





Panel discussion:

BSM and LFV searches with τ 's

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IJCLab (Orsay)

Belle II Physics Week, October, 2024

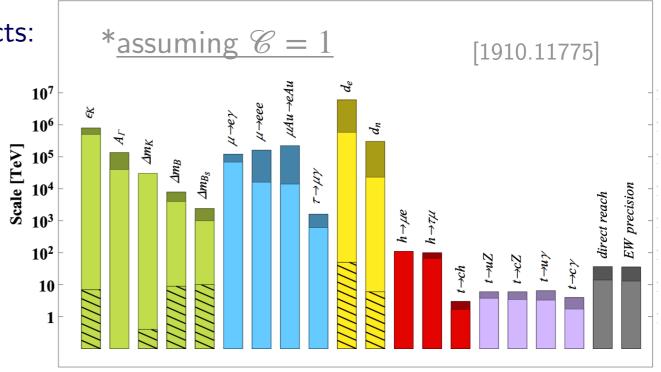


Laboratoire de Physique des 2 Infinis

The Precision Frontier

• Powerful indirect probe of New Physics effects:

EFT constraints on the scale of New Physics (i.e., $\Lambda/g_{\rm NP}$)



• What is the scale of New Physics?

Observable

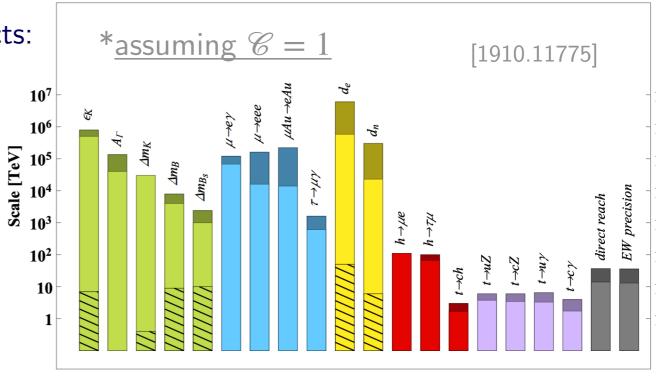
- If we <u>naively</u> assume NP to be flavor blind, then constraints from $\mu \to e$ processes or $K^0 \overline{K}^0$ mixing would be far more constraining than any τ observables...
- However, the SM Yukawa sector has a very peculiar structure (Why? \Rightarrow Flavor Problem).

$$M_{u,d,\ell} = \begin{pmatrix} \bullet & \bullet & \bullet \\ \bullet & \bullet \\ \bullet & \bullet \end{pmatrix} \qquad V_{\rm CKM} = \begin{pmatrix} \bullet & \bullet & \bullet \\ \bullet & \bullet & \bullet \\ \bullet & \bullet & \bullet \end{pmatrix}$$

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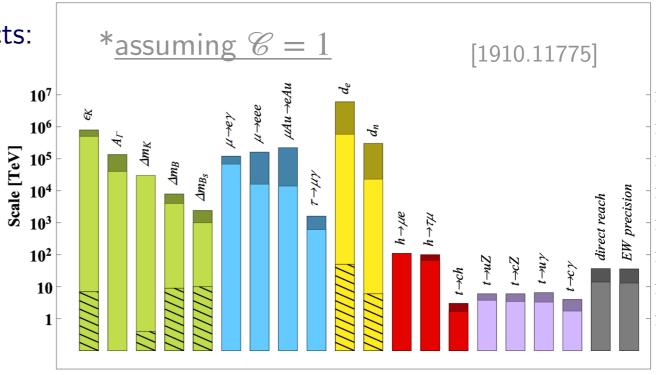
NP could also have a **hierarchical flavor structure!** that would **suppress flavor-changing processes** with **light generations** (e.g., in scenarios with an approximate $U(2)^5$ symmetry).

[Barbieri et al. '11], [Faroughy et al. '20], [Allwicher et al. '23]

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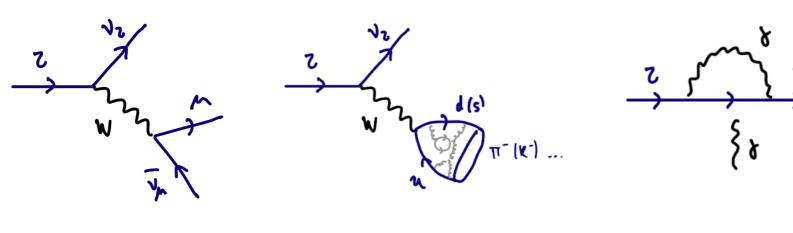
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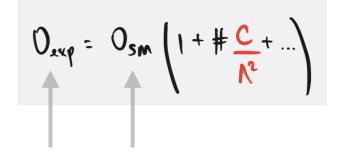
[Barbieri et al. '11], [Faroughy et al. '20], [Allwicher et al. '23]

More generally, exploring *τ*-lepton physics is fundamental to probe the (unknown) flavor structure of physics beyond the SM!

Precision physics with τ 's

• Search for deviations w.r.t. the SM predictions:





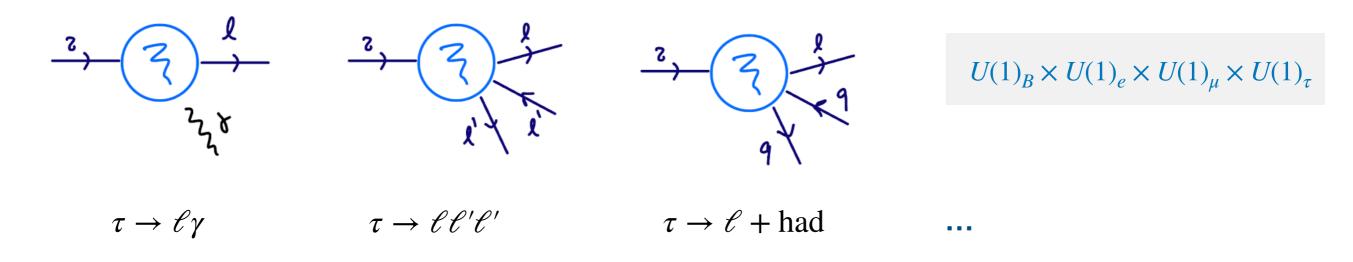
High precision/accuracy needed!

 $\tau \to \ell \nu \bar{\nu} \qquad \qquad \tau \to \nu P$

au-dipoles

See talks by Roig, Bruno, Hoferichter and Passemar

• Search for forbidden SM processes (by accidental symmetries):

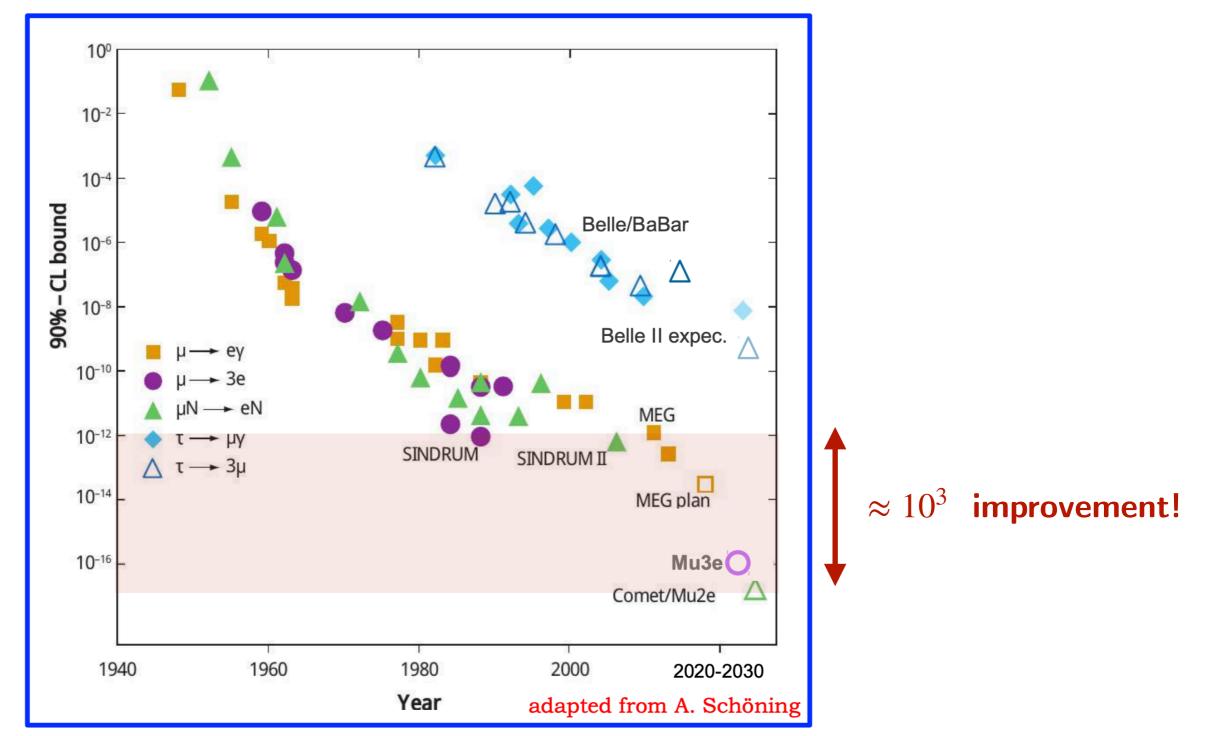


⇒ Very clean probes of New Physics!

See talks by Ardu, Calibbi and Zupan

Charged LFV searches

[From Calibbi's talk]



There will be a **huge improvement** in sensitivity in $\mu \rightarrow e$ experiments — it is fundamental to improve τ -data too (complementarity)!

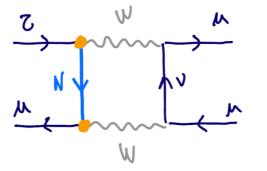
Complementarity

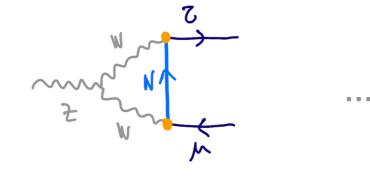
Concrete models predict **correlations** between different observables:

• Heavy neutrinos:

• Leptoquarks:

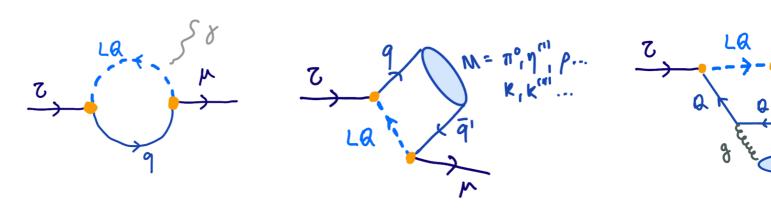
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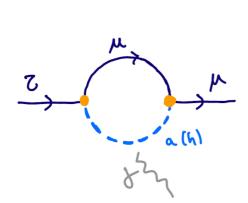
[llakovac et al. '94], [Abada et al. '14]...

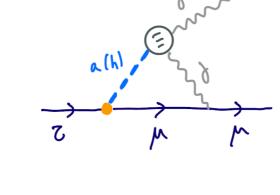
M=n°, y



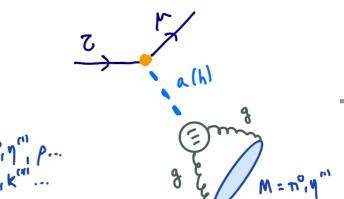
• ALPs (or Higgs):

[Petrov et al. '13], [Becirevic et al. (OS), '16], [Dorsner et al. '16]...





za(h)q $M = \pi^{o} \eta^{c'}$

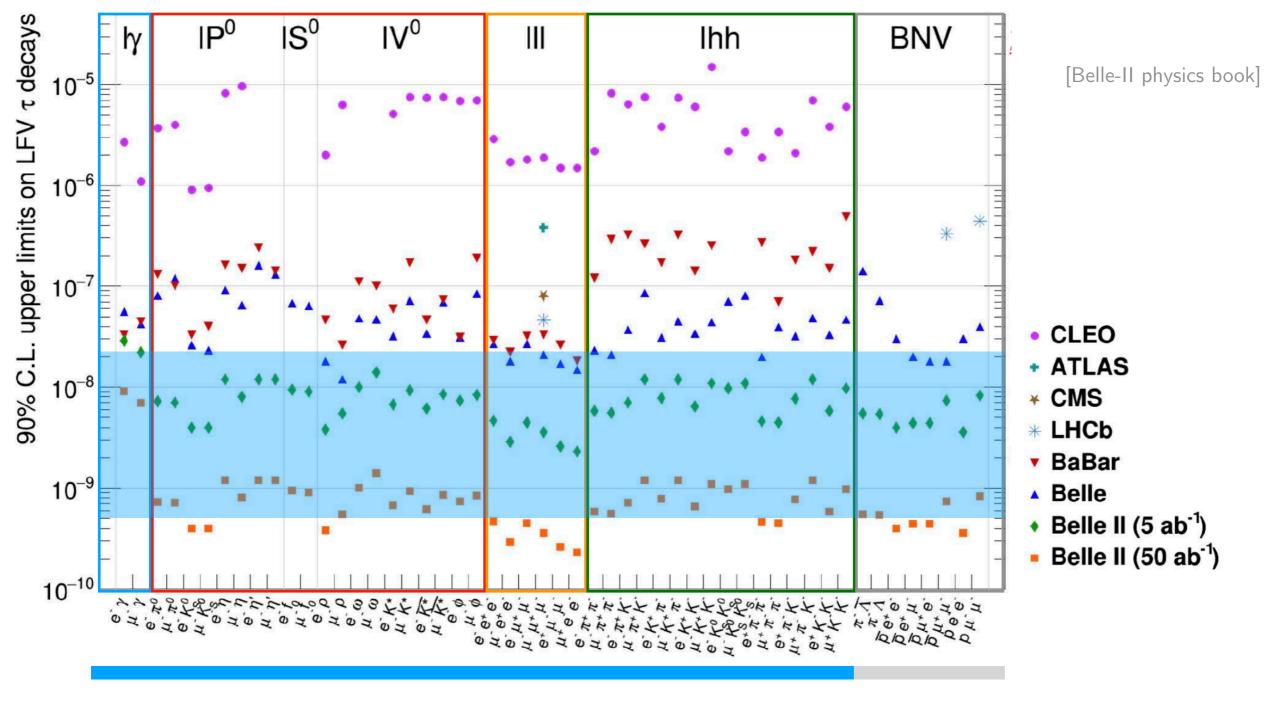


[Han et al. '08], [Harnik et al. '12], [Blankenburg et al. '13]...

[Cornella, Paradisi, **OS**. '19]...

Belle-II and τ -decays

Belle-II will **improve** the **sensitivity** on $\tau \to e$ and $\tau \to \mu$ decays by a **factor** $\mathcal{O}(10)$:



LFV decays

BNV decays

Discussion

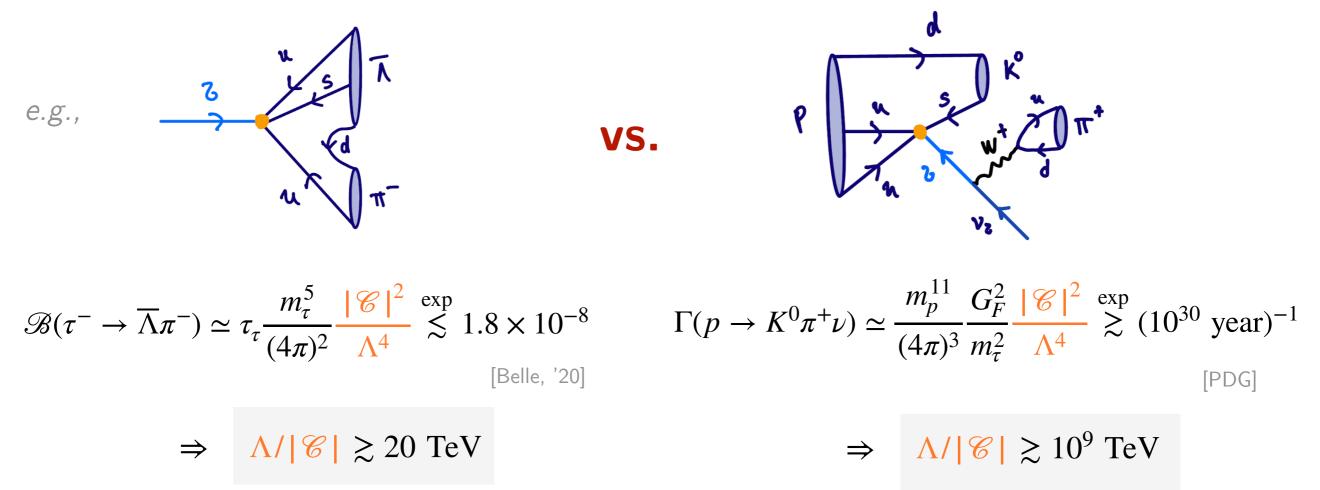
I. Should we look for BNV in τ -decays?

The lowest-order operators that violate B in the SM appear at d = 6:

[Weinberg, '79]

 $O_1 \sim QQQL$ $O_2 \sim QQue$ $O_3 \sim Qu_R d_R L$ $O_4 \sim u_R u_R d_R e_R$

They could in principle induce BNV τ -decays depending on their flavor content:



<u>Caution</u>: the same operators that generate BNV τ -decays may also induce the proton decay (via an insertion of G_F or EW loops) — potentially much more constraining!

O. Sumensari

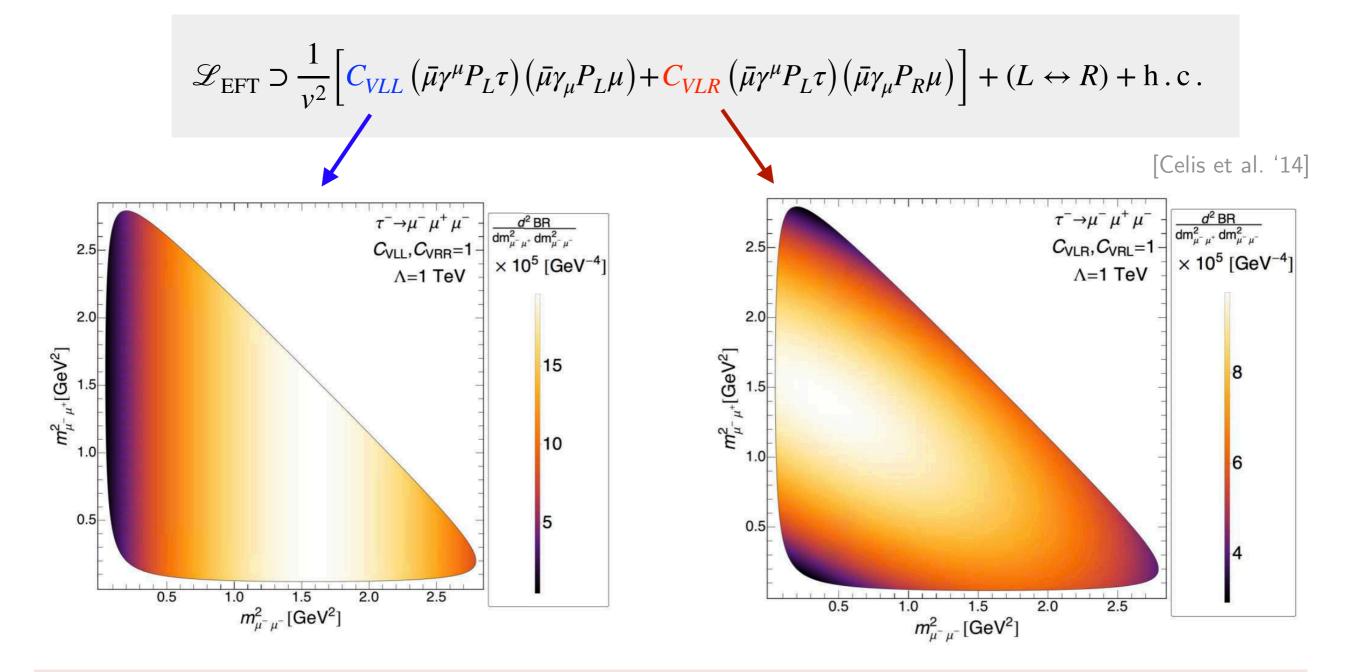
see [Beneke et al. '24] for a discussion of *B*-meson BNV decays

II. EFTs and Dalitz plot

See talk by Ardu!

Which Lorentz structure?

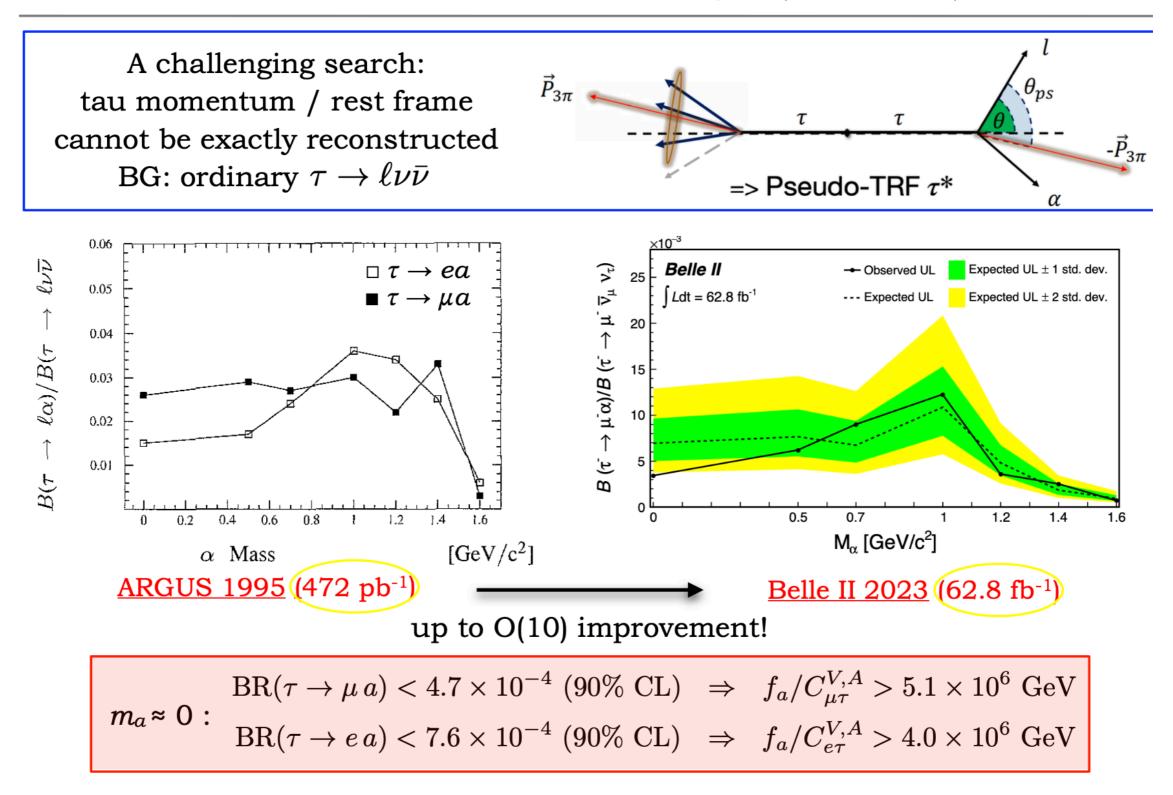
[Dassinger et al. '07]



Question: Which **EFT scenario** is used in the **MC simulation**? Can/should we further exploit this information experimentally?

III. Beyond the SMEFT

Present limits on $\tau \rightarrow e a$, $\tau \rightarrow \mu a$ (invisible *a*)

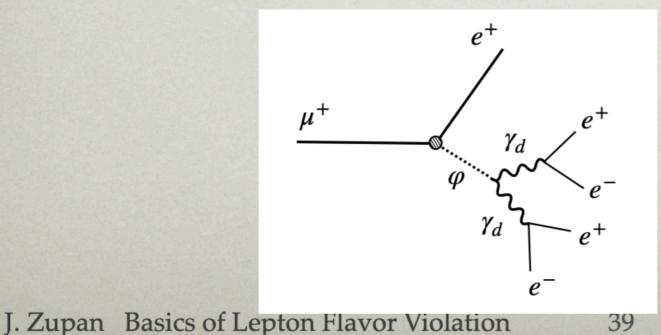


Other channels to be explored?

IV. New ideas

 $\mu \rightarrow 5e$

- if $\frac{m_{\mu}}{\Lambda} \phi(\bar{e}\mu)$ coupling \Rightarrow mediates $\mu \rightarrow e\phi$
 - if φ QCD axion \Rightarrow escapes the detector $\mu \rightarrow e + inv$
 - MEG-II, Mu3e, Mu2e-X, COMET-X can search for it
 - if φ can decay \Rightarrow sensitivity to even higher scales
 - example: $\mu \rightarrow 5e$ can probe $f_a \gtrsim 10^{13} \text{GeV}$



Hostert, Menzo, Pospelov, JZ, 2306.1563

Can we do the same for τ 's?

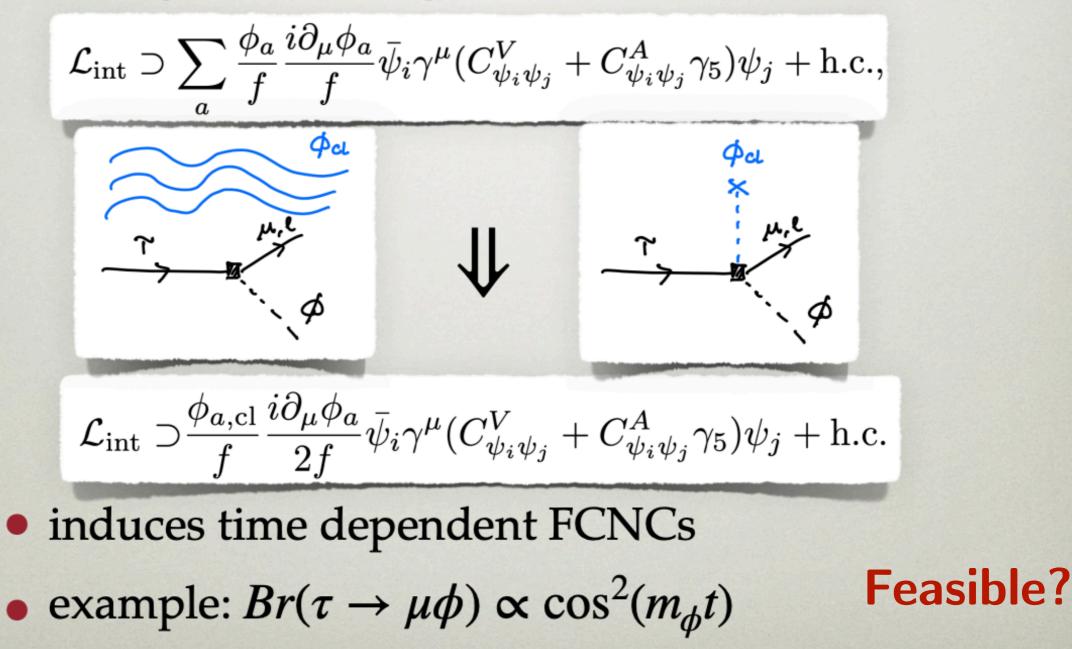
KEK, Oct 17 2024

IV. New ideas

See talk by Jure!

NON-ÅBELIAN PNGB

in the light DM background



J. Zupan Challenging BSM searches with taus

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Take-home

Improving the sensitivity of LFV τ decays is fundamental to constraining the flavor structure beyond the SM and is a powerful approach to probe high-energy scales (otherwise inaccessible...) — in particular, in models that <u>couple hierarchically</u> to leptons.

These decays also provide **key information** to **indirectly constrain potential** NP explanations of discrepancies in low-energy data (e.g., *B-physics*) and to probe scenarios with **light particles**.

Questions:

- I. Does it really make sense to look for **BNV decays of** τ 's? How to avoid *p*-decay bounds?
- II. Which model is used to simulate the signal in $\tau \rightarrow \ell \ell \ell \ell$ searches? Can we further exploit the **kinematical distributions**?
- III. Are there searches for **displaced vertices** $(\tau \rightarrow \ell \varphi (\rightarrow \gamma \gamma, \ell \ell))$ in LFV decays?
- IV. **Time-modulation** of $\tau \to \ell \varphi$: Which are the main challenges experimentally?
- V. What can we learn if the e^+e^- beam is polarized both in models with light/heavy particles?

Back-up

EFT for $\tau \to \ell \gamma$ and $\tau \to \ell \ell \ell$

See talk by Ardu!

$$\mathscr{L}_{\text{LEFT}} \supset \frac{1}{v^2} \sum_{a} C_a O_a \quad \blacktriangleleft$$

Operators invariant under $SU(3)_c \times U(1)_{em}$

• Dipoles:

$$O_{D_L}^{\ell} = e m_{\tau} \bar{\ell}_R \sigma_{\mu\nu} \tau_L F^{\mu\nu}$$

... and $(L \leftrightarrow R)$

• Four-leptons:

$$O_{V_{LL}}^{\ell} = (\bar{\ell} \gamma^{\mu} P_L \tau) (\bar{\ell} \gamma_{\mu} P_L \ell)$$

$$O_{V_{LR}}^{\ell} = (\bar{\ell} \gamma^{\mu} P_L \tau) (\bar{\ell} \gamma_{\mu} P_R \ell)$$

$$O_{V_{LR}}^{\ell} = (\bar{\ell} \gamma^{\mu} P_L \tau) (\bar{\ell} \gamma_{\mu} P_R \ell)$$

$$O_{LR}^{\ell} = (\bar{\ell} \sigma^{\mu\nu} P_L \tau) (\bar{\ell} \sigma_{\mu\nu} P_L \ell)$$

SMEFT for $\tau \rightarrow \ell \gamma$ and $\tau \rightarrow \ell \ell \ell$

See talk by Ardu!

• Dipoles:

$$\mathcal{O}_{B} = \left(\bar{L}_{i}\sigma^{\mu\nu}e_{Rj}\right)HB_{\mu\nu}$$
$$\mathcal{O}_{W} = \left(\bar{L}_{i}\tau^{I}\sigma^{\mu\nu}e_{Rj}\right)HW_{\mu\nu}^{I}$$

• Four-leptons:

$$\mathcal{O}_{ll} = (\bar{L}_i \gamma^{\mu} L_j) (\bar{L}_k \gamma_{\mu} L_l)$$
$$\mathcal{O}_{le}_{ijkl} = (\bar{L}_i \gamma^{\mu} L_j) (\bar{e}_{Rk} \gamma_{\mu} e_{Rl})$$
$$\mathcal{O}_{ee}_{ijkl} = (\bar{e}_{Ri} \gamma^{\mu} e_{Rj}) (\bar{e}_{Rk} \gamma_{\mu} e_{Rl})$$

$$\mathscr{L}_{\text{SMEFT}} \supset \frac{1}{\Lambda^2} \sum_{a} \mathscr{C}_a \mathscr{O}_a \quad \blacktriangleleft$$

[Buchmuller et al. '86, Grzadkowski et al. '10]

Operators invariant under $SU(3)_c \times SU(2)_L \times U(1)_Y$

*Flavor indices omitted... Matching @ tree-level: $\int C_{D_L} \propto \frac{v^2}{\Lambda^2} \left[\cos \theta_W \mathscr{C}_{eW} - \sin \theta_W \mathscr{C}_{eB} \right]$ $C_{V_{LL}} \propto \frac{v^2}{\Lambda^2} \mathscr{C}_{ll} + \dots$ $C_{V_{LR}} \propto \frac{v^2}{\Lambda^2} \mathscr{C}_{le} + \dots$... and $(L \leftrightarrow R)$ $\begin{array}{c} \label{eq:constraint} \mathbf{C}_{S_{LL}} = \mathcal{O}(\Lambda^{-4}) \\ C_{S_{LR}} = \mathcal{O}(\Lambda^{-4}) \end{array} \begin{array}{c} \mathbf{C}_{T_L} = \mathcal{O}(\Lambda^{-4}) \\ \mathbf{C}_{S_{LR}} = \mathcal{O}(\Lambda^{-4}) \end{array}$

Suppressed in the SMEFT!

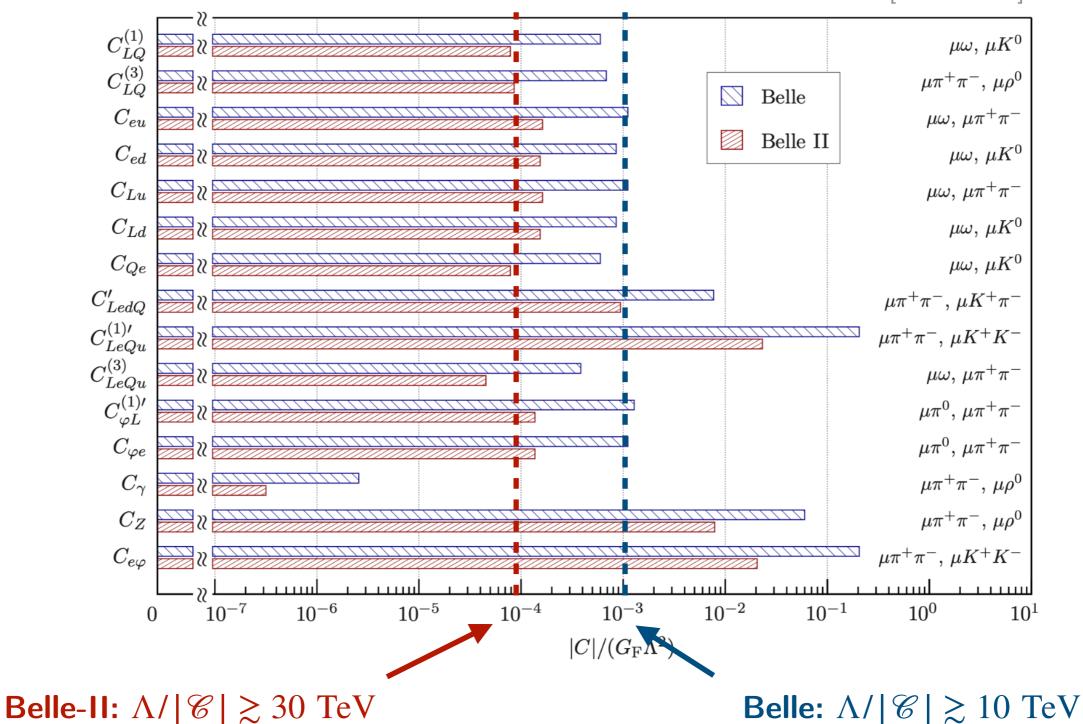
Gauge invariance is useful for constraining the possible operators

EFT for $\tau \rightarrow \ell + had$



New Physics scale of SMEFT operators probed by **hadronic** *τ*-**decays**:

[2203.14919]

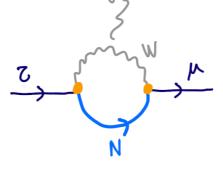


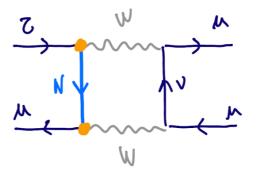
From EFTs to concrete models

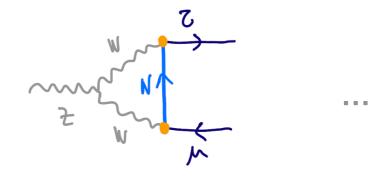
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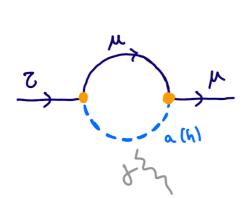


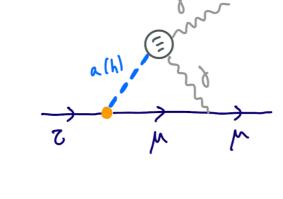


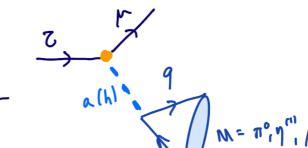
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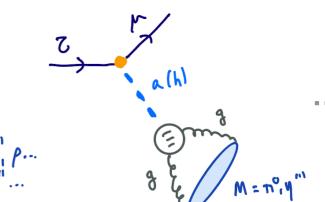
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M=nº, y

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