# Dark sector search with hadrons -Wish list beyond Belle II Physics Book-

### Oct 14, 2024 [Belle II Physics Week, KEK] Kohsaku Tobioka [Tobi] Florida State University, KEK Theory center

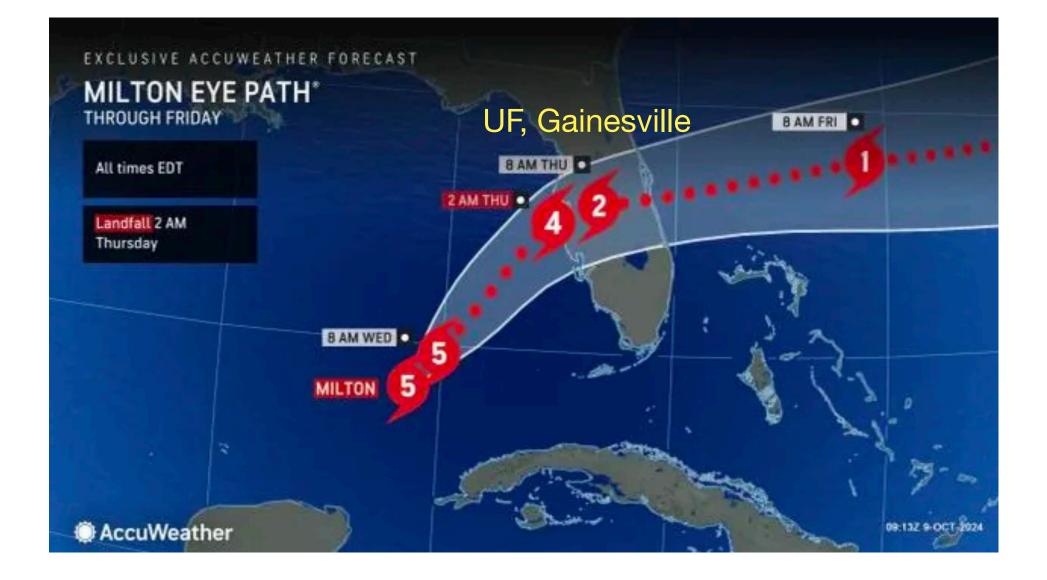




### Before start...

### Degeneracy in Florida





### **Before start...**

 Originally I wanted to overview many proposals for Belle II, but I realize I already have many for given time. [Apologies!]

->My wish list for young exp colleagues. [5 modes]

### **B-factory is quite unique**

ktobioka@fsu.edu not @ufl.edu...



https://www.irasutoya.com/

 I'm attending only Mon/Tue/(Wed?), but feel free to interact/email or ZOOM. Many collaborations with experimentalists started over corridor chats (for me).

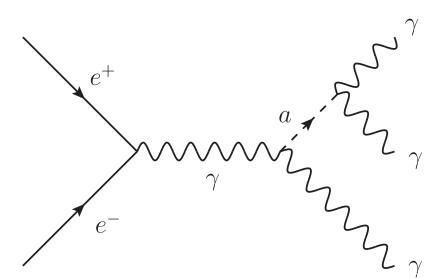


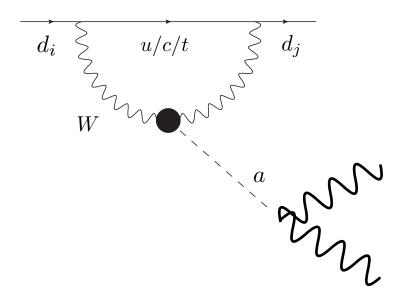
# **BSM wishlist beyond Belle II Physics Book**

- 16 Dark Sectors and Light Higgs
  - 16.1 Theory
  - Experiment: Scattering Processes 16.2
    - 16.2.1 Search for Dark Photons decaying into Light Dark Matter ("Single-photon search")
    - 16.2.2 Search for Axion-like particles
    - 16.2.3 Search for Dark Photons decaying into charged leptons and charged hadrons
    - 16.2.4 Search for Dark Photons decaying into Light Dark Matter in  $e^+e^- \to A'\ell^+\ell^-$
  - 16.3 Experiment: Quarkonium Decay
    - 16.3.1 Searches for BSM physics in invisible  $\Upsilon(1S)$  decays 16.3.2 Probe of new light CP even Higgs bosons from bottomonium  $\chi_{b0}$  decay
    - 16.3.3 Search for a CP-odd Higgs boson in radiative  $\Upsilon(3S)$  decays
    - 16.3.4 Prospects for lepton universality tests in  $\Upsilon(1S)$  decays
  - Conclusions 16.4

### This list will be expanded [all the afternoon talks!]

Only electroweak boson couplings!!





Many interesting channels with hadronic/gluon couplings **Heavy QCD axion** 

[S. Gori's talk for lepton couplings.]



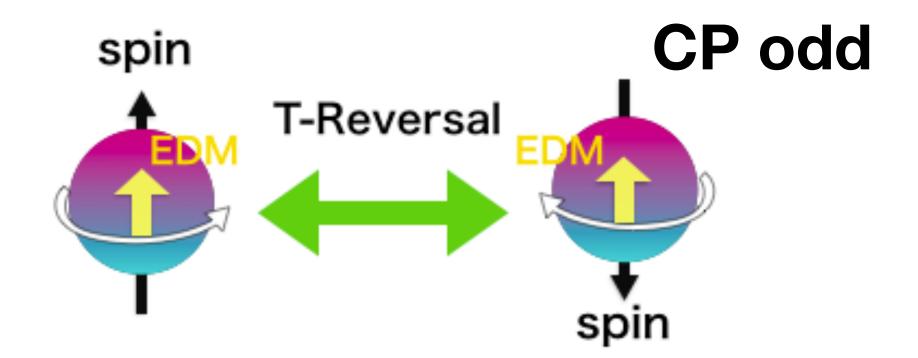


# **Strong CP problem and QCD Axion**

### The strong CP problem

- The unknown of the SM: CP phase in the strong sector
- Neutron EDM sets a very stringent upper bound:  $\bar{\theta} \lesssim 10^{-10}$
- We expect the order 1 number because CP is violated in the CKM matrix (thanks to Belle and Babar).
   If we start from CP conserving theory for θ, we need to break it to explain CKM.

 $\frac{\alpha_s \bar{\theta}}{G} G^{a\mu
u} \tilde{G}^{a\mu
u}$  $8\pi$ 



# Strong CP problem and QCD Axion

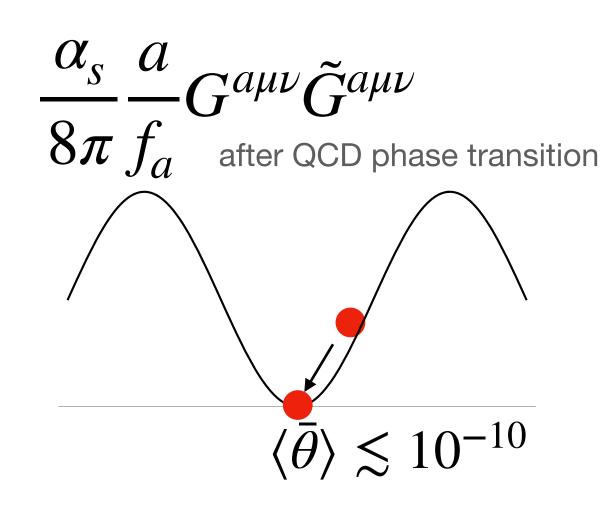
### The strong CP problem

- The unknown of the SM: CP phase in the strong sector
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### **QCD** Axion solution

- Promote  $\theta$  to a field  $a/f_a$  dynamically settles the CP phase to the minimum.
- Peccei-Quinn symmetry: Global U(1) that generates the axion as a Nambu-Goldstone boson.  $f_a$  is the breaking scale.
- Attractive dark matter candidate, typically ma<meV.</li>

$$\frac{\alpha_s\bar{\theta}}{8\pi}G^{a\mu\nu}\tilde{G}^{a\mu\nu}$$



# From QCD Axion to Heavy QCD Axion

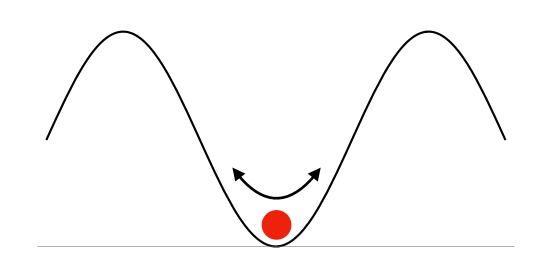
Standard QCD axion has mass prediction

$$m_a \sim \frac{m_\pi f_\pi}{f_a} \sim 0.1 \mathrm{MeV}$$

Heavy QCD axion, axion heavier than "standard mass"

 $m_a \sim \left( m_0^2 + \frac{m_\pi f_\pi}{f_a} \right)^{n/2} \gg \frac{m_\pi f_\pi}{f_a}$ 

 $V\left(\frac{100 \text{GeV}}{f}\right)$ 



Models: additional QCD SU(3)' to raise ma Berezhiani et al('01); Hook('04); Fukuda et al('04). Dimopoulos et al('16); Hook et al('19); Valenti ('22)... Another class: Agrawal and Howe ('17)...

# From QCD Axion to Heavy QCD Axion

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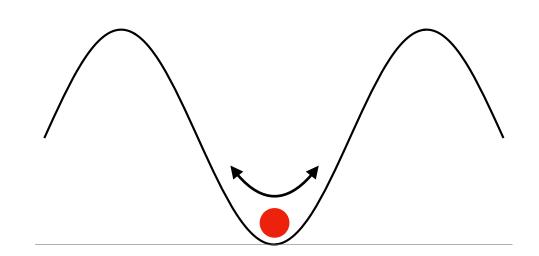
$$m_a \sim \left( m_0^2 + \frac{m_\pi f_\pi}{f_a} \right)^{1/2}$$

Why interesting?

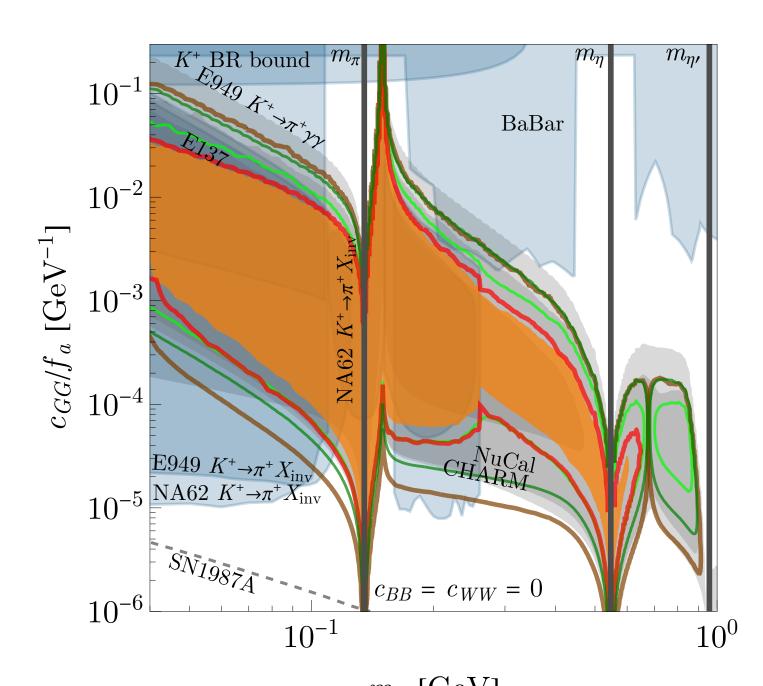
- 1. Viable with lower  $f_{a}$ . Should cover  $m_a > 100 MeV$ .
- 2. Better quality of PQ symmetry

**★**It it not dark matter because it decays in cosmological time scale

100GeV



Models: additional QCD SU(3)' to raise ma Berezhiani et al('01); Hook('04); Fukuda et al('04). Dimopoulos et al('16); Hook et al('19); Valenti ('22)... Another class: Agrawal and Howe ('17)...



7

### ALPs: Axion-like Particles vs Heavy QCD Axion

### Typical ALPs



$$\frac{\alpha_W c_2}{8\pi} a W \tilde{W} \quad \frac{\alpha_1 c_1}{8\pi} a B \tilde{B}$$

$$\frac{\alpha_{EM}c_{\gamma}}{8\pi}aF\tilde{F} = \frac{g_{a\gamma\gamma}}{4}aF\tilde{F}$$



Completely free

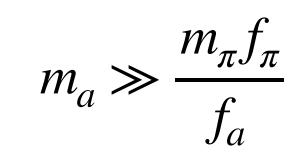
Heavy QCD Axion

Must:

$$\frac{c_g \alpha_s}{8\pi} \frac{a}{f_a} G^{a\mu\nu} \tilde{G}^{a\mu\nu}$$

Optional: *aBÃ*, *aWŴ*, *aFÃ* 

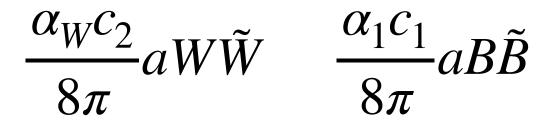
Heavier than standard one  $m_a$ 



# **ALPs:**Axion-like Particles vs Heavy QCD Axion

### Typical ALPs

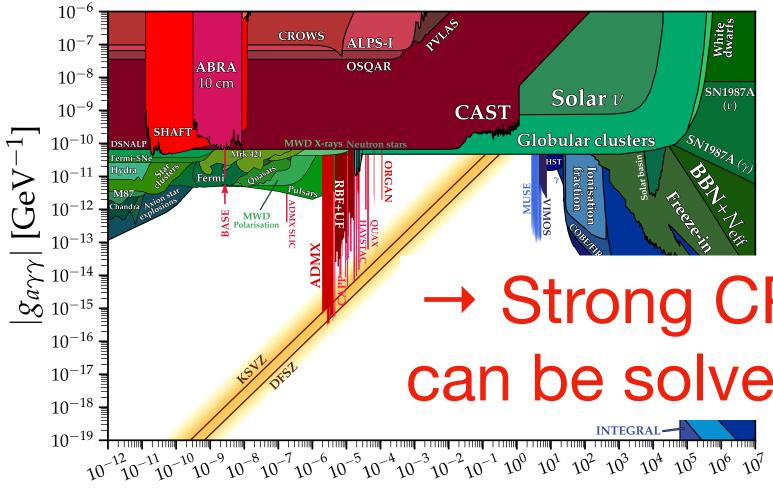
### • couplings



 $\frac{\alpha_{EM}c_{\gamma}}{8\pi}aF\tilde{F} = \frac{g_{a\gamma\gamma}}{4}aF\tilde{F}$ 



Completely free



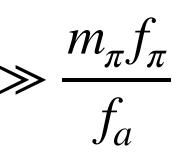
 $m_a \,[\mathrm{eV}]$ 

### Heavy QCD Axion

Must:  $\frac{c_g \alpha_s}{8\pi} \frac{a}{f_a} G^{a\mu\nu} \tilde{G}^{a\mu\nu}$ 

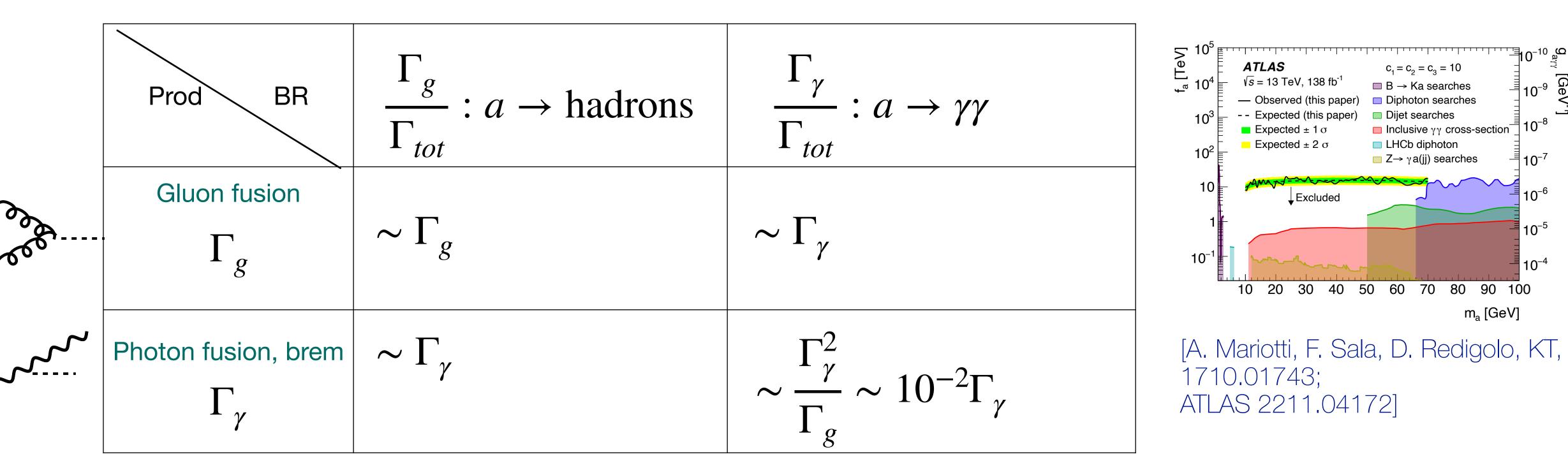
Optional:  $aB\tilde{B}, aW\tilde{W}, aF\tilde{F}$ 

Heavier than standard one  $m_a \gg \frac{m_{\pi} f_{\pi}}{f_{\pi}}$ 

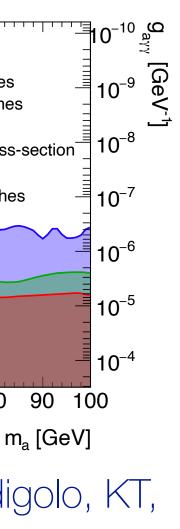


### → Strong CP can be solved

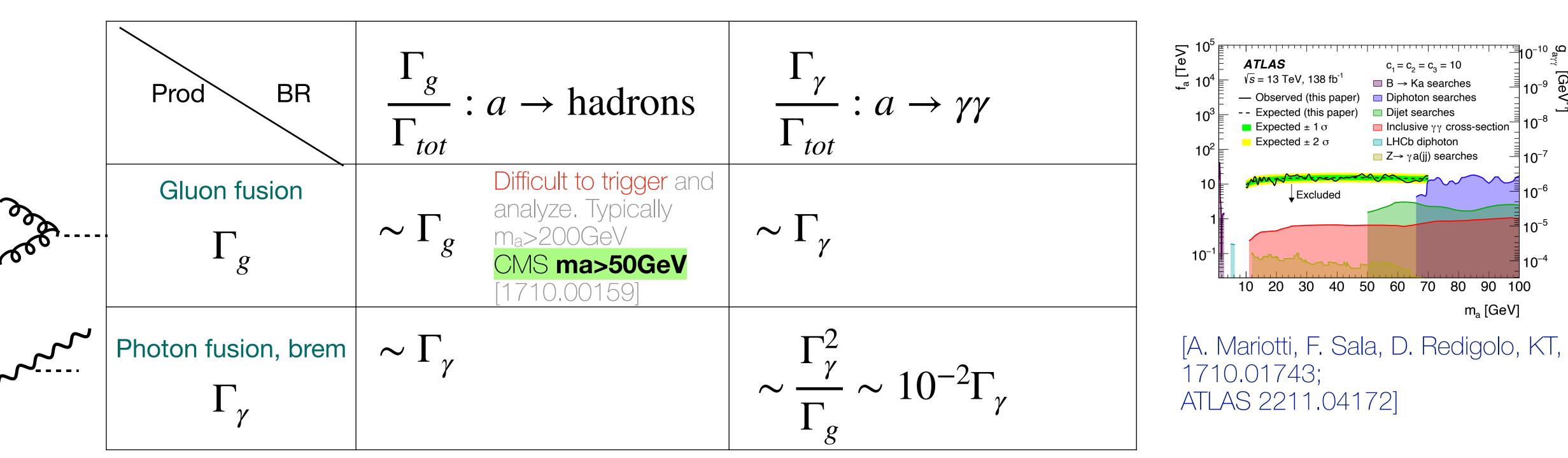
• To get ideas, consider resonance search, from LHC to lower mass Using narrow width approximation with  $c_v \sim c_g$ :  $\sigma \cdot BR \sim \frac{\Gamma_{\text{prod}}}{m_{\alpha} E_{CM}^2} \frac{\Gamma_{\text{decay}}}{\Gamma_{\text{tot}}} \qquad \alpha_s^2 \gg \alpha_{EM}^2$ 



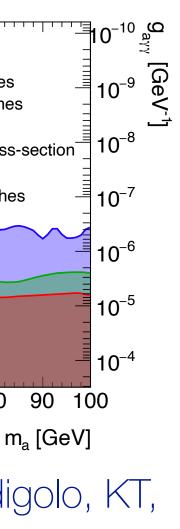
$$\rightarrow \Gamma_{tot} \sim \Gamma_g \gg \Gamma_{\gamma}$$



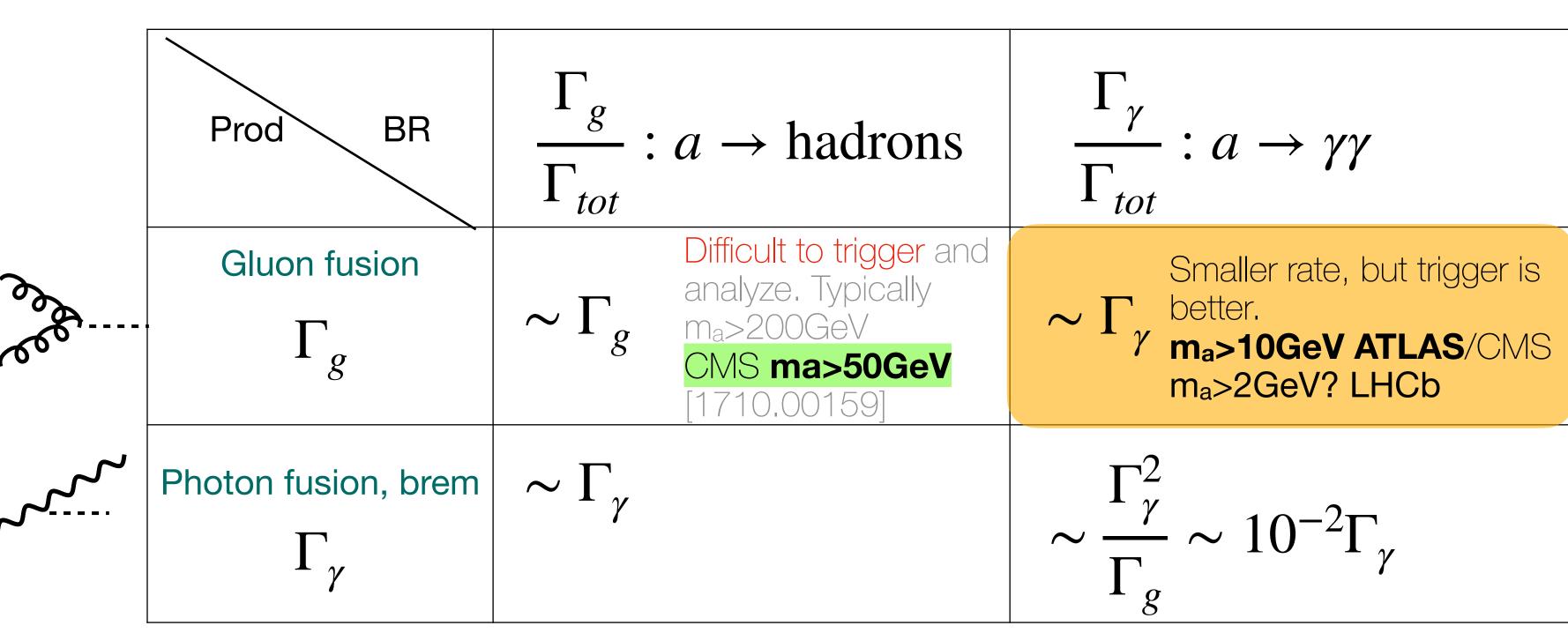
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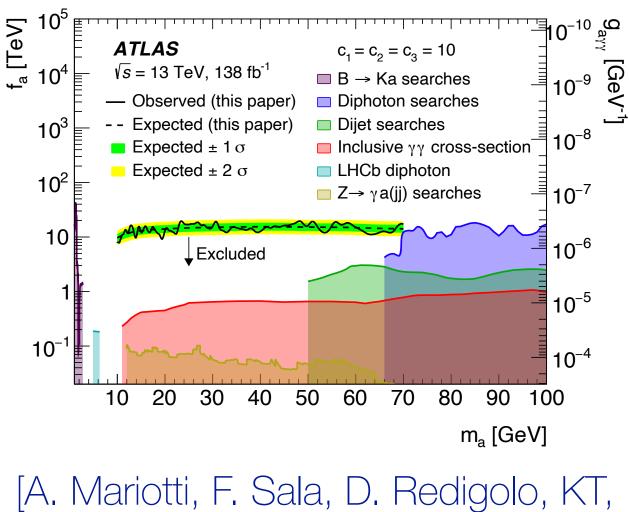
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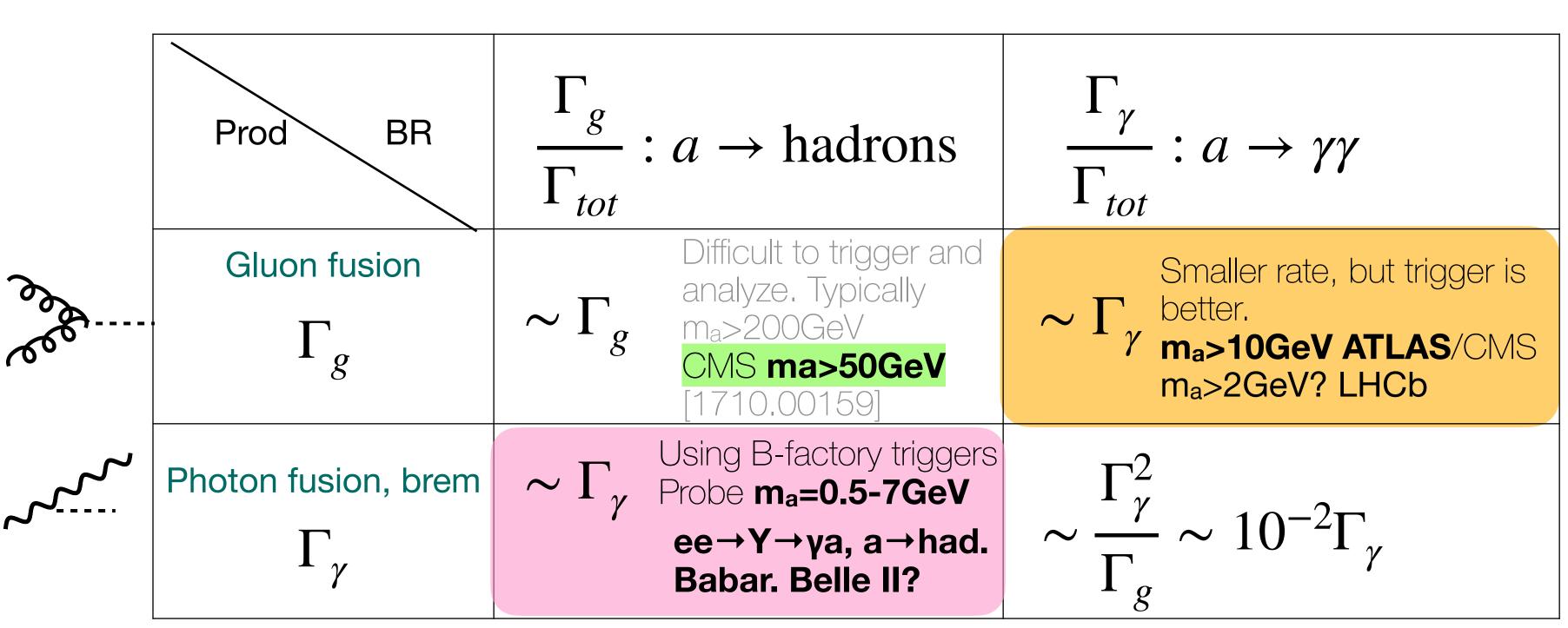


 $\sigma \cdot \mathrm{BR} \sim \frac{\Gamma_{\mathrm{prod}}}{m_a E_{CM}^2} \frac{\Gamma_{\mathrm{decay}}}{\Gamma_{\mathrm{tot}}} \qquad \alpha_s^2 \gg \alpha_{EM}^2 \to \Gamma_{tot} \sim \Gamma_g \gg \Gamma_{\gamma}$ 

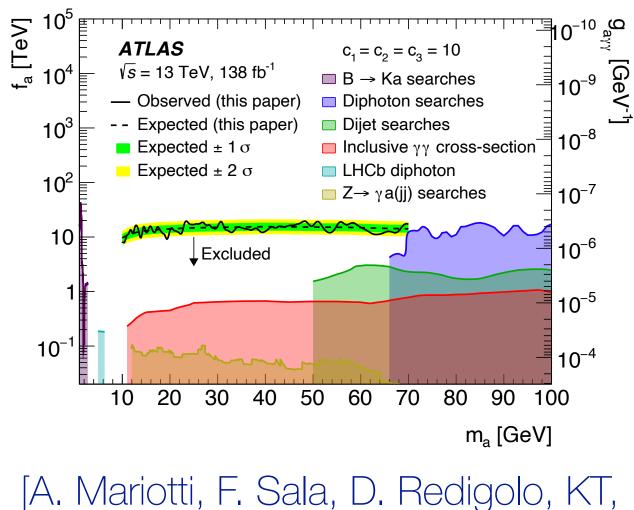


[A. Mariotti, F. Sala, D. Redig 1710.01743; ATLAS 2211.04172]

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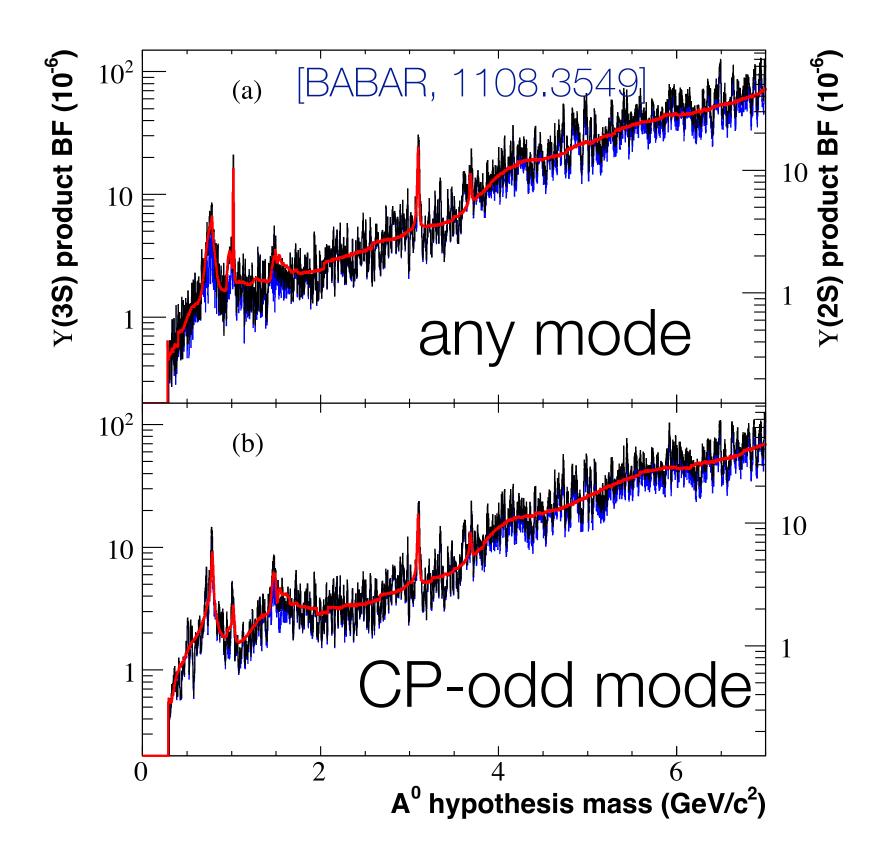


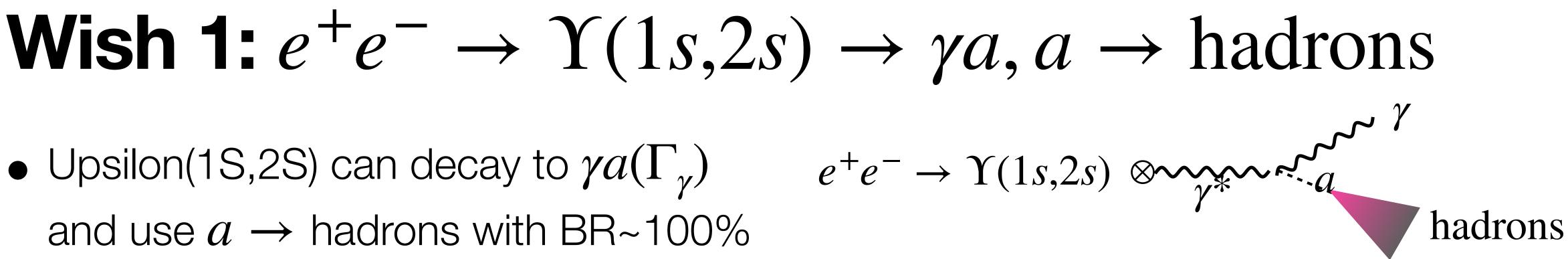
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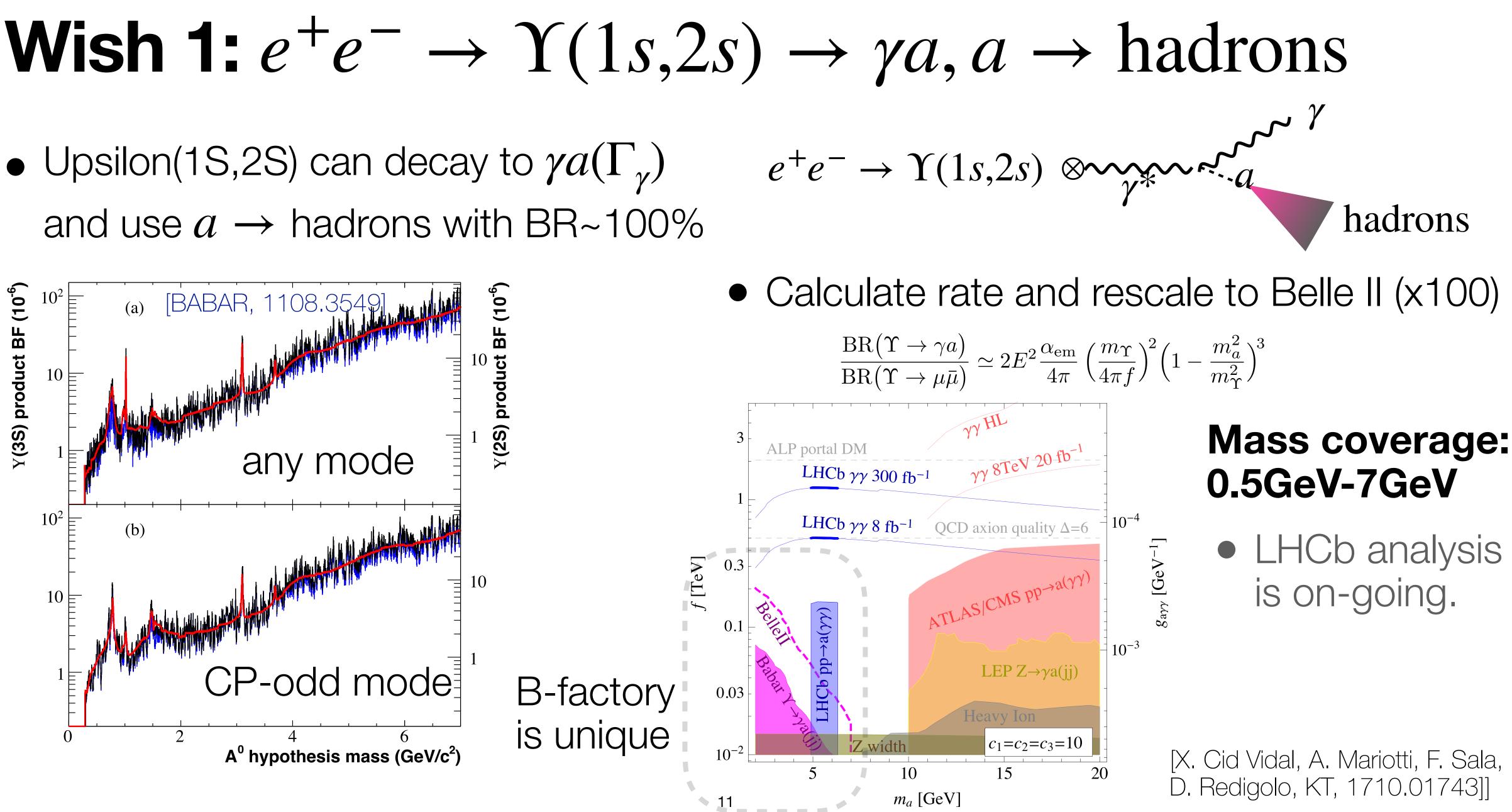
[A. Mariotti, F. Sala, D. Redig 1710.01743; ATLAS 2211.04172]

and use  $a \rightarrow$  hadrons with BR~100%



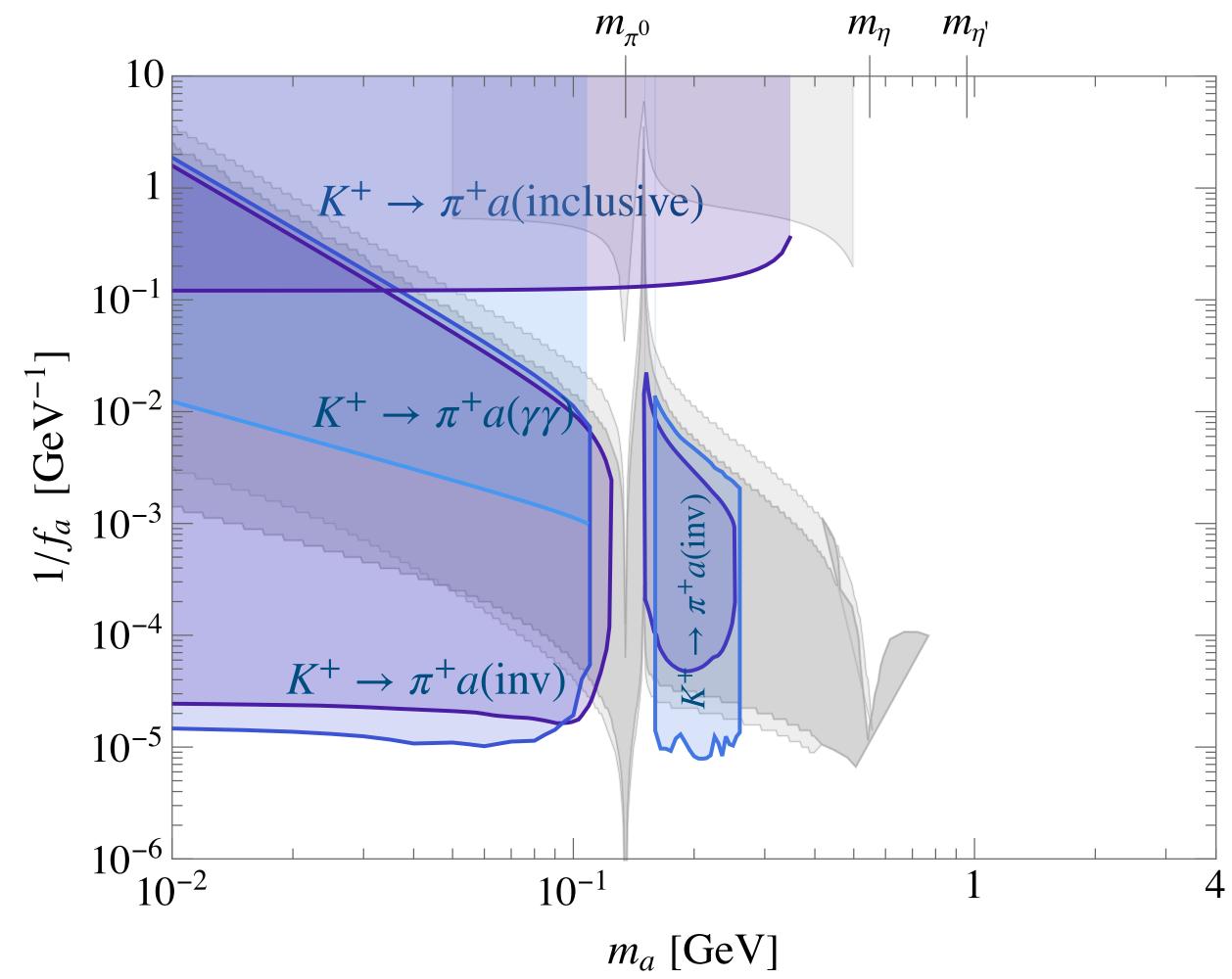






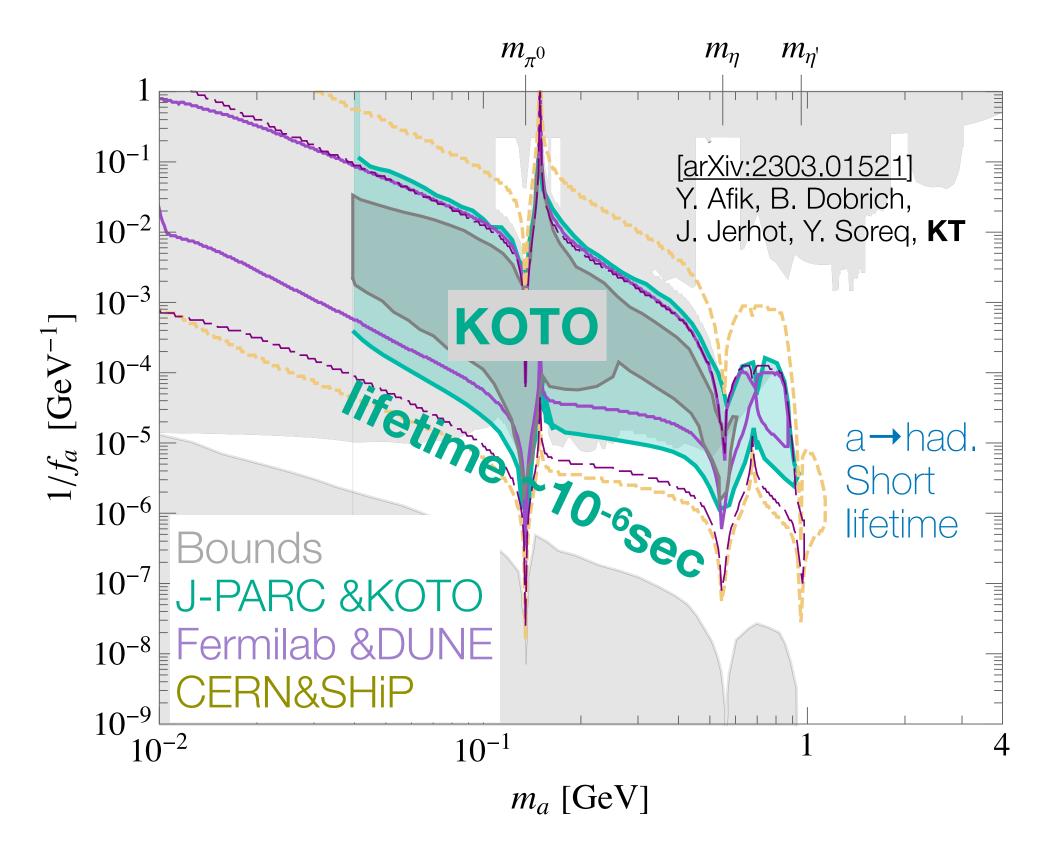
### MeV-GeV range, kaon and proton beam-dump

• LHC down to 10GeV(2GeV).  $K^+ \rightarrow \pi^+ a$  constraints m<sub>a</sub><m<sub>K</sub>-m<sub> $\pi$ </sub>~0.35GeV

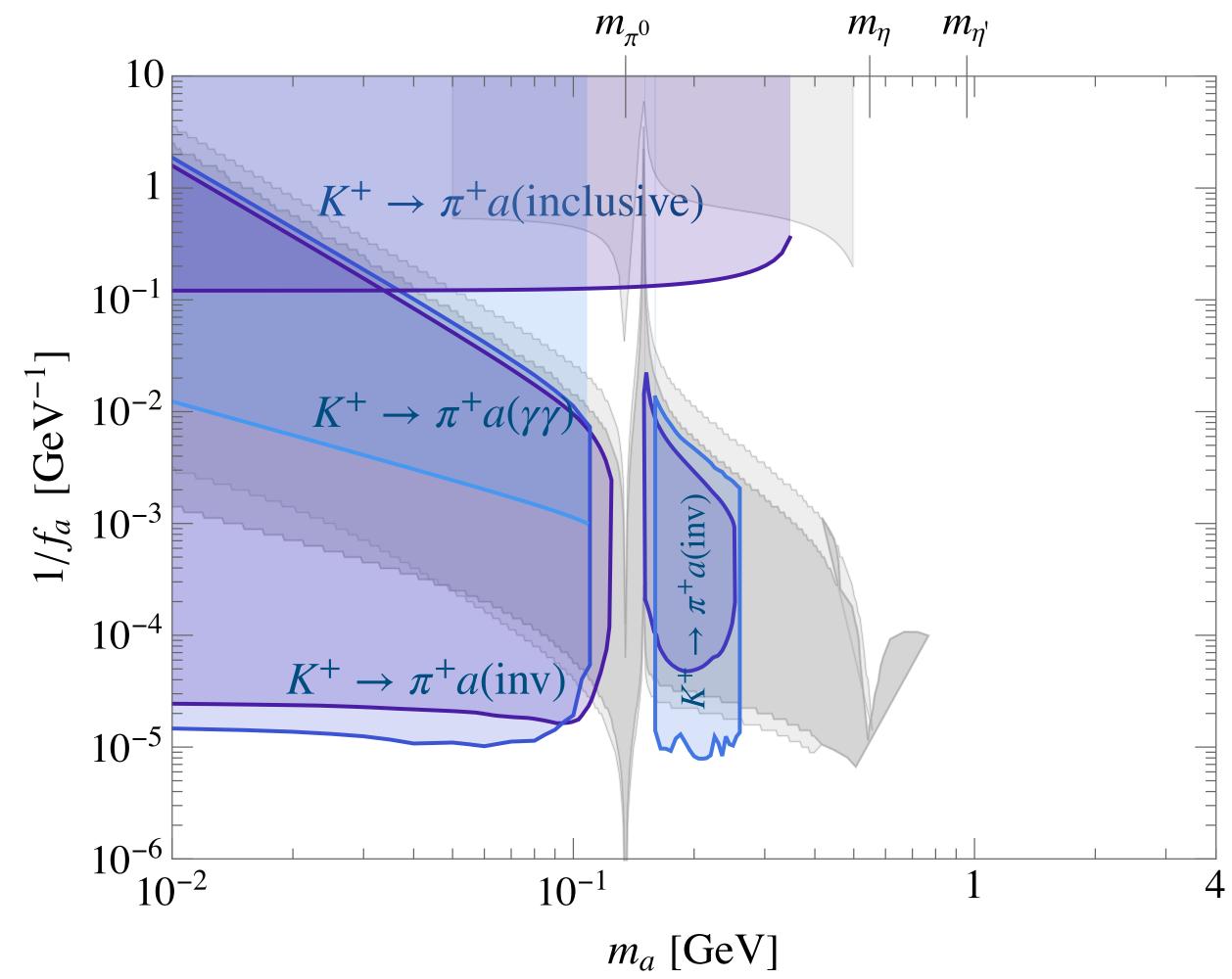


### MeV-GeV range, kaon and proton beam-dump

- Beam-dump search is limited to m<sub>a</sub><1GeV due to shorter lifetime.

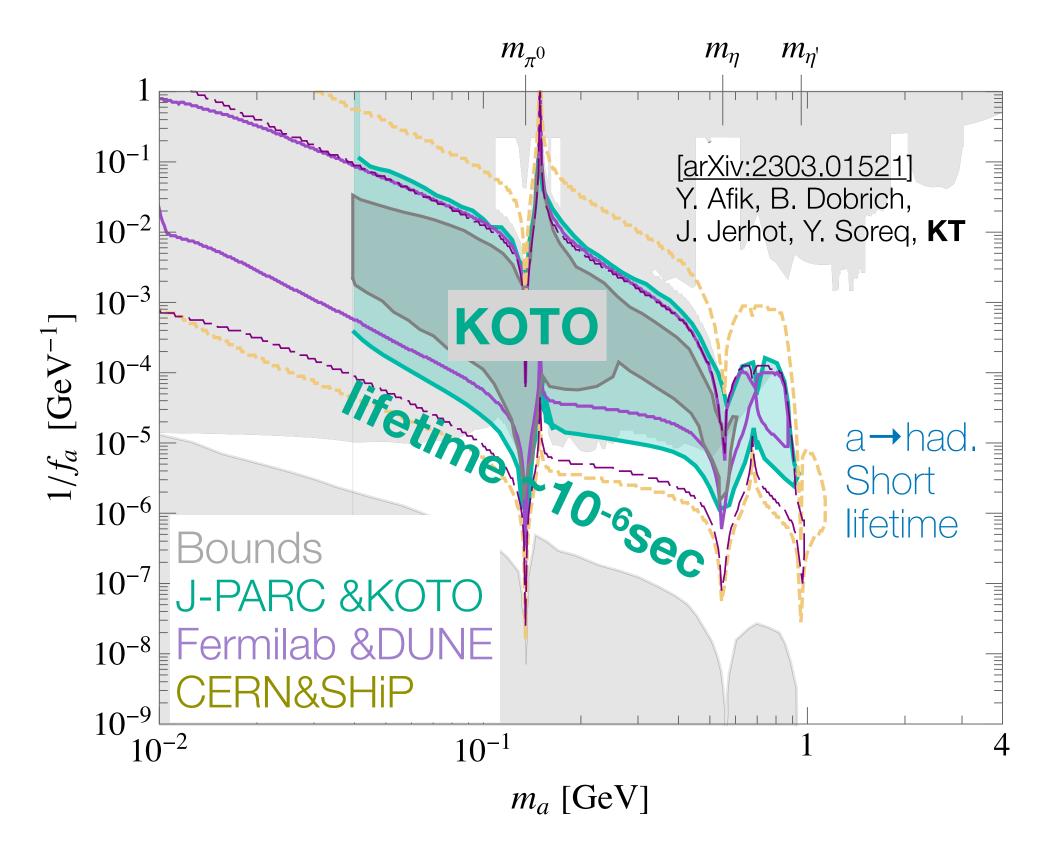


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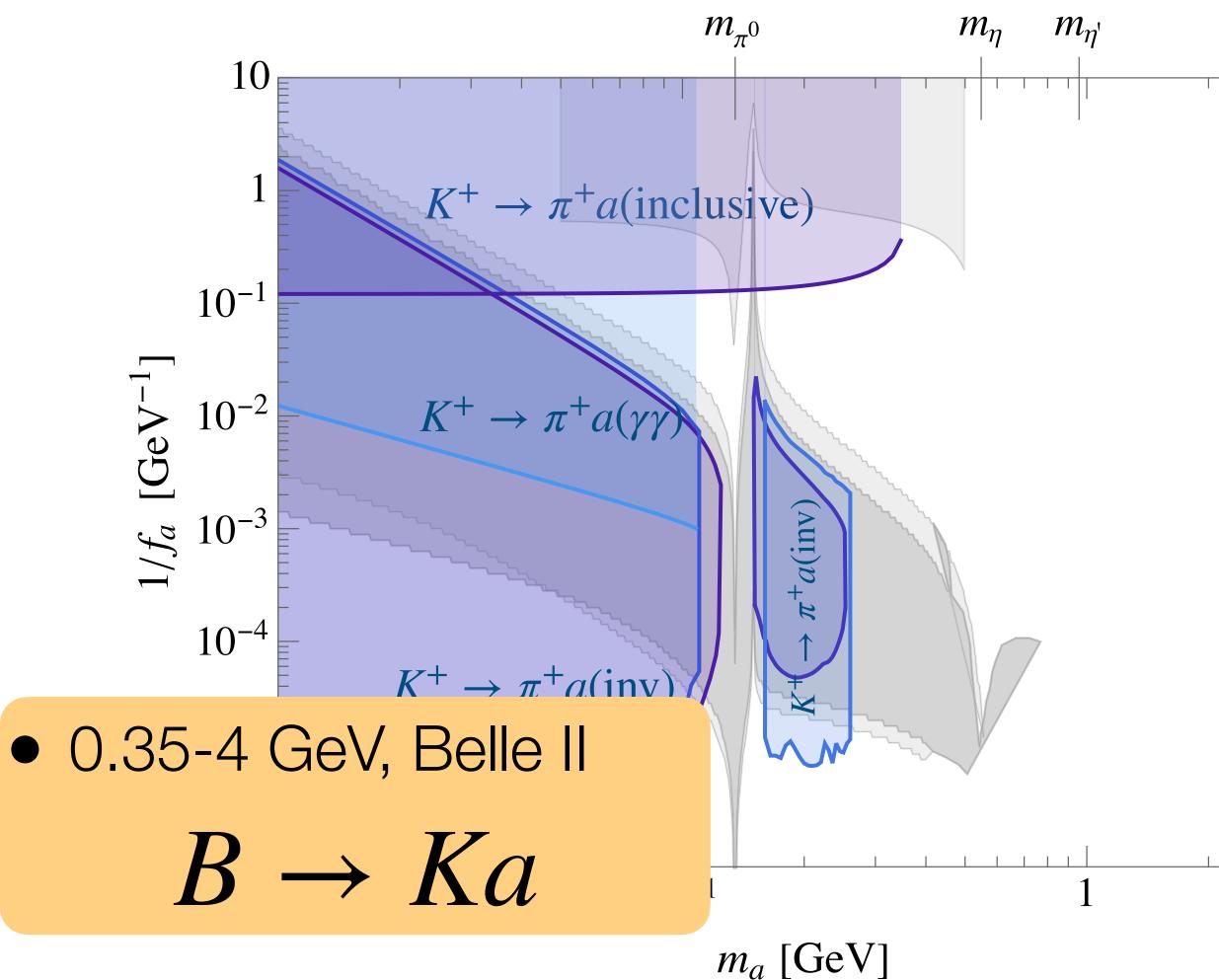


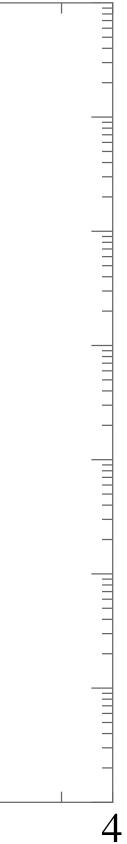
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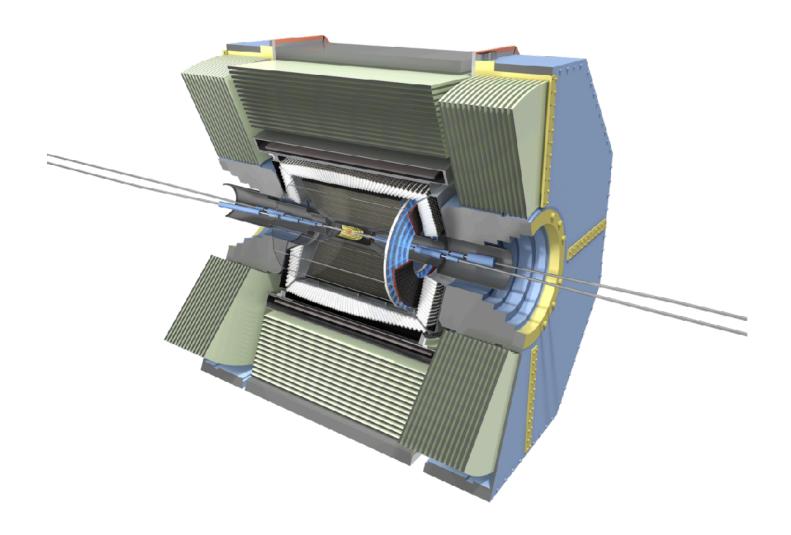
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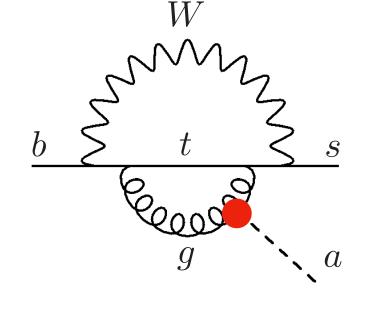


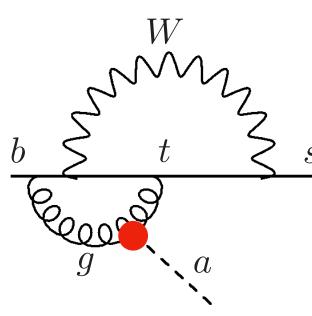


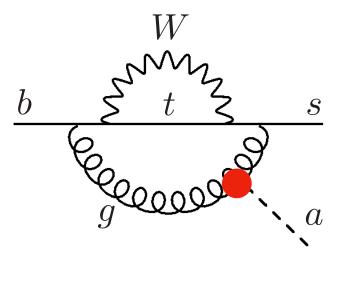
# $B \rightarrow Ka \text{ for } m_a \sim GeV$

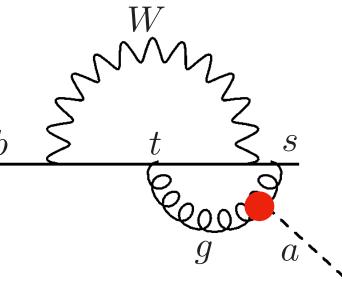
### **Production rate of** $B \rightarrow Ka$

- If the optional aWW coupling exists, there is 1-loop contribution [finite due to GIM, diverge at 2-loop].
- Robust production is from gluon coupling: leading is at 2-loop!





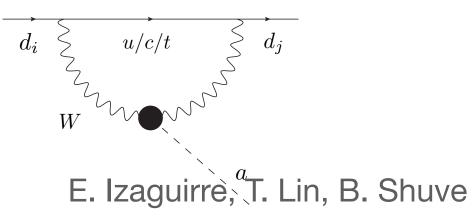




Need 2-loop to generate  $b \rightarrow sa$  from aGG

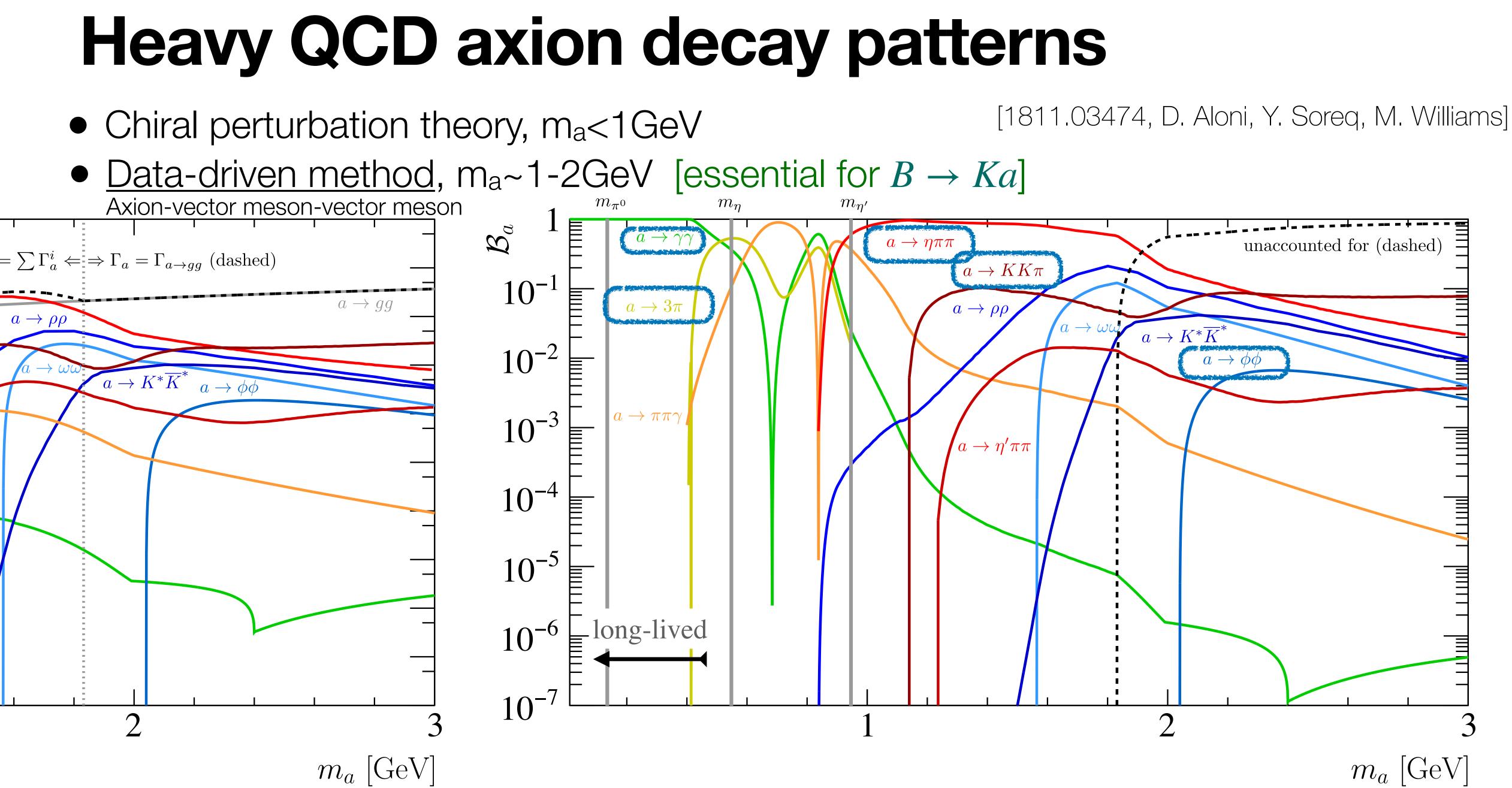
- 1-loop QCD for agg
- 1-loop with W-boson for flavor changing





PRD 104 (**2021**) 055036 S. Chakraborty, M. Kraus, V. Loladze, T. Okui, KT  $\Lambda_{UV} \text{ EFT}(\Lambda_{UV})$  $\mathcal{O}_{gg} = \frac{1}{8\pi} \frac{a}{f_a} G^a_{\mu\nu} \tilde{G}^{a\mu\nu}$  $\mathcal{O}_{qq} = \sum \frac{\partial_{\mu} a}{f} \, \bar{q} \gamma^{\mu} \gamma_5 q \; ,$ Renormalization Group Evolution@2loop  $\mathcal{O}_{bs} = \frac{\partial_{\mu}a}{f_{\sigma}} \,\bar{s}_{\mathrm{L}} \gamma^{\mu} \gamma_{5} b_{\mathrm{L}} + \mathrm{h.c.} \,.$ Matching to weak-scale EFT @2loop  $\mathcal{L}_{bsa} = C_W \frac{\partial_\mu a}{f_a} \bar{s}_{\mathrm{L}} \gamma^\mu \gamma_5 b_{\mathrm{L}}$ mw BR[B $\rightarrow$ Ka]~10<sup>-5</sup> (f<sub>a</sub>/100GeV)<sup>-2</sup>

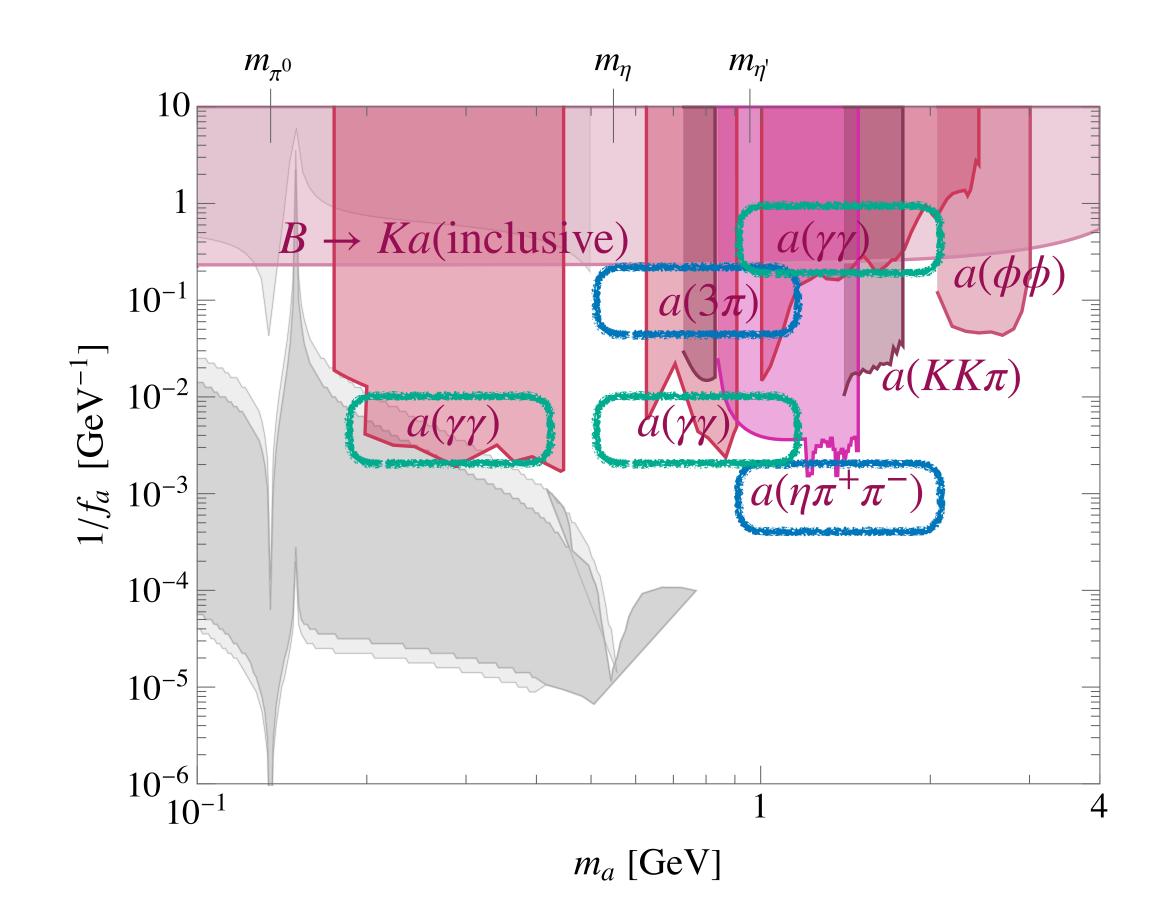




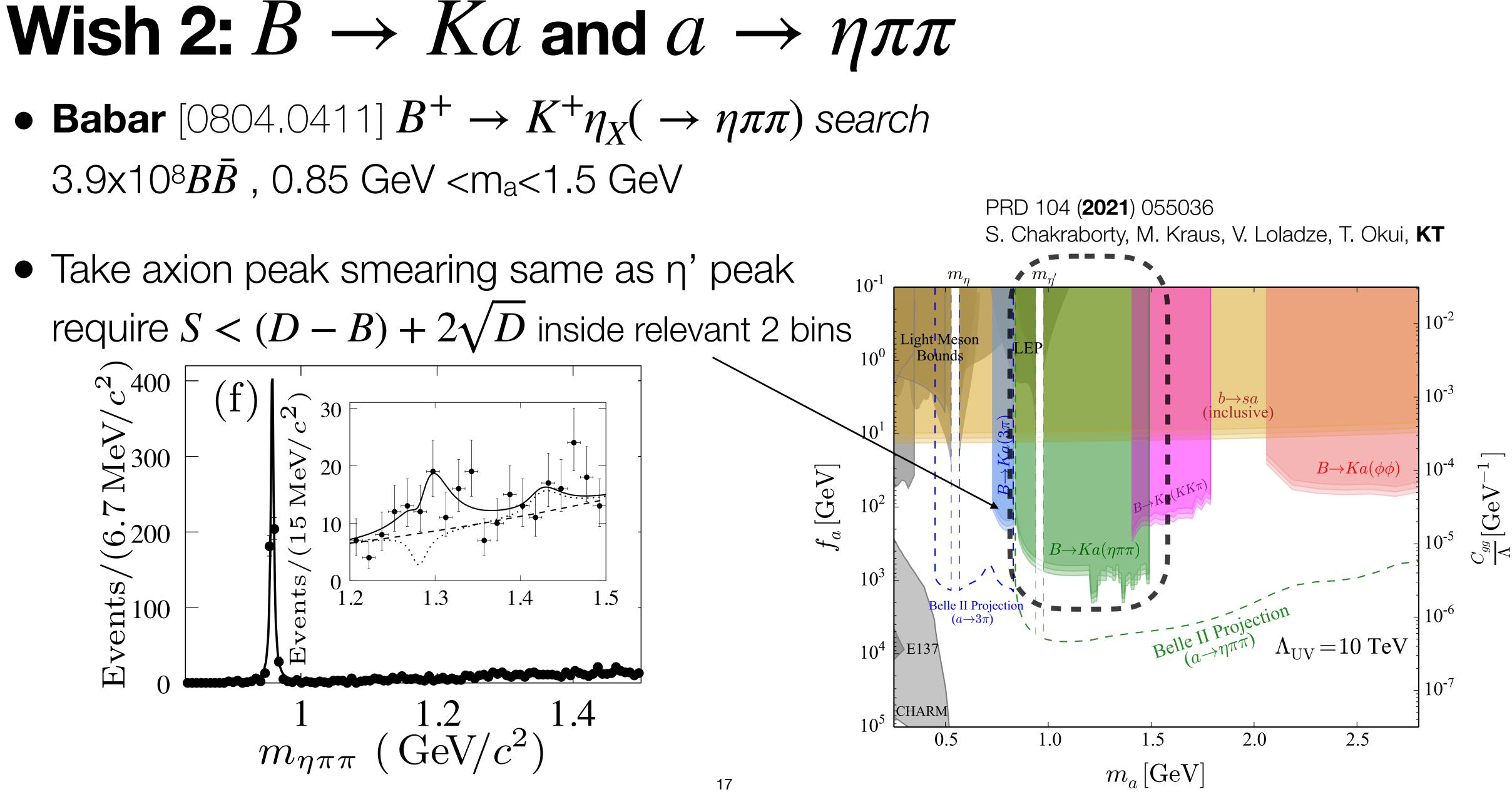
# Heavy QCD axion decay patterns

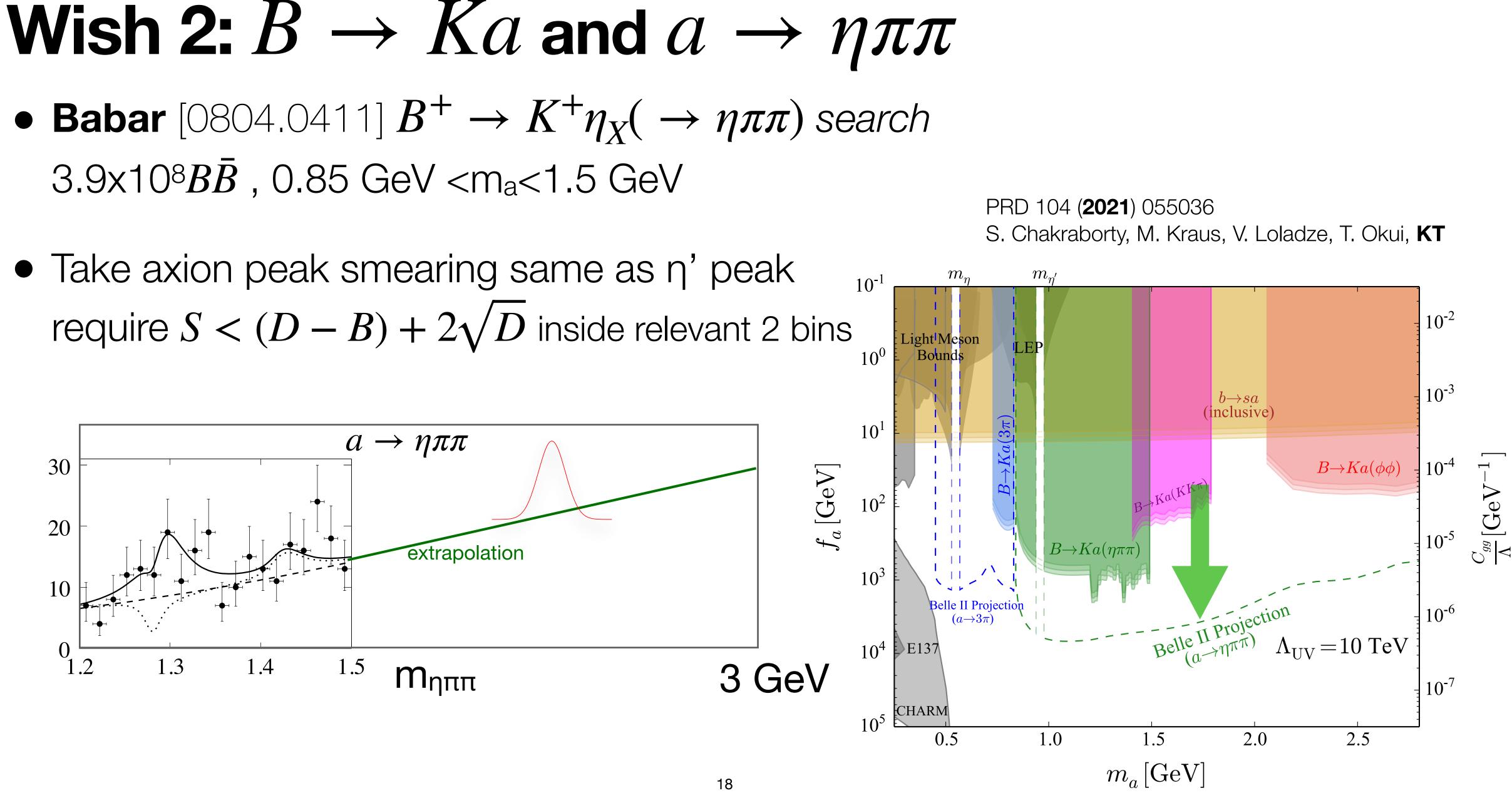
- Chiral perturbation theory, m<sub>a</sub><1GeV</li>
- <u>Data-driven method</u>,  $m_a \sim 1-2 \text{GeV}$  [essential for  $B \rightarrow Ka$ ] Axion-vector meson-vector meson

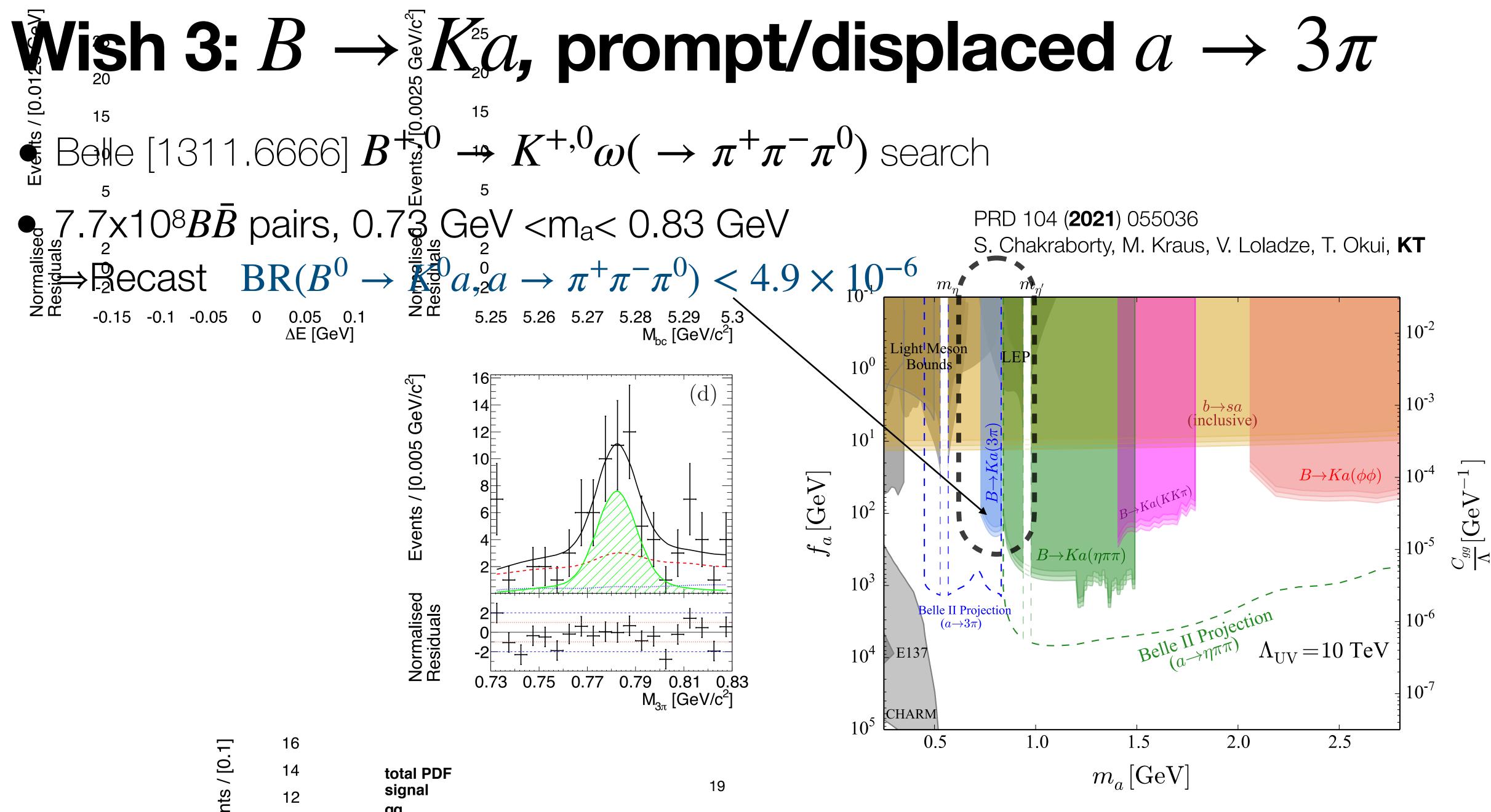
- Recast Old (Babar) and new (Belle II) analyses  $B \to KX(\pi^0 \pi^+ \pi^-, \eta \pi^+ \pi^-, KK\pi, \phi \phi, \gamma \gamma),$
- Need similar/dedicated analyses at Belle II Also new searches.

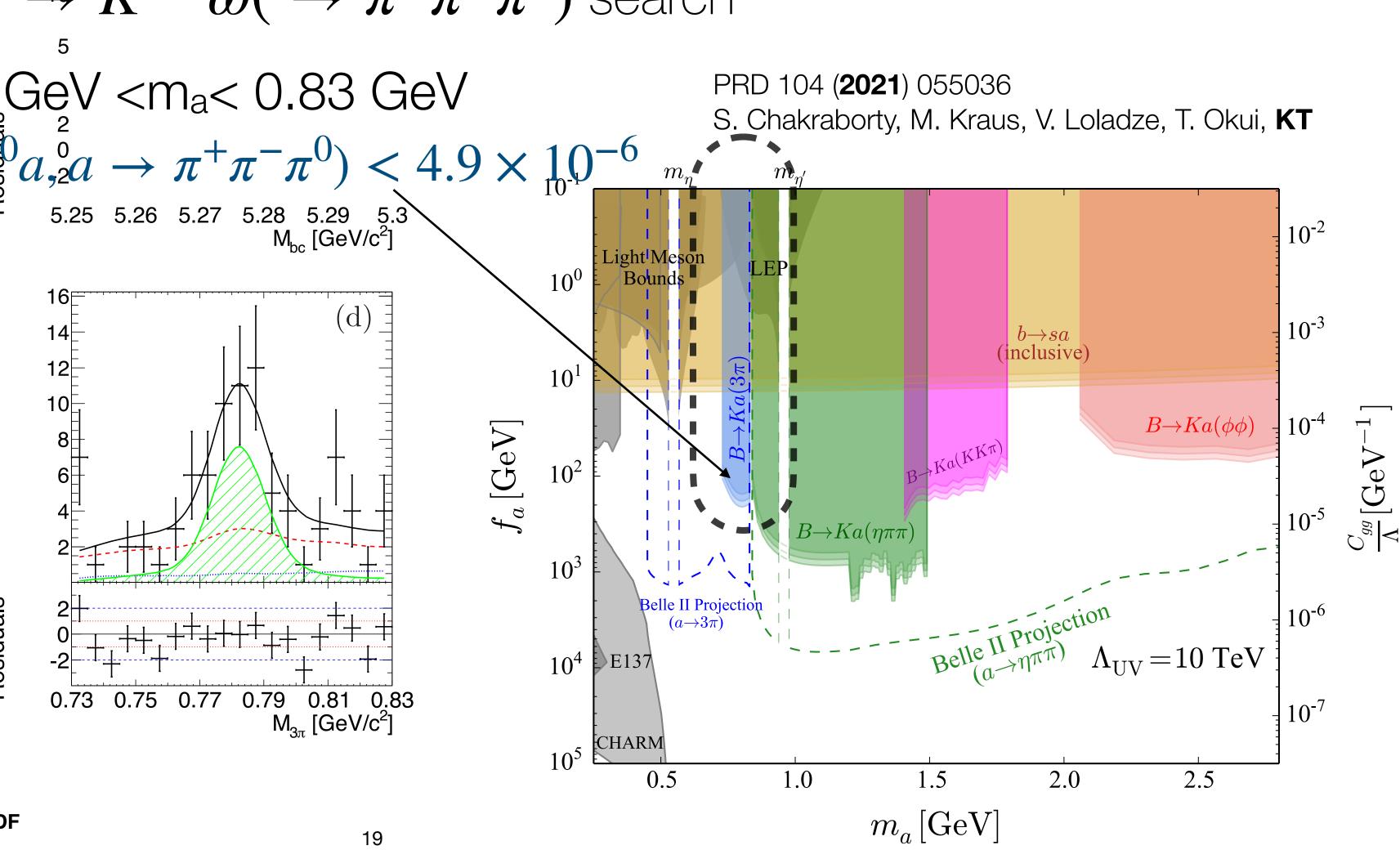


# $3.9 \times 10^8 B\bar{B}$ , 0.85 GeV <m<sub>a</sub><1.5 GeV

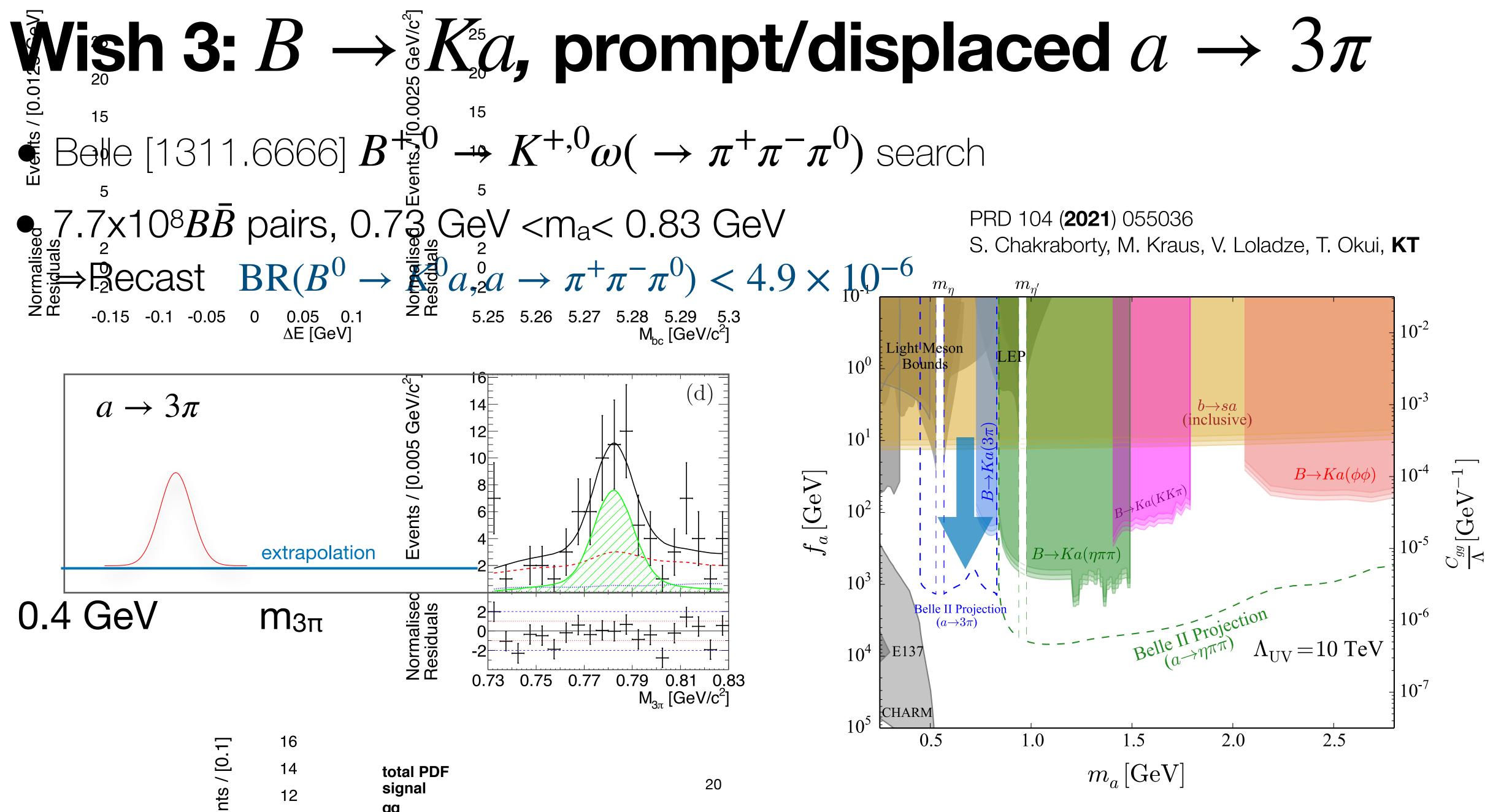






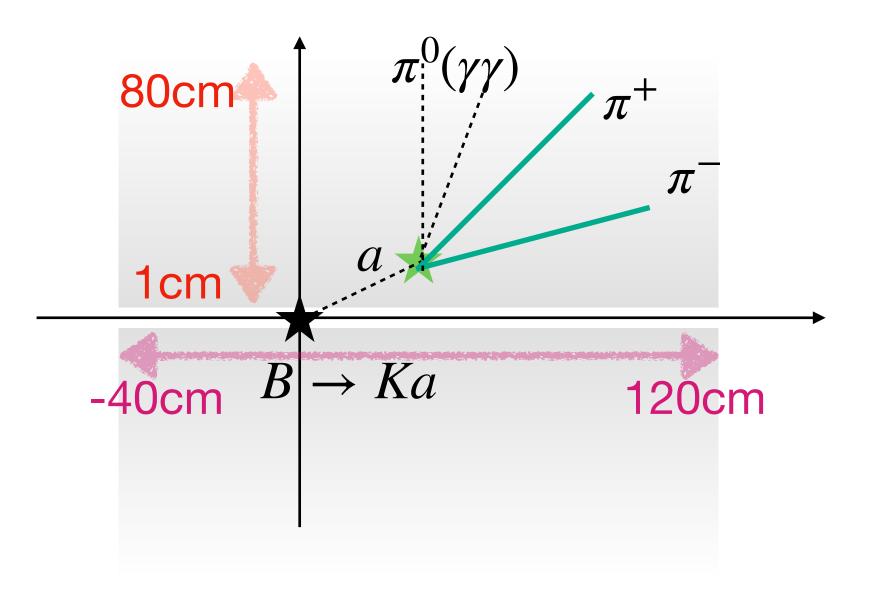


signal aa



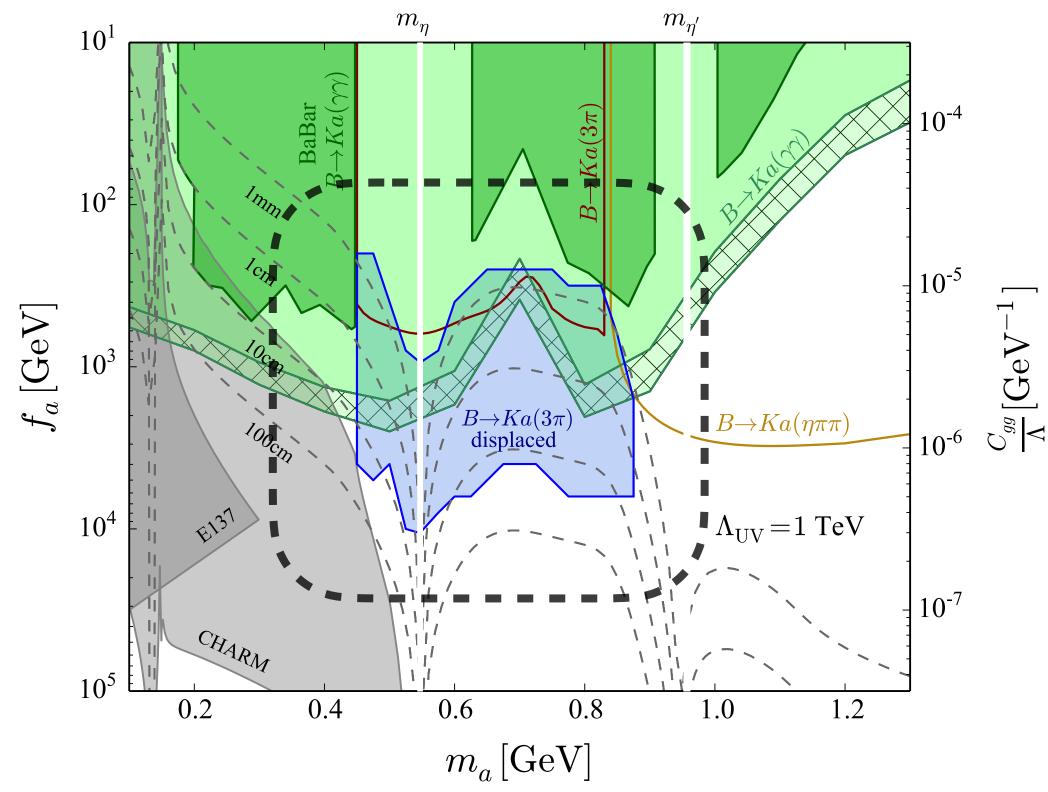
# Wish 3: $B \rightarrow Ka$ , prompt/displaced $a \rightarrow 3\pi$

- Very low background due to >1cm DV.  $B \to KK^*(\to \pi^0 K_I, K_I \to \pi^+\pi^-)$



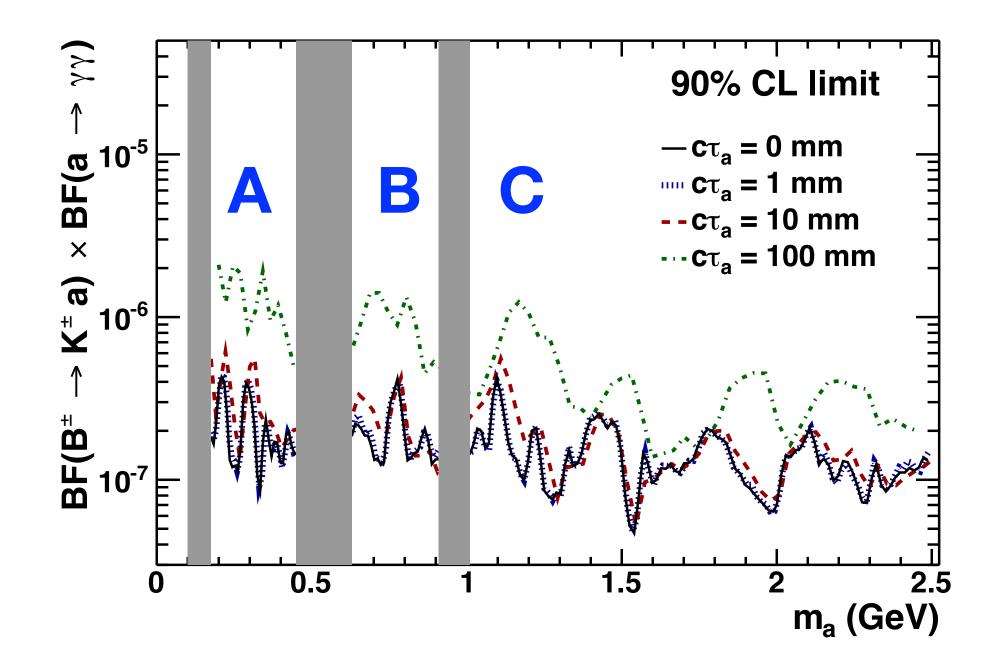
• **Displaced decay** is also possible because 2 charged pions reconstruct vertex.

Phys.Rev.D 105 (2022) L071701 E. Bertholet, S. Chakraborty, V. Loladze, T. Okui, A. Soffer, KT

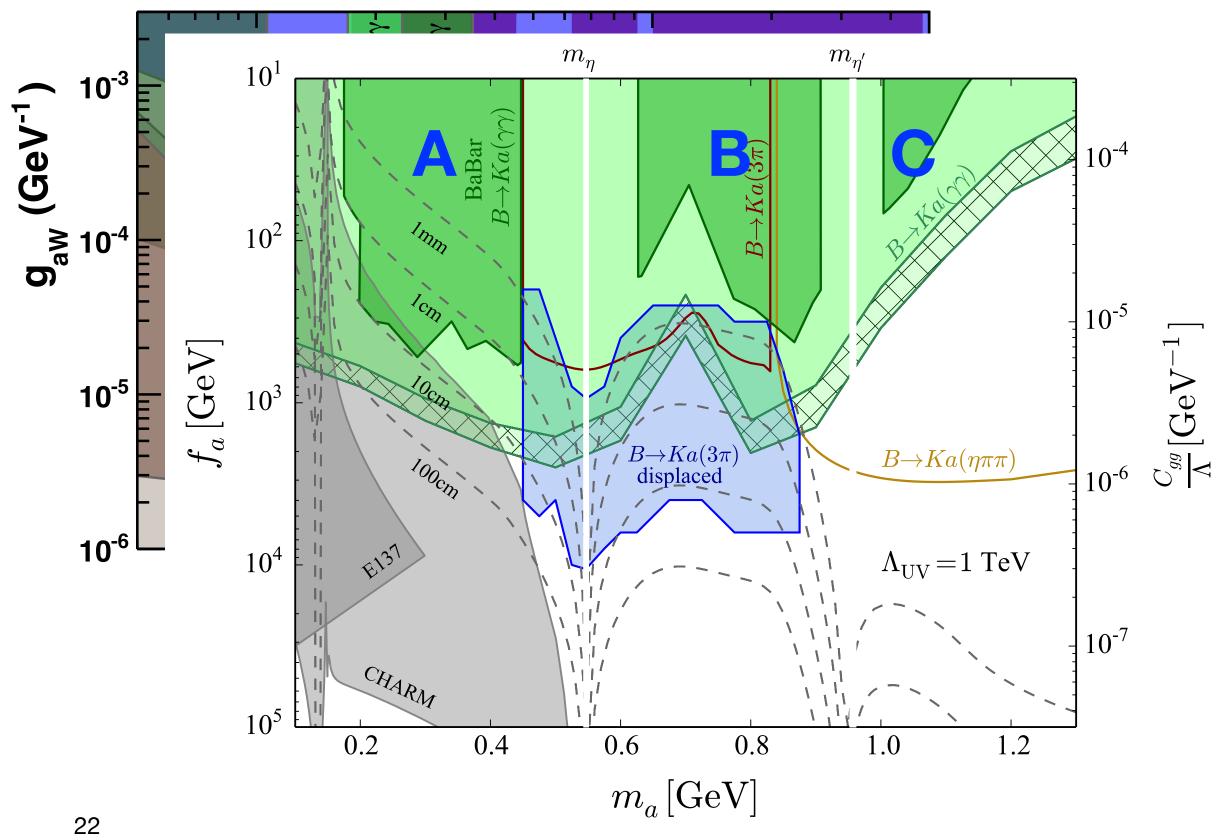


### Wish 4: $B \rightarrow Ka$ and $a \rightarrow \gamma\gamma$

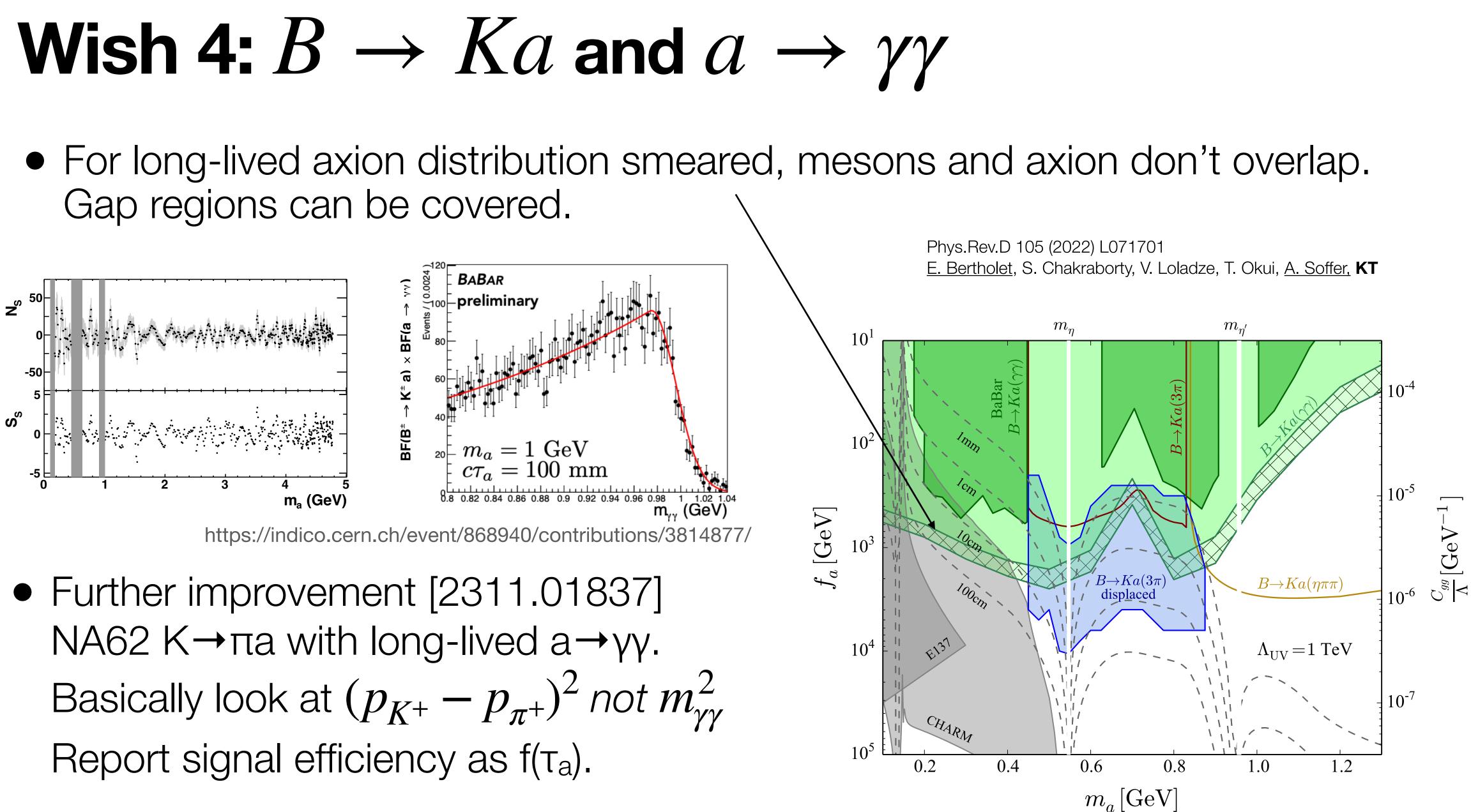
- Babar [2111.01800] (B. Shuve) dedicated ALP search. Reinterpret.
- Same/better analysis can be done at Belle II.



Phys.Rev.D 105 (2022) L071701 E. Bertholet, S. Chakraborty, V. Loladze, T. Okui, A. Soffer, KT

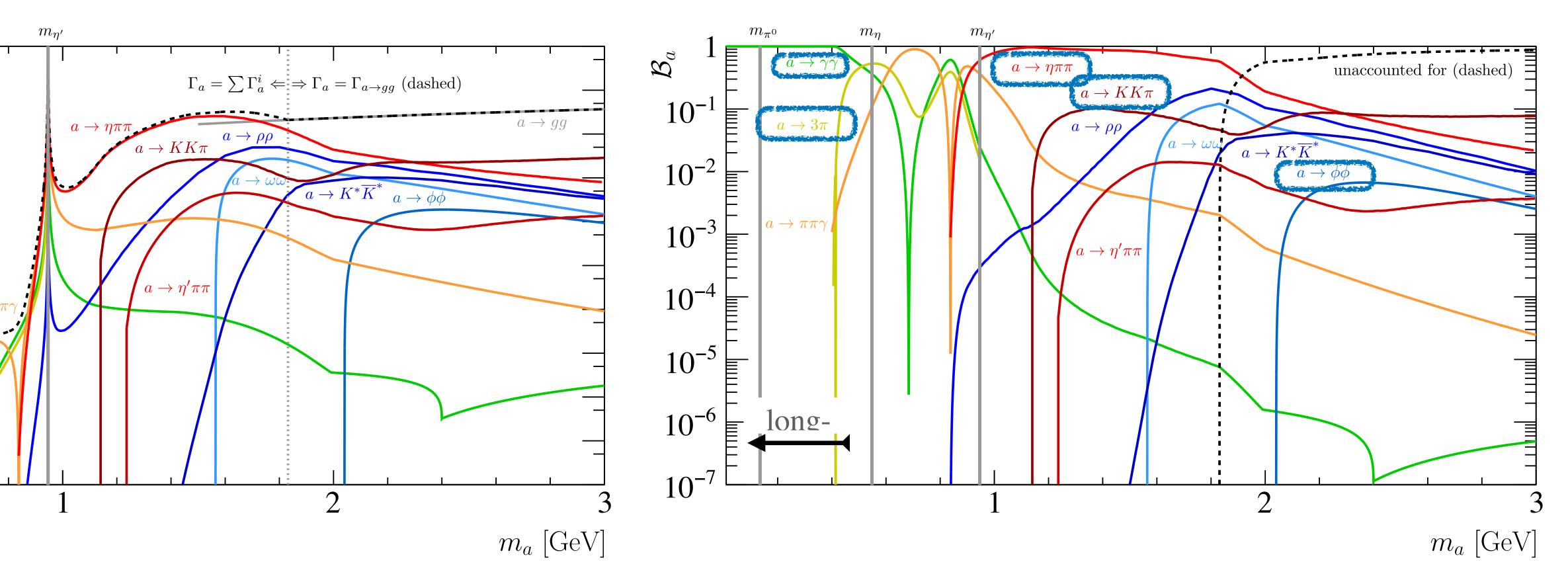


Gap regions can be covered.



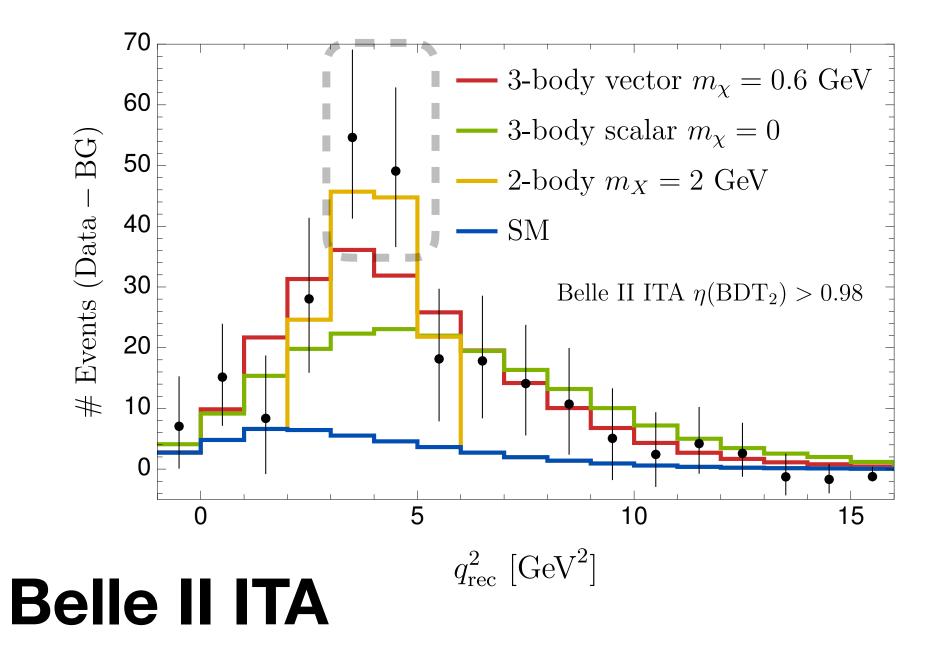
# Wish 5: $B \rightarrow Ka$ , $a \rightarrow$ hadrons

### • Exclusive BR calculation is unstable for ma>2GeV but inclusive mode can be used. Similar to $e^+e^- \rightarrow \Upsilon \rightarrow \gamma a (\rightarrow hadrons)$



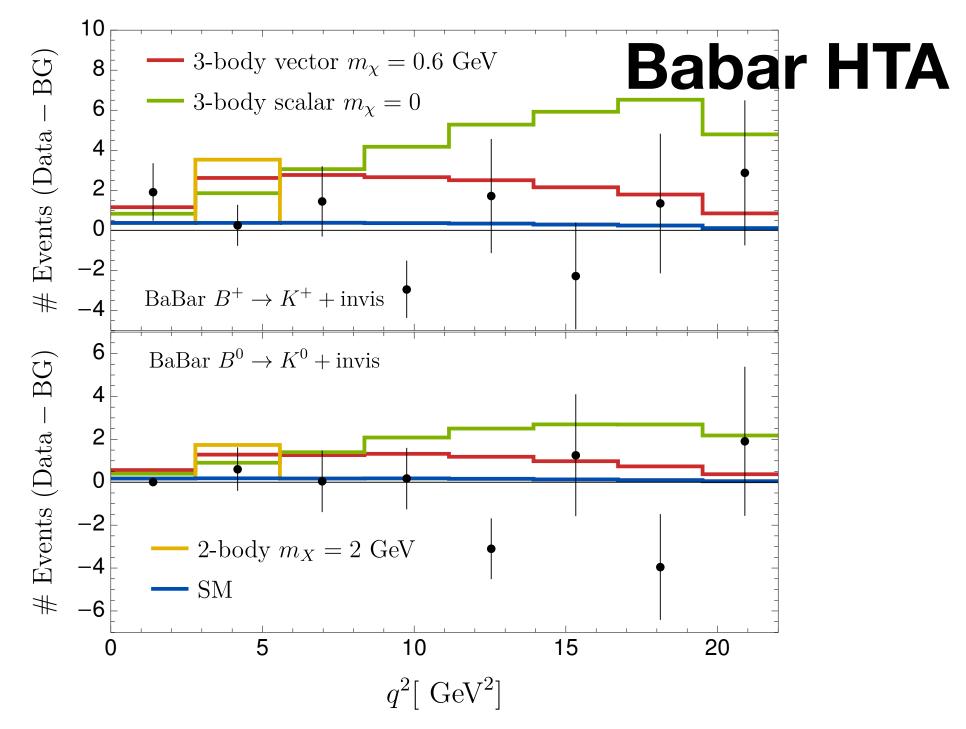
### **Excess in** $B \rightarrow K + inv$ and inclusive tagging

- Latest **Belle II** data shows an excess around m<sub>inv</sub>~2GeV [2311.14647]
- But combining past Babar analysis disfavors 2-body kinematics.



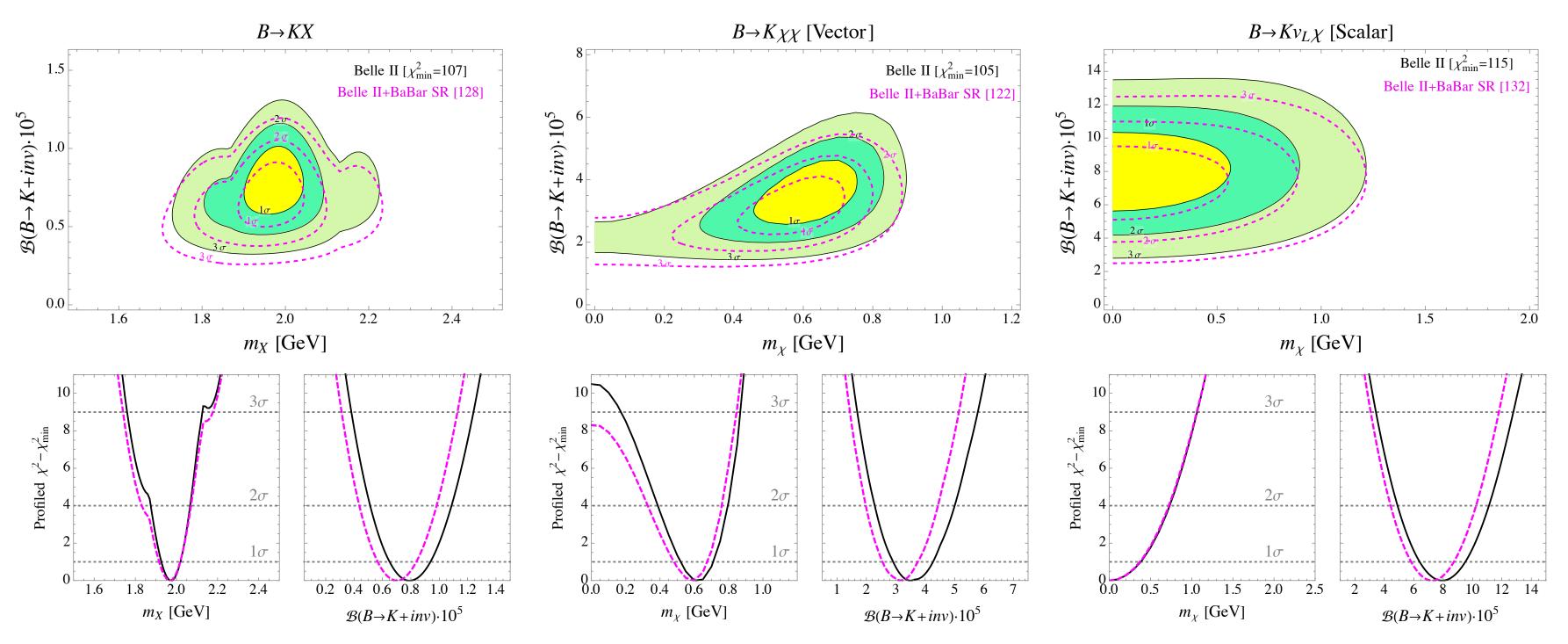
K. Fridell, M. Ghosh, T. Okui, **KT** ('23)

With only this data 3-body decay (vector current) and 2-body are reasonable.



### **Excess in** $B \rightarrow K + inv$ and inclusive tagging

- Latest Belle II data shows an excess around minv~2GeV
- But combining past Babar analysis disfavor 2-body kinematics.



With only this data 3-body decay (vector current) and 2-body are reasonable.

| $\chi^2_{\rm min} - 100$     | 2b   | V    | V'   | S    | Т    | SM    |
|------------------------------|------|------|------|------|------|-------|
| Belle II                     | 6.8  | 15.2 | 4.7  | 15.1 | 11.9 | 44.6  |
| + BaBar SR                   | 27.6 | 30.4 | 22.1 | 31.8 | 29.8 | 61.0  |
| $+ \text{ BaBar } s_B < 0.8$ | 73.3 | 78.8 | 72.9 | 90.2 | 86.9 | 106.7 |

K. Fridell, M. Ghosh, T. Okui, **KT** ('23)

### Summary

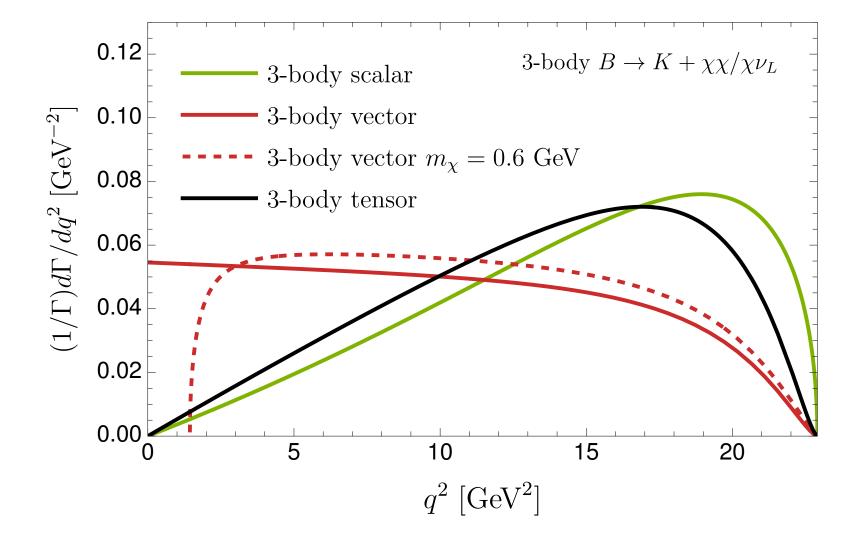
- Wishlist
- 1.  $e^+e^- \rightarrow \Upsilon \rightarrow \gamma a$ ,  $a \rightarrow hadrons$  [previous work: Babar recast] 2.  $B \rightarrow Ka, a \rightarrow \eta \pi \pi$  [Babar recast] 3.  $B \to Ka, a \to \pi^+ \pi^- \pi^0$  prompt [Belle recast] / displaced [new] 4.  $B \rightarrow Ka$ ,  $a \rightarrow \gamma\gamma$  [Babar dedicated analysis, possible to improve further] 5.  $B \rightarrow Ka, a \rightarrow hadrons [new]$ • CP even scalar  $\rightarrow \pi\pi$ , KK would be interesting too.

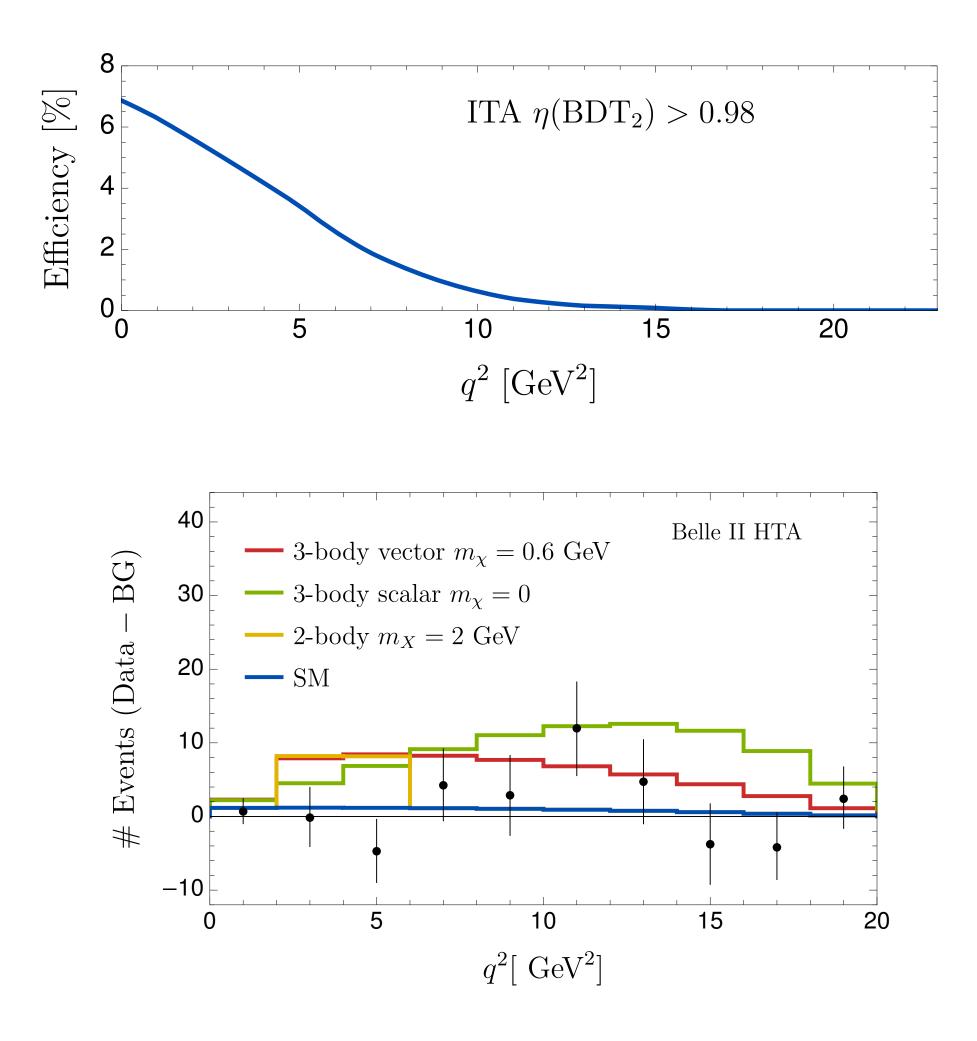
### Unique possibilities to probe Heavy QCD axion @ Belle II based on hadrons.

Thank you!

Backup

### **Excess in** $B \rightarrow K + inv$ and inclusive tagging

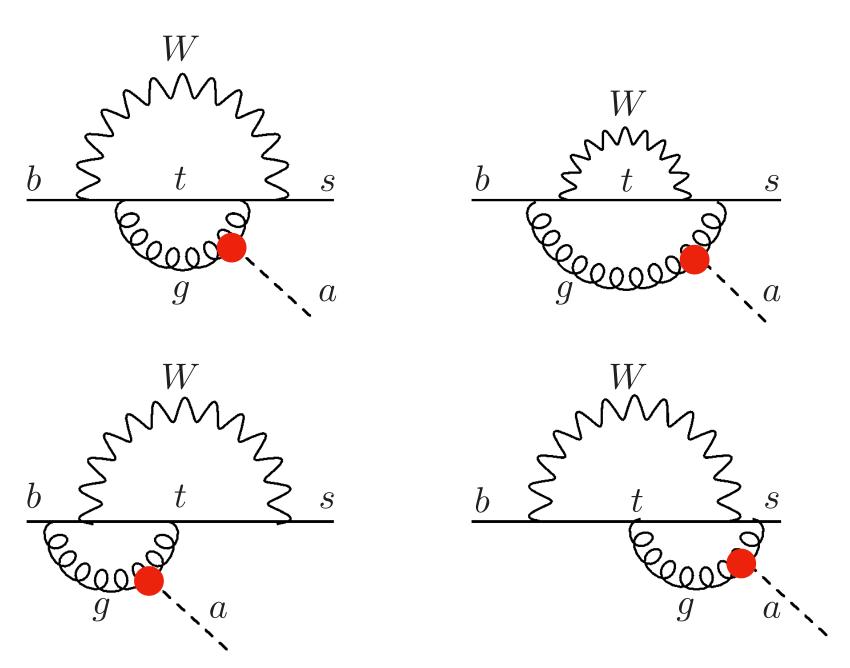






# **Two-loop calculation from the EFT**

\*two-loop is the leading order!



Need 2-loop to generate  $b \rightarrow sa$  from aGG

- 1-loop QCD for aqq
- 1-loop with W-boson for flavor changing

Previous estimate: Ref[Yotam, etc]

#### <sup>Kohsaku</sup> For more detail, KEK Theory seminar by Chakraborty, Mar 15 (16) 9am EST (11 am JST)

 $\Lambda_{UV}$  EFT( $\Lambda_{UV}$ )

$$\begin{split} \mathcal{O}_{gg} &= \frac{1}{8\pi} \frac{a}{f_a} G^a_{\mu\nu} \tilde{G}^{a\mu\nu} ,\\ \mathcal{O}_{qq} &= \sum_q \frac{\partial_\mu a}{f_a} \,\bar{q} \gamma^\mu \gamma_5 q ,\\ \mathcal{O}_{bs} &= \frac{\partial_\mu a}{f_a} \,\bar{s}_{\mathrm{L}} \gamma^\mu \gamma_5 b_{\mathrm{L}} + \mathrm{h.c.} \end{split}$$

Renormalization Group Evolution

> Matching to weak-scale EFT @2loop MW  $\mathcal{L}_{bsa} = C_W \frac{\partial_{\mu}a}{f_a} \bar{s}_{\mathrm{L}} \gamma^{\mu} \gamma_5 b_{\mathrm{L}}$

RGE is essential  $C_W \sim \frac{\alpha_w}{4\pi} \left(\frac{\alpha_s}{4\pi}\right)^2 V_{tb} V_{ts}^* \frac{m_t^2}{m_W^2} \ln[\Lambda_{\rm UV}^2/M_W^2]$ 

Two boundary conditions at  $\Lambda_{UV}$ 

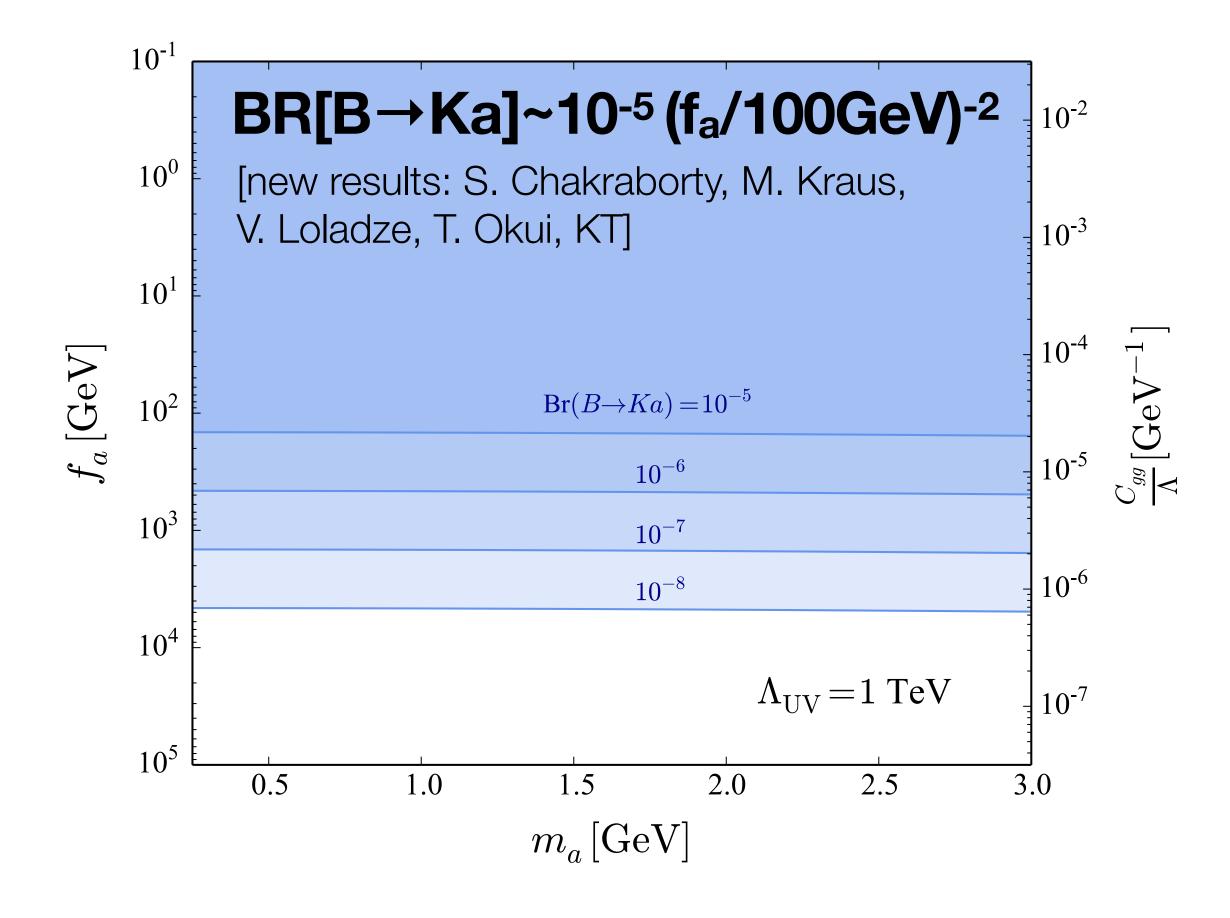
- 1. Size of  $\Lambda_{UV}$  (>  $4\pi f_a$ )  $\rightarrow$  require  $\Lambda_{UV}$  > TeV
- 2. Size of  $C_{qq}(\Lambda_{UV}), C_{bs}(\Lambda_{UV}) \rightarrow \text{fix as big as the counter term}$



### **Axion Production Rate from** $B \rightarrow Ka$

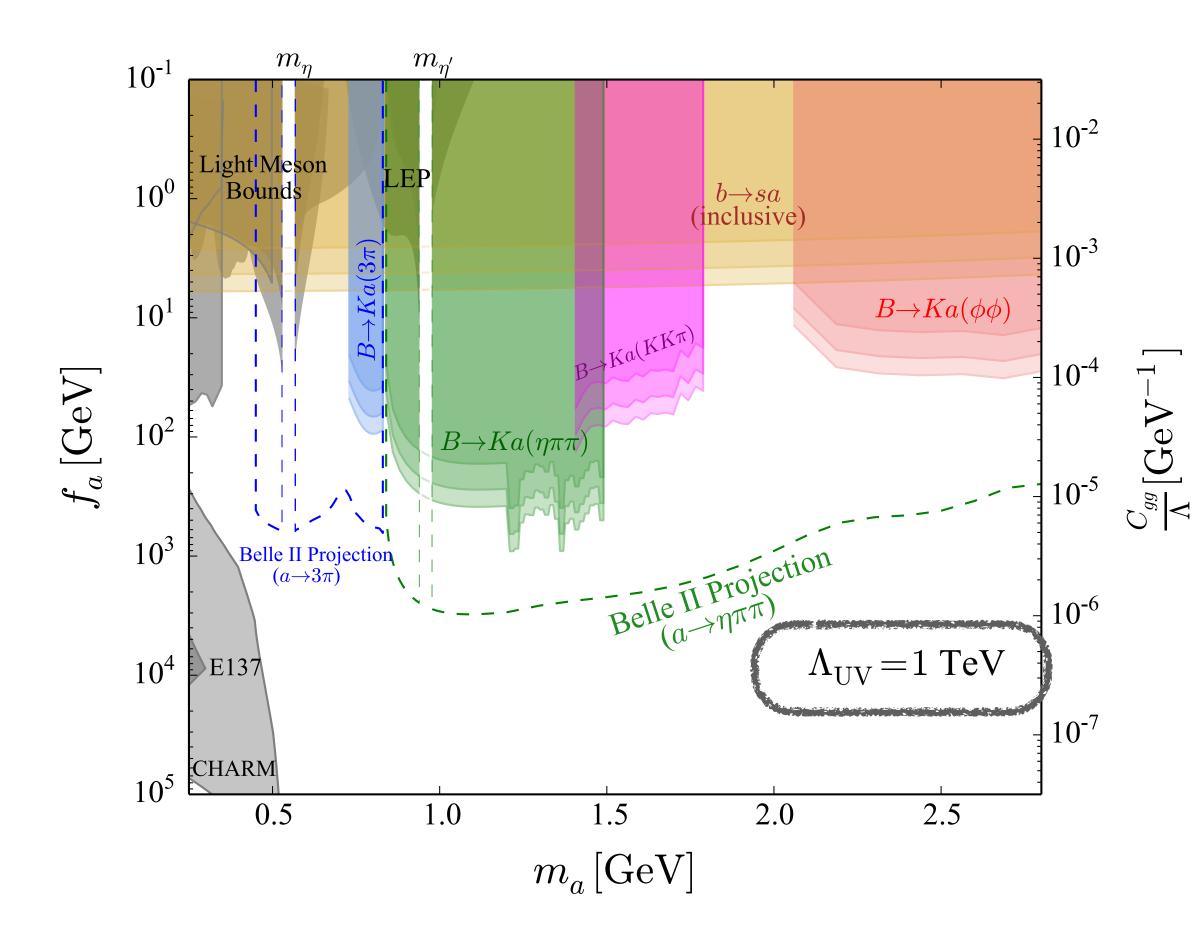
$$\Gamma_{B\to Ka} = |C_W|^2 \frac{m_B^3}{64\pi f_a^2} \left(1 - \frac{m_K^2}{m_B^2}\right)^2 \lambda_{Ka} \left[f_0(m_a^2)\right]^2$$
our result
• Form factor
$$f_0(m_a^2) = \frac{0.330}{1 - m_a^2/37.5 \text{ GeV}^2}$$
• Phase space
$$\lambda_{Ka} = \left[\left(1 - \frac{(m_K + m_a)^2}{m_B^2}\right)\left(1 - \frac{(m_K - m_a)^2}{m_B^2}\right)\right]^{\frac{1}{2}}$$

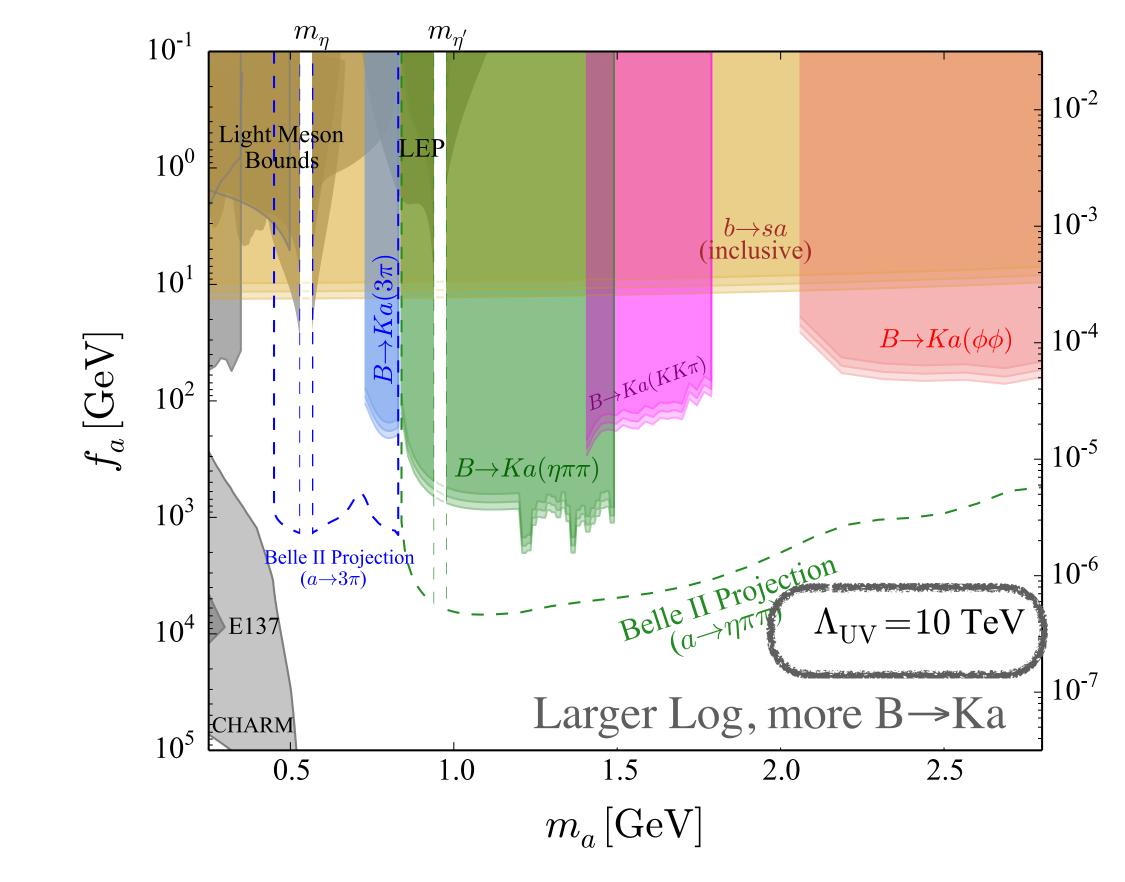
Two boundary conditions 1.  $\Lambda_{UV}$  1TeV. Larger  $\Lambda_{UV}$ , more BR[B->Ka] 2. Size of  $C_{qq}(\Lambda_{UV}), C_{bs}(\Lambda_{UV}) \rightarrow$  here set to 0





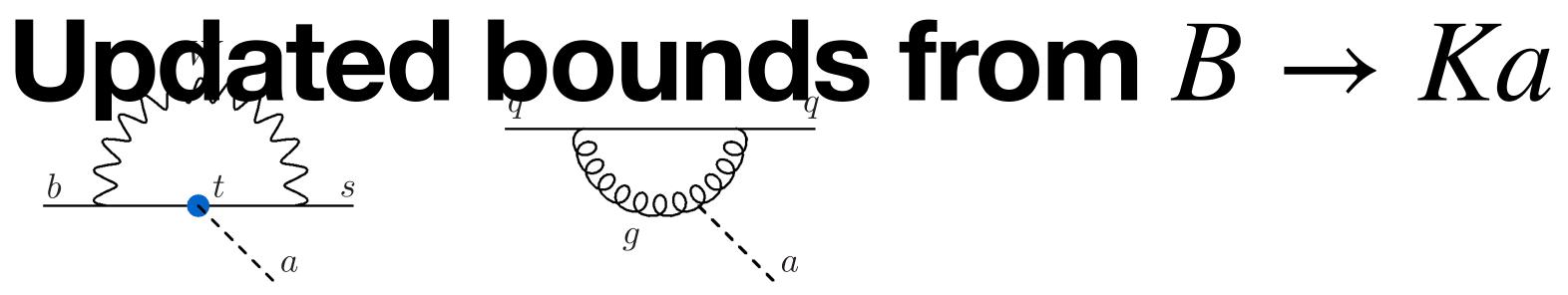
### **Updated bounds from** $B \rightarrow Ka$





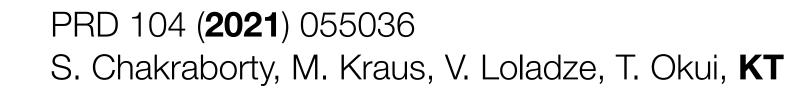


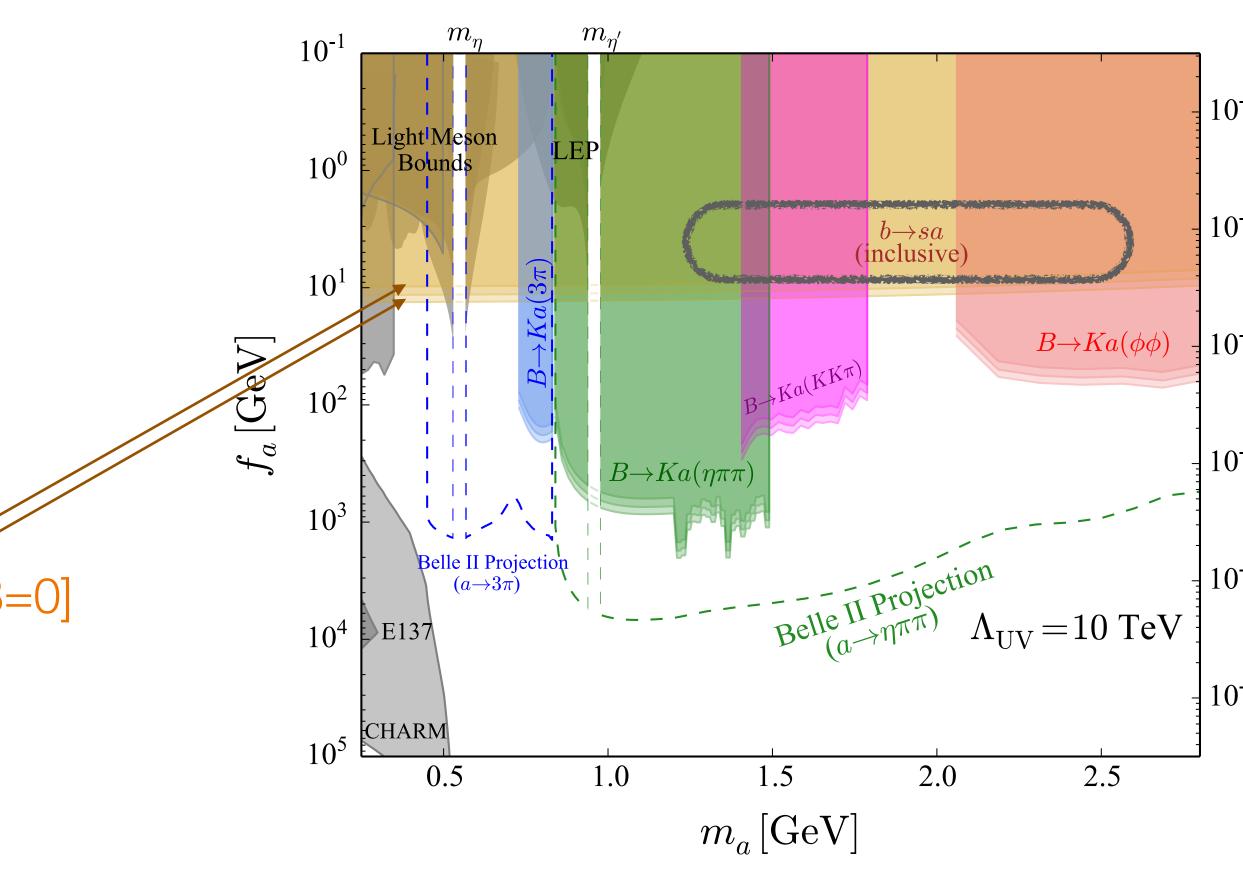


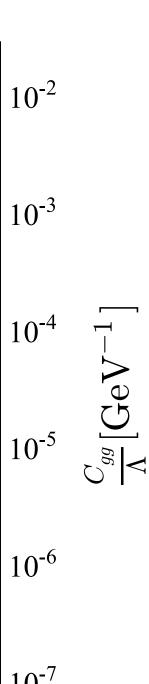


$$BR(b \to sa) \simeq \frac{1}{\Gamma_B f_a^2} \frac{1}{32\pi m_B}$$

 Change UV boundary conditions  $-3 < A, B < 3 \Rightarrow 3$  [ines] optimistic, pesmistic, A=B=0]  $C_{qq}(\Lambda_{\rm UV}) \equiv AC_F\left(\frac{\alpha_s}{4\pi}\right)^2$  $\langle 4\pi \rangle$  $C_{bs}(\Lambda_{\rm UV}) \equiv BC_F \left(\frac{\alpha_s}{4\pi}\right)^2 \frac{\alpha_w}{4\pi} \sum_{k} V_{ik} V_{kj}^* \xi_k$ 

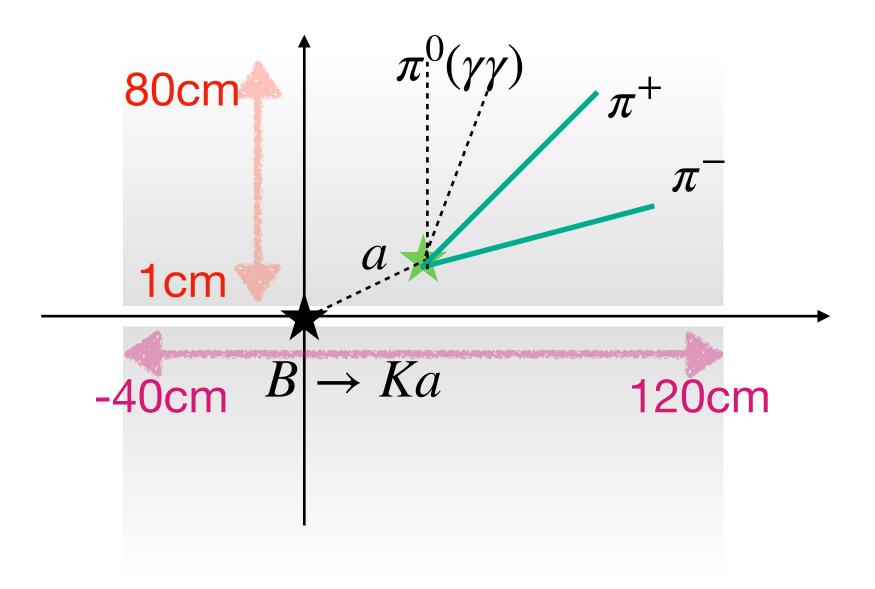


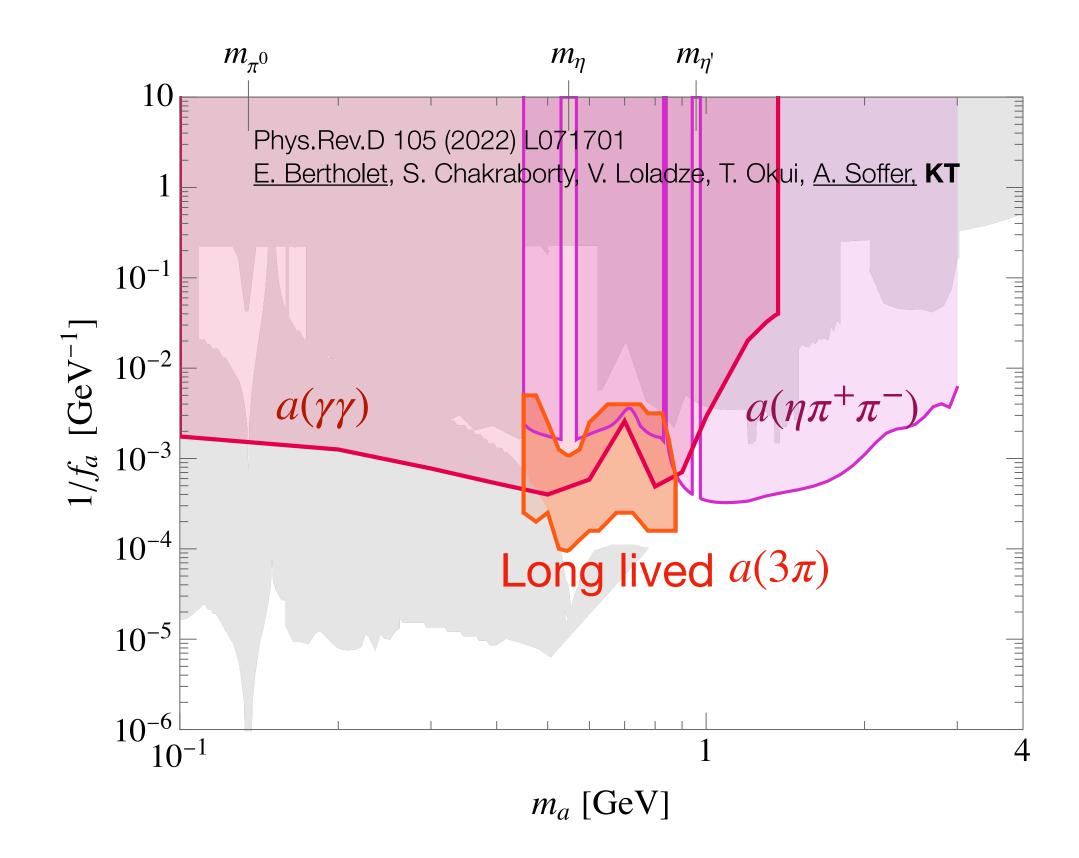






### More on B decays at Belle II for prompt decays and displaced decay (nanosecond lifetime)





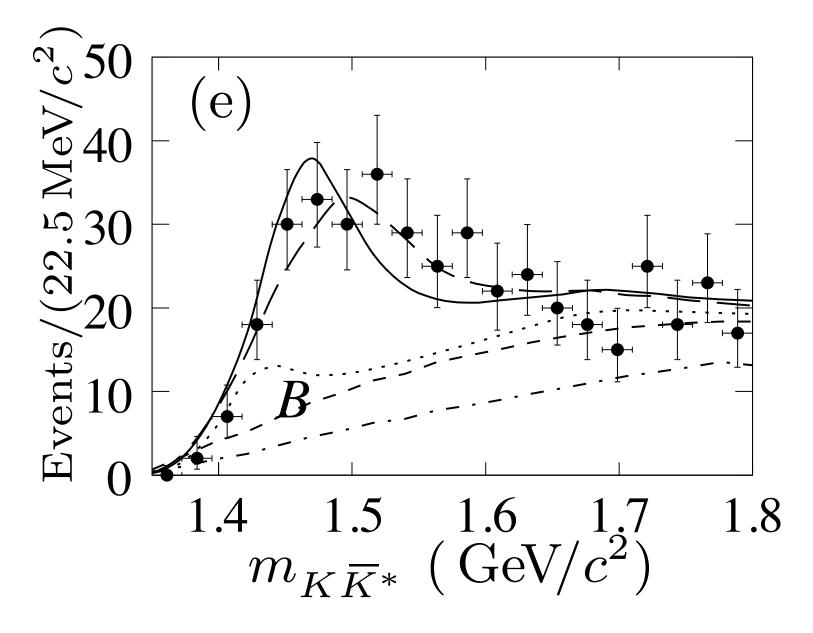


### **Other channels**

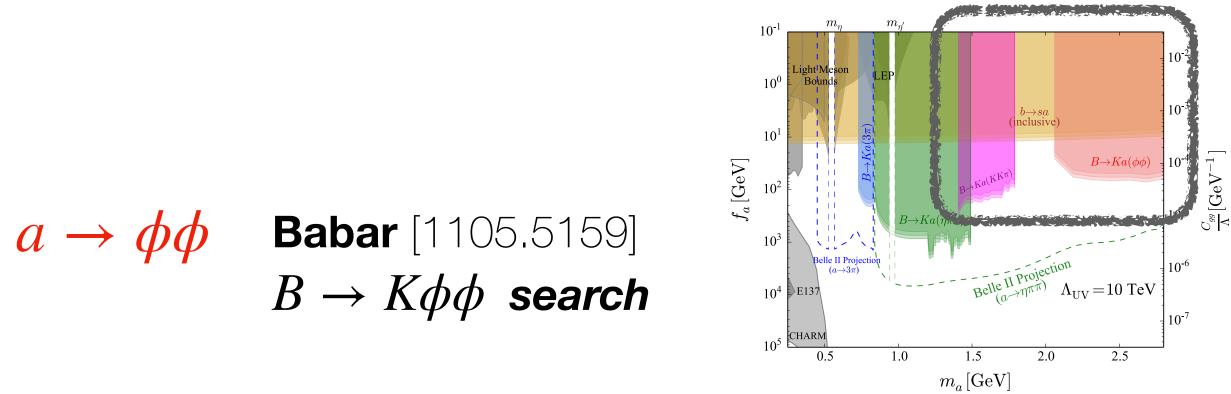
 $a \rightarrow KK\pi$  Babar [0804.0411]  $B^+ \to K^+ \eta_X (\to KK^* \to KK\pi)$  search

- 3.9x10<sup>8</sup>BB

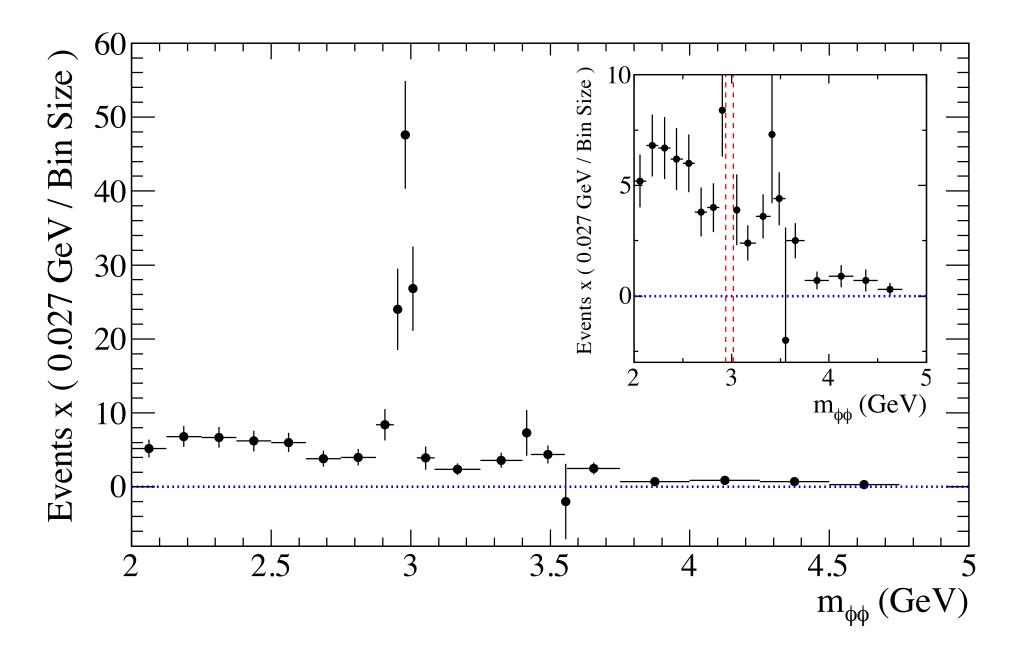
   1.4 GeV <ma < 1.8 GeV</li>
- take 2 bins around the axion mass (45MeV) take 1 bin around the axion mass (125MeV)  $S < (D - B) + 2\sqrt{D}$ require  $2\sqrt{D}$  inside relevant bins



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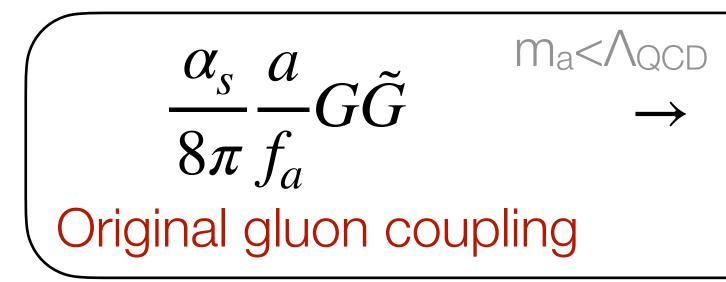
• 4.6x10<sup>8</sup> $B\bar{B}$ , 2 GeV <m<sub>a</sub><3 GeV







### 1. Photon coupling is probed in most searches.



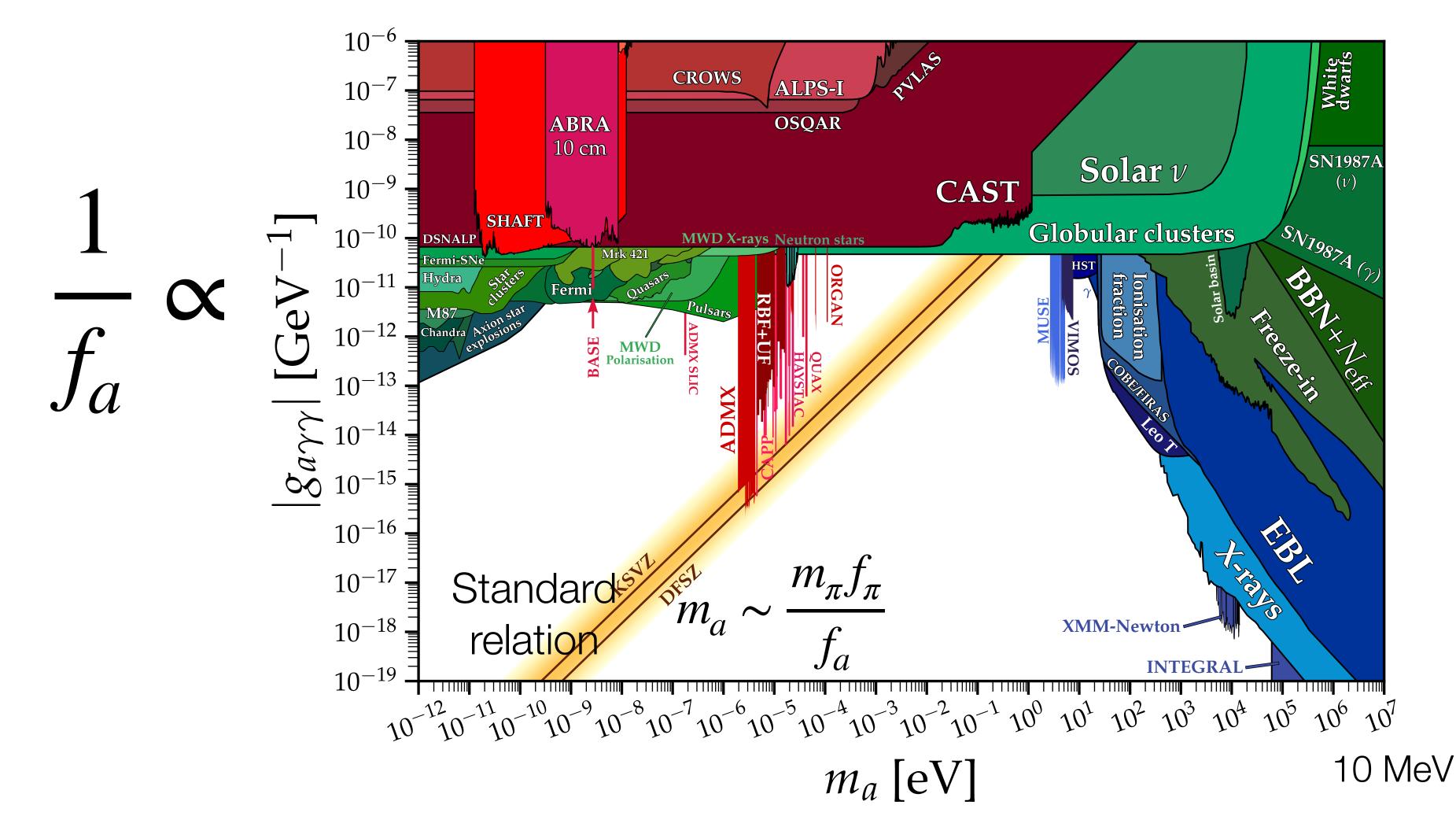
2. Axion dark matter is the most popular. ∧<sub>CP</sub>~∧<sub>DM</sub> Often assume it's 100% DM and extremely light.

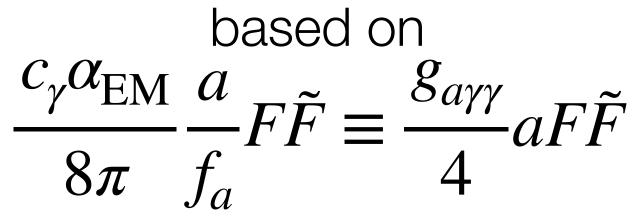
#### Mass m<sub>a</sub>~10<sup>-4</sup>-10<sup>-6</sup> eV

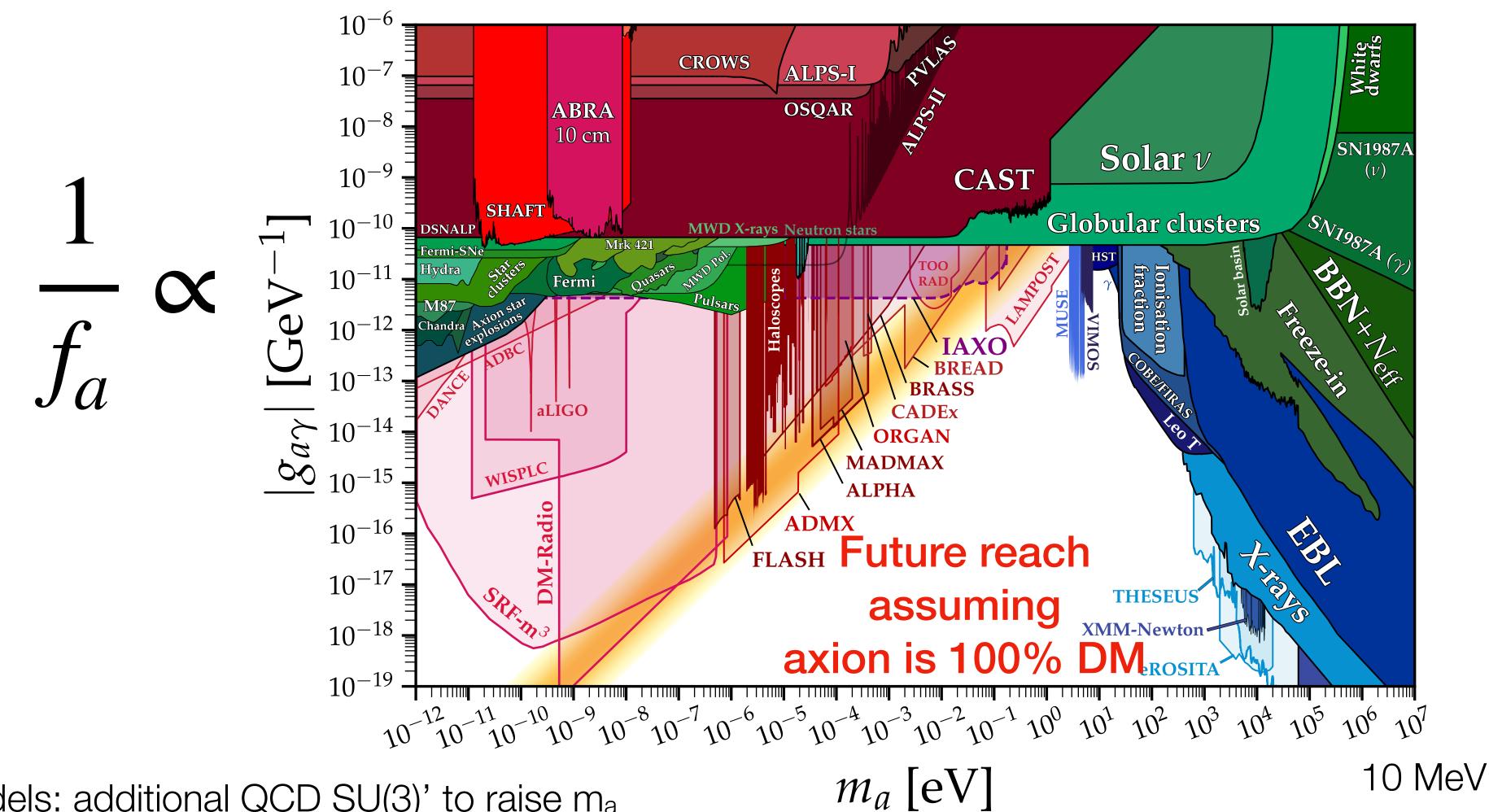
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$$\frac{c_{\gamma}\alpha}{8\pi}\frac{a}{f_a}F\tilde{F} + (a\text{-hadron couplings})$$
**Photon coupling**

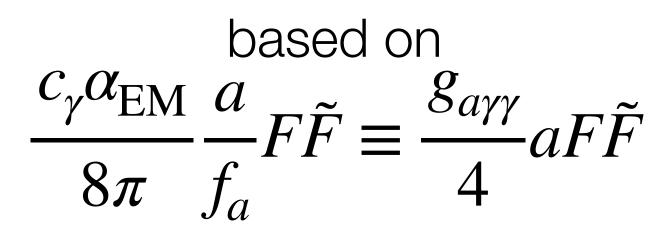
Hadronic coupling is often ignored.

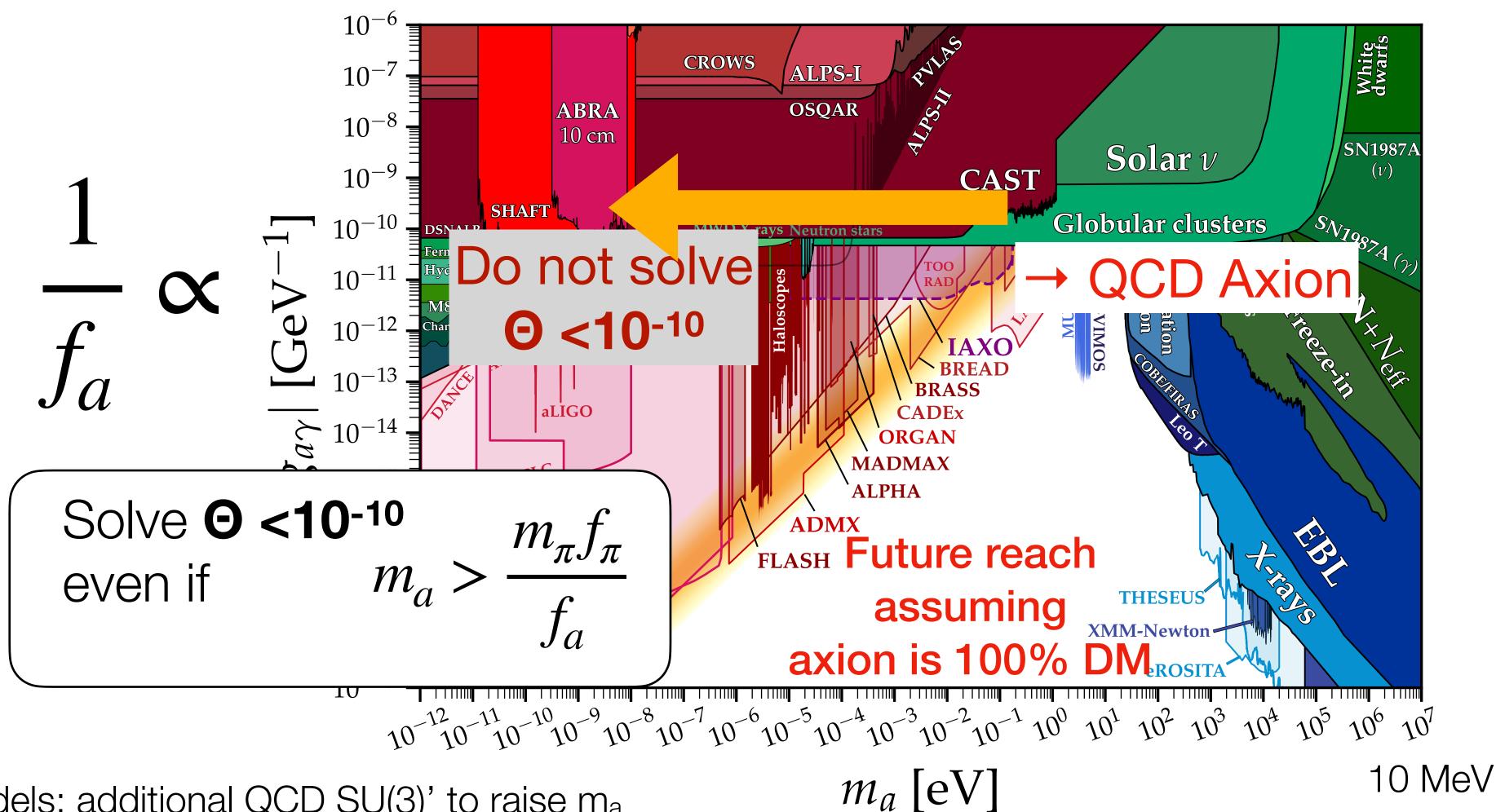




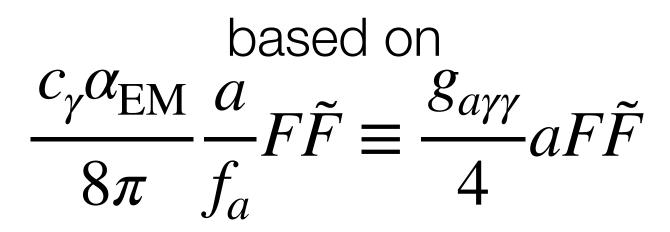


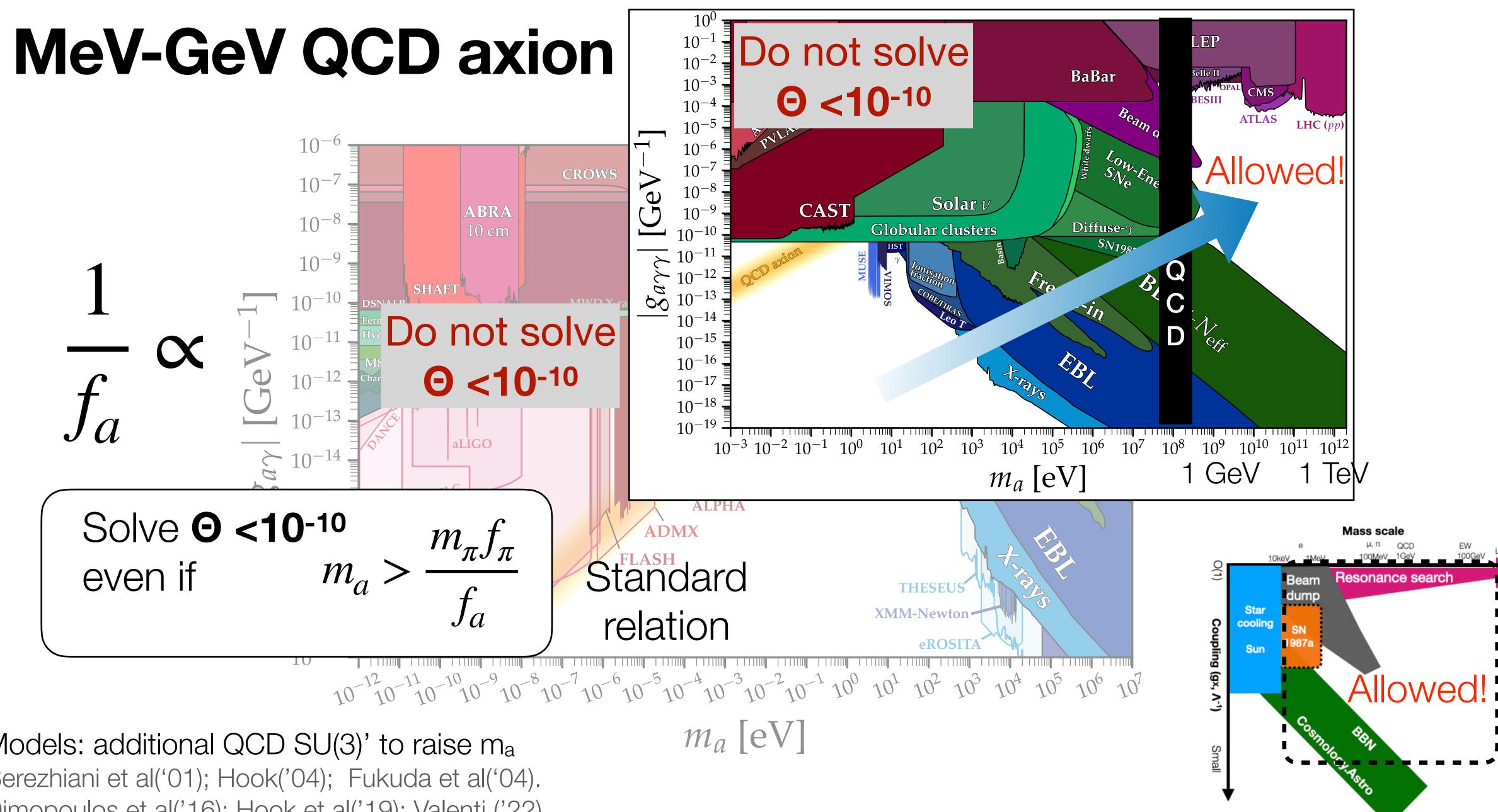
Models: additional QCD SU(3)' to raise m<sub>a</sub> Berezhiani et al('01); Hook('04); Fukuda et al('04). Dimopoulos et al('16); Hook et al('19); Valenti ('22)... Another class: Agrawal and Howe ('17)...





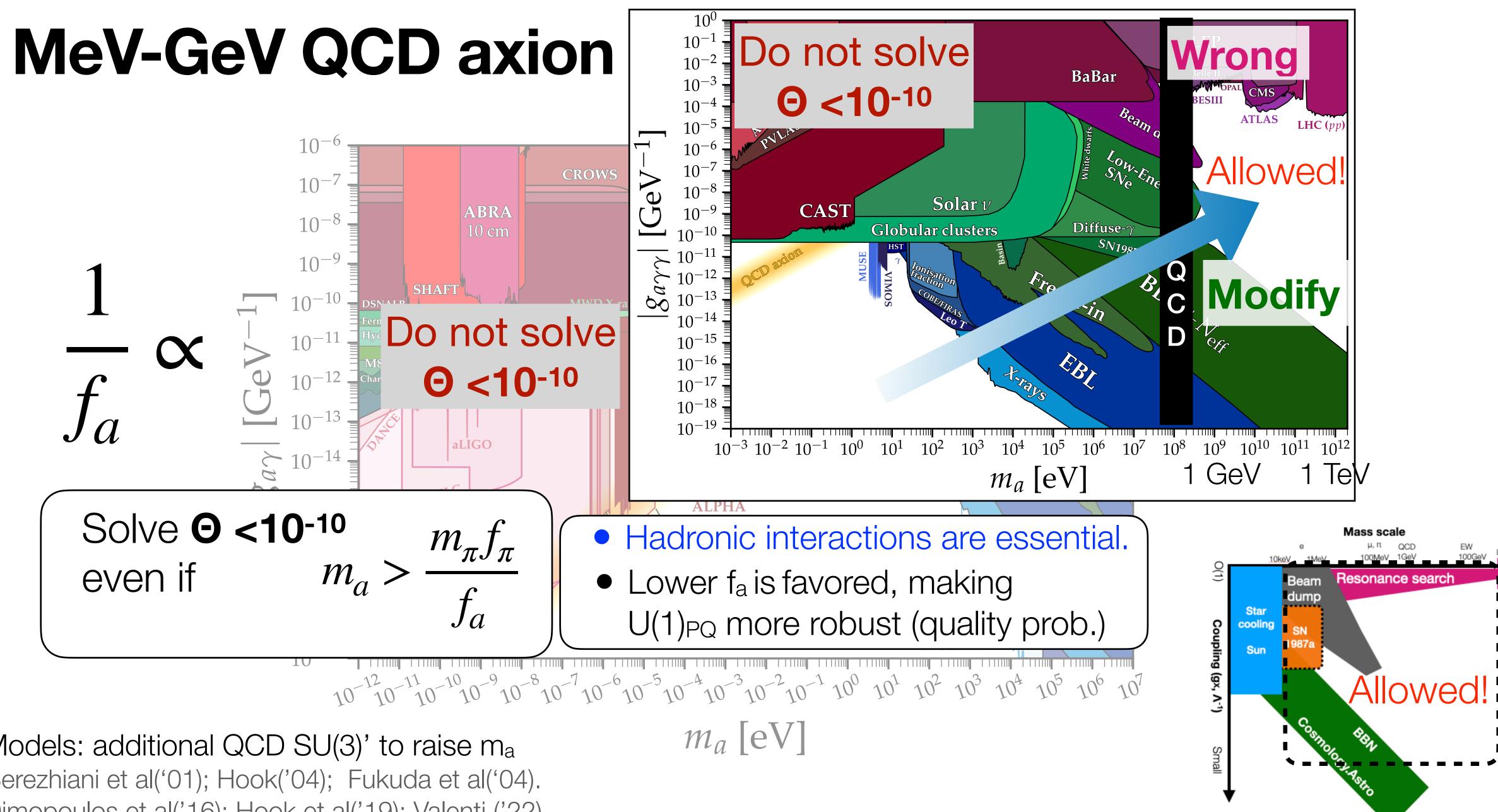
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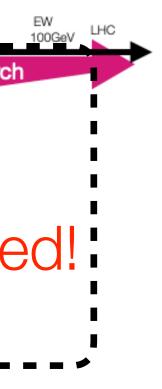


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## Impact of other couplings

- We've considered aGG coupling.  $B \rightarrow Ka$  is from 2-loop (our work).
- $\frac{\alpha_w}{8\pi} \frac{a}{f_a} W \tilde{W}$

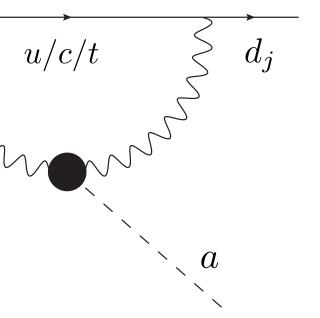
• As long as the strong CP problem is concerned, can't drop aGG $\Rightarrow$ Hadronic decay modes are still dominant ( $a \rightarrow \gamma \gamma$  can be enhanced).

\*Many ALPs can't solve the strong CP because aGG is omitted.

• Lepton-axion couplings are optional. Cleaner signal as  $B \rightarrow Ka( \rightarrow \mu\mu)$ but other bounds become more stringent too (e.g. LEP, LHC)

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### • Other reasonable couplings $\rightarrow$ Production from 1-loop, more signal rate.



1611.09355, E. Izaguirre, T. Lin, B. Shuve

For axion-quark coupling example, see 2002.04623 J. Martin Camalich, M. Pospelov, P. N. H. Vuong, R. Ziegler, J. Zupan

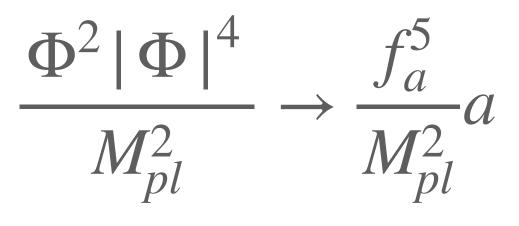




# **Additional motivation: Quality problem**

- Parameter for axion DM: f<sub>a</sub>=10<sup>9-</sup>10<sup>13</sup> GeV
- With such high f<sub>a</sub>, known theoretical issue: "*axion (PQ) quality problem*"
- Any global symmetry including  $U(1)_{PQ}$  expected to be broken by the gravity





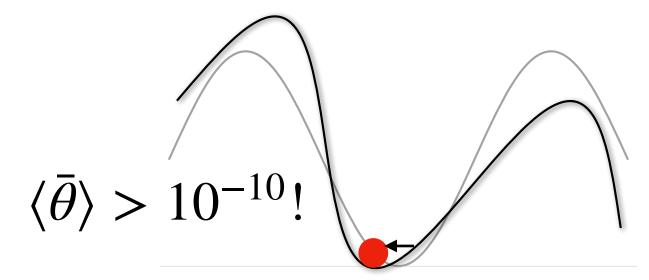
Need model building→going to non-minimal is necessary

The heavy QCD axion has no issue because f<sub>a</sub> is much lower. f<sub>a</sub> below 10TeV is generically OK.

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Refs

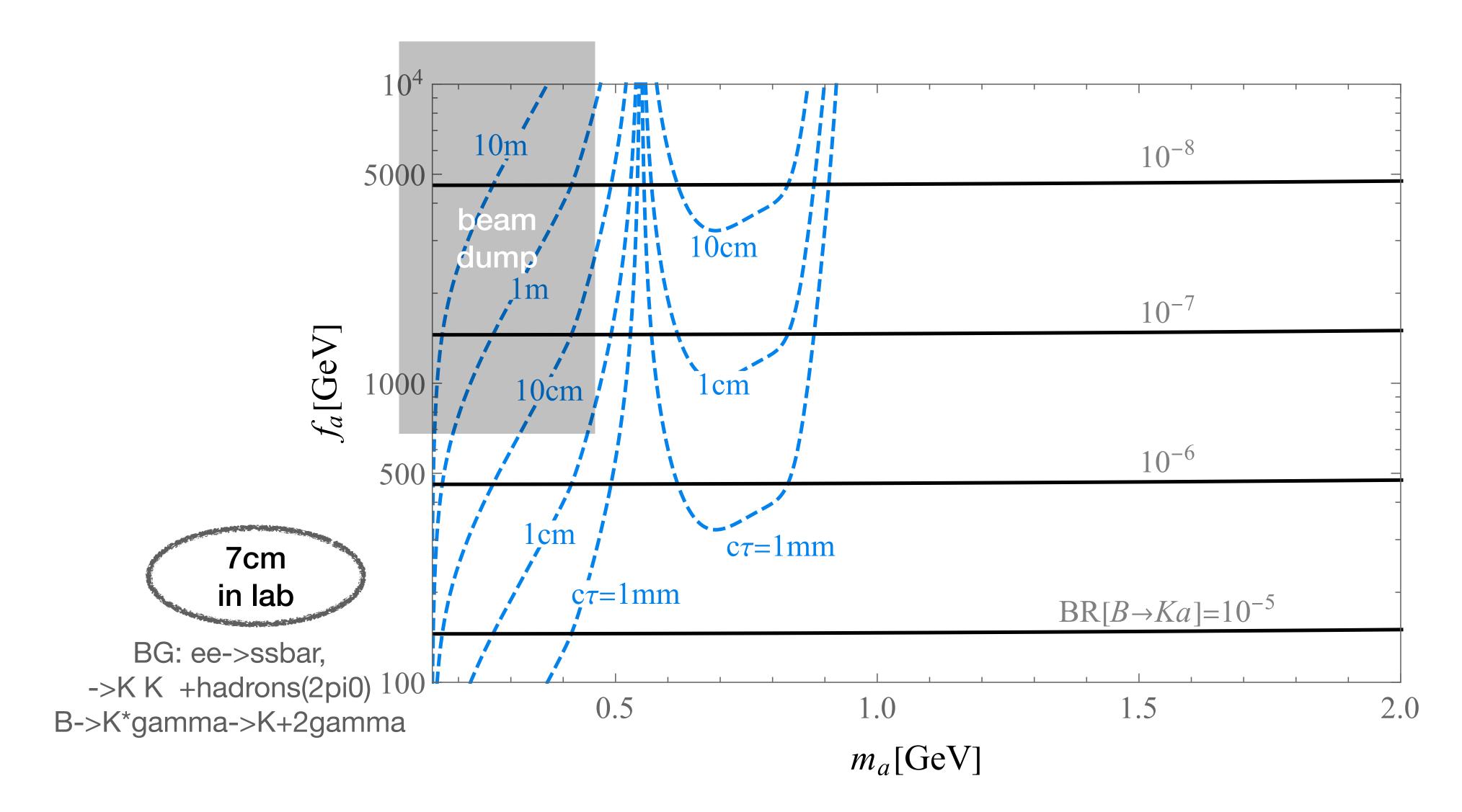
easily ruins axion-solution for the strong CP problem







### Lifetime





### Decay modes

 $10^7$  $10^{3}$ 

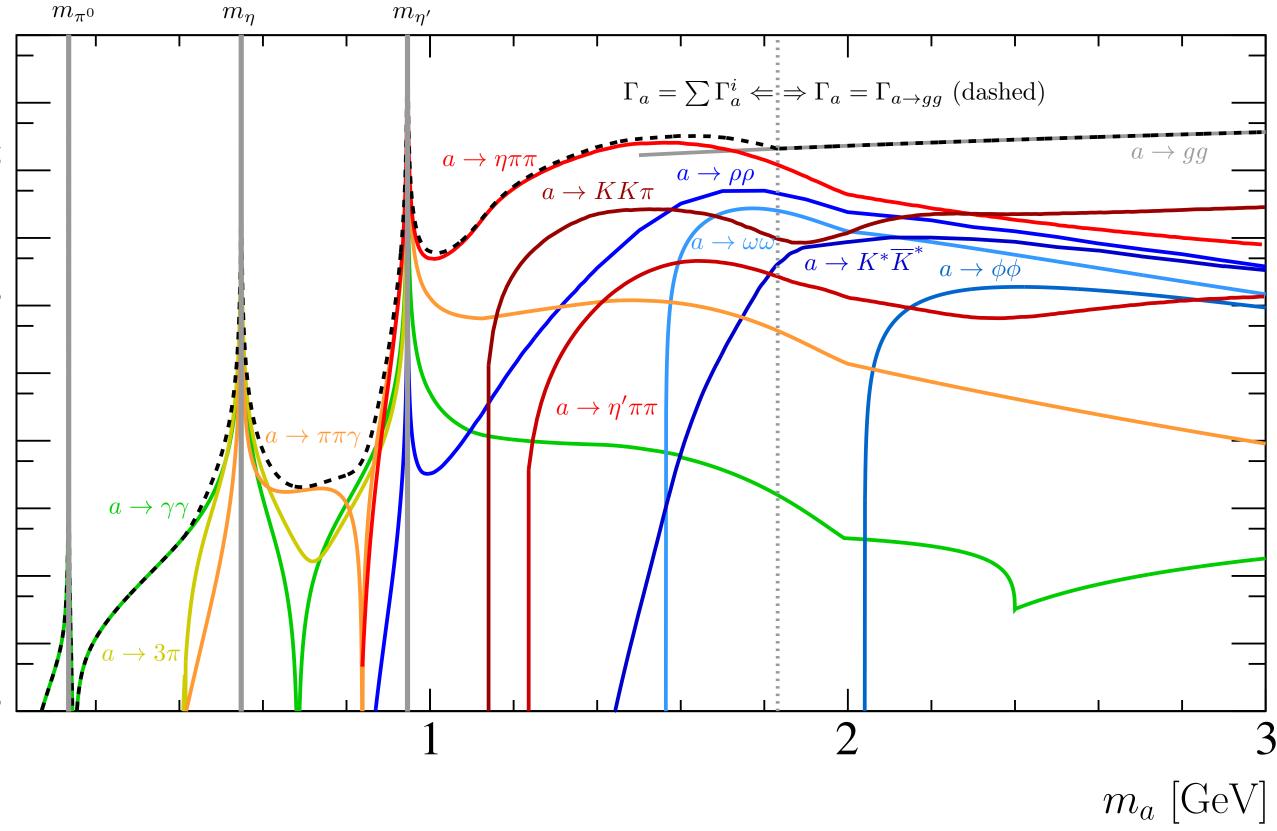
10

 $10^{-1}$ 

 $10^{-3}$ 

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1811.03474, D. Aloni, Y. Soreq, M. Williams





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