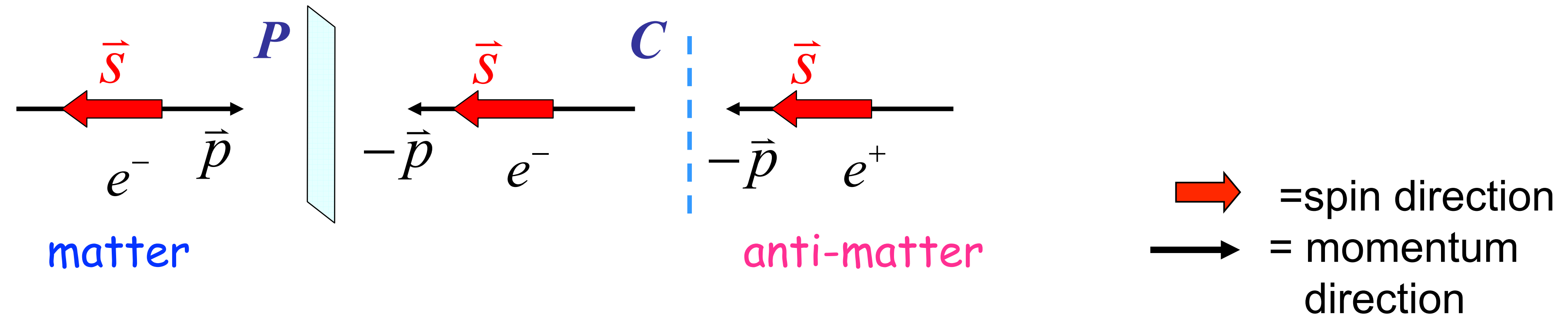


Why do we study flavor physics?

- Obtain deeper understanding of the fundamental flavor structure of the nature.
- Charge Parity (CP) violation and its connection to the matter dominated Universe.
(Matter-antimatter asymmetry of universe requires CP -violating interactions (Sakharov 1967))
- Discover effects of new particles and forces beyond the Standard Model – even particles too massive to be produced at the high energy colliders.

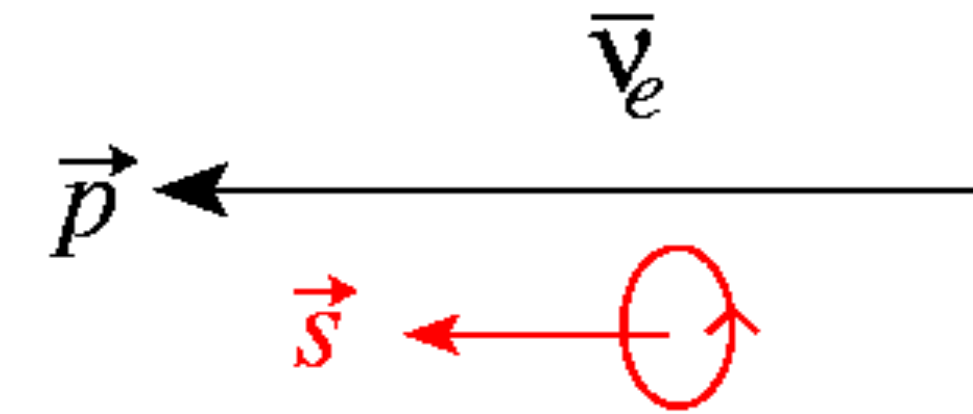
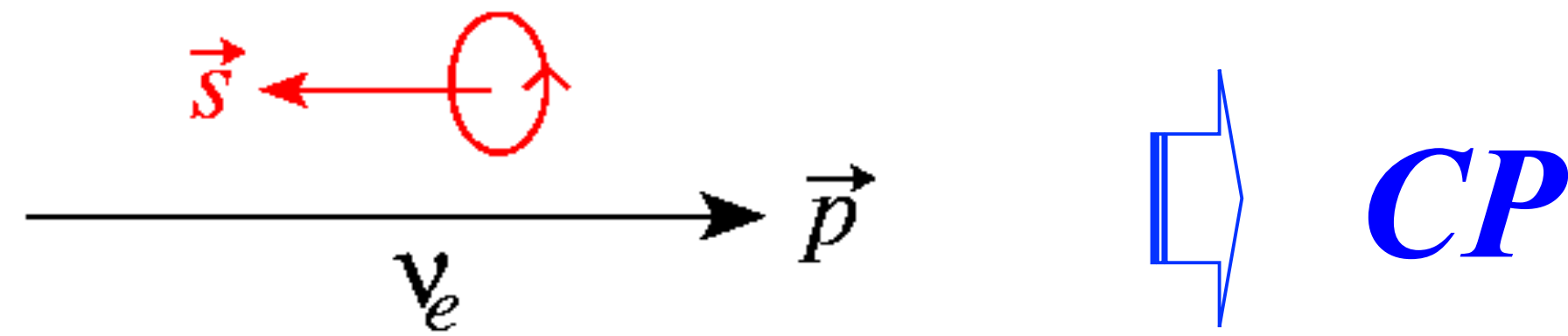


What is CP?



- Parity, P
 - $(\vec{r} \rightarrow -\vec{r}, \vec{p} \rightarrow -\vec{p}, \vec{L} \rightarrow \vec{L})$
- Charge Conjugation, C
 - $e^- \rightarrow e^+, K^+ \rightarrow K^-$
- CP is the product of the C and P operators.
 - CP transformation converts a particle into its anti-particle.

CP Symmetry, Particles - Antiparticles



CP

“CP-mirror,” a device that returns the mirror image of the anti-matter.

CP is violated IF particles and anti-particles behave differently!

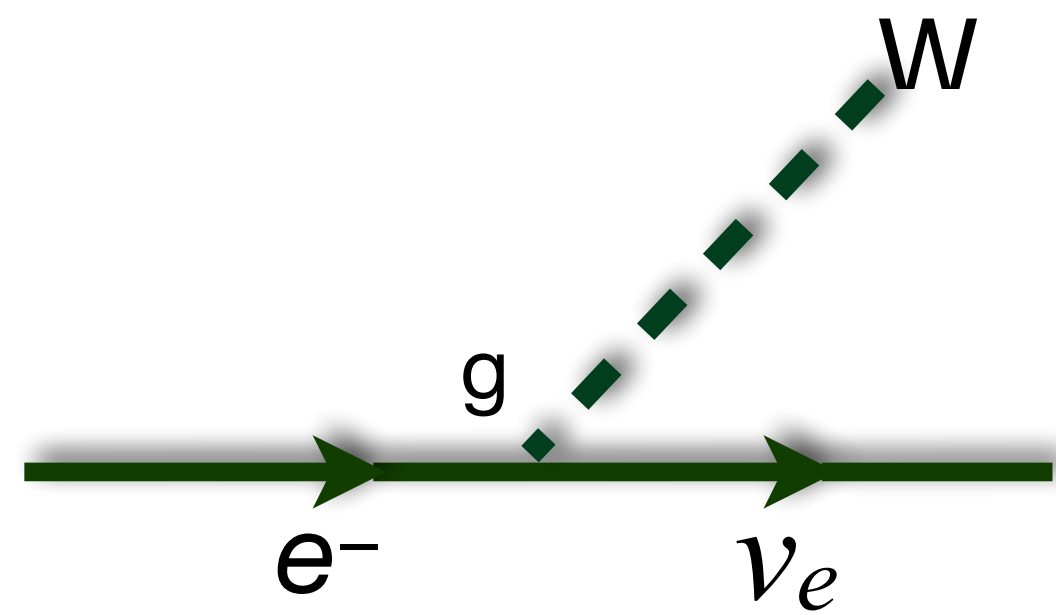


What is different about the weak interaction?

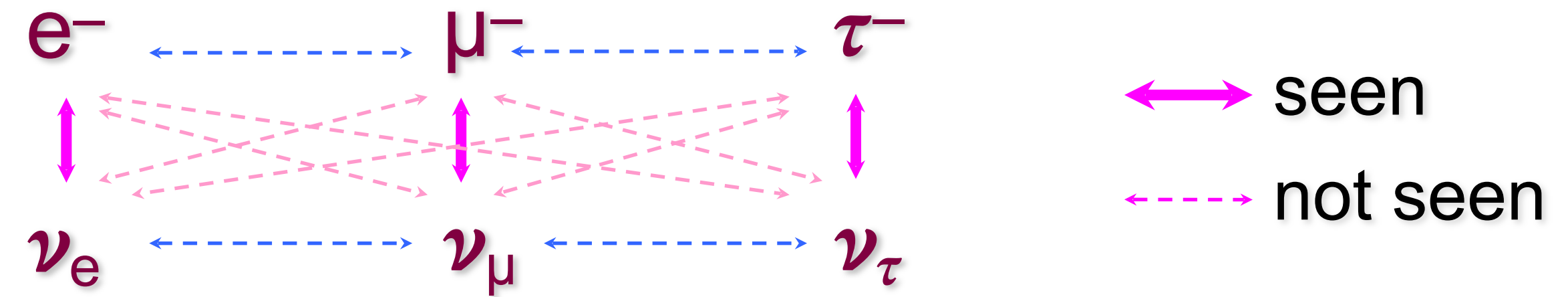
Only the weak interaction allows flavor change

CP is violated in weak interactions

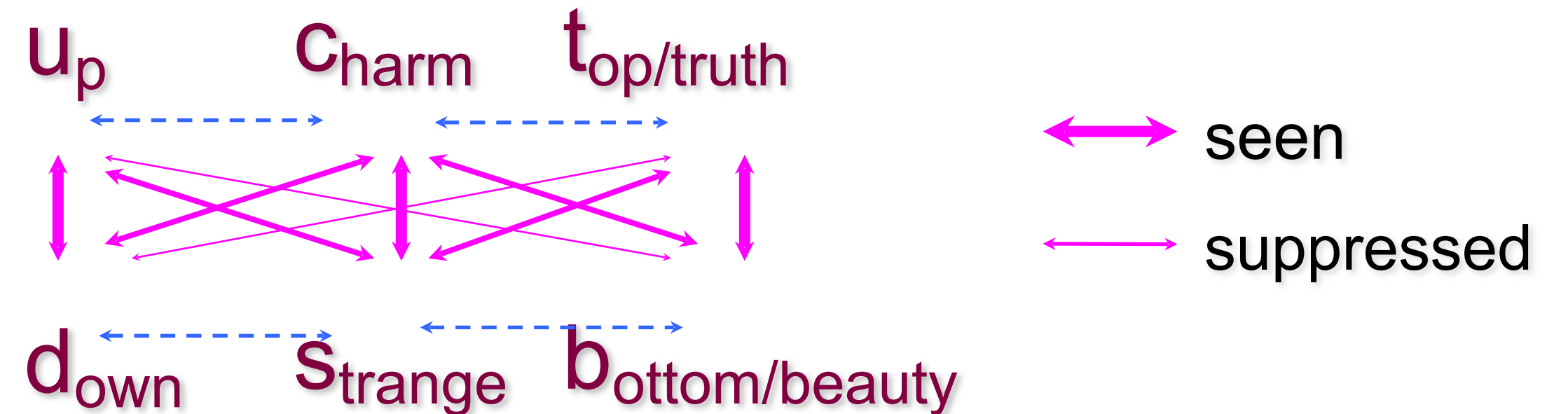
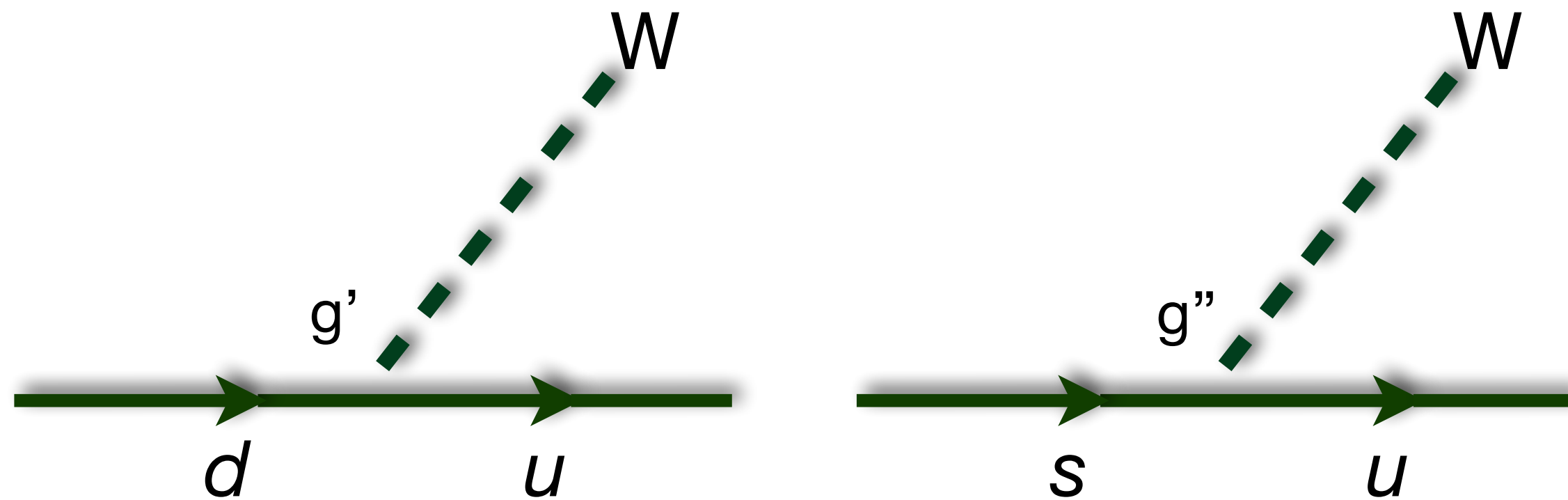
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Quarks: generation crossing

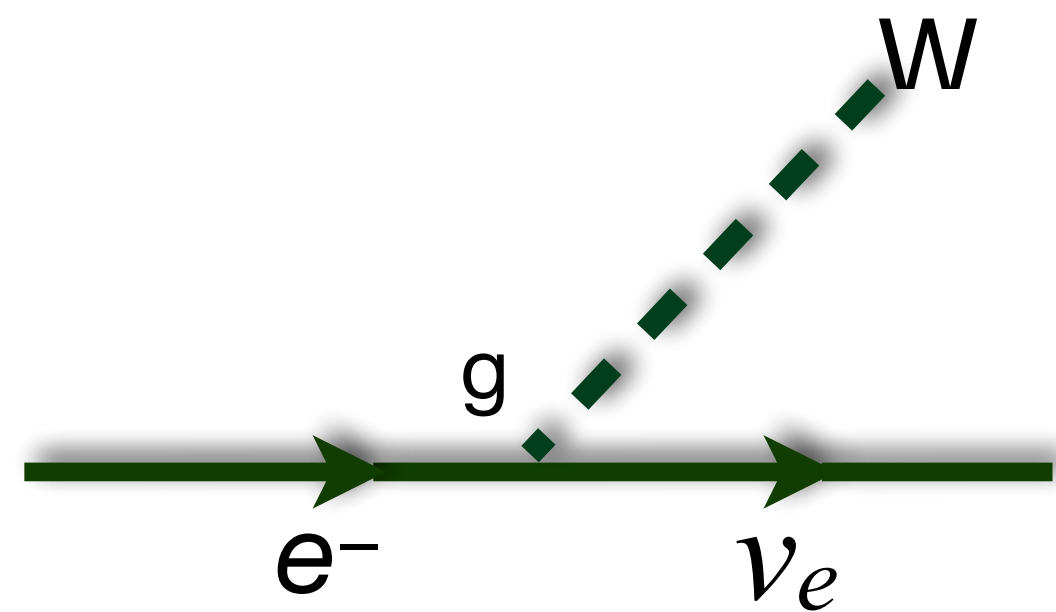


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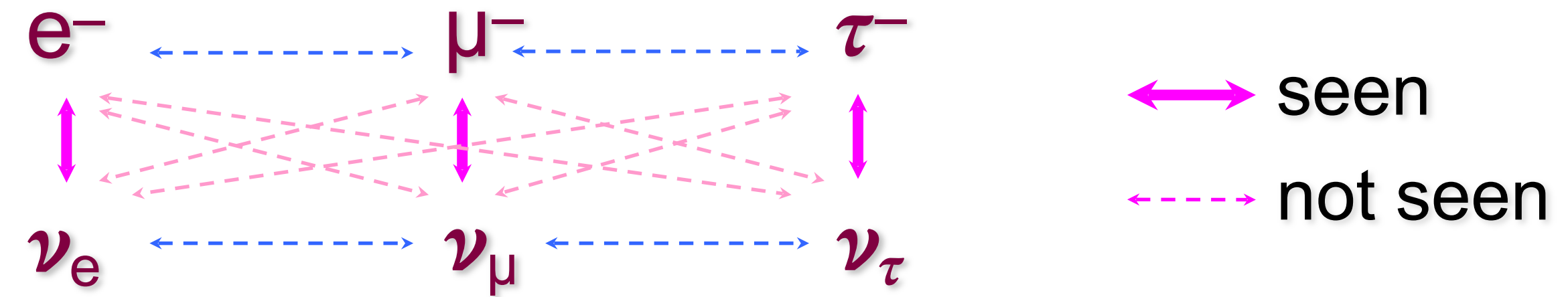
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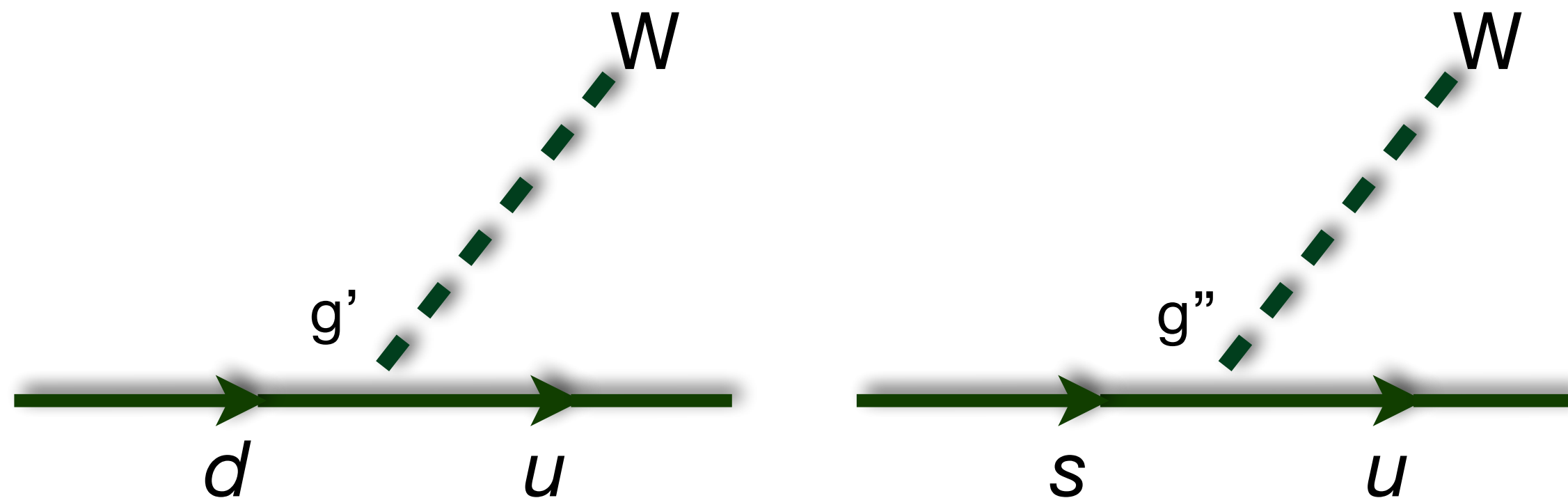
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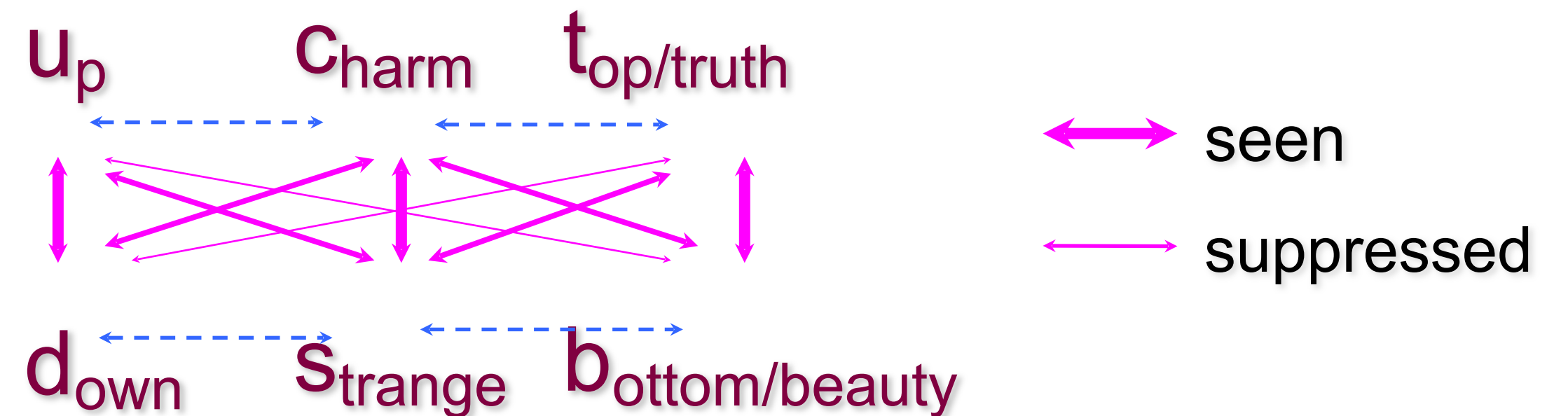
W^\pm "charged current"



Quarks: generation crossing



the pattern was not initially understood



Quark Mixing

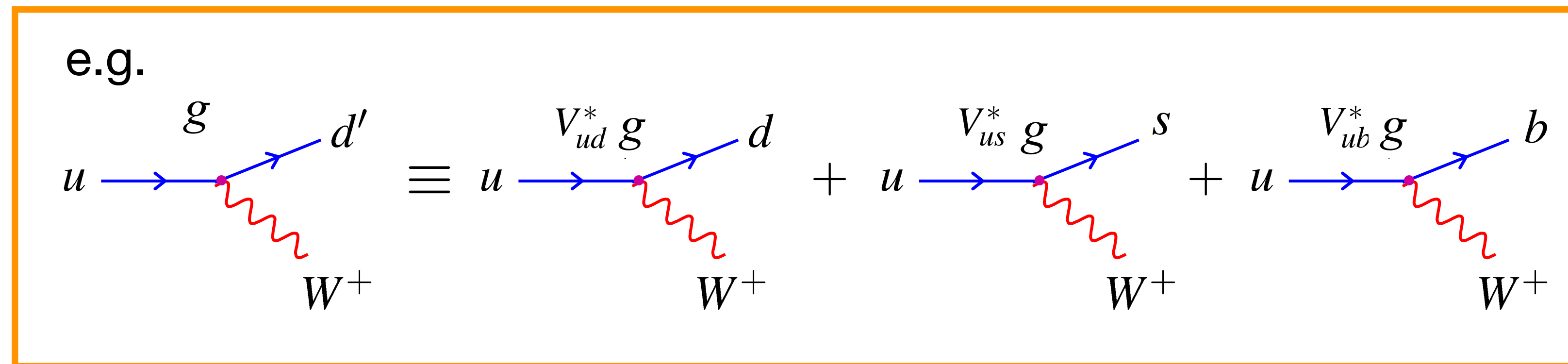
The solution was suggested by Cabibbo in 1963, perfected by Glashow, Iliopoulos, and Maiani (GIM) in 1970, and an explanation for CP violation is proposed by Kobayashi and Maskawa (KM) in 1973:

Instead of $\begin{pmatrix} u \\ d \end{pmatrix}, \begin{pmatrix} c \\ s \end{pmatrix}, \begin{pmatrix} t \\ b \end{pmatrix}$ the weak force couples the pairs $\begin{pmatrix} u \\ d' \end{pmatrix}, \begin{pmatrix} c \\ s' \end{pmatrix}, \begin{pmatrix} t \\ b' \end{pmatrix}$

where $d', s',$ and b' are linear combinations of the physical quarks $d, s,$ and b

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \underbrace{\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}}_{V_{CKM}} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

Cabibbo **K**obayashi **M**askawa
(CKM) Matrix



Universal weak coupling g must be multiplied by element of CKM matrix V_{ij} .

(Kobayashi-Maskawa **1973**)

Kobayashi and Maskawa proposed the explanation for CP violation in weak interactions prior to discovery of bottom and top quarks.



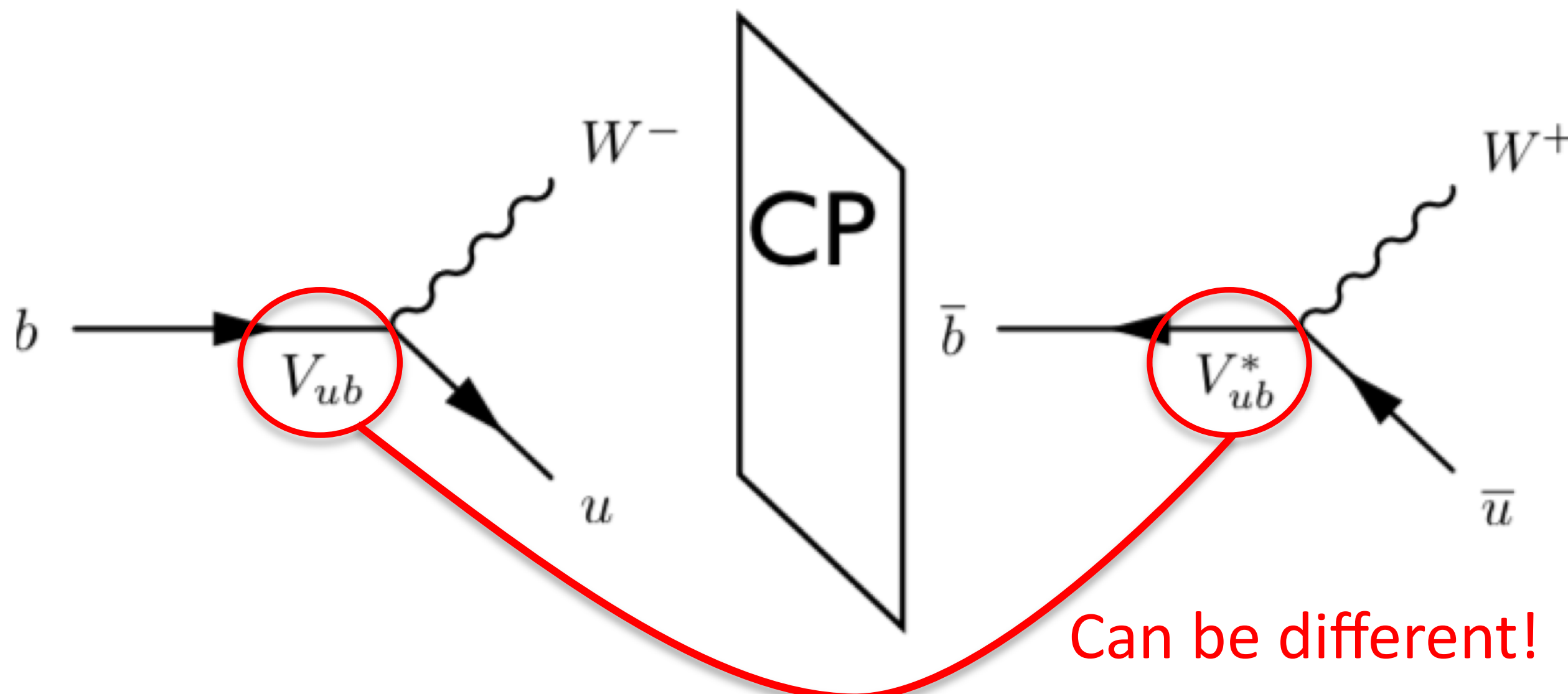
Explicit parameterization (Wolfenstein)

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} = \begin{bmatrix} 1 - \frac{\lambda^2}{2} & \lambda & \lambda^3 A(\rho - i\eta) \\ -\lambda & 1 - \frac{\lambda^2}{2} & \lambda^2 A \\ \lambda^3 A(1 - \rho - i\eta) & -\lambda^2 A & 1 \end{bmatrix} + \mathcal{O}(\lambda^4)$$

V_{CKM}

irreducibly complex

\Rightarrow **CP Violation!**



CP-violation measured in B-decays: 2002



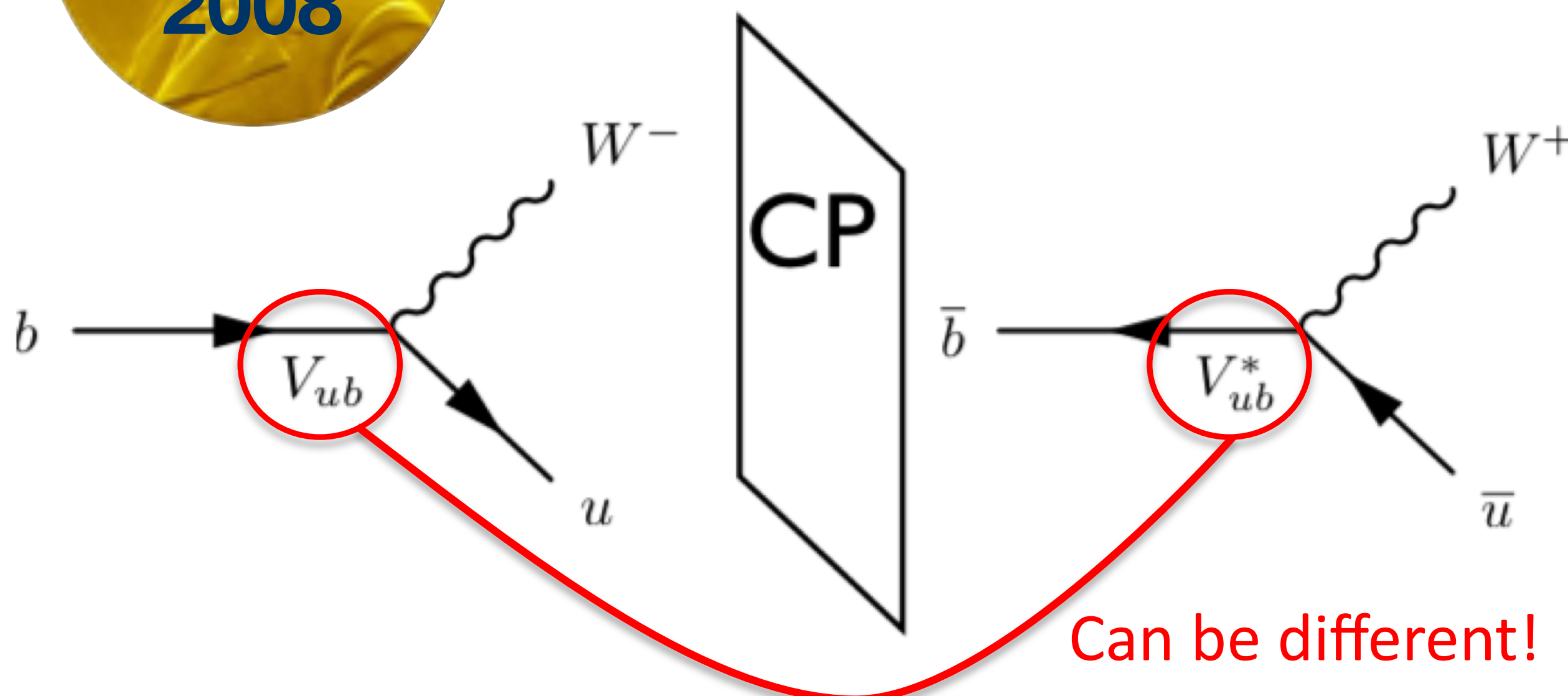
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The CKM Matrix and “Unitarity Triangle(s)”

- Unitarity condition: $(V_{CKM})^\dagger V_{CKM} = \mathbf{1} \Leftrightarrow \sum_j V_{ij} V_{jk}^* = \delta_{ik}$

- The 9 unitarity conditions of the 3×3 generations CKM matrix:

$$|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 1$$

$$|V_{cd}|^2 + |V_{cs}|^2 + |V_{cb}|^2 = 1$$

$$|V_{td}|^2 + |V_{ts}|^2 + |V_{tb}|^2 = 1$$

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

$$V_{ud} V_{us}^* + V_{cd} V_{cs}^* + V_{td} V_{ts}^* = 0$$

$$V_{ud} V_{ub}^* + V_{cd} V_{cb}^* + V_{td} V_{tb}^* = 0$$

$$V_{us} V_{ub}^* + V_{cs} V_{cb}^* + V_{ts} V_{tb}^* = 0$$

$$V_{ud} V_{cd}^* + V_{us} V_{cs}^* + V_{ub} V_{cb}^* = 0$$

$$V_{ud} V_{td}^* + V_{us} V_{ts}^* + V_{ub} V_{tb}^* = 0$$

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The CKM Matrix and “Unitarity Triangle(s)”

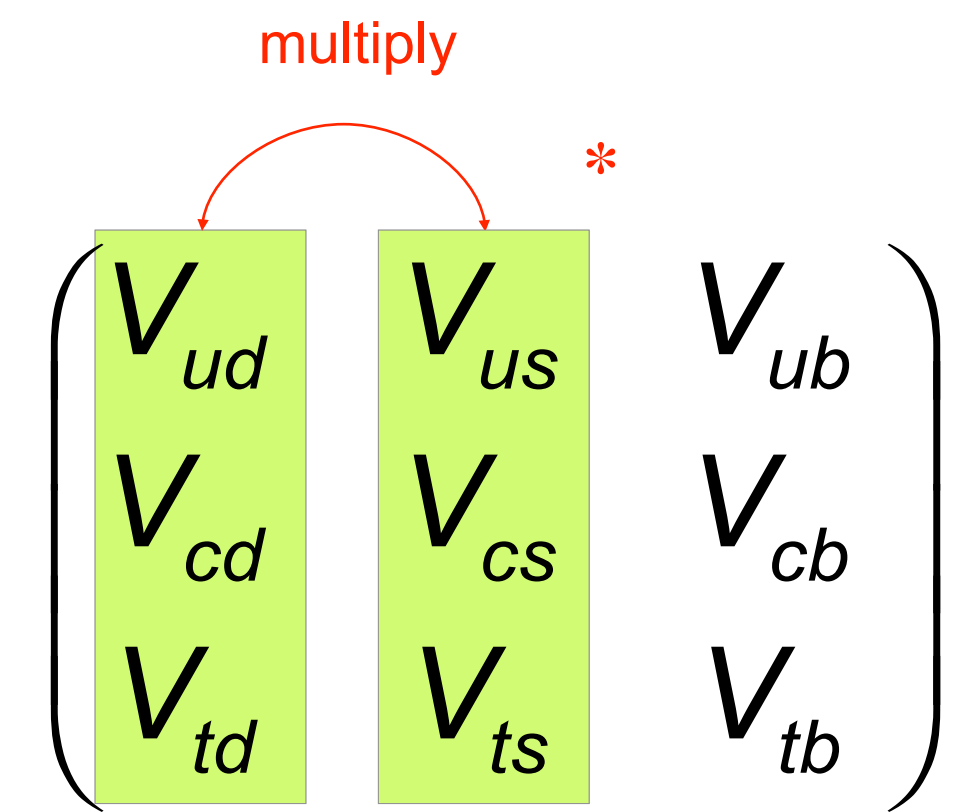
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- The 6 complex “**Unitarity Triangles**” involve different physics processes

$$V_{ud} V_{us}^* + V_{cd} V_{cs}^* + V_{td} V_{ts}^* = 0$$

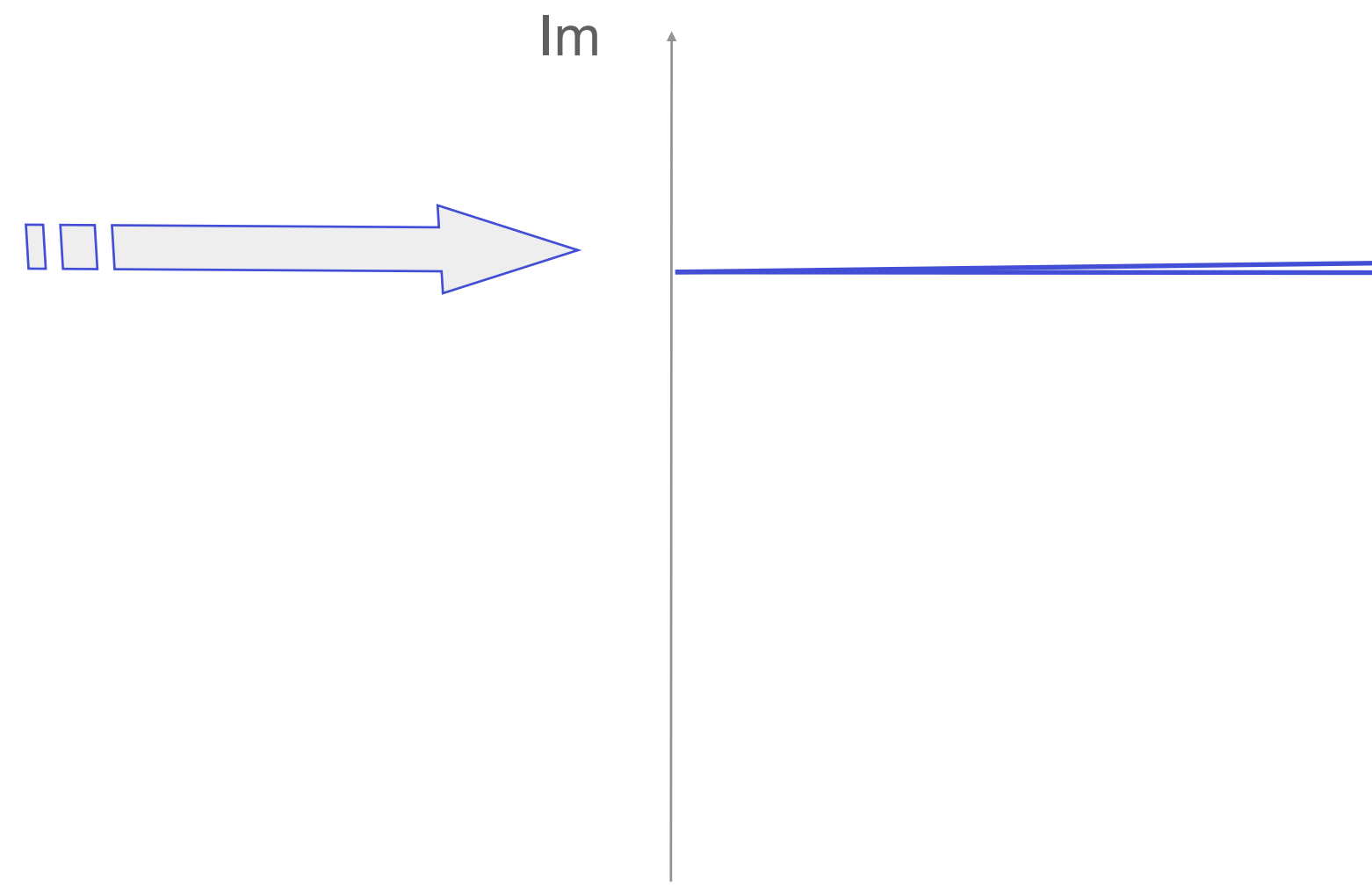
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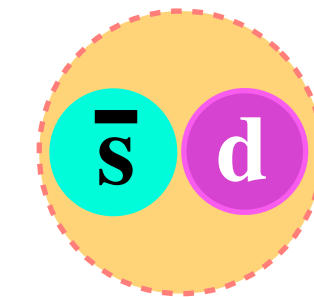
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K^0 system



The CKM Matrix and “Unitarity Triangle(s)”

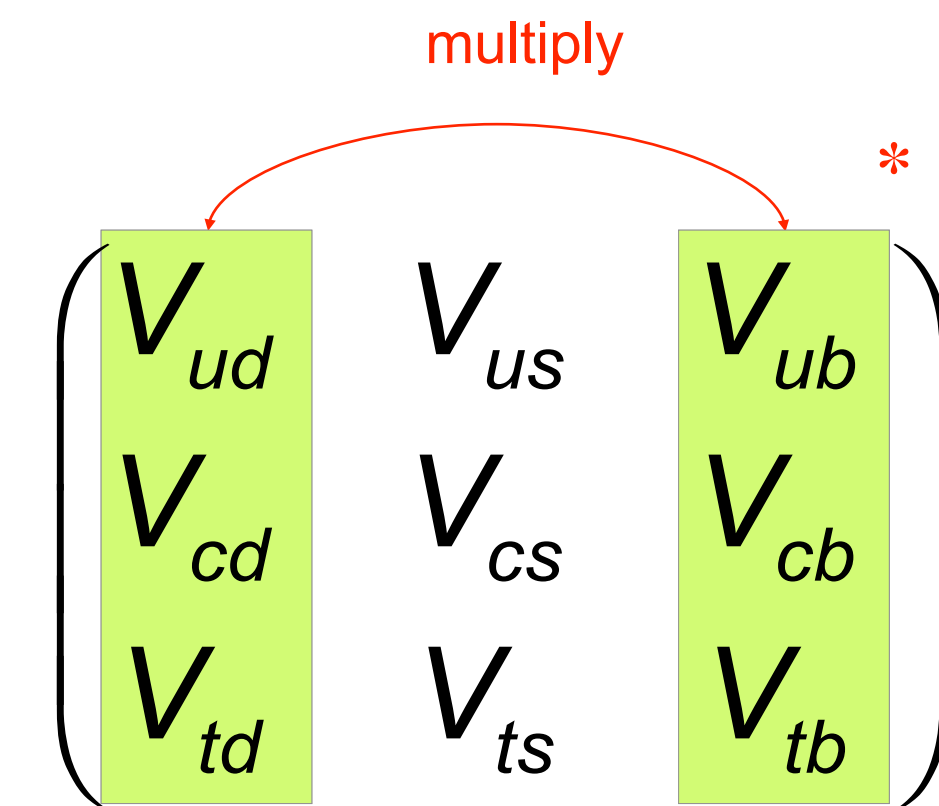
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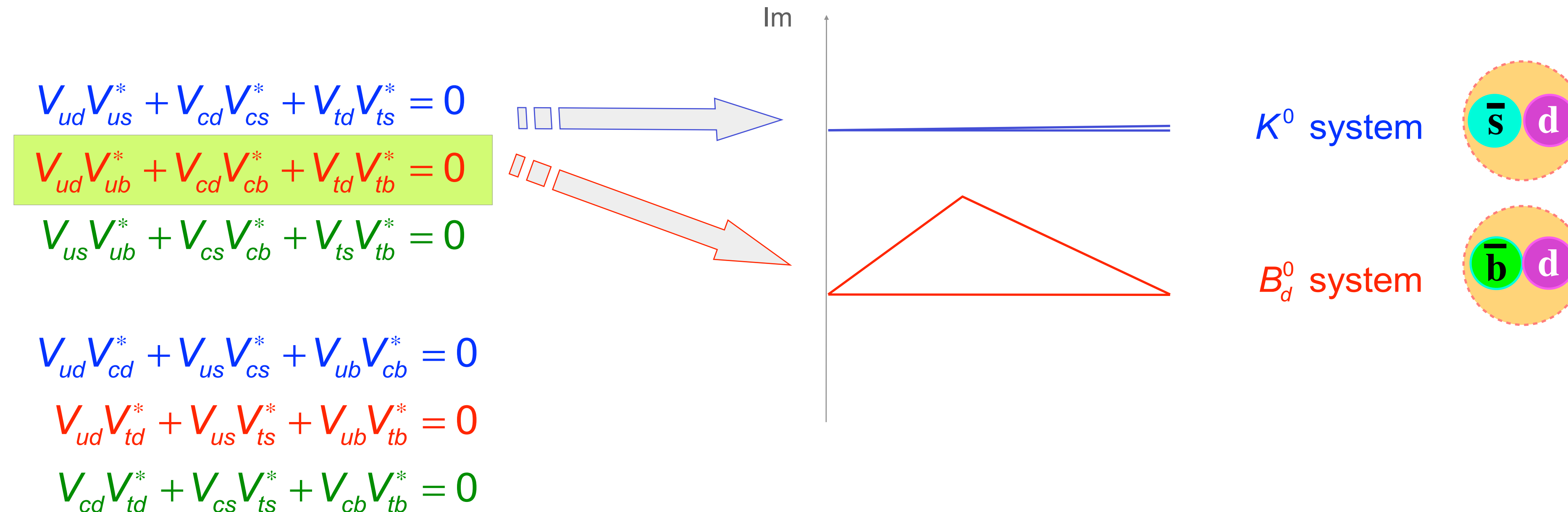
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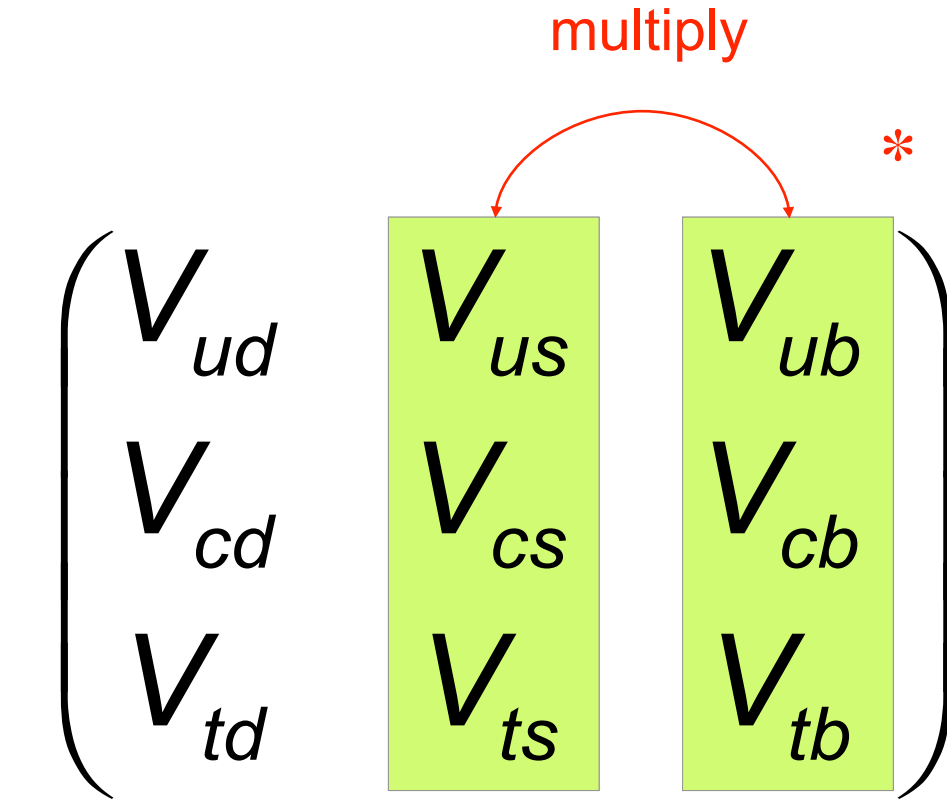
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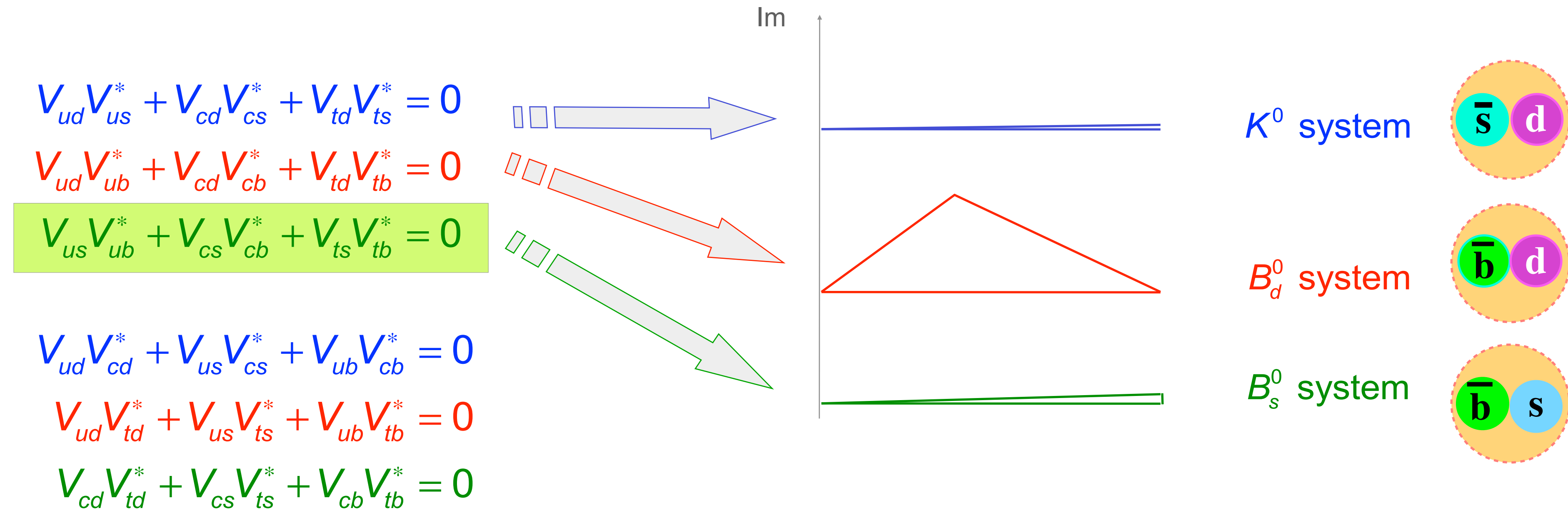
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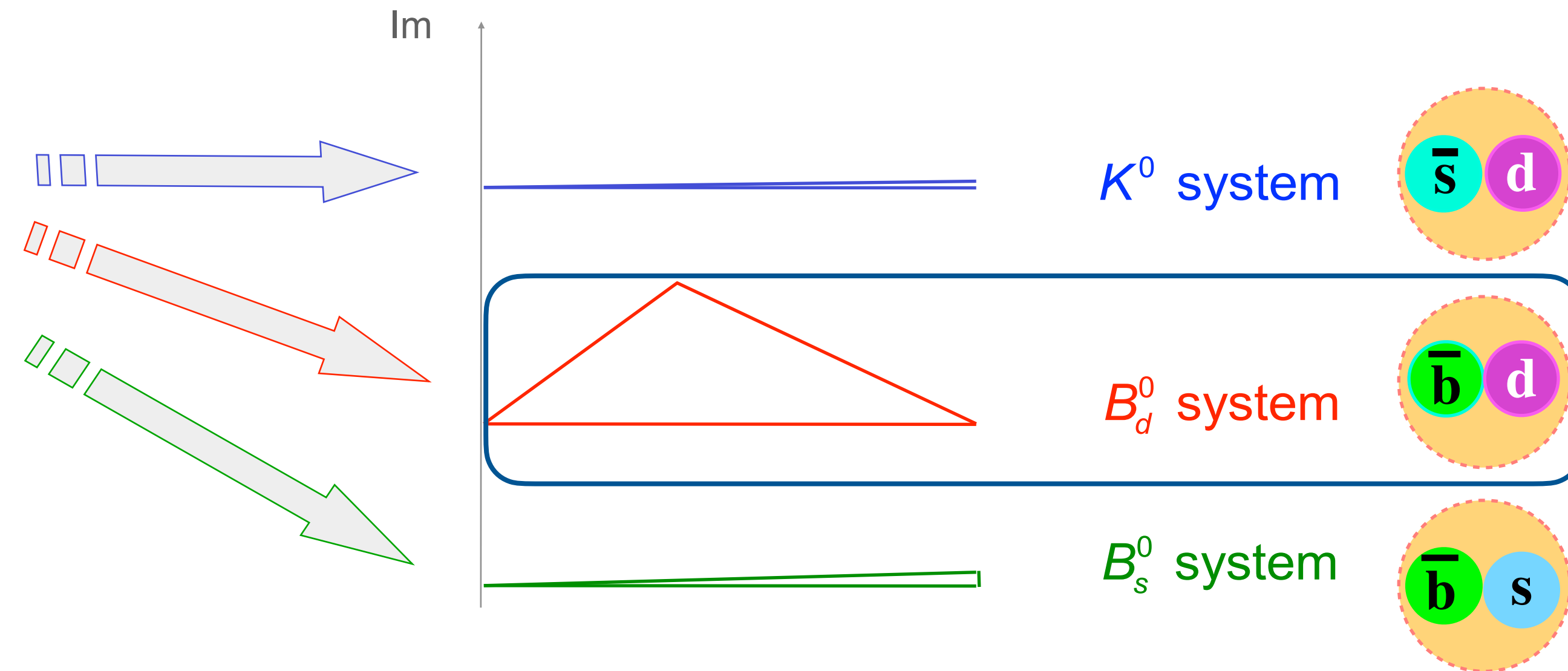
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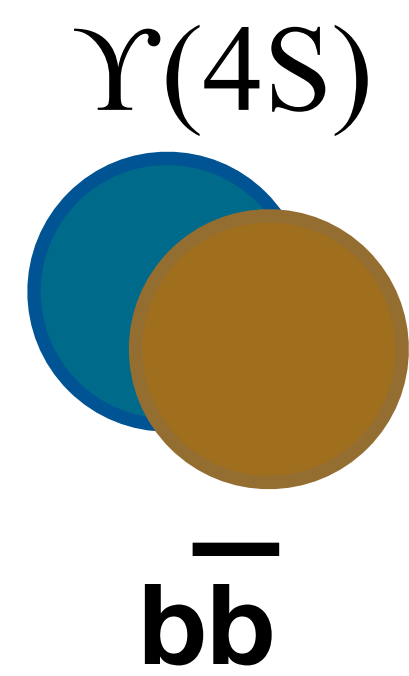
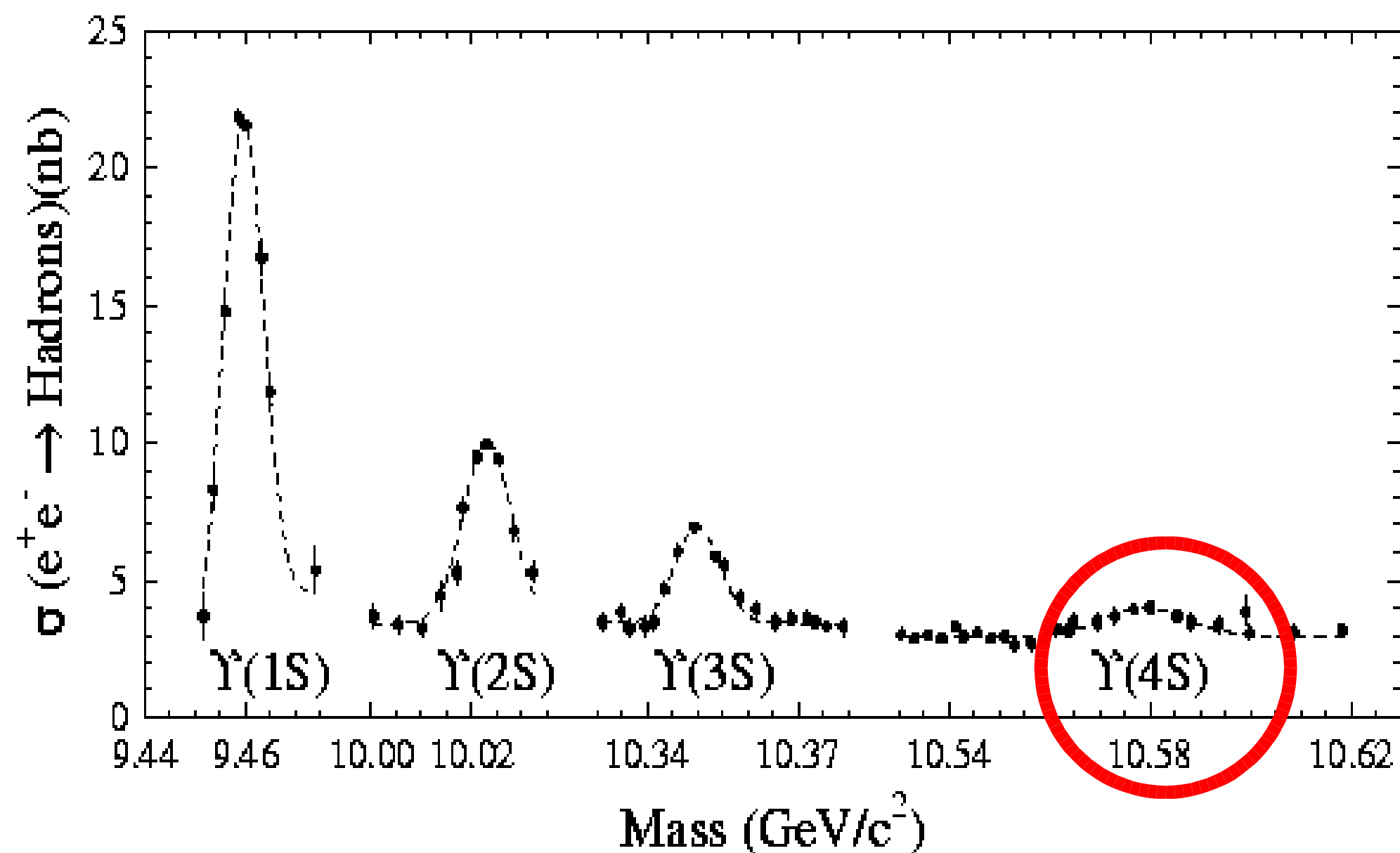


height measures the amount of CP violation

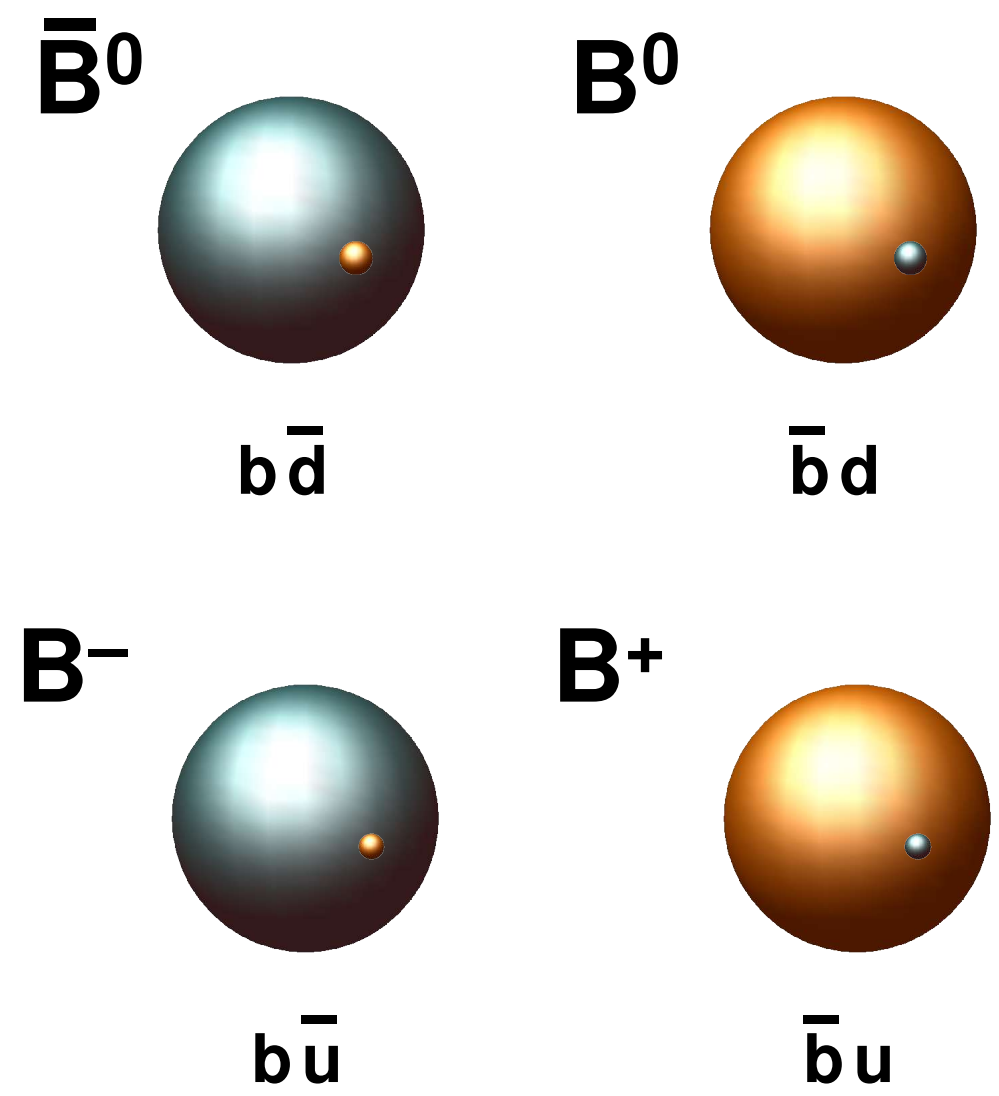
One Way to Produce B 's

Use e^+e^- annihilations at **Upsilon (4S)** [$\Upsilon(4S)$] particle which decays to B mesons

$$e^+e^- \rightarrow \Upsilon(4S) \rightarrow B\bar{B}$$



or



$$m_{\Upsilon(4S)} \sim 10.58 \text{ GeV}$$

$$m_{B^0} \sim m_{B^-} \sim 5.28 \text{ GeV}$$

The $\Upsilon(4S)$ - a clean source of B meson *pairs*

Equal amounts of matter and anti-matter

Use e^+e^- annihilations at Upsilon (4S) [$\Upsilon(4S)$] particle which decays to B mesons

B-Factory experiments [Belle at KEK(Japan), BaBar at SLAC (Stanford)] (1999-2009)

- CP asymmetry observed in diverse processes in B decay
- > many measurements, (over)constrain CKM, confirm unitarity

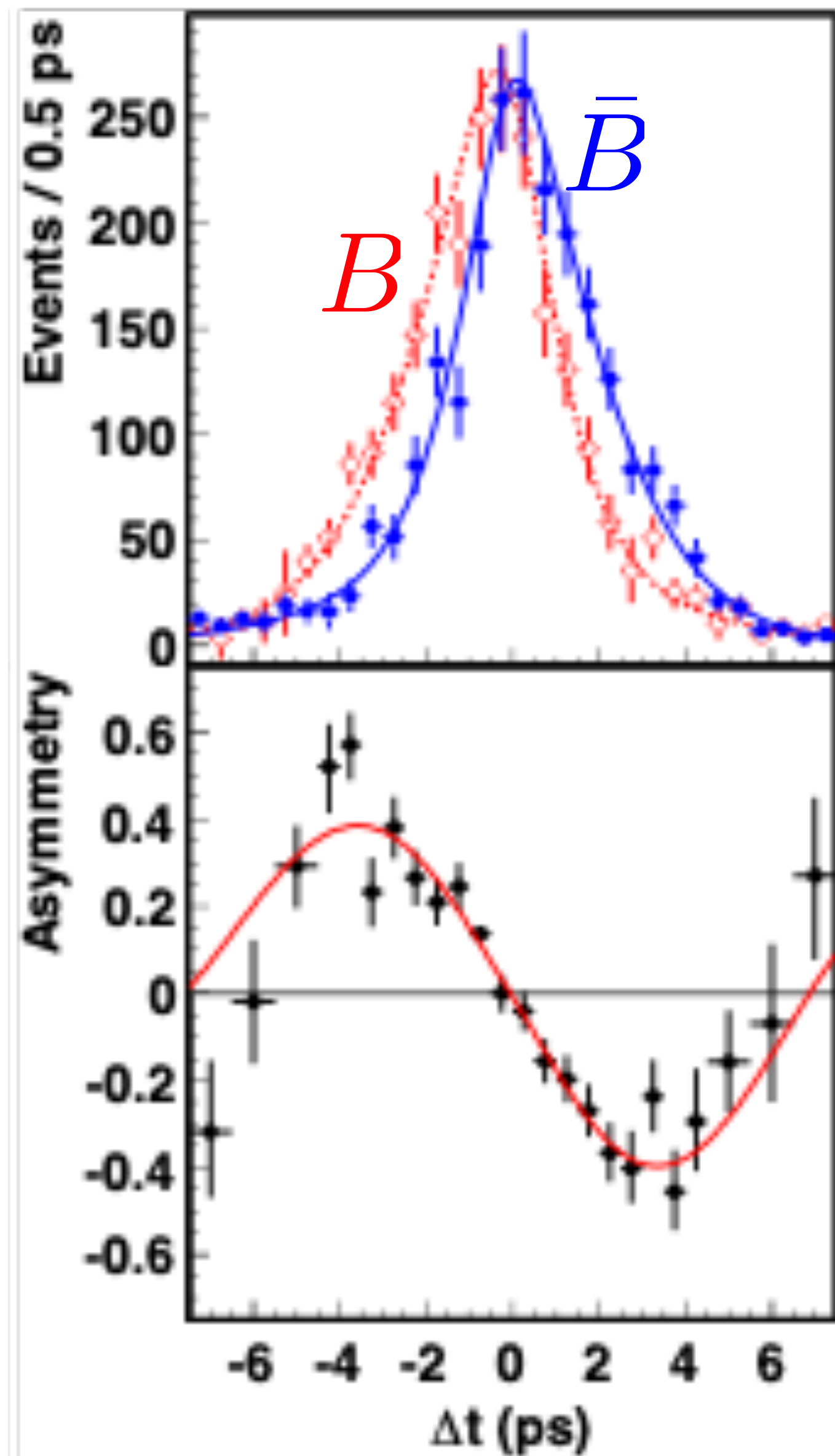
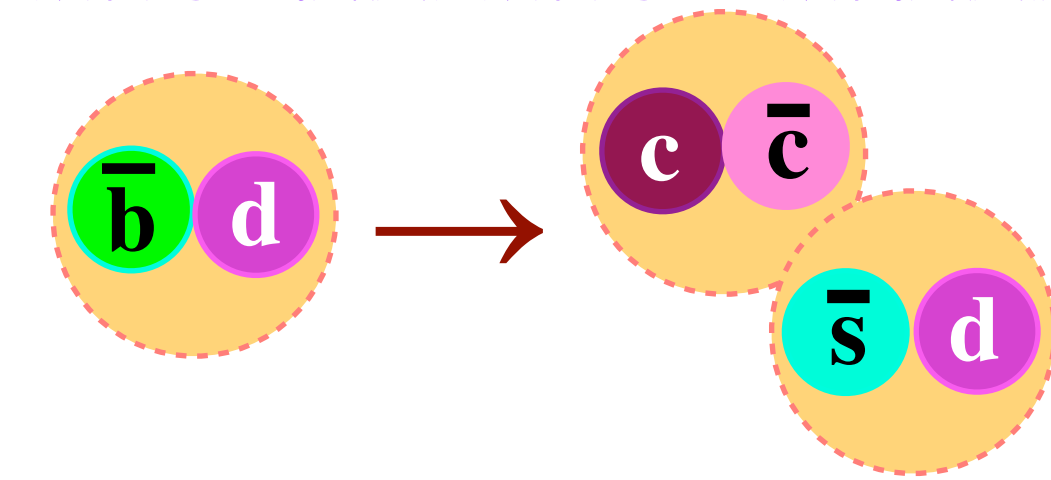
Next generation B factory, BELLE II detector at Super KEKB will be covered by Dr. Seema Choudhury

The $\Upsilon(4S)$ - a clean source of B meson *pairs*

Equal amounts of matter and anti-matter

Measurement of CP Violation in $B^0 \rightarrow J/\psi K_s^0$

PRL 108, 171802 (2012)



• Top plot shows the decay time distribution for B^0 and anti- B^0 . Matter and anti-matter are essentially different.

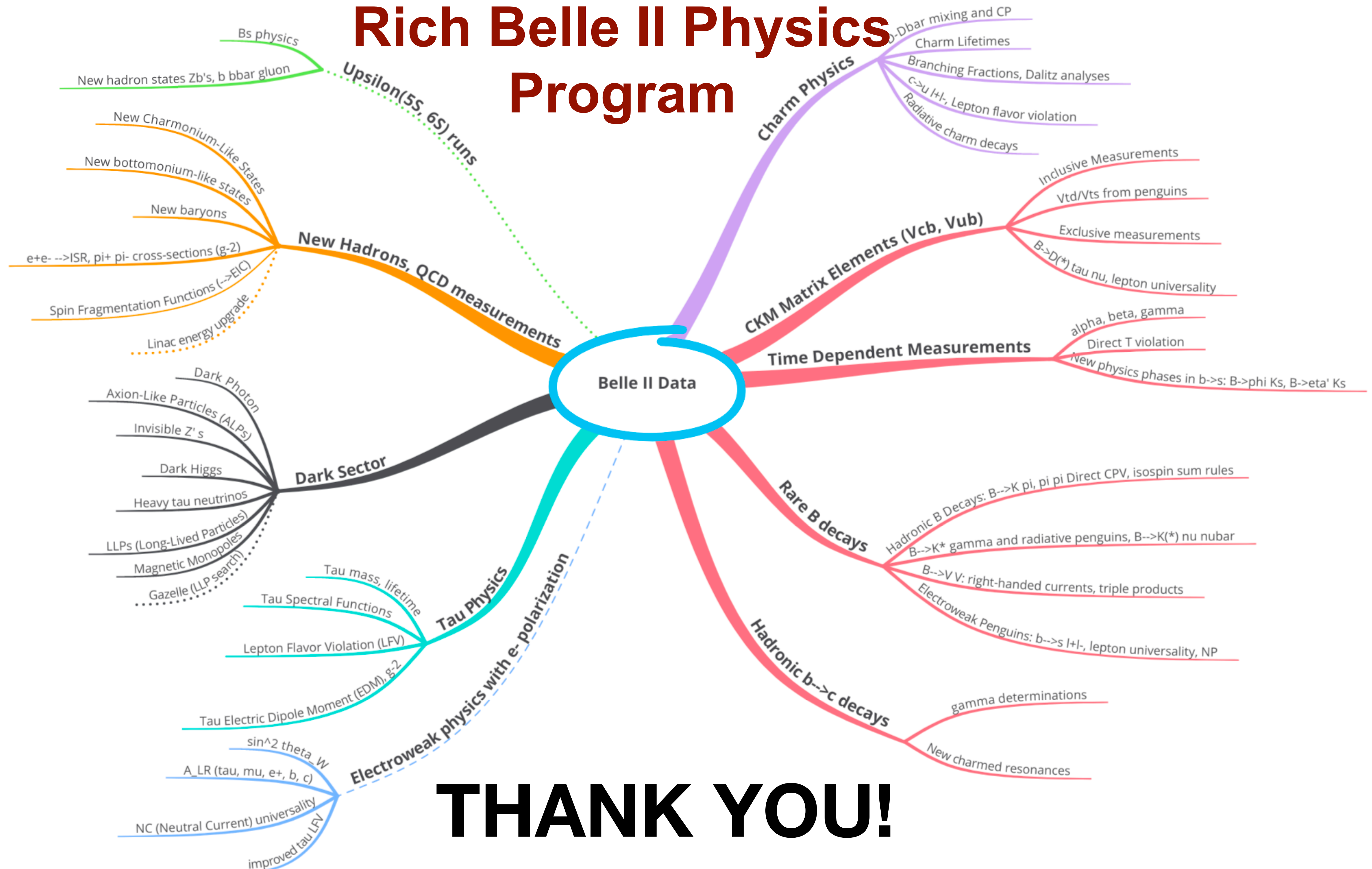
• If CP were conserved in nature, the bottom distribution would be completely flat.

Belle collected 7.7×10^8 B events for this plot

- However, the Standard Model sources of CP violation cannot explain the matter dominance of the universe.
- We have to search for new sources of CP violation
 ⇒ New physics beyond the Standard Model.



Rich Belle II Physics Program

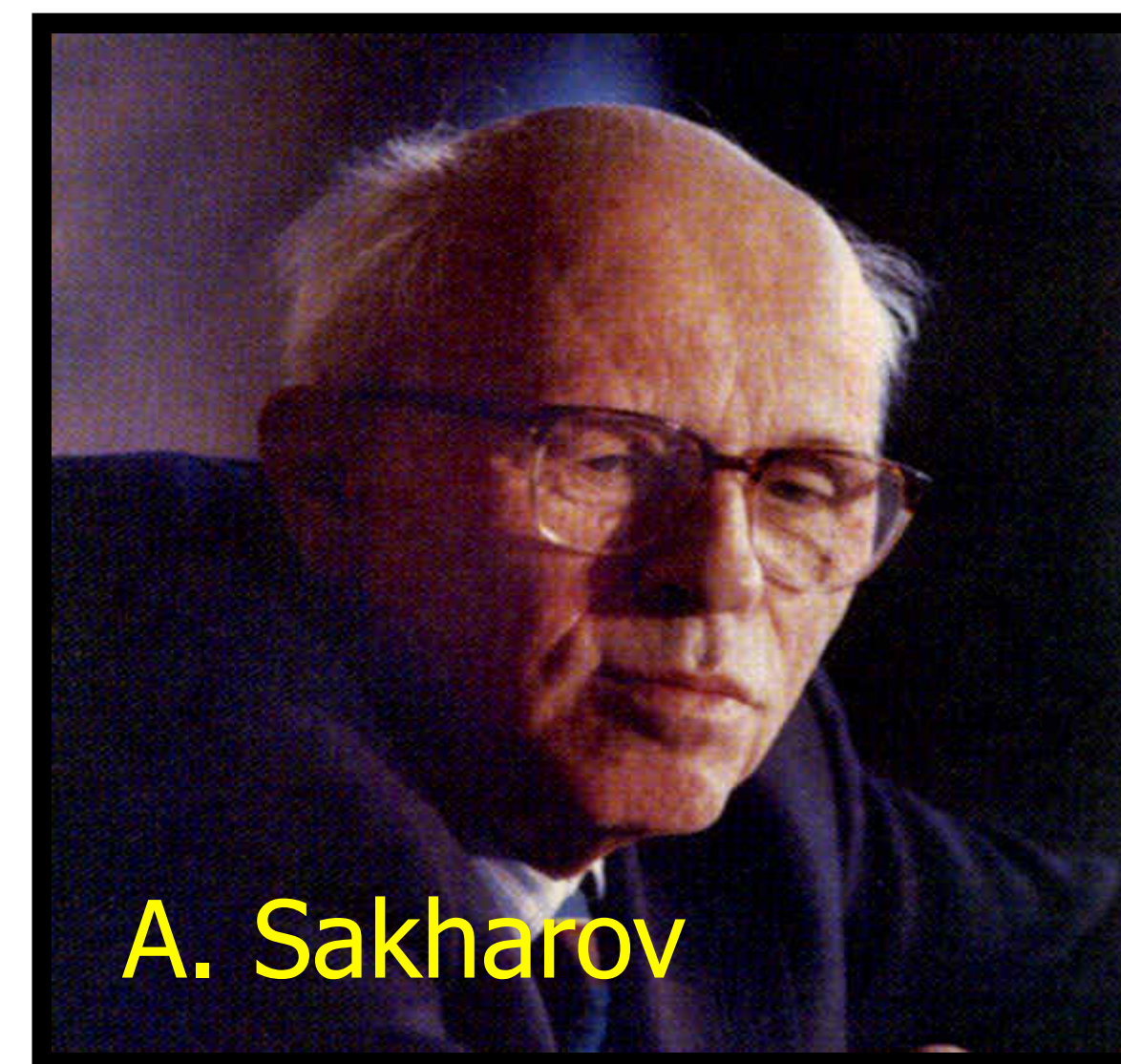
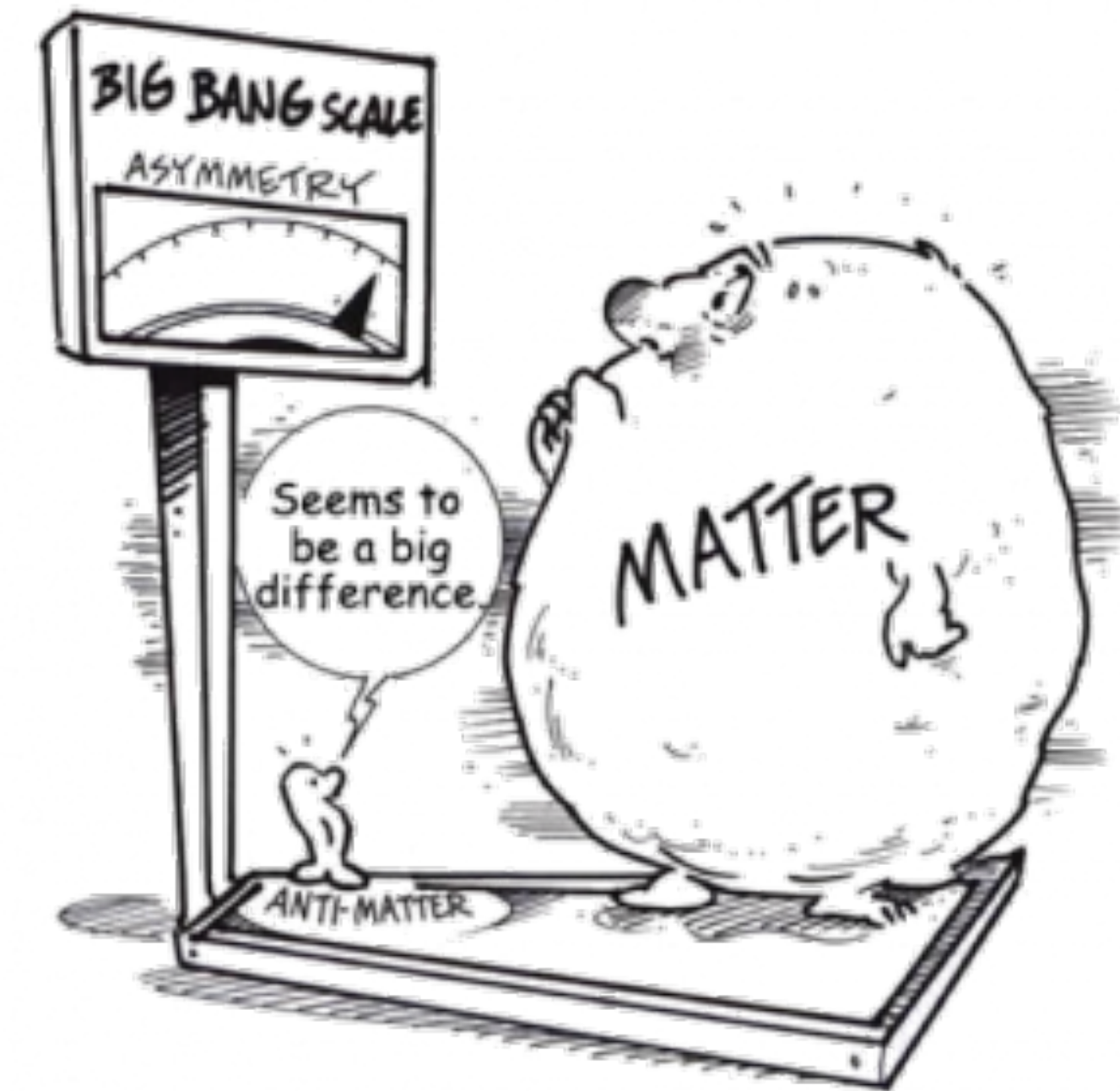


THANK YOU!

Extra

To generate this initial asymmetry three conditions must be met (Sakharov, 1967):

- Baryon number violation
- *Charge and Charge Parity violation*
- Departure from thermal equilibrium; in thermal equilibrium any baryon number violating process will be balanced by the inverse reaction



A. Sakharov

Symmetries in Physics

There are two main types of symmetries in physics: continuous and discrete.

- For every continuous symmetry there exists a corresponding conservation law.

Noether's Theorem

Laws of physics invariant under

Conservation of

Spatial Translations



Momentum

Time Translations



Energy

Rotations



Angular Momentum

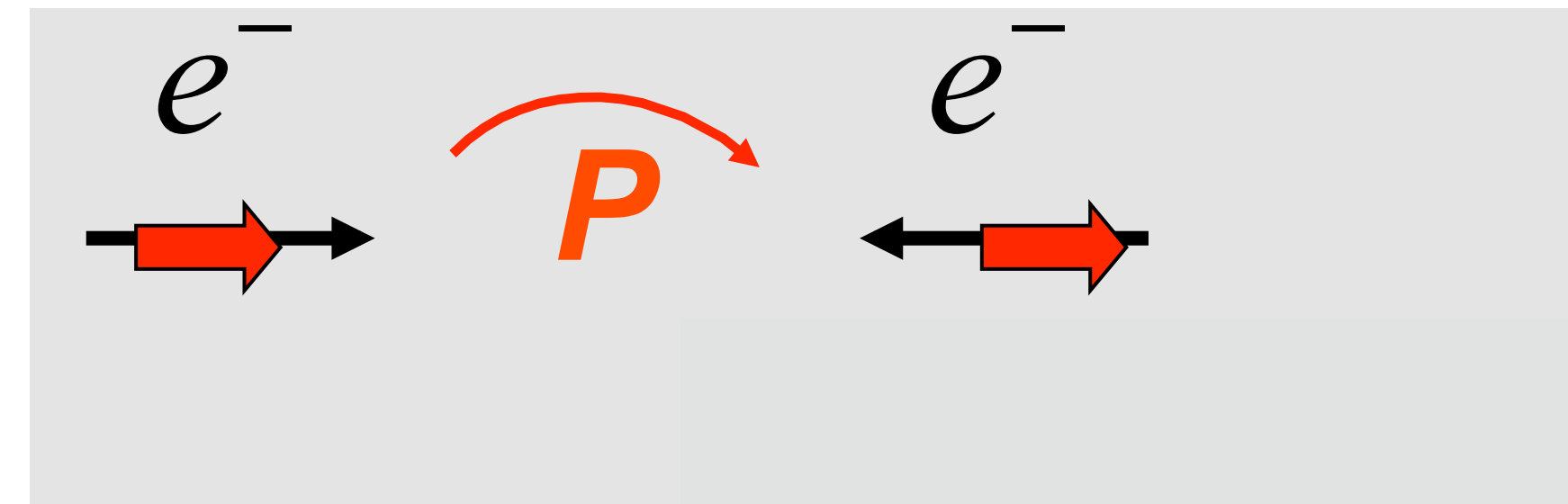


Emmy Noether
(1882-1935)

- Discrete symmetries are very important in particle physics!

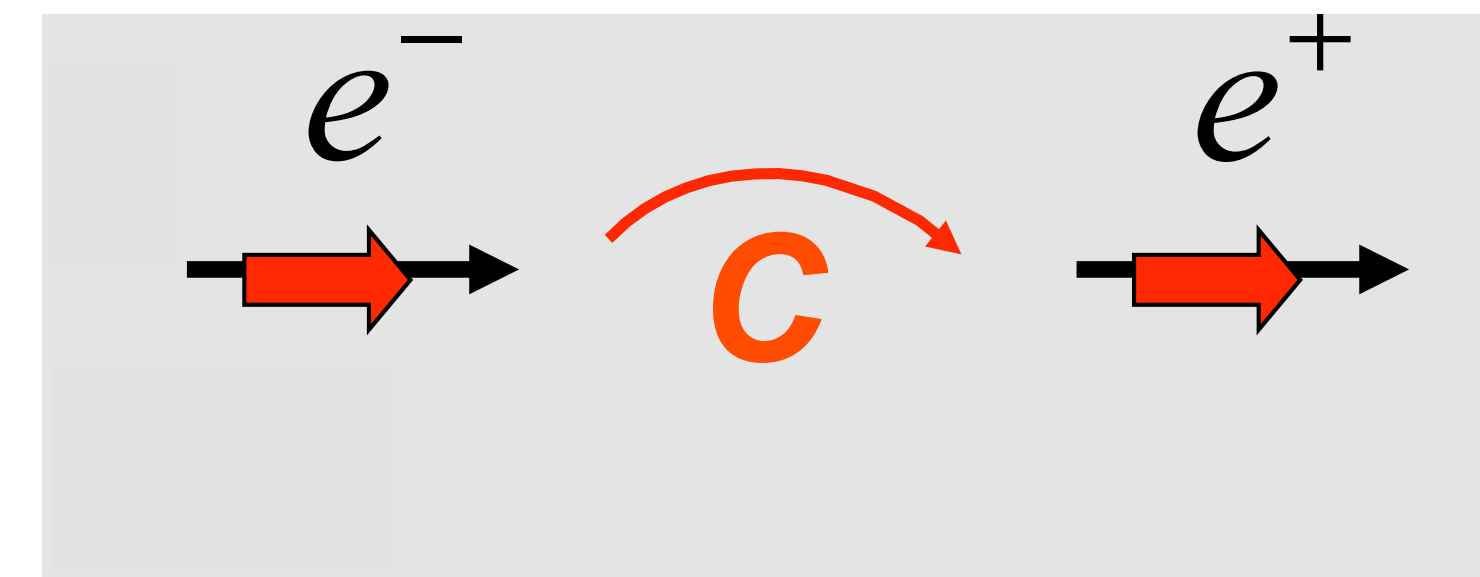
- Parity, **P**
- Parity reflects a system through the origin.

- $(\vec{r} \rightarrow -\vec{r}, \vec{p} \rightarrow -\vec{p}, \vec{L} \rightarrow \vec{L})$



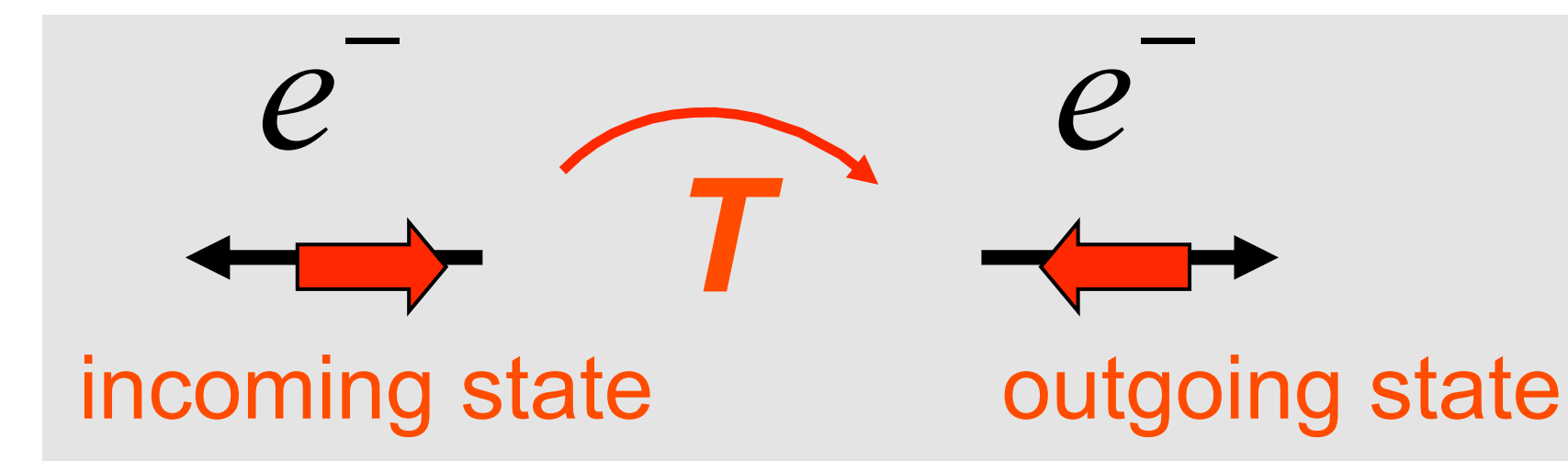
- Charge Conjugation, **C**
- **C** changes the sign of all the “internal” quantum numbers.

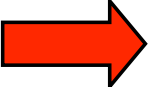
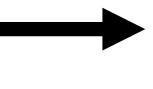
- $e^- \rightarrow e^+, K^+ \rightarrow K^-, p \rightarrow \bar{p}, \gamma \rightarrow \gamma$



- Time reversal, **T**
- **T** reverse any motion in system.

- $t \rightarrow -t, \vec{p} \rightarrow -\vec{p}, \vec{L} \rightarrow -\vec{L}$



 = spin direction
 = momentum direction

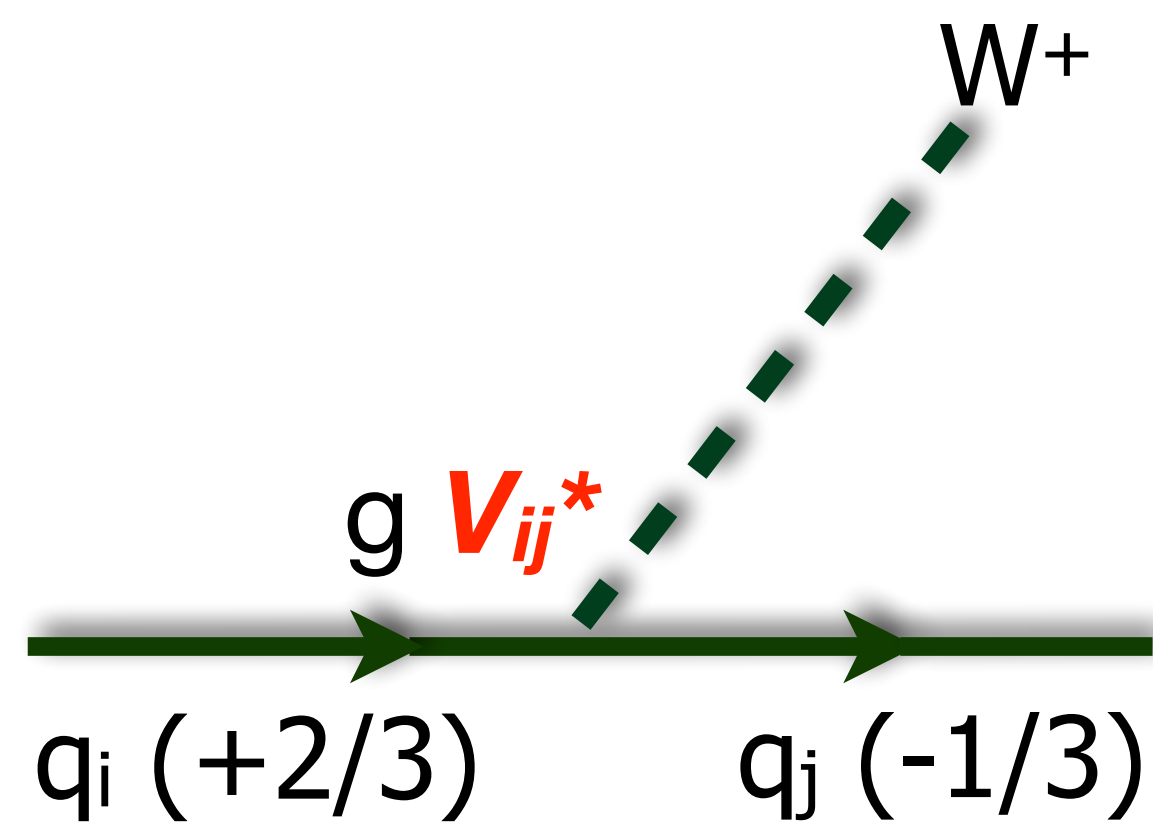
Quark Mixing

(Kobayashi-Maskawa 1973)



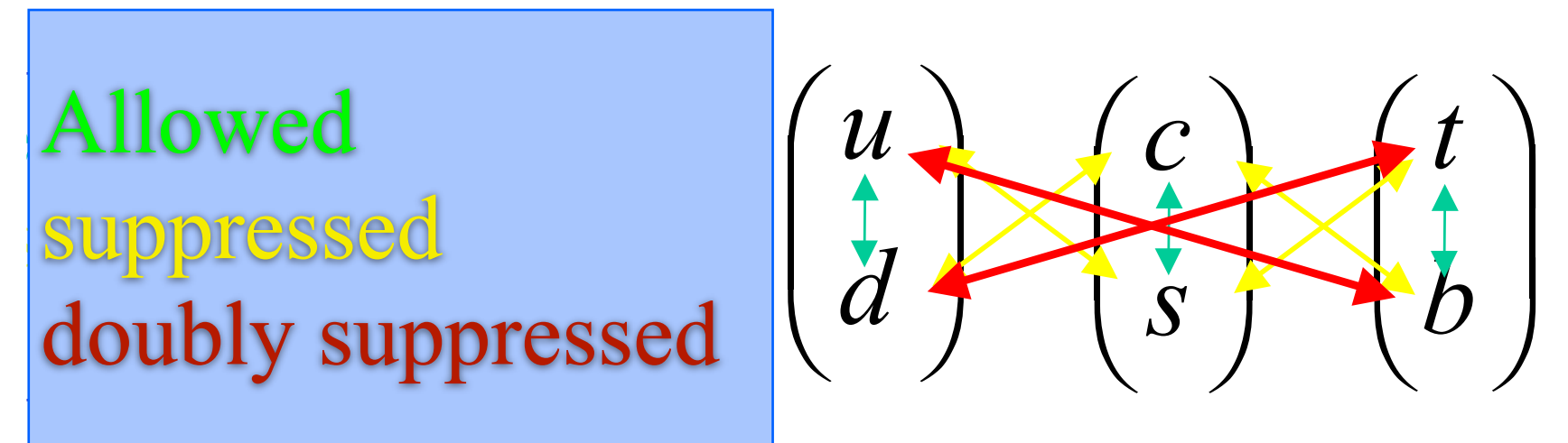
Kobayashi and Maskawa proposed a daring explanation for CP violation in weak interactions.

- Need at least **3 Generation of Quarks** (then not known).
- Get a CP-violating phase $e^{i\delta}$, if one assumes three families.



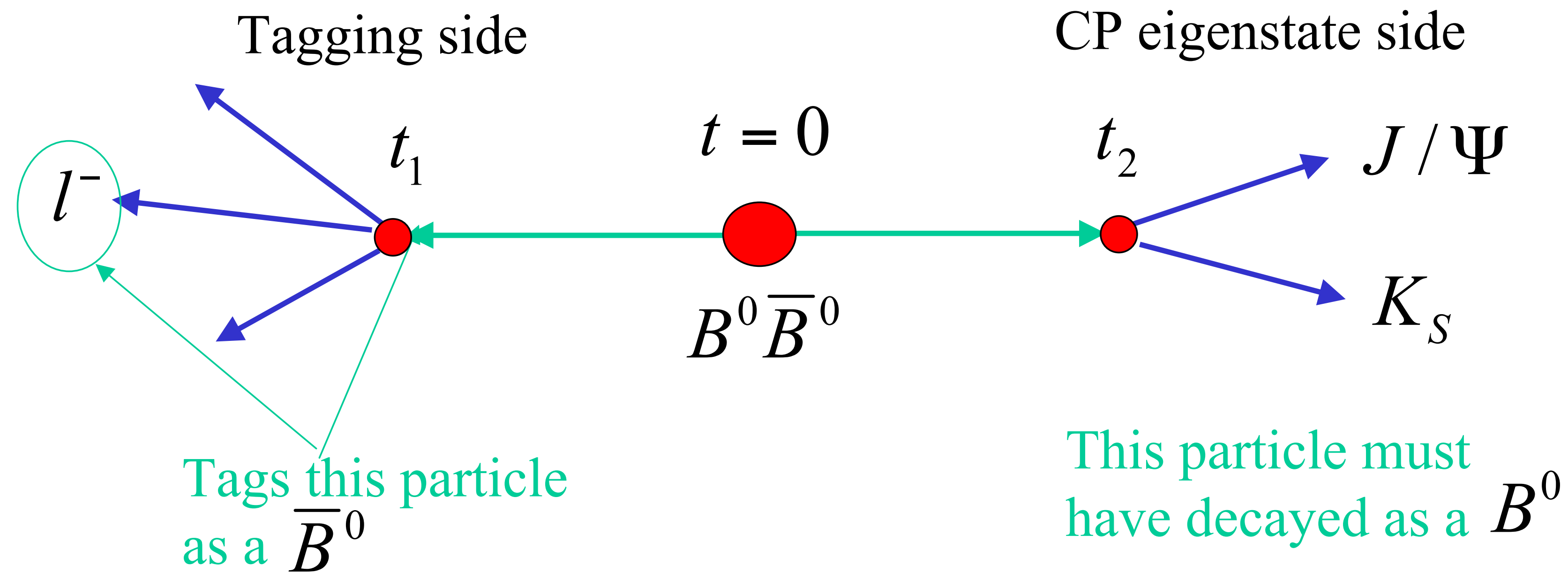
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Cabibbo Kobayashi Maskawa
(CKM)
matrix



$$|V_{ij}| = \begin{pmatrix} \text{large} & \text{small} & \text{tiny} \\ \text{small} & \text{large} & \text{small} \\ \text{tiny} & \text{small} & \text{large} \end{pmatrix}$$

relative magnitudes



If $t_1 = t_2$, then the particle on the CP eigenstate must be a B^0 . Note that the tagging information is communicated across space instantaneously despite the fact that the B 's could be separated by a finite distance (a few hundred microns). This is an example of quantum entanglement (the EPR paradox).

SuperKEKB - Belle II

$e^- : 7.0 \text{ GeV}, e^+ : 4.0 \text{ GeV}$

