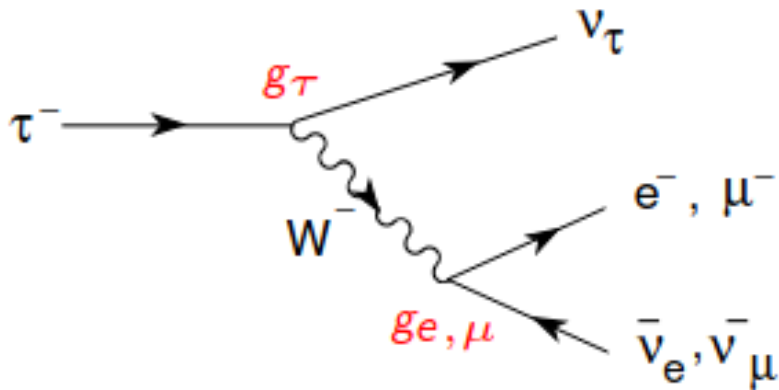


# Opportunities with Tau Leptonic decays at Belle-II



Cinvestav

Pablo Roig (Cinvestav, México)



'Fundamental'  
Theory

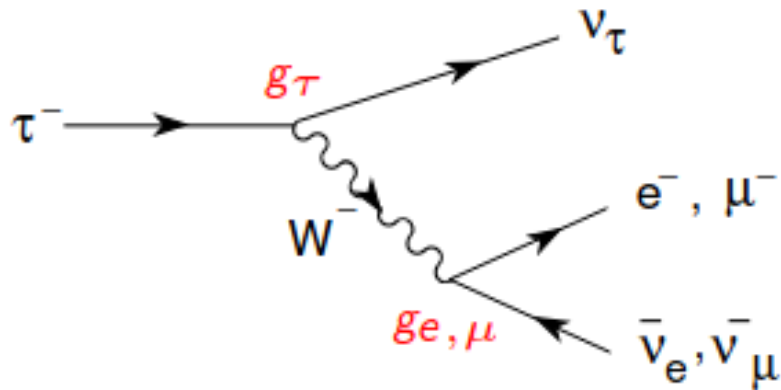
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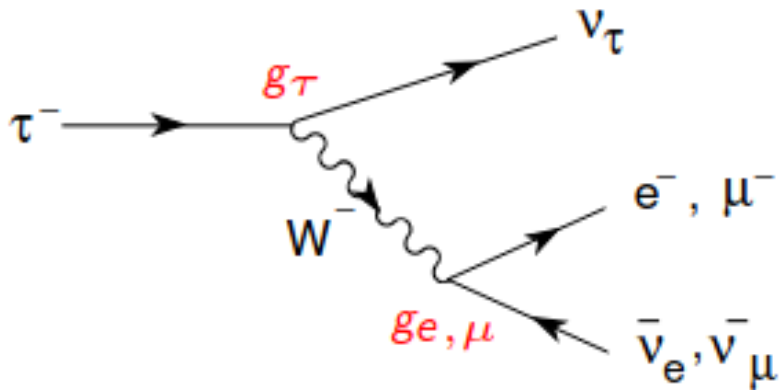
2024 Belle-II Physics Week KEK, Oct. 14-18

# Opportunities with Tau Leptonic decays at Belle-II



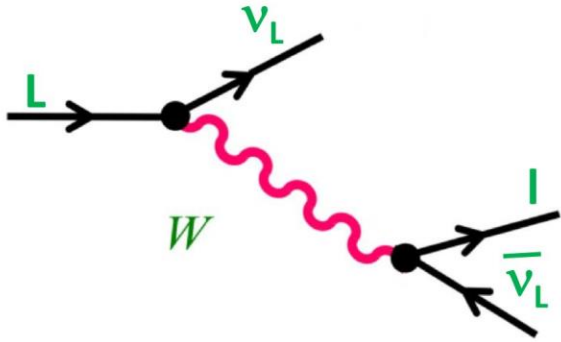
- Lepton Universality
- Lorentz structure of the charged current

# Opportunities with Tau Leptonic decays at Belle-II



- Lepton Universality
- Lorentz structure of the charged current (including possible heavy sterile neutrinos)

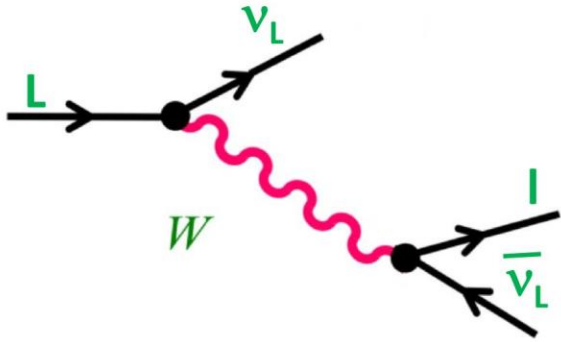
# Lepton Universality



$$T(l \rightarrow \nu_l l' \bar{\nu}_{l'}) \sim \frac{g^2}{M_W^2 - q^2} \xrightarrow{q^2 \ll M_W^2} \frac{g^2}{M_W^2} = 4\sqrt{2} G_F$$

$$\Gamma \sim G_F^2 m_L^5$$

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The observable:

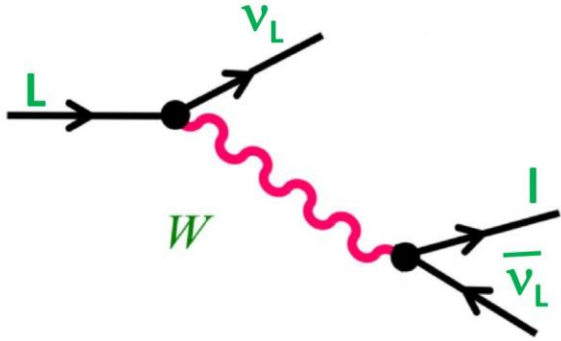
$$R_\mu = \frac{\mathcal{B}(\tau^- \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau)}{\mathcal{B}(\tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau)}$$

We extract

$$\left| \frac{g_\mu}{g_e} \right|_\tau = \sqrt{R_\mu \frac{f(m_e^2/m_\tau^2)}{f(m_\mu^2/m_\tau^2)}} \quad f(x) = 1 - 8x + 8x^3 - x^4 - 12x^2 \ln x$$

It must be 1 in the SM, as a consequence of gauge symmetry

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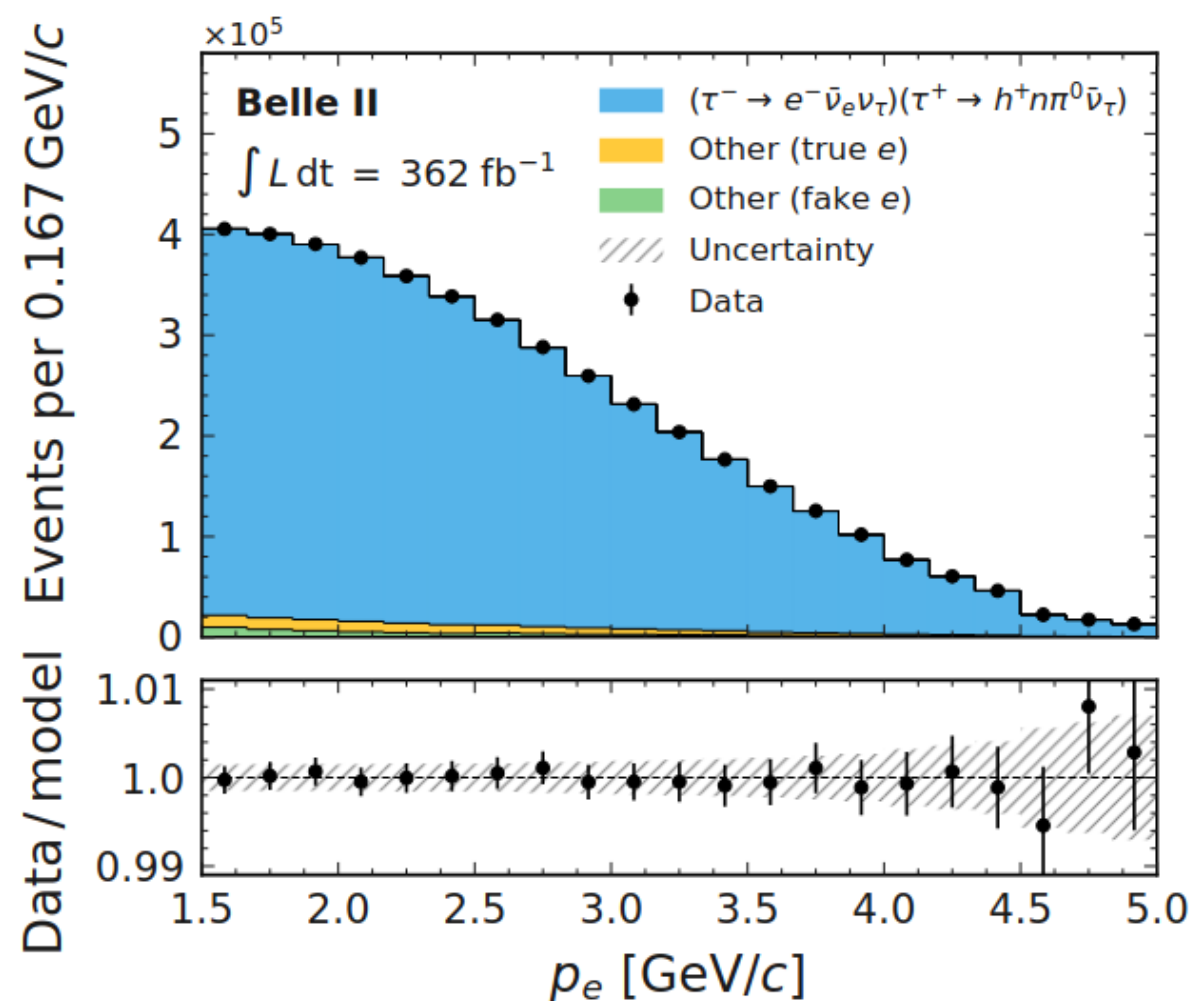
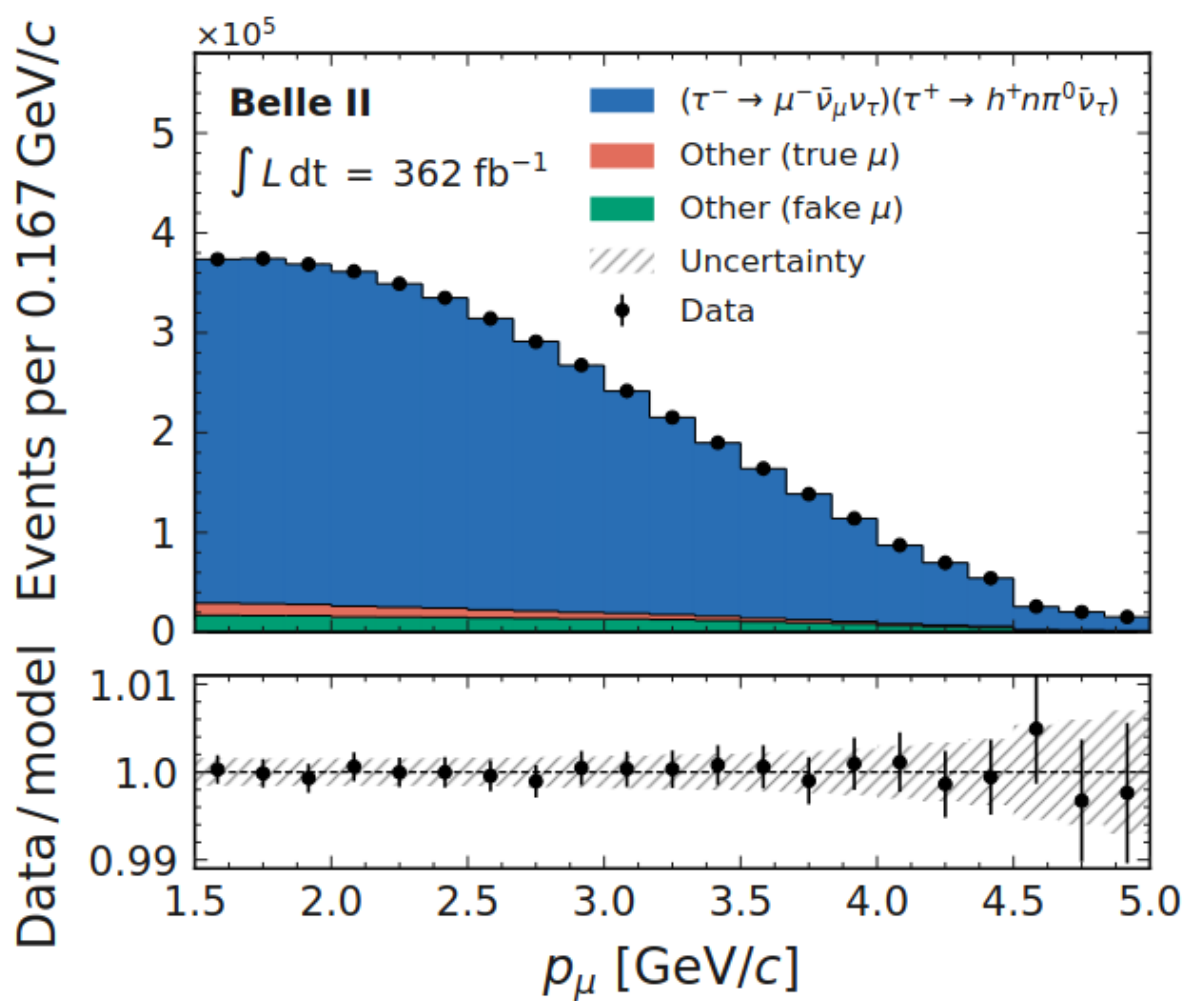
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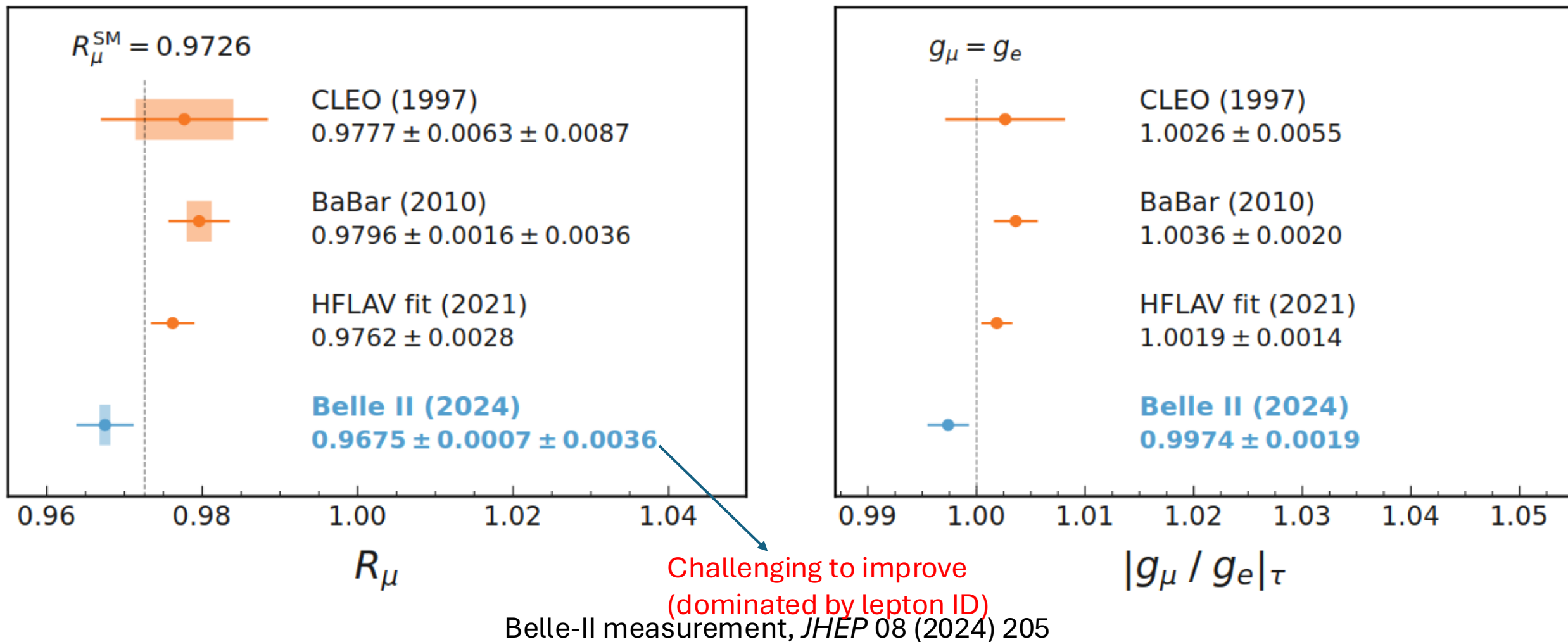


Belle-II measurement, *JHEP* 08 (2024) 205



Belle-II measurement, *JHEP* 08 (2024) 205

**Figure 6.** Observed momentum distribution for muon (left) and electron (right) candidates with fit results overlaid. The lower panel shows the ratio between data and fit results. The hatched area indicates the possible variation of the fitted yields due to systematic effects, with the constraints of the nuisance parameters reduced to their fit uncertainties and correlations taken into account.



**Figure 7.** Determinations of  $R_\mu$  (left) and  $|g_\mu/g_e|_\tau$  (right) from previous individual measurements [11, 12] and the fit from the Heavy Flavor Averaging Group [15], compared with the result of this work. The shaded areas represent the statistical uncertainties, while the error bars indicate the total uncertainties. The vertical dashed line indicates the SM prediction, including mass effects.



# This result in the context of other purely leptonic tests of LU

$$R_{\mu/e}^{\tau} = \frac{\text{Br}[\tau^{-} \rightarrow \mu^{-} \bar{\nu}_{\mu} \nu_{\tau}]}{\text{Br}[\tau^{-} \rightarrow e^{-} \bar{\nu}_{e} \nu_{\tau}]},$$

Belle-II '24 =>  $g_{\mu}/g_e=0.9974(19)$  1.4 $\sigma$

$$R_{\tau/\mu}^{\tau} = \frac{\text{Br}[\tau^{-} \rightarrow e^{-} \bar{\nu}_{e} \nu_{\tau}]}{\text{Br}[\mu^{-} \rightarrow e^{-} \bar{\nu}_{e} \nu_{\mu}]}, \text{ and}$$

$g_{\tau}/g_{\mu}=1.0010(14)$  0.7 $\sigma$

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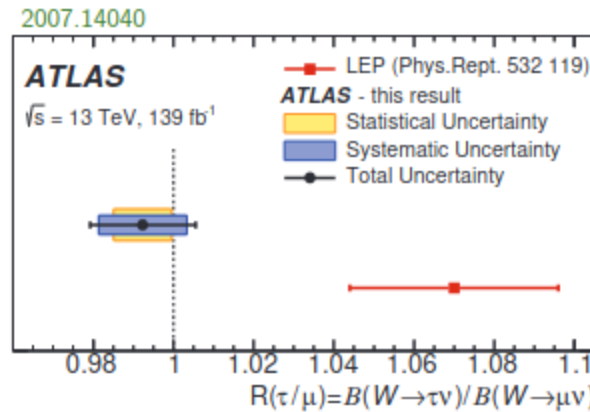
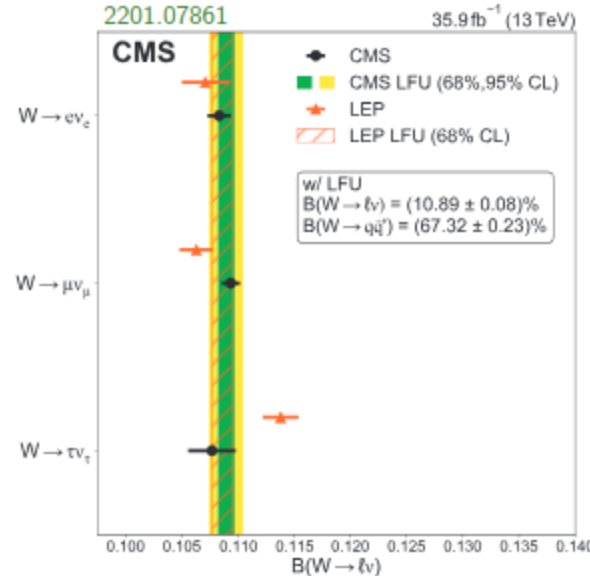
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From W leptonic decays:

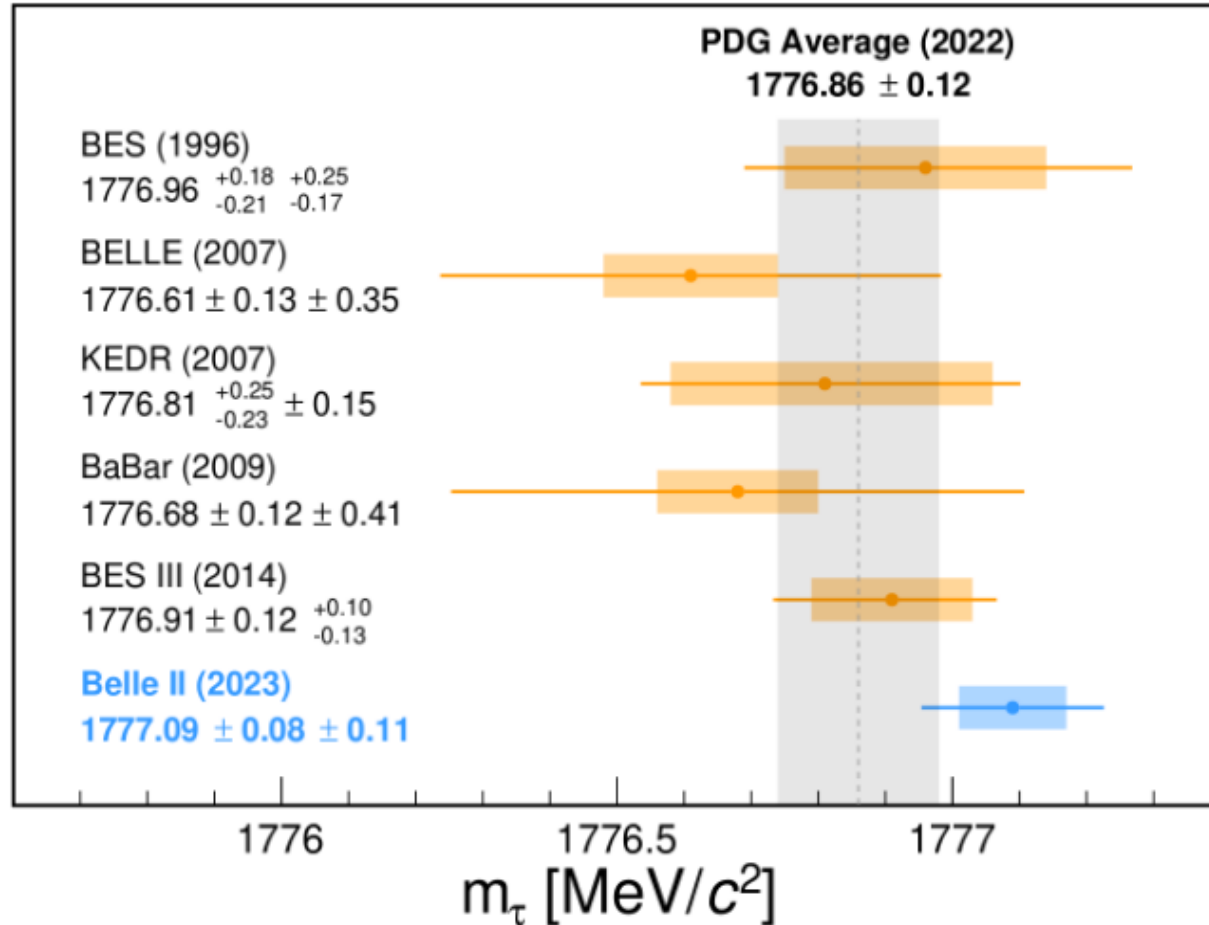


NP constrained at few TeV level

# Lepton Universality

$$\Gamma \sim G_F^2 m_L^5$$

It can also be checked benefitting from the Belle-II  $\tau$  mass measurement, using the leptonic BR determination:



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Opportunities with Tau Leptonic decays at Belle-II

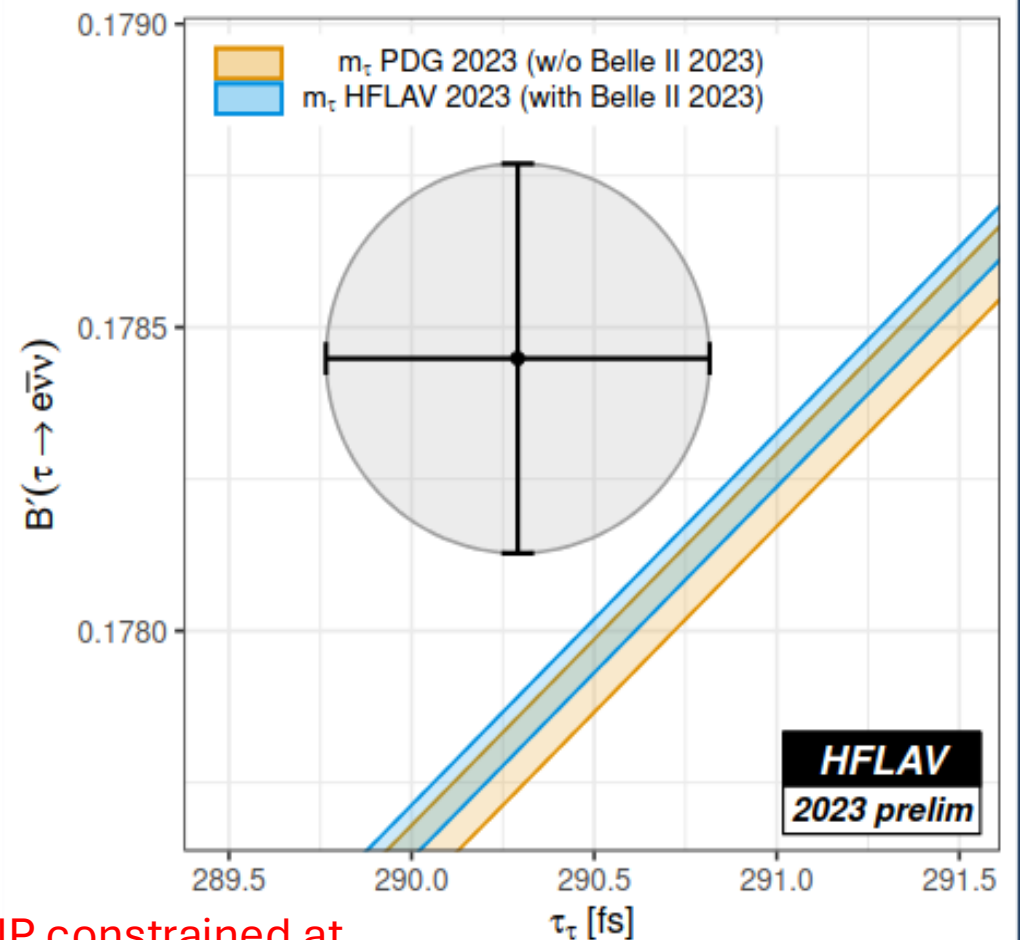
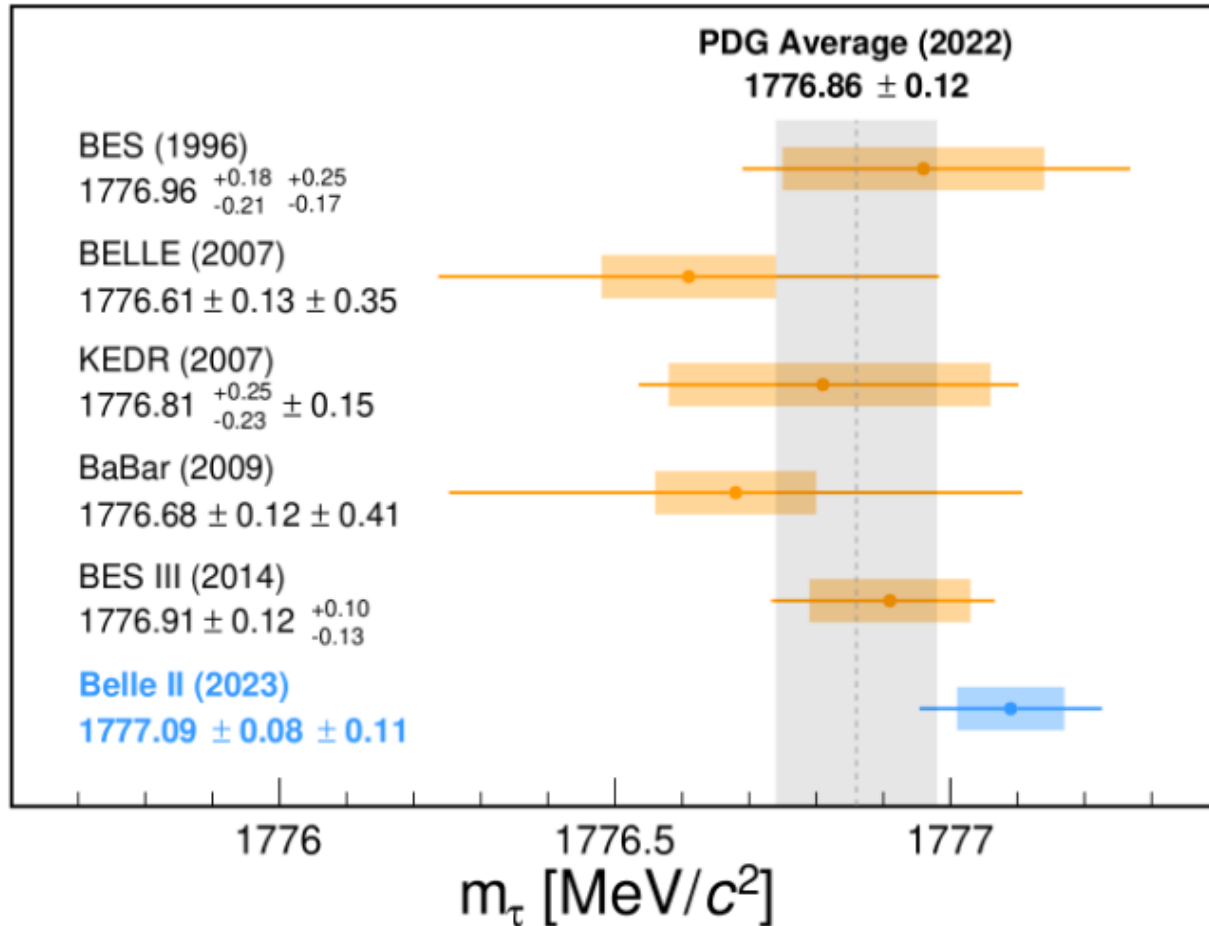
Pablo Roig (Cinvestav, Mexico)

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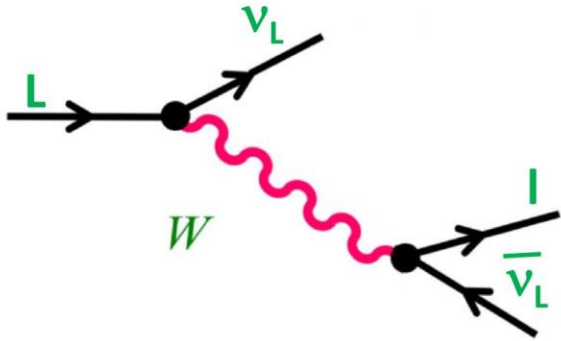


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Opportunities with Tau Leptonic decays at Belle-II

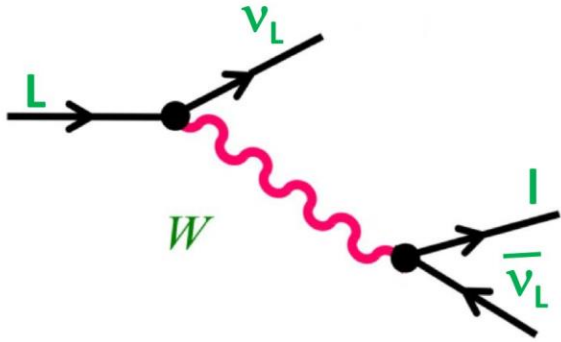
# Lorentz structure of the charged current



$$\mathcal{H} = 4 \frac{G_{\ell\ell}}{\sqrt{2}} \sum_{n,\epsilon,\omega} g_{\epsilon\omega}^n \left[ \bar{\ell}'_{\epsilon} \Gamma^n (\nu_{\ell'})_{\sigma} \right] \left[ (\nu_{\ell})_{\lambda} \Gamma_n l_{\omega} \right],$$

It depends on 4S, 4V & 2T complex couplings labelled also by charged lepton chiralities.

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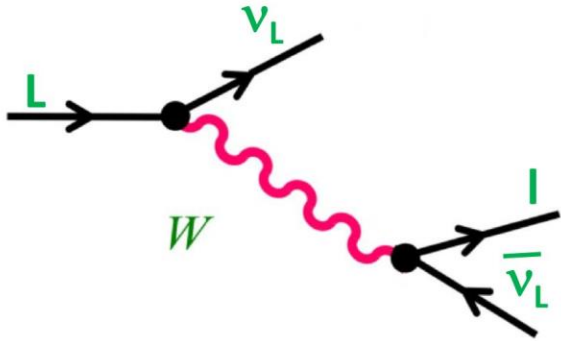


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$$1 = \frac{1}{4} (|g_{RR}^S|^2 + |g_{RL}^S|^2 + |g_{LR}^S|^2 + |g_{LL}^S|^2) + 3 (|g_{RL}^T|^2 + |g_{LR}^T|^2) + (|g_{RR}^V|^2 + |g_{RL}^V|^2 + |g_{LR}^V|^2 + |g_{LL}^V|^2) .$$

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$$\frac{d^2\Gamma_{\ell \rightarrow \ell'}}{dx d\cos\theta} = \frac{m_{\ell}\omega^4}{2\pi^3} G_{\ell\ell}^2 \sqrt{x^2 - x_0^2} \left\{ F(x) - \frac{\xi}{3} \mathcal{P}_{\ell} \sqrt{x^2 - x_0^2} \cos\theta A(x) \right\}$$

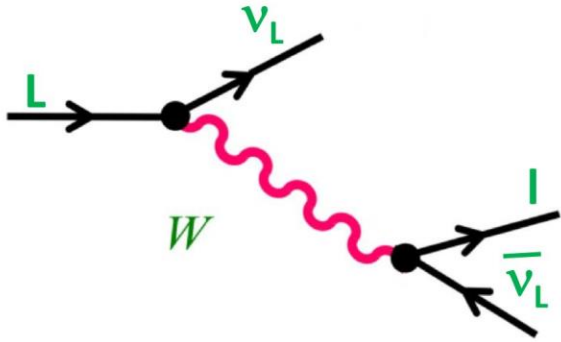
$$\omega \equiv (m_{\ell}^2 + m_{\ell'}^2)/2m_{\ell}$$

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Five additional parameters if the  $l'$  polarization is also measured.



# Lorentz structure of the charged current

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SM		$\mu^- \rightarrow e^- \bar{\nu}_e \nu_\mu$	$\tau^- \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau$	$\tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau$	$\tau^- \rightarrow \ell^- \bar{\nu}_\ell \nu_\tau$
3/4	$\rho$	$0.74979 \pm 0.00026$	$0.763 \pm 0.020$	$0.747 \pm 0.010$	$0.745 \pm 0.008$
0	$\eta$	$0.057 \pm 0.034$	$0.094 \pm 0.073$	—	$0.013 \pm 0.020$
1	$\xi$	$1.0009 \begin{smallmatrix} + 0.0016 \\ - 0.0007 \end{smallmatrix}$	$1.030 \pm 0.059$	$0.994 \pm 0.040$	$0.985 \pm 0.030$
3/4	$\xi\delta$	$0.7511 \begin{smallmatrix} + 0.0012 \\ - 0.0006 \end{smallmatrix}$	$0.778 \pm 0.037$	$0.734 \pm 0.028$	$0.746 \pm 0.021$
1	$\xi'$	$1.00 \pm 0.04$	—	—	—
1	$\xi''$	$0.65 \pm 0.36$	—	—	—

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1	$\xi''$	$0.65 \pm 0.36$	<i>Belle Phys.Rev.Lett.</i> 131 (2023) 2, 021801 <i>Phys.Rev.D</i> 108 (2023) 1, 012003		

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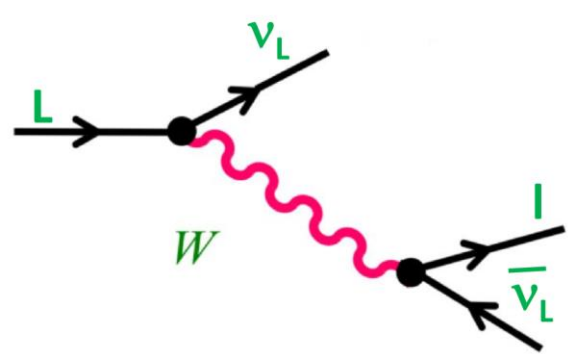
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1	$\xi$	$1.0009$	$\eta = \frac{1}{2} \text{Re} [g_{LL}^V g_{RR}^{S*} + g_{RR}^V g_{LL}^{S*} + g_{LR}^V (g_{RL}^{S*} + 6g_{RL}^{T*}) + g_{RL}^V (g_{LR}^{S*} + 6g_{LR}^{T*})]$	
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# Michel parameters in presence of massive Dirac or Majorana neutrinos

(Juanma Márquez, Gabriel López Castro & P. R., JHEP11(2022)117)

$$\frac{d\Gamma}{dx d\cos\theta} = \frac{m_1}{4\pi^3} \omega^4 G_{II'}^2 \sqrt{x^2 - x_0^2} \left( F(x) - \frac{\xi}{3} \mathcal{P} \sqrt{x^2 - x_0^2} \cos\theta A(x) \right) \times [1 + \hat{\zeta} \cdot \vec{\mathcal{P}}_{I'}(x, \theta)],$$

Massless case

**'Fundamental' Theory**

Λ

$$\frac{d\Gamma}{dx d\cos\theta} = \sum_{j,k} \frac{m_1}{4\pi^3} \omega^4 G_{II'}^2 \sqrt{x^2 - x_0^2} \times \left( (F_{IS}(x) + F'_{IS}(x) + F''_{IS}(x)) - \mathcal{P} \cos\theta (F_{AS}(x) + F'_{AS}(x) + F''_{AS}(x)) \right) \times [1 + \hat{\zeta} \cdot \vec{\mathcal{P}}_{I'}(x, \theta)],$$

Linear in  $\nu$  masses

Quadratic in  $\nu$  masses

Massive case

**EFT**

E

## What is new:

- We write our expressions in the PDG parametrization form, in a way that complements all previous results, facilitating their application to model-dependent scenarios.
- We classify the Dirac and Majorana contributions with the help of a flag parameter  $\epsilon = 0, 1$ , making easier to distinguish between Dirac and Majorana nature of neutrinos.
- We discuss their main differences, together with some examples of its application to model-dependent theories.
- We also introduced and discussed the leading W-boson propagator correction to the differential decay rate including the final charged-lepton polarization.

— Dirac  
— Majorana

Study suggested by Denis Epifanov (generalizing earlier work by Michel, Bouchiat-Michel, Fetscher-Gerber-Johnson, Langacker-London, Shrock, Doi-Kotani-Takasugi, etc.)

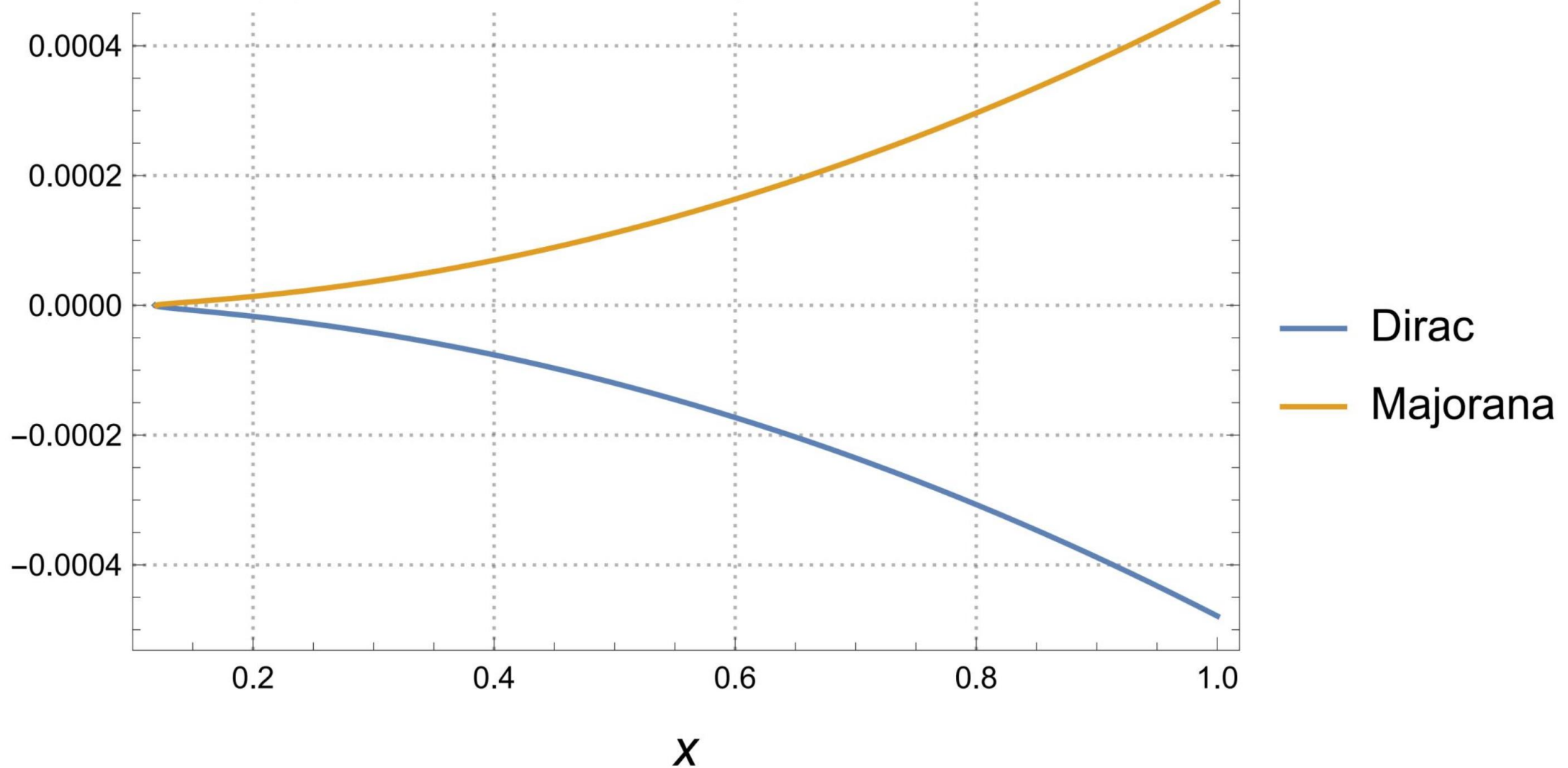
JHEP 11 (2022) 117

Figure 6: Neutrino mass contribution to Dirac and Majorana distributions.

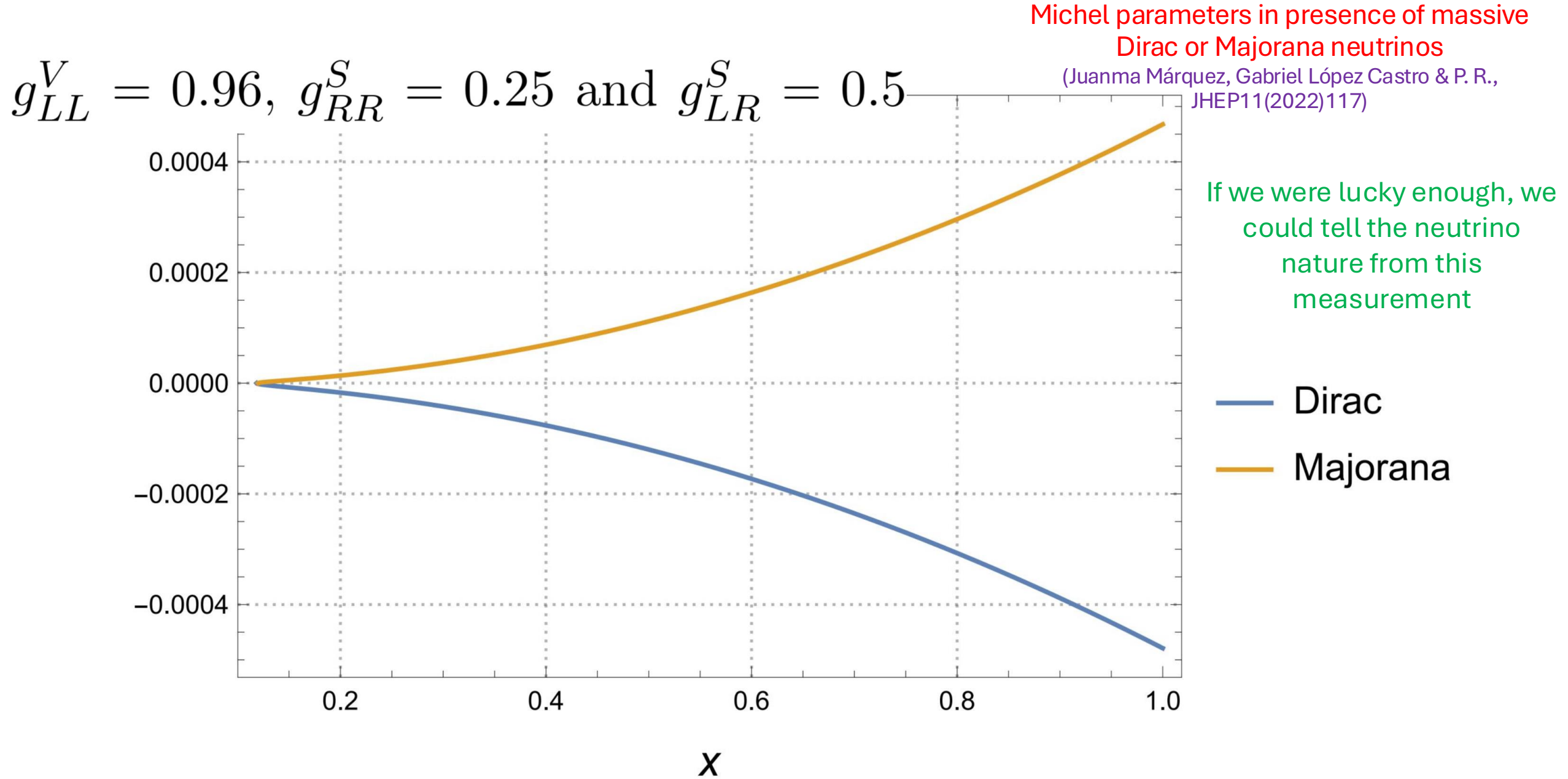
Michel parameters in presence of massive Dirac or Majorana neutrinos

(Juanma Márquez, Gabriel López Castro & P. R., JHEP11(2022)117)

$$g_{LL}^V = 0.96, g_{RR}^S = 0.25 \text{ and } g_{LR}^S = 0.5$$

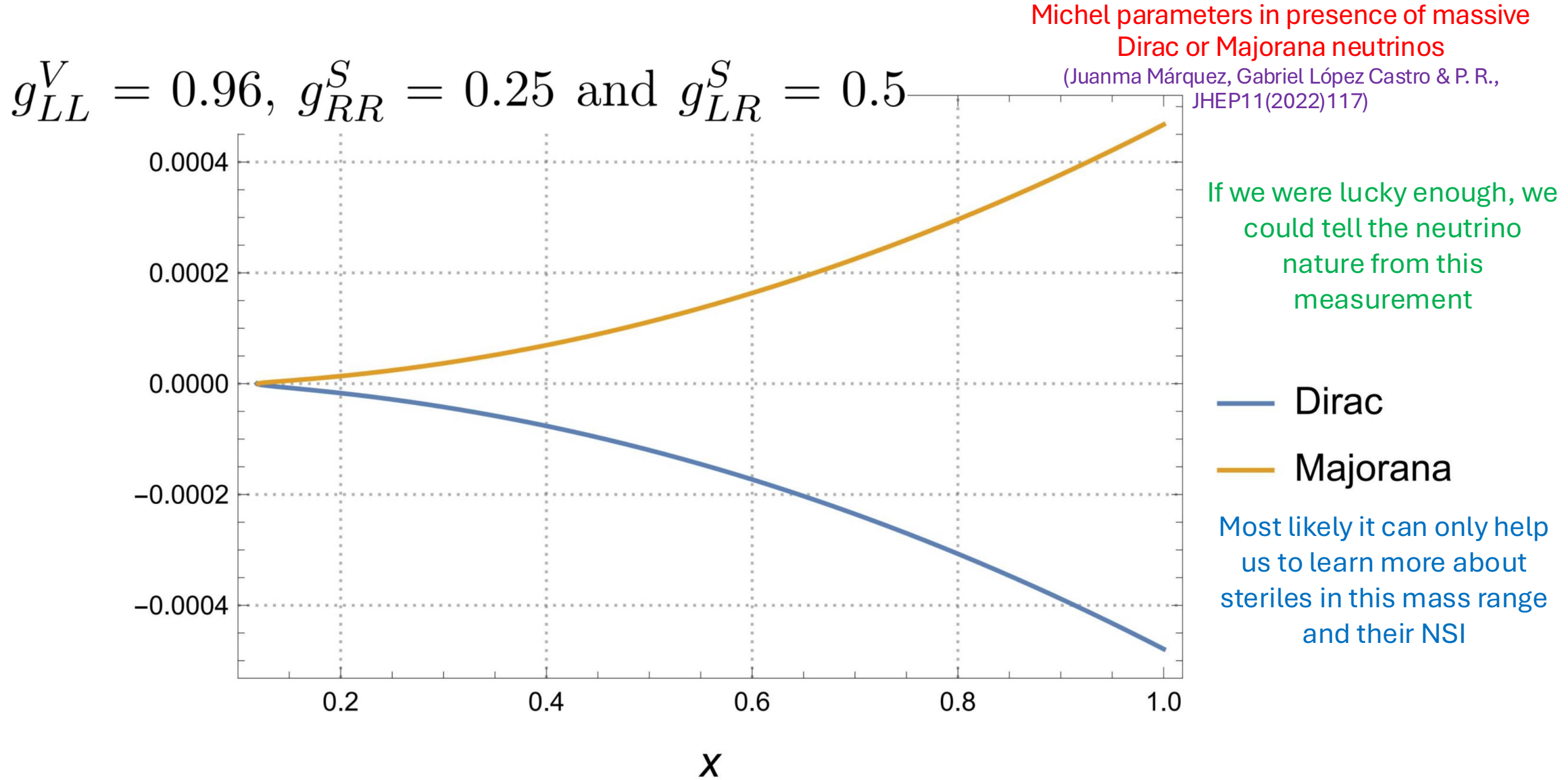


**Figure 6:** Neutrino mass contribution to Dirac and Majorana distributions.

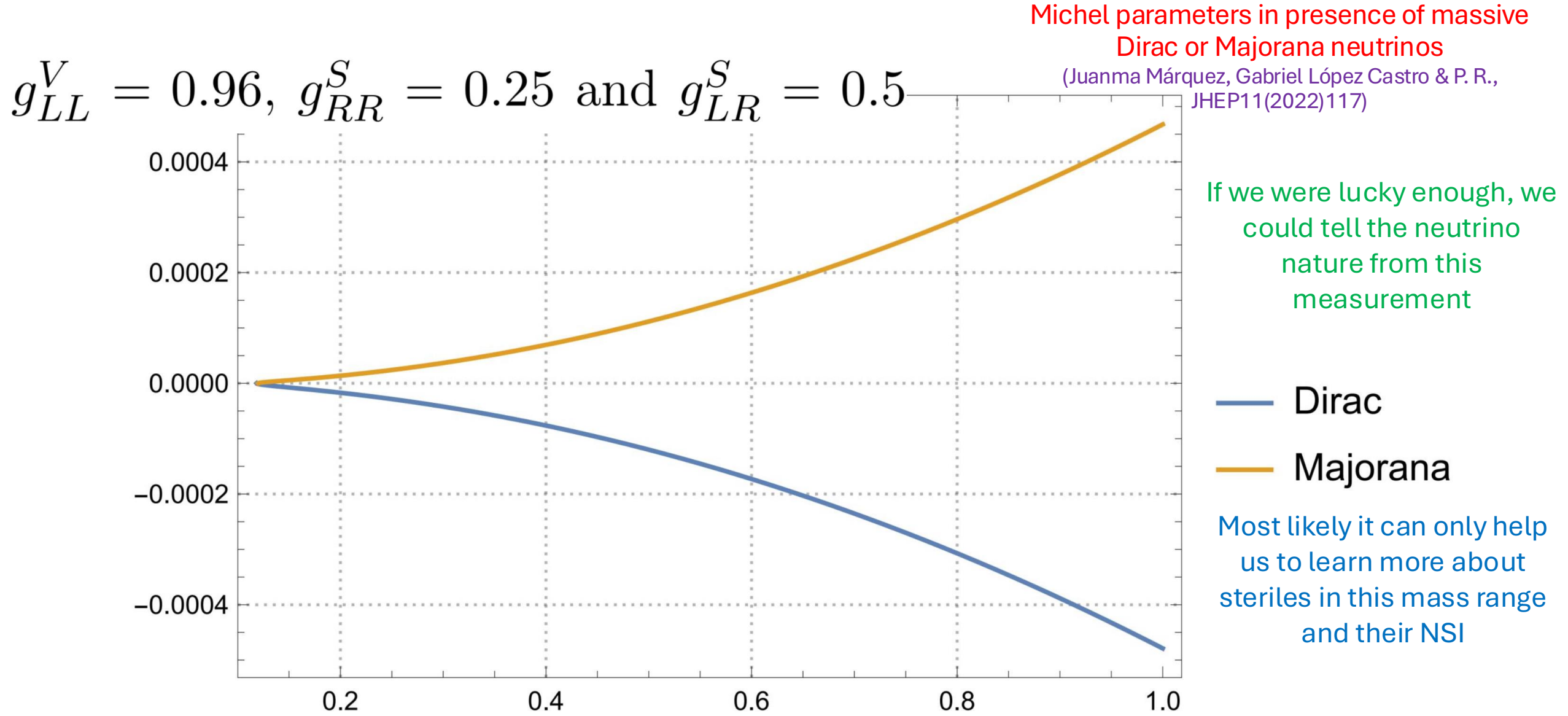


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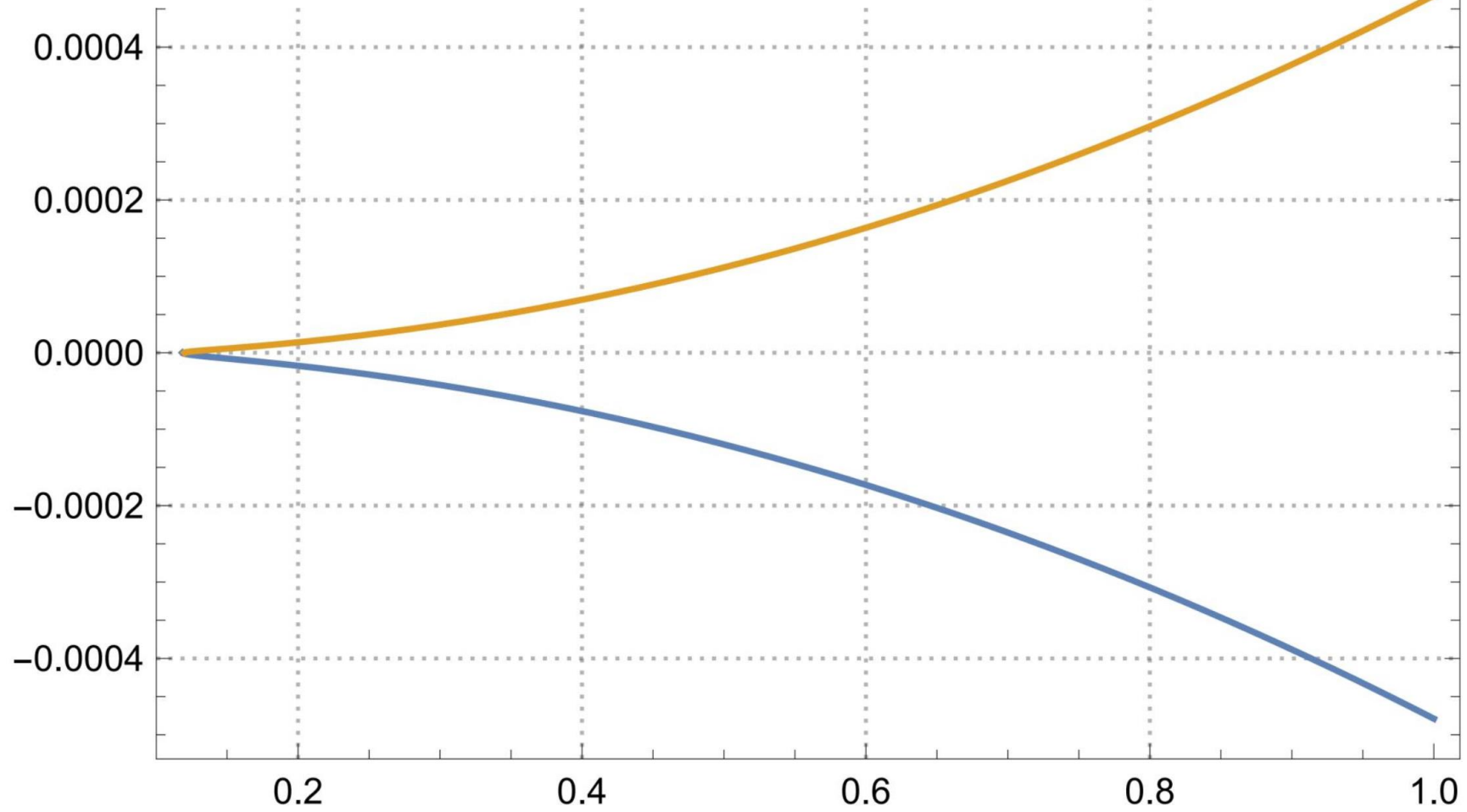
Also interesting radiative processes: Belle'18, Arbuzov-Kopilova, Flores Tlalpa- López Castro-R., giving access to additional Michel parameters.

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Michel parameters in presence of massive Dirac or Majorana neutrinos

(Juanma Márquez, Gabriel López Castro & P. R., JHEP11(2022)117)

$$g_{LL}^V = 0.96, g_{RR}^S = 0.25 \text{ and } g_{LR}^S = 0.5$$



If we were lucky enough, we could tell the neutrino nature from this measurement

Dirac  
Majorana

Most likely it can only help us to learn more about steriles in this mass range and their NSI

Also interesting radiative processes: Belle'18, Arbuzov-Kopilova, Flores Tlalpa- López Castro-R., giving access to additional Michel parameters

X

Great benefits in Chiral Belle-II!!

Figure 6: Neutrino mass contribution to Dirac and Majorana distributions.

# Inferring the nature of active neutrinos: Dirac or Majorana

(C.S. Kim, M.V.M. Murthy & D. Sahoo, PRD105(2022)11,113006)

This paper apparently avoids the Kayser's confusion theorem 'Any property differentiating Dirac/Majorana neutrinos will be suppressed by active neutrino masses, with neutrinos coupling to the SM's  $SU(2)_L$ '.

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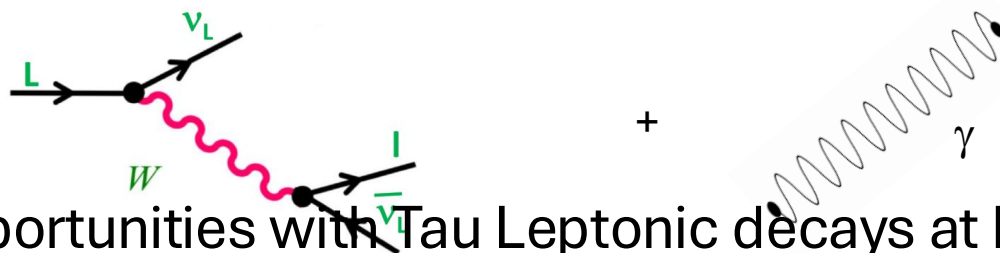
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Looking in depth (Juanma Márquez, Diego Portillo & P. R., PRD109 (2024) 3, 033005), there is a loophole in their derivation. When corrected, it yields observables which are orders of magnitude smaller than initially thought, likely preventing the observation of this effect.



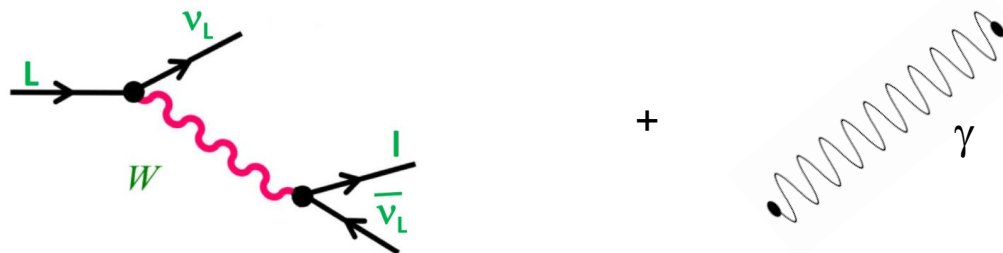
See also [Akhmedov&Trautner, JHEP109 \(2024\) 3, 033005](#)

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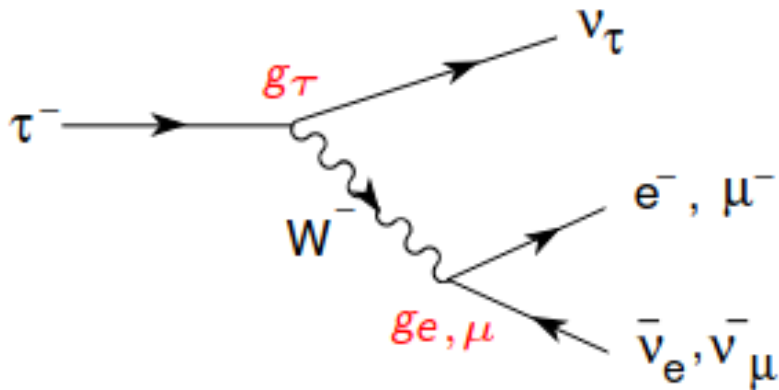
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In fact, applying Kim's paper one should get BRs  $\sim 10^{-4}$  for the b2b configuration, while we predict rates  $\sim 10^{-10}$ . **Belle-II can disprove the former prediction!**



See also Akhmedov&Trautner,  
JHEP109 (2024) 3, 033005

# Opportunities with Tau Leptonic decays at Belle-II



- Lepton Universality
- Lorentz structure of the charged current (including possible heavy sterile neutrinos)