

# The CDF-II Anomaly and Potential New Physics Explanations

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**IIT Hyderabad**  
**21<sup>th</sup> Oct. 2024**

# Outline

- The CDF-II Anomaly
- W Boson Mass in SM
- The Oblique S,T,U Parameters
- New Physics Scenarios
  - Models with Triplet Scalars
  - Models with Doublet Scalars
  - Models with modified Z boson mass
- Conclusions

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  - **Also 7-sigma away from SM prediction !!**

# How you measure $W$ mass?

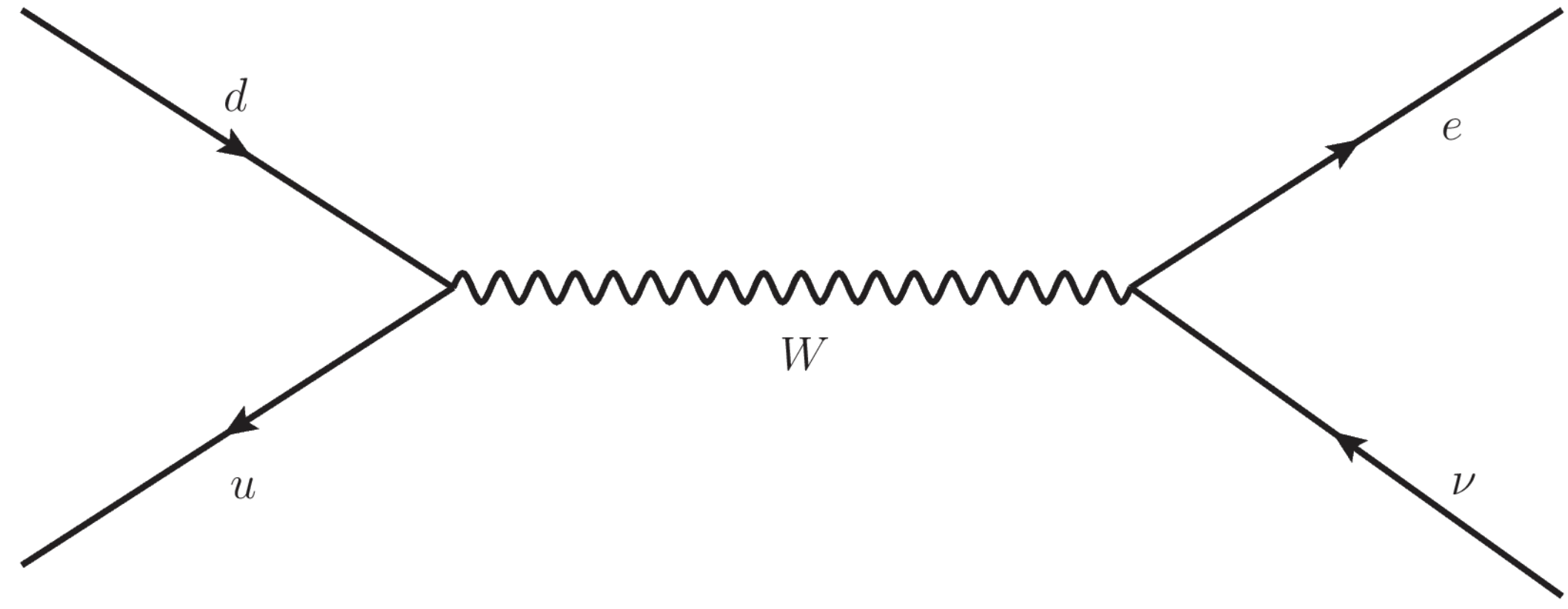
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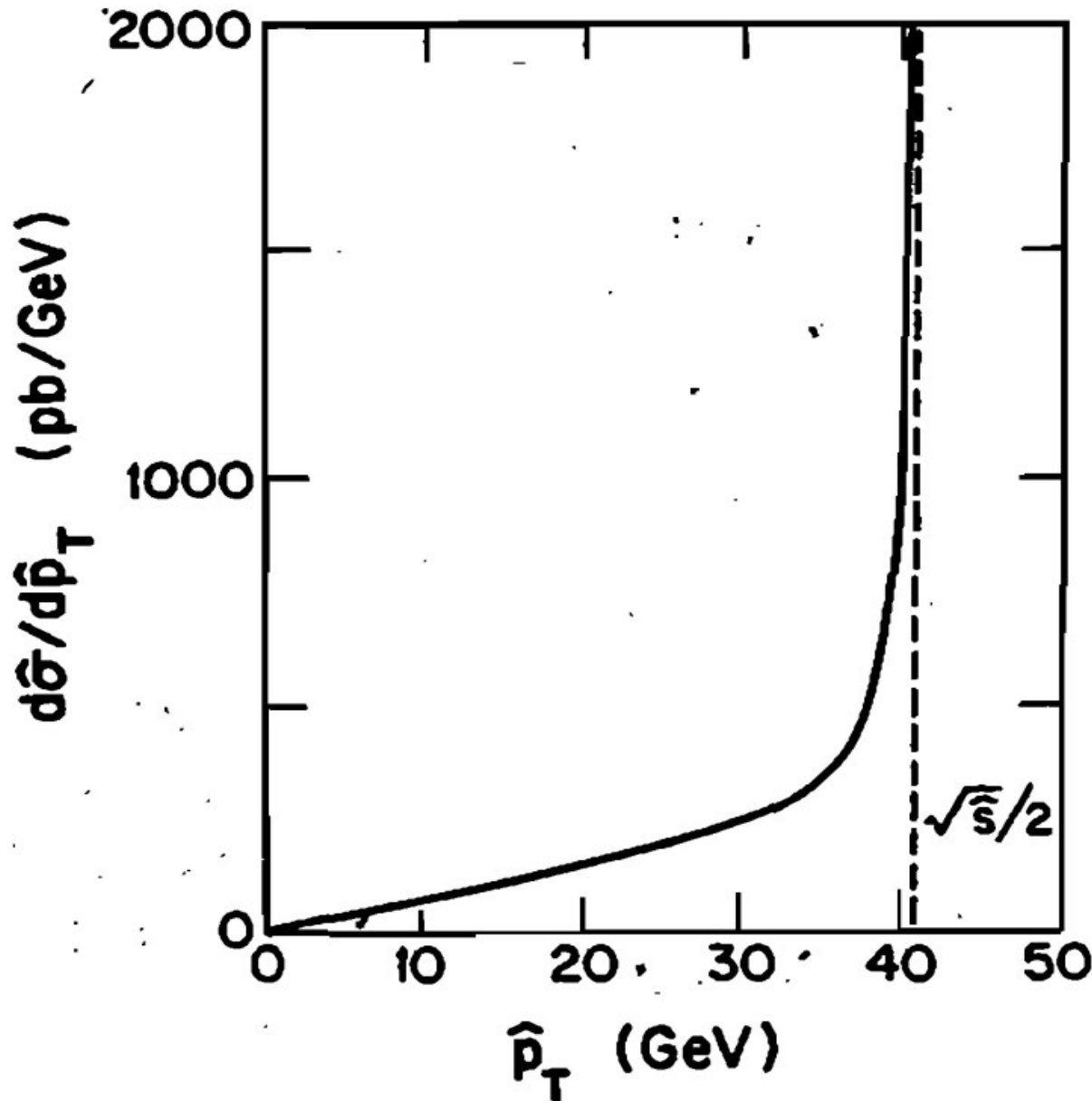
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  - Use the transverse momentum (pT) information of the charged lepton

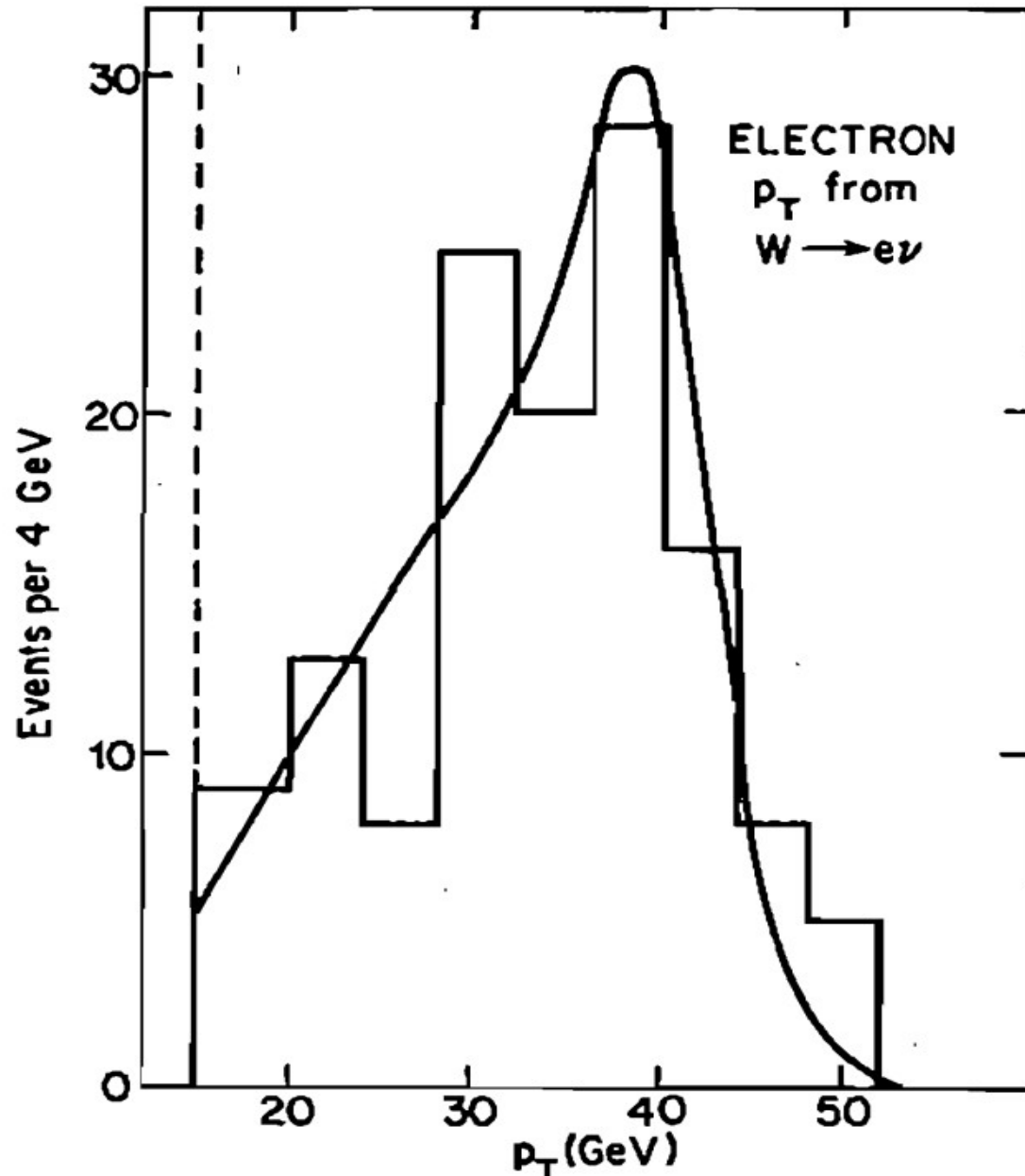


# How you measure $W$ mass?



**Fig. 8.9.** Jacobian peak  $p_T$ -distribution in the  $u\bar{d} \rightarrow e^+\nu$  subprocess at fixed  $\hat{s}$ .

# How you measure $W$ mass?



**Fig. 8.10.** Distribution of electron  $p_T$  from  $p\bar{p} \rightarrow W \rightarrow e\nu$  events at  $\sqrt{s} = 630$  GeV (UA1 data) compared with a calculation folding in the smearing from  $p_T(W)$ .

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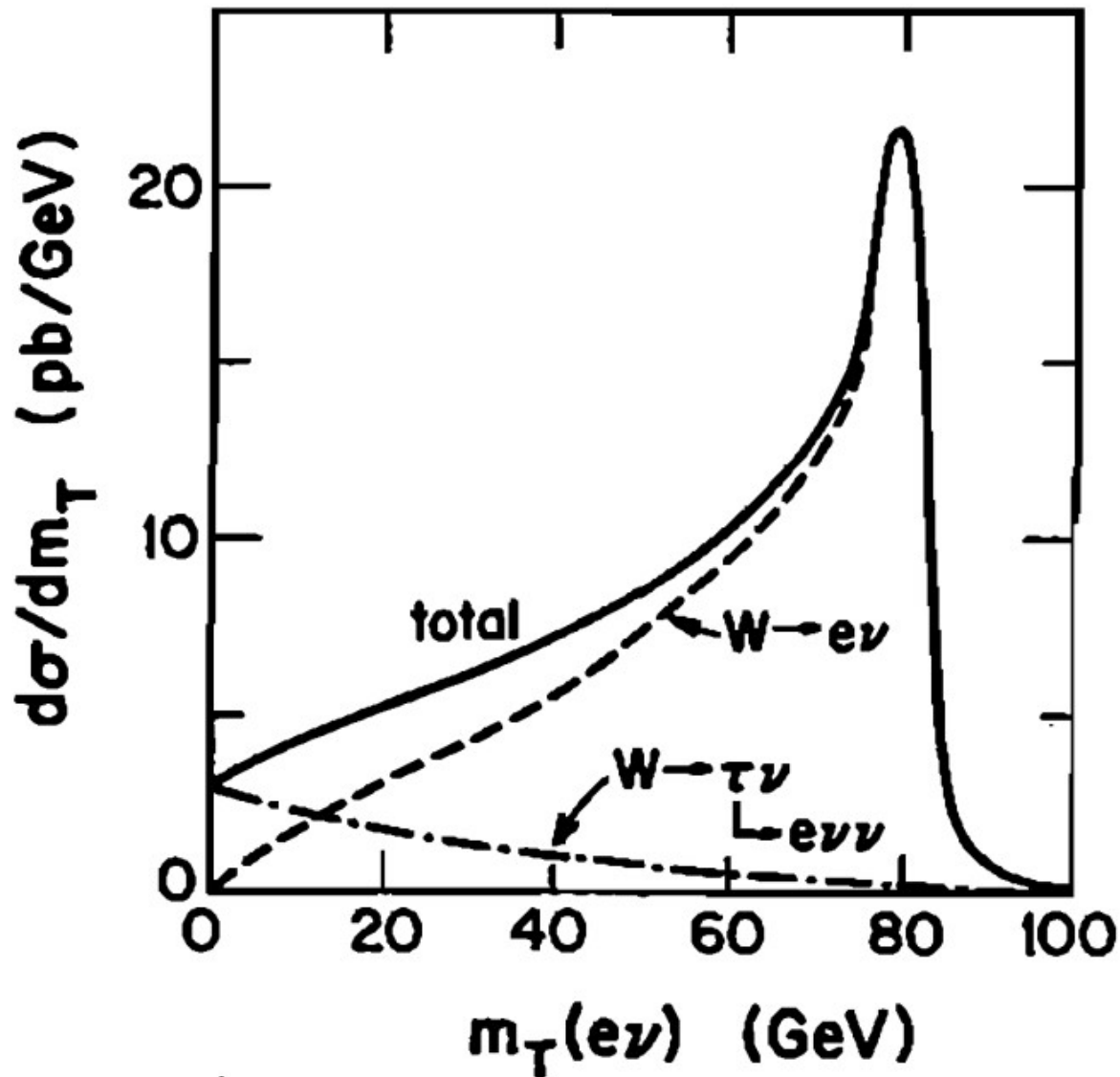
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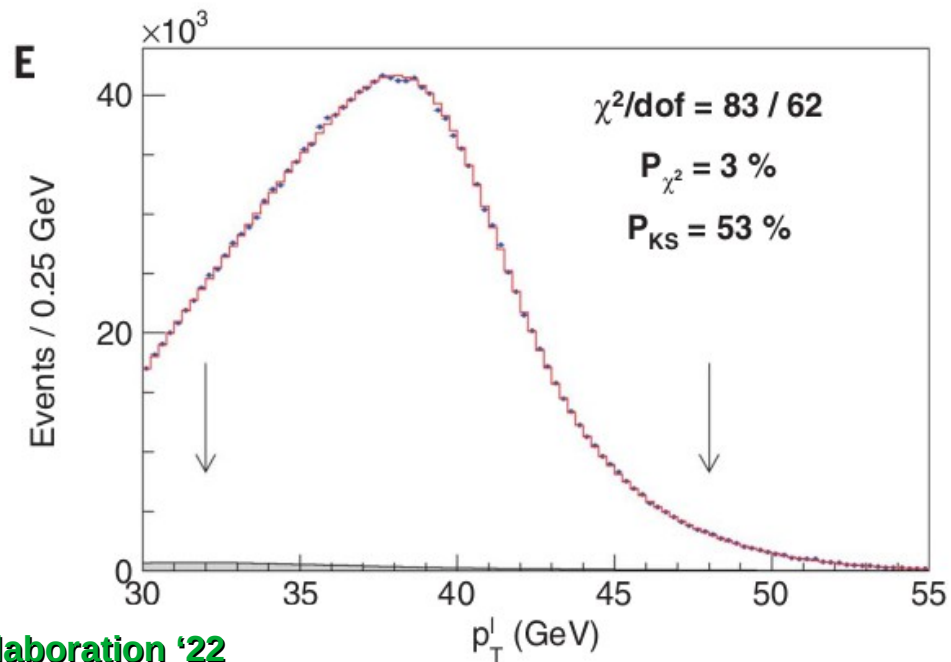
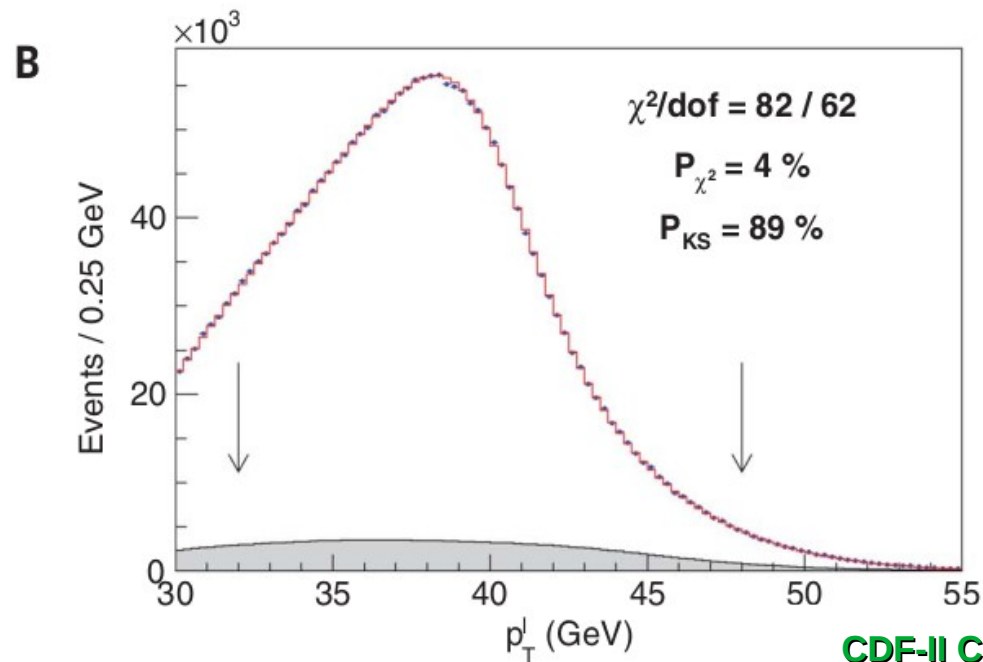
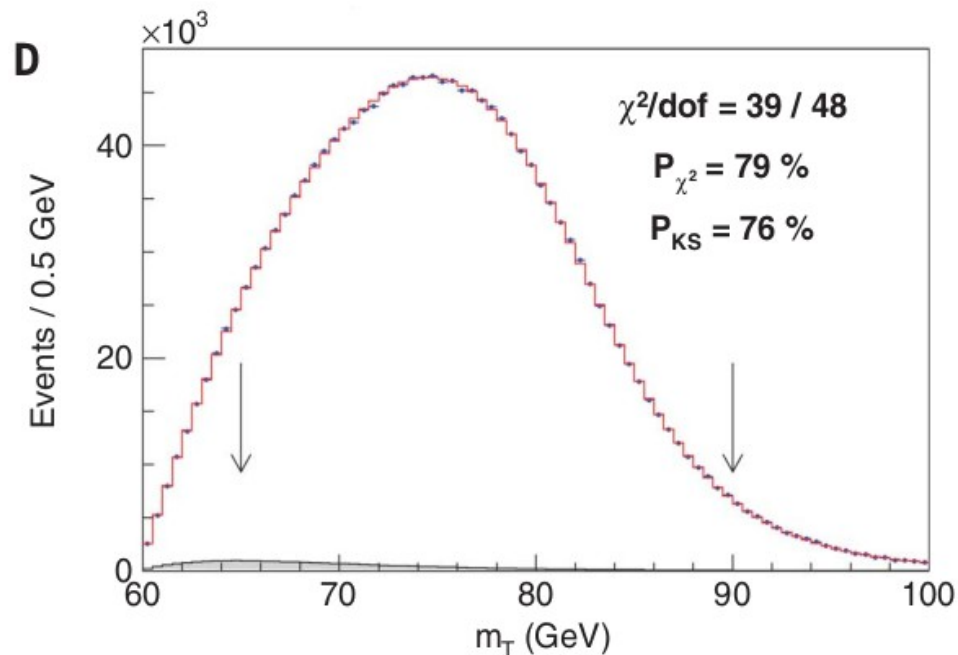
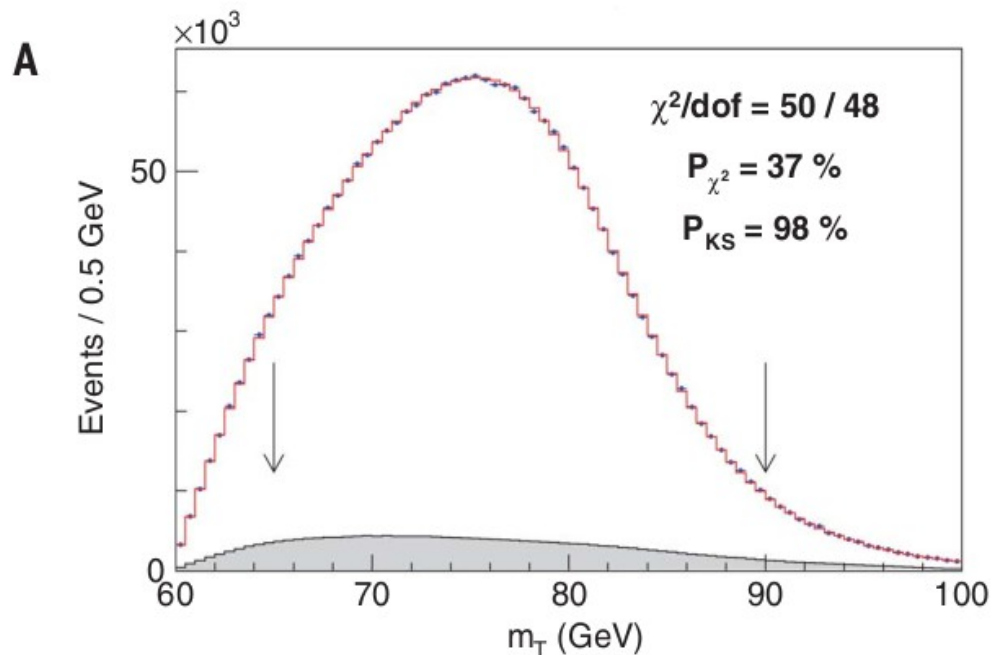
$$m_T^2(e, \nu) = (|\mathbf{p}_{eT}| + |\mathbf{p}_{\nu T}|)^2 - (\mathbf{p}_{eT} + \mathbf{p}_{\nu T})^2$$

# How you measure W mass?



**Fig. 8.11.** Predicted distribution of transverse mass  $m_T(e\nu)$  for  $p\bar{p} \rightarrow W^\pm \rightarrow e\nu$  collisions at  $\sqrt{s} = 630$  GeV.

# The CDF-II W mass measurement



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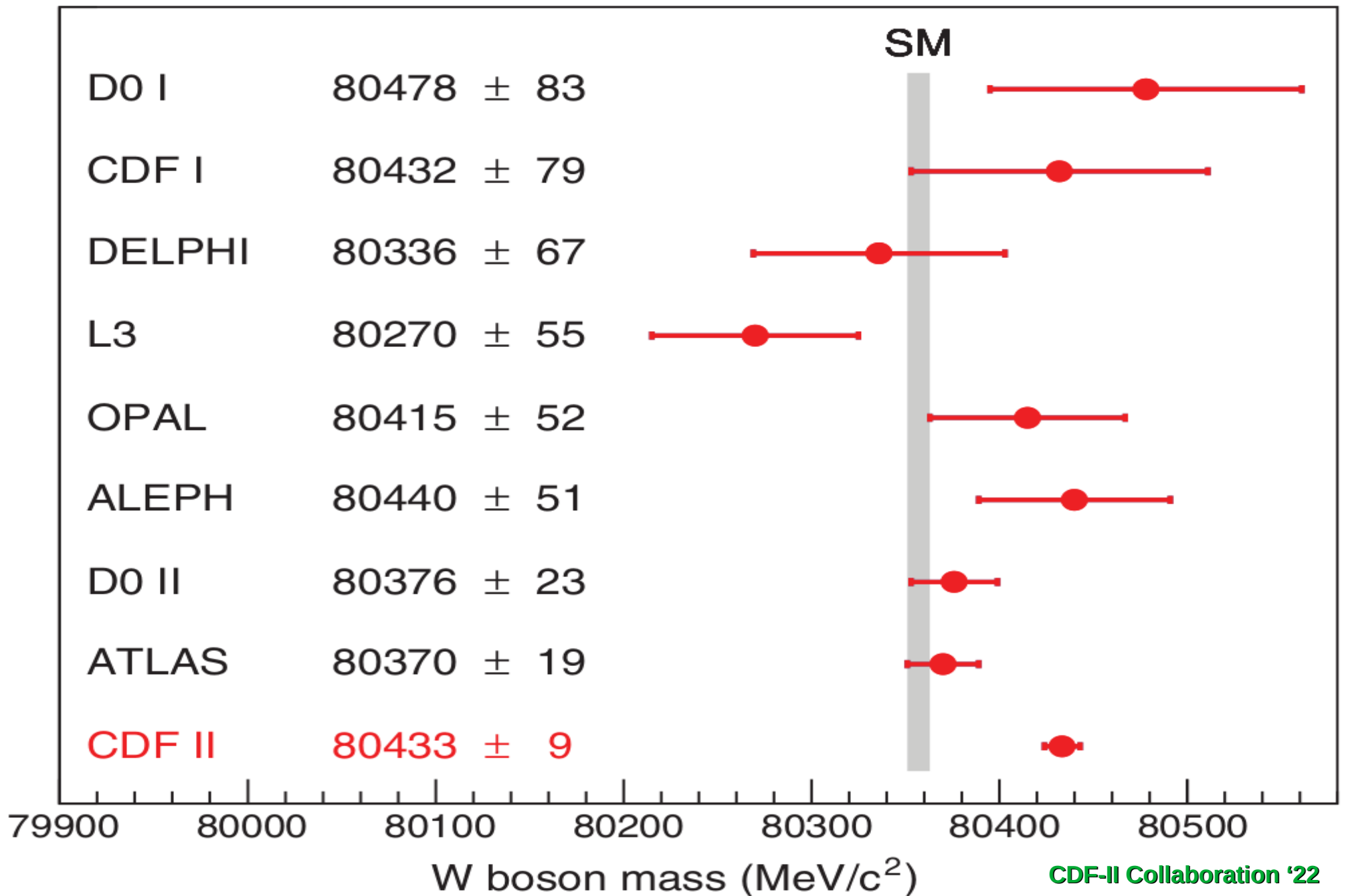
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# The CDF-II Anomaly



$$m_W = 80360.2 \pm 9.9 \text{ MeV}$$

LEP combination

Phys. Rep. 532 (2013) 119

D0

PRL 108 (2012) 151804

CDF

Science 376 (2022) 6589

LHCb

JHEP 01 (2022) 036

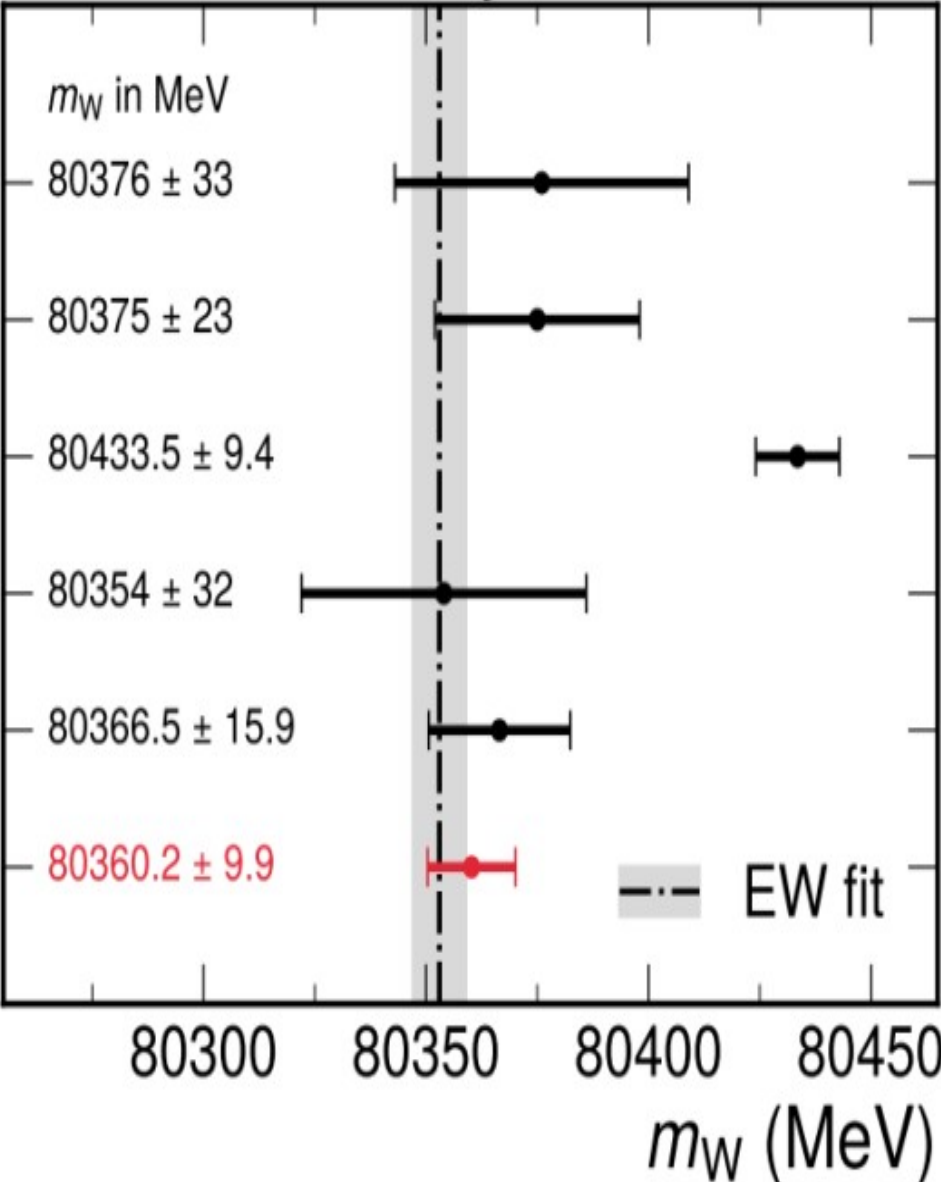
ATLAS

arxiv:2403.15085, subm. to EPJC

**CMS**

*This Work*

# CMS Preliminary



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  - **I will also discuss their connections to other shortcomings of the SM**

# W boson Mass in SM

- In SM the W, Z bosons get their mass after SSB via Higgs mechanism
  - The mass of W, Z are proportional to the Higgs vev and the gauge couplings
  - They are also tied together by the rho-parameter (on-shell renormalization scheme)

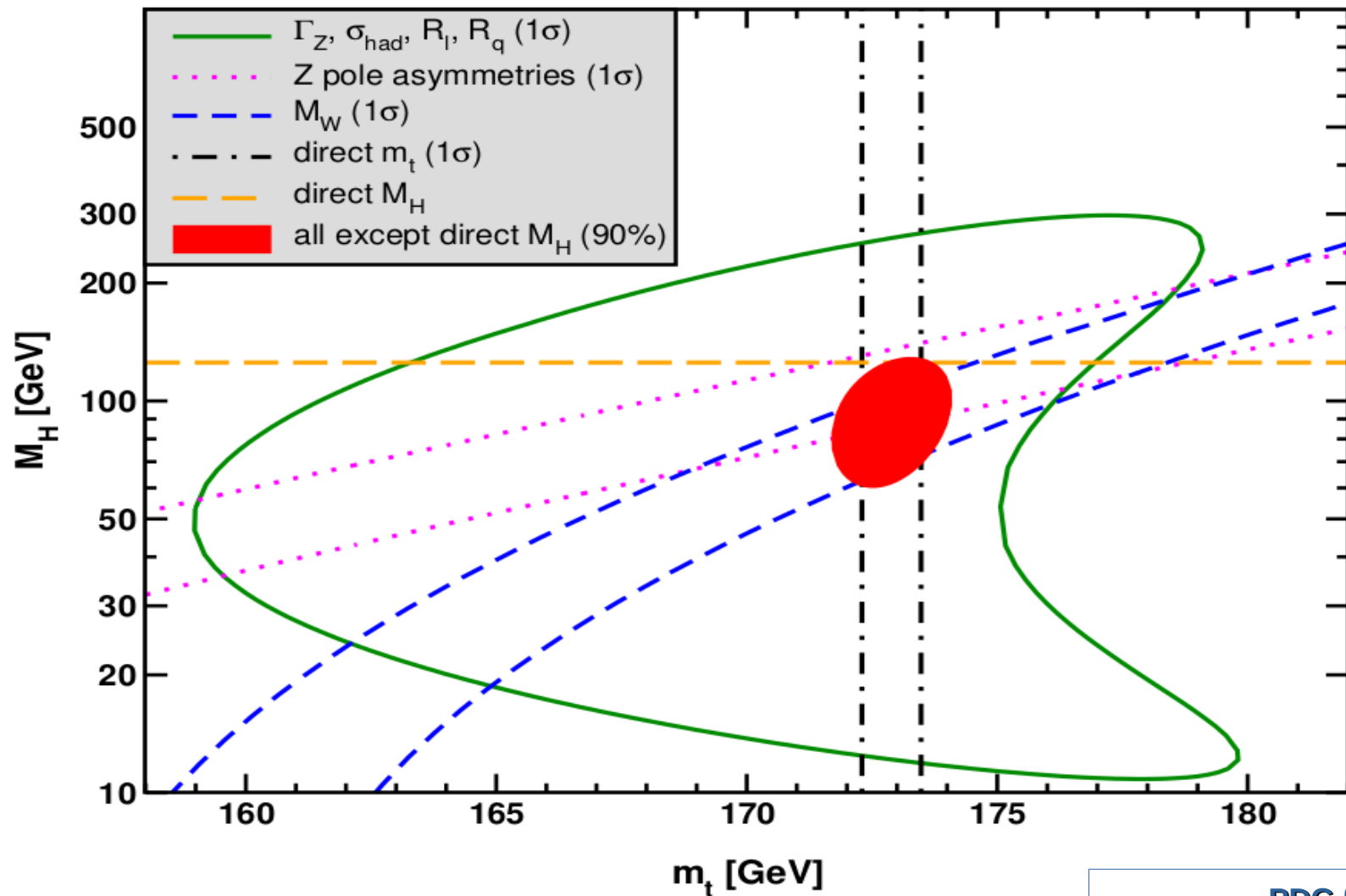
$$\rho = \frac{m_W^2}{m_Z^2 \cos^2 \theta_W} \quad \text{in SM : } \rho = 1$$

- So if you know mass of one gauge boson then you can compute the SM expectation for the other
  - Typically Z mass can be easily and precisely measured
  - Z mass used as input for W mass expectation



# SM Global Fit

- You actually do a global fit



# The Oblique S, T, U Parameters

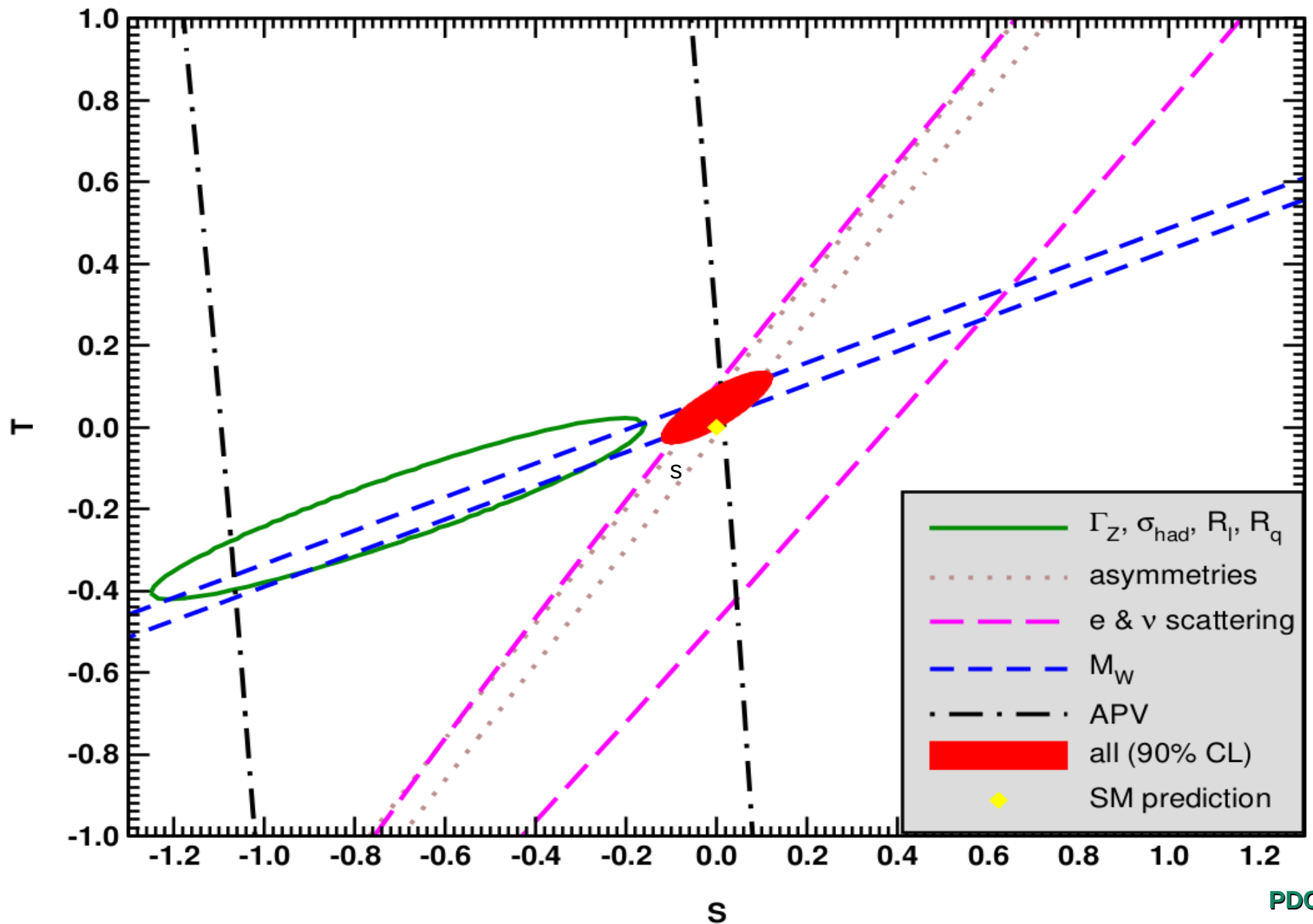
- New physics effects can change the SM relations
- Model independent parametrization: The Oblique S, T, U parameters

$$\hat{\alpha}(M_Z)T \equiv \frac{\Pi_{WW}^{\text{new}}(0)}{M_W^2} - \frac{\Pi_{ZZ}^{\text{new}}(0)}{M_Z^2},$$

$$\begin{aligned} \frac{\hat{\alpha}(M_Z)}{4\hat{s}_Z^2\hat{c}_Z^2}S \equiv & \frac{\Pi_{ZZ}^{\text{new}}(M_Z^2) - \Pi_{ZZ}^{\text{new}}(0)}{M_Z^2} - \frac{\hat{c}_Z^2 - \hat{s}_Z^2}{\hat{c}_Z\hat{s}_Z} \frac{\Pi_{Z\gamma}^{\text{new}}(M_Z^2)}{M_Z^2} \\ & - \frac{\Pi_{\gamma\gamma}^{\text{new}}(M_Z^2)}{M_Z^2}, \end{aligned}$$

$$\begin{aligned} \frac{\hat{\alpha}(M_Z)}{4\hat{s}_Z^2}(S + U) \equiv & \frac{\Pi_{WW}^{\text{new}}(M_W^2) - \Pi_{WW}^{\text{new}}(0)}{M_W^2} - \frac{\hat{c}_Z}{\hat{s}_Z} \frac{\Pi_{Z\gamma}^{\text{new}}(M_Z^2)}{M_Z^2} \\ & - \frac{\Pi_{\gamma\gamma}^{\text{new}}(M_Z^2)}{M_Z^2}. \end{aligned}$$

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- **New physics effects that can explain CDF-II anomaly fall into some general categories**
  - **Triplet scalars changing W-mass at tree level**
  - **Doublet scalars changing W-mass at one loop level**
  - **A new U(1) gauge symmetry and the Z' boson changing Z-boson mass at tree level**

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$$M_W = \frac{g}{2} \sqrt{v_\Phi^2 + 4v_\Omega^2} \text{ and } M_Z = \frac{\sqrt{g^2 + g'^2}}{2} v_\Phi$$

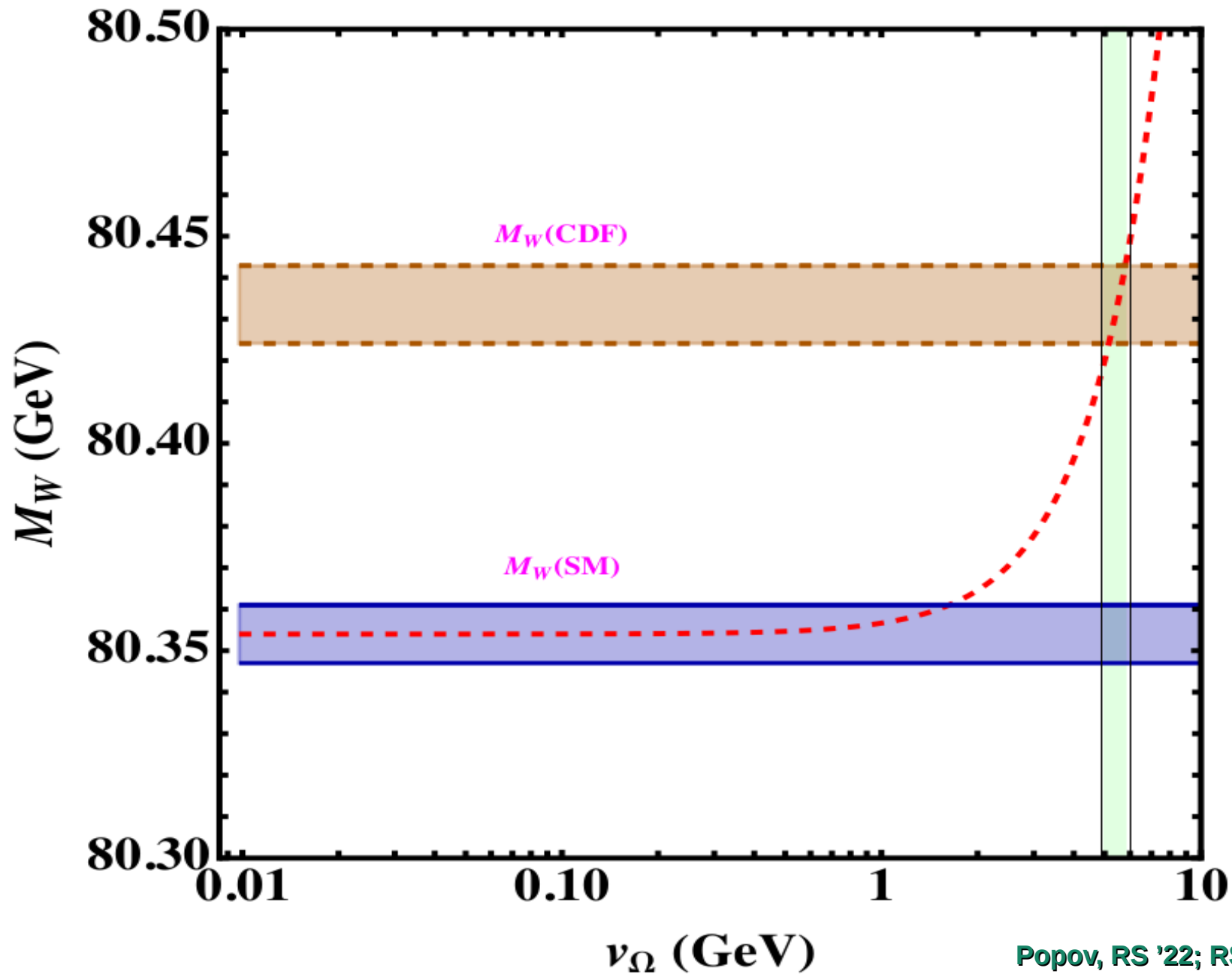
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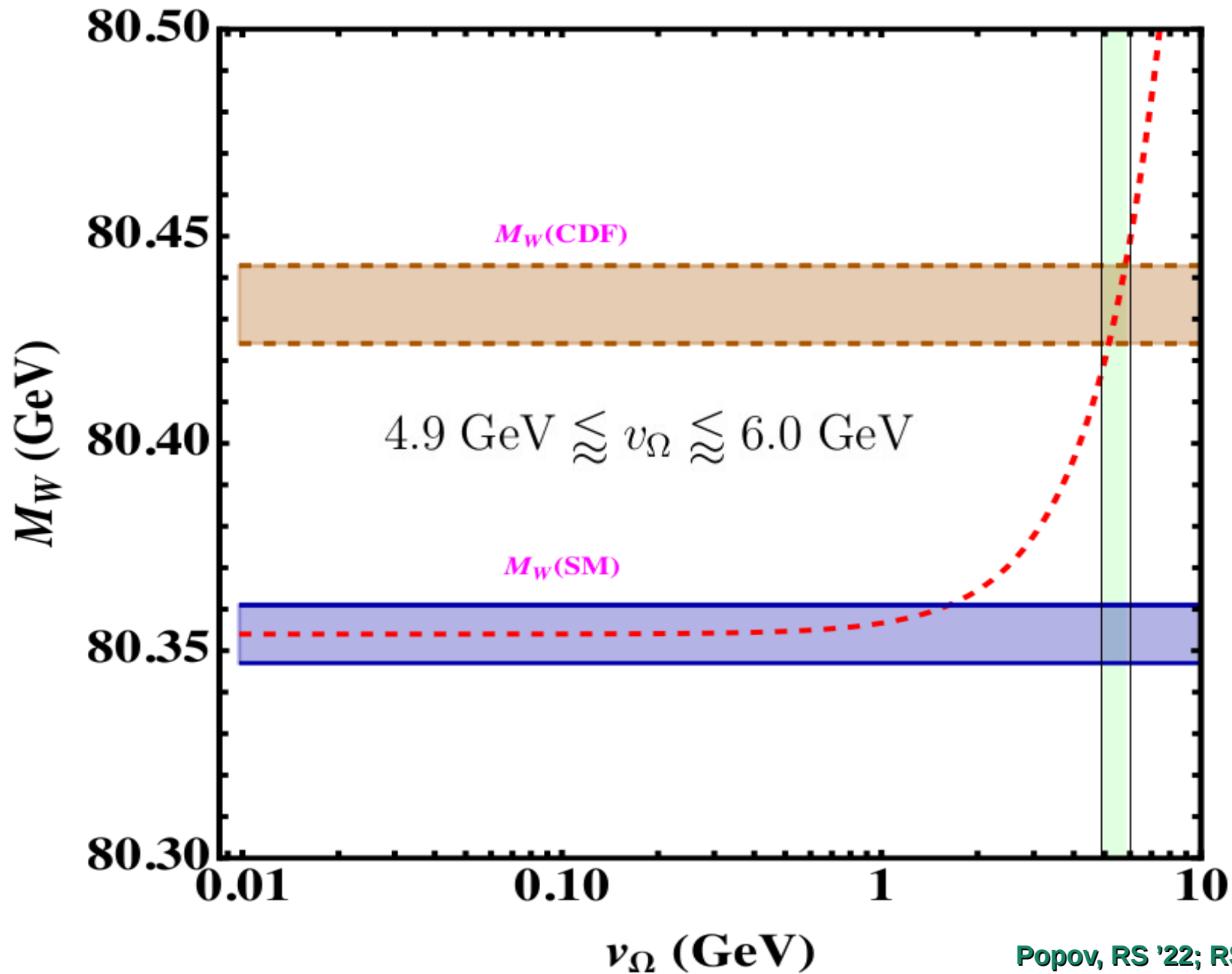
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$$\rho = \frac{\sqrt{v_\Phi^2 + 4v_\Omega^2}}{v_\Phi} \approx 1 + 2 \frac{v_\Omega^2}{v_\Phi^2} \quad M_{W,\text{CDF}}^2 - M_{W,\text{SM}}^2 = g^2 v_\Omega^2$$

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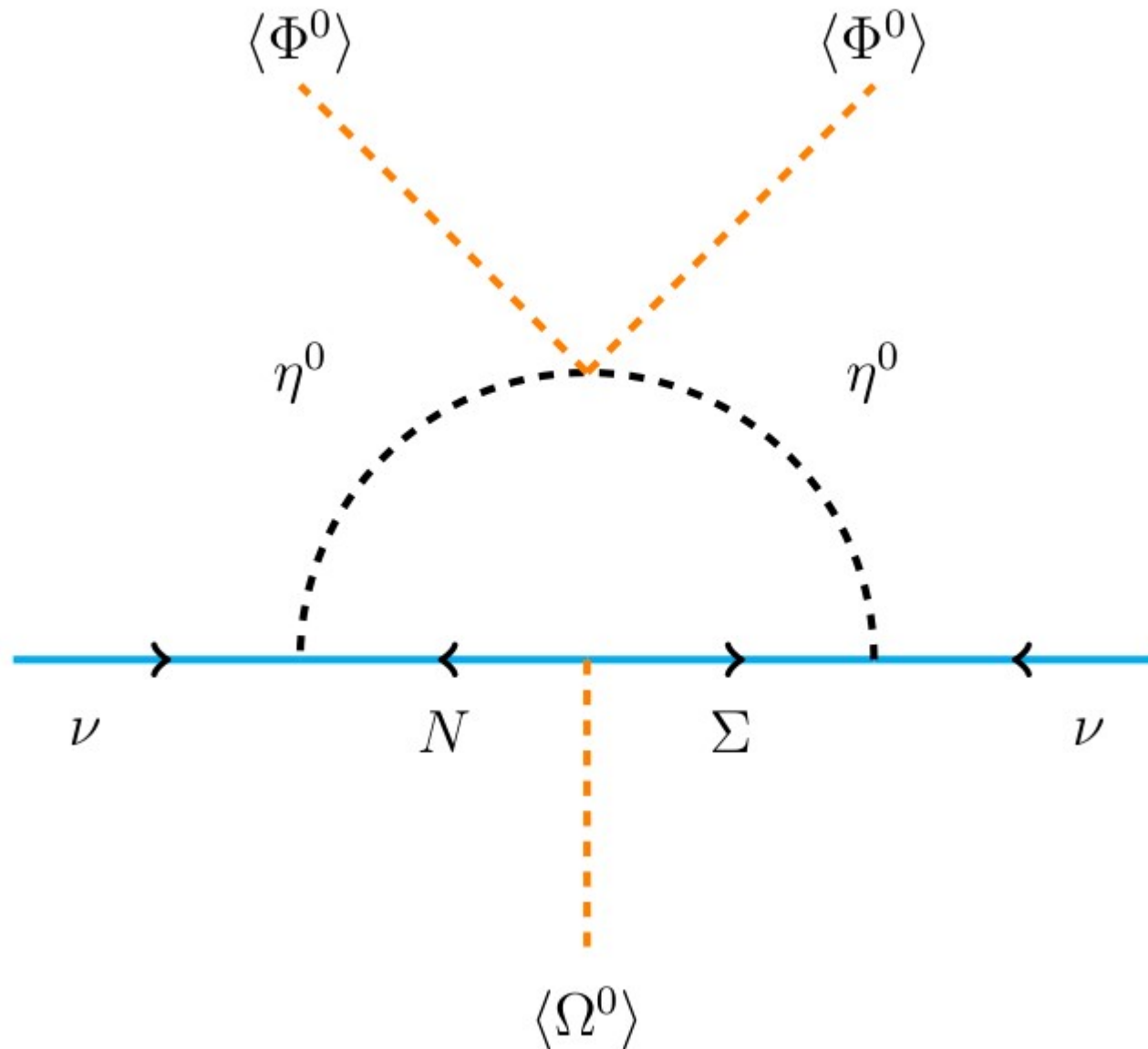


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  - Example: Singlet-Triplet Scotogenic model [Ma '09, Valle et.al. '09]

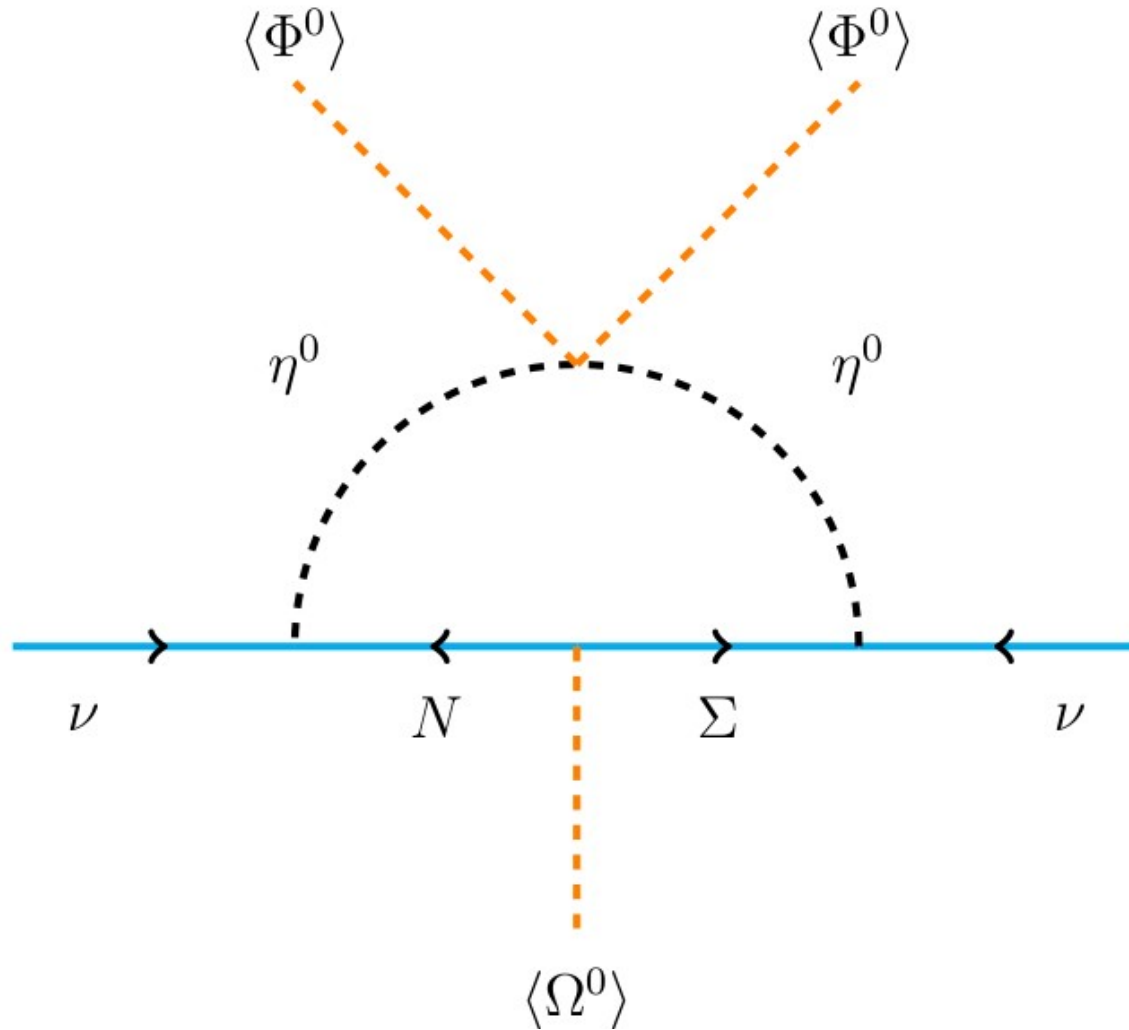
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$$(\mathcal{M}_\nu)_{\alpha\beta} = \sum_{\sigma=1}^2 \left( \frac{ih_{\alpha\sigma}}{\sqrt{2}} \right) \left( \frac{-ih_{\beta\sigma}}{\sqrt{2}} \right) \frac{m_{\chi_\sigma}}{(4\pi)^2} \left[ \frac{m_{\eta_R}^2 \ln \left( \frac{m_{\chi_\sigma}^2}{m_{\eta_R}^2} \right)}{m_{\chi_\sigma}^2 - m_{\eta_R}^2} - \frac{m_{\eta_I}^2 \ln \left( \frac{m_{\chi_\sigma}^2}{m_{\eta_I}^2} \right)}{m_{\chi_\sigma}^2 - m_{\eta_I}^2} \right]$$

# Triplet Scalars: W mass

- Including loop corrections the W mass: Tree level correction + loop corrections parametrized in terms of S, T, U

$$M_W = M_W^{\text{SM}} \left[ \frac{\sqrt{v_\Phi^2 + 4v_\Omega^2}}{v_\Phi} - \frac{\alpha_{\text{em}}}{4(c_W^2 - s_W^2)} (S - 1.55T - 1.24U) \right]$$

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$$S \simeq \underbrace{\frac{1}{12\pi} \log \left( \frac{m_{\eta^0}^2}{m_{\eta^+}^2} \right)}_{\text{Scalar doublet contribution}} + \frac{1}{18\pi}$$

Scalar doublet contribution

$$T \simeq \underbrace{\frac{1}{6\pi} \frac{1}{\sin^2(\theta_W) \cos^2(\theta_W)} \frac{\Delta M}{M_Z^2}}_{\text{Scalar triplet contribution}} + \underbrace{\frac{2\sqrt{2}G_F}{(4\pi)^2 \alpha_{\text{em}}} \left[ \frac{m_{\eta^0}^2 + m_{\eta^+}^2}{2} - \frac{m_{\eta^0}^2 m_{\eta^+}^2}{m_{\eta^+}^2 - m_{\eta^0}^2} \log \left( \frac{m_{\eta^+}^2}{m_{\eta^0}^2} \right) \right]}_{\text{Scalar doublet contribution}}$$

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$$U \simeq \underbrace{\frac{\Delta M}{3\pi M_H^\pm}}_{\text{Scalar triplet contribution}} \quad (37)$$

Scalar triplet contribution

$$+ \underbrace{\frac{1}{12\pi} \left[ -\frac{5m_{\eta^+}^4 - 22m_{\eta^+}^2 m_{\eta^0}^2 + 5m_{\eta^0}^4}{3(m_{\eta^+}^2 - m_{\eta^0}^2)^2} + \frac{(m_{\eta^+}^2 + m_{\eta^0}^2)(m_{\eta^+}^4 - 4m_{\eta^+}^2 m_{\eta^0}^2 + m_{\eta^0}^4)}{(m_{\eta^+}^2 - m_{\eta^0}^2)^3} \log \left( \frac{m_{\eta^+}^2}{m_{\eta^0}^2} \right) \right]}_{\text{Scalar doublet contribution}}$$

Scalar doublet contribution

# Triplet Scalars: $W$ mass

- Global fits for  $S, T, U$  including CDF-II results:

$$S = 0.06 \pm 0.10, \quad T = 0.11 \pm 0.12 \quad \text{and} \quad U = 0.14 \pm 0.09$$

Lu et.al. '22

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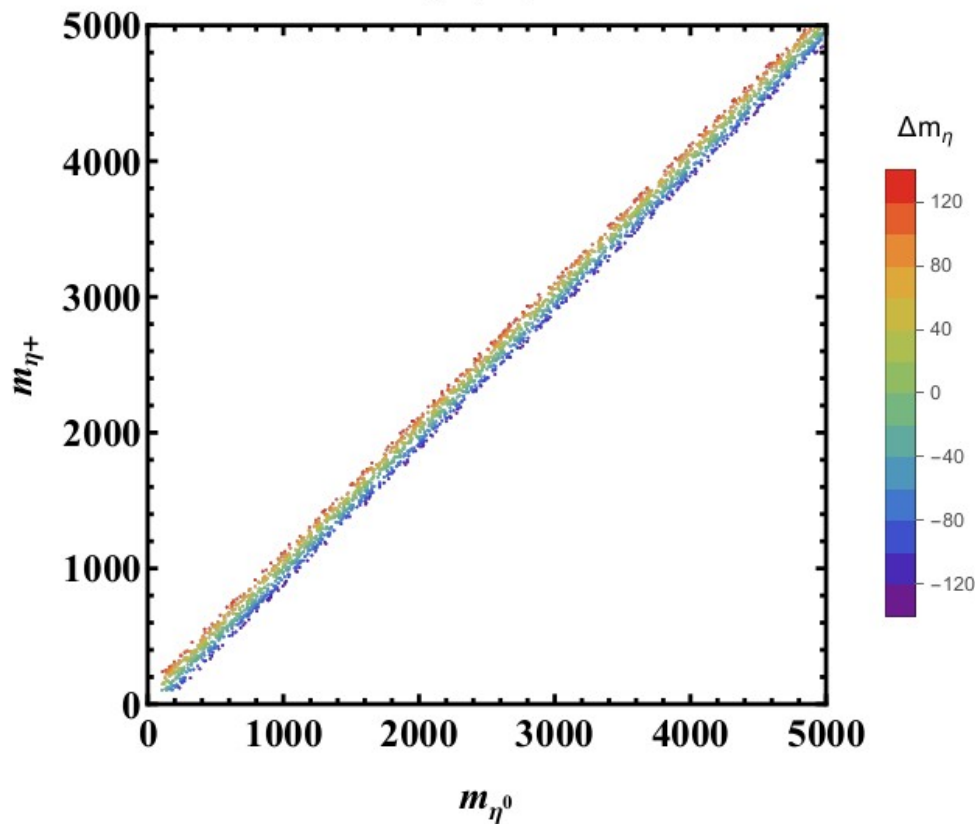
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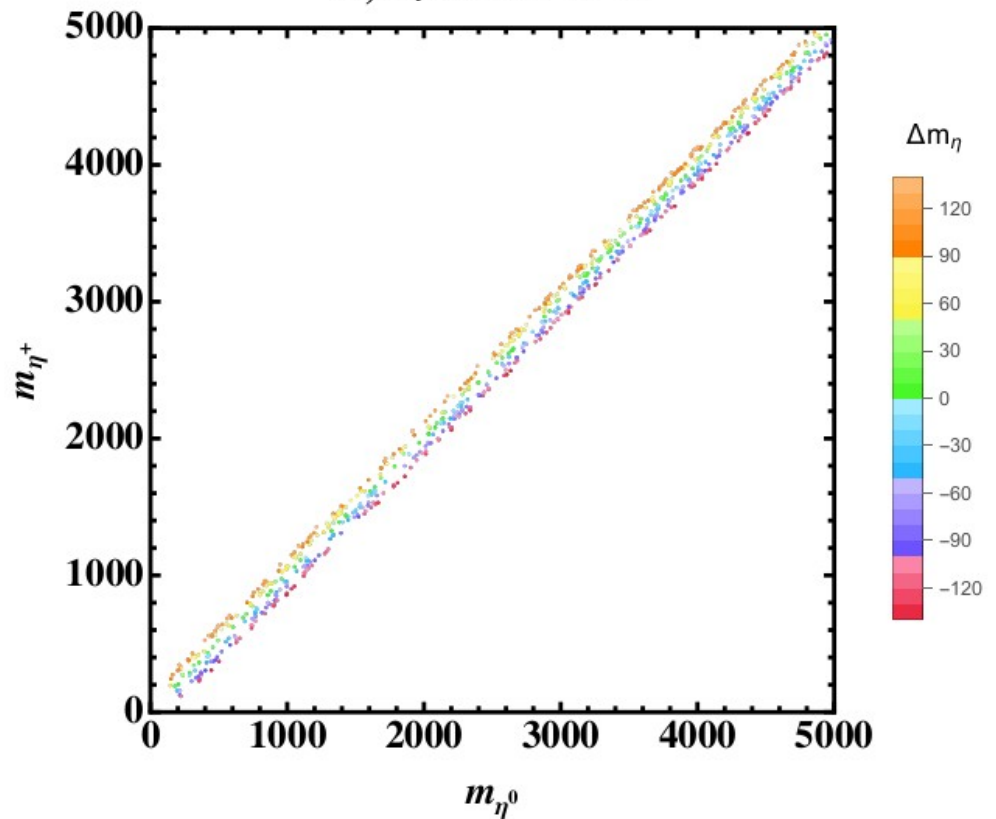
Lu et.al. '22

- Compressed spectrum for the scalars in the model RS et.al. '22

S,T,U≠0



S,T≠0 and U=0



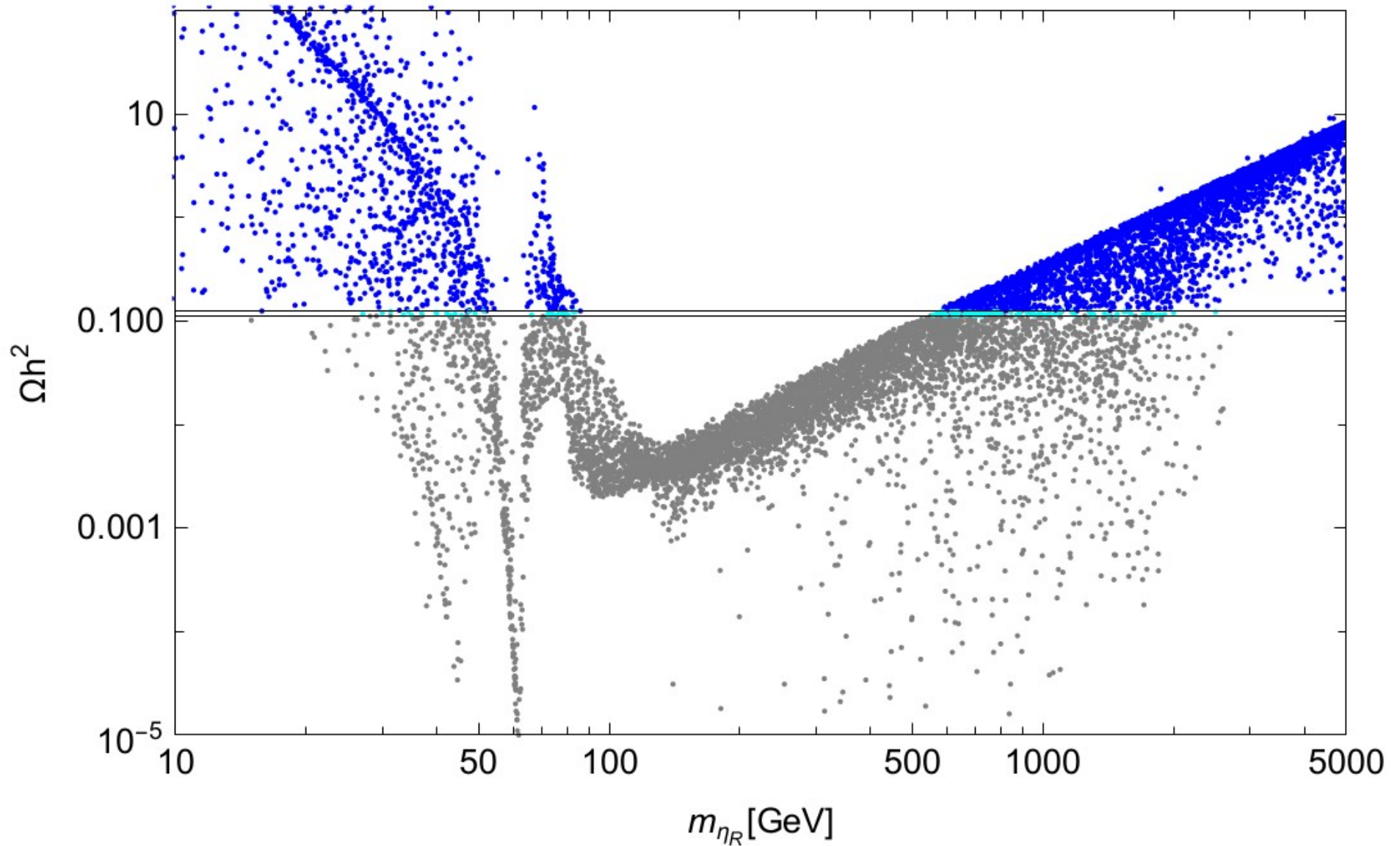
# Triplet Scalars: Dark Matter

- Doublet scalar can be good dark matter candidate [RS et.al. '22](#)



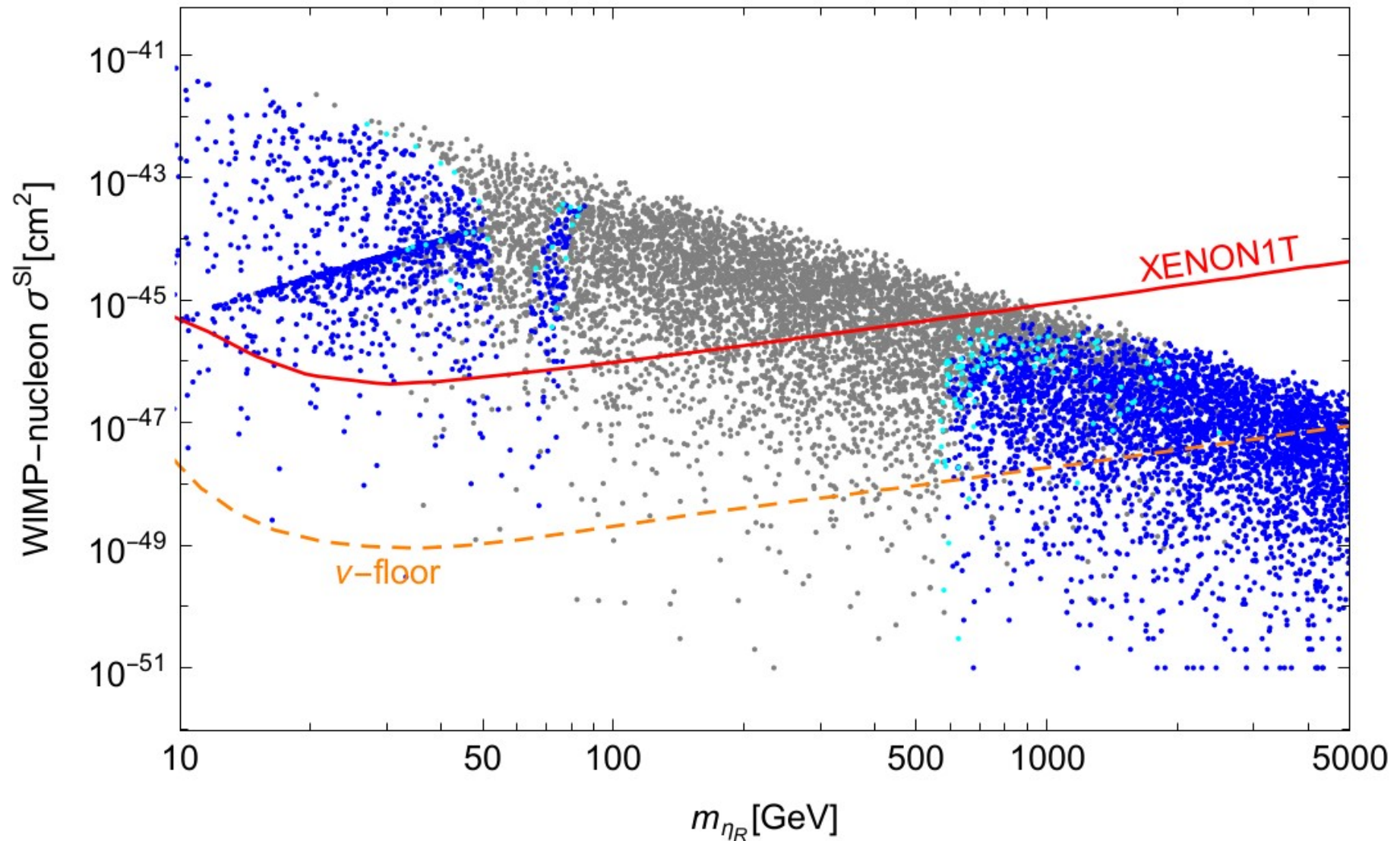
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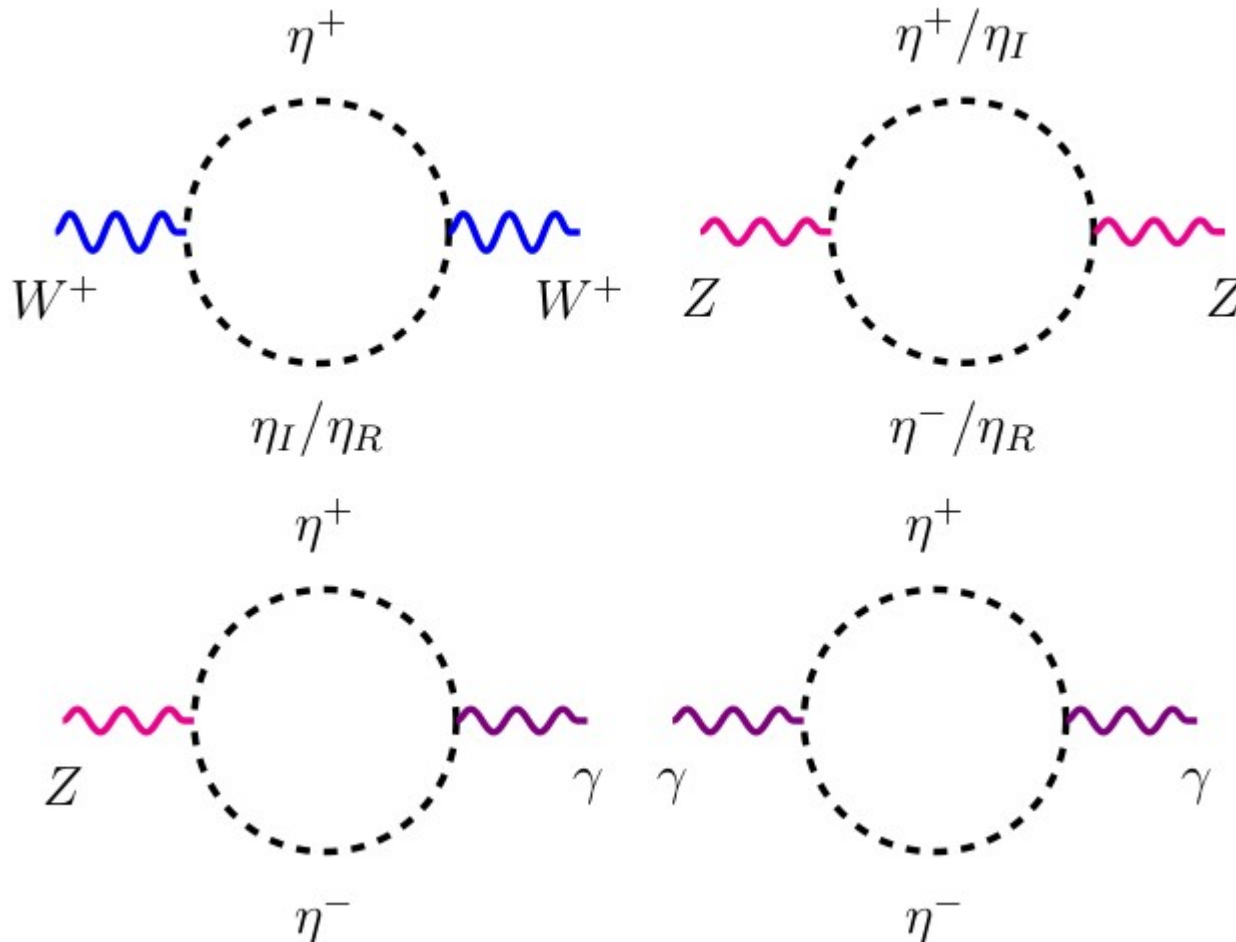
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RS. et. al. '22, '22, '22

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$$U = \frac{1}{12\pi} \left[ -\frac{5m_{\eta^+}^4 - 22m_{\eta^+}^2 m_{\eta^0}^2 + 5m_{\eta^0}^4}{3(m_{\eta^+}^2 - m_{\eta^0}^2)^2} + \frac{(m_{\eta^+}^2 + m_{\eta^0}^2)(m_{\eta^+}^4 - 4m_{\eta^+}^2 m_{\eta^0}^2 + m_{\eta^0}^4)}{(m_{\eta^+}^2 - m_{\eta^0}^2)^3} \log \left( \frac{m_{\eta^+}^2}{m_{\eta^0}^2} \right) \right]$$

$$m_W = m_W^{\text{SM}} \left[ 1 - \frac{\alpha_{em}}{4(c_W^2 - s_W^2)} (S - 1.55T - 1.24U) \right]$$

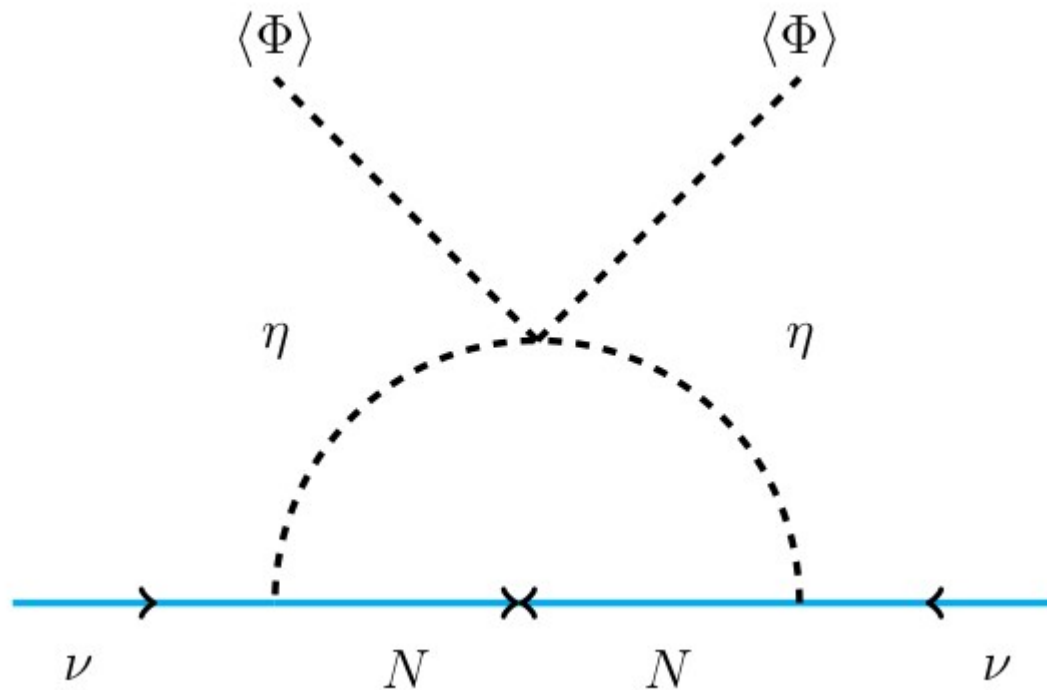
# Doublet Scalar: Scotogenic Model

- Scotogenic model: A simple model which explain both neutrino mass and dark matter E. Ma '06
  - Dark matter generates neutrino mass at one loop



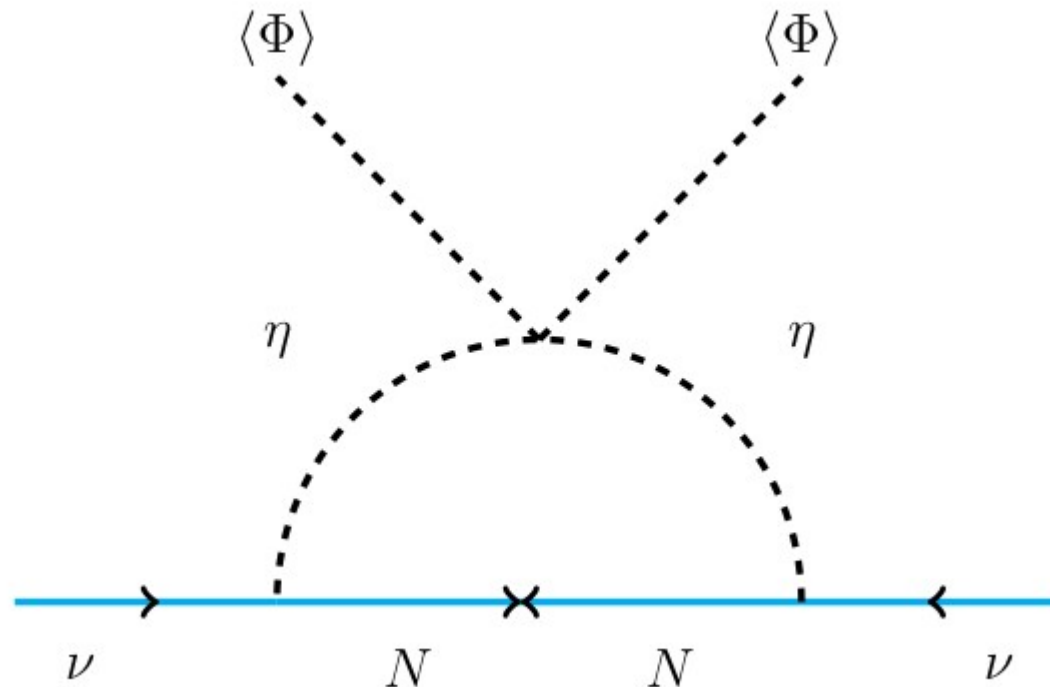
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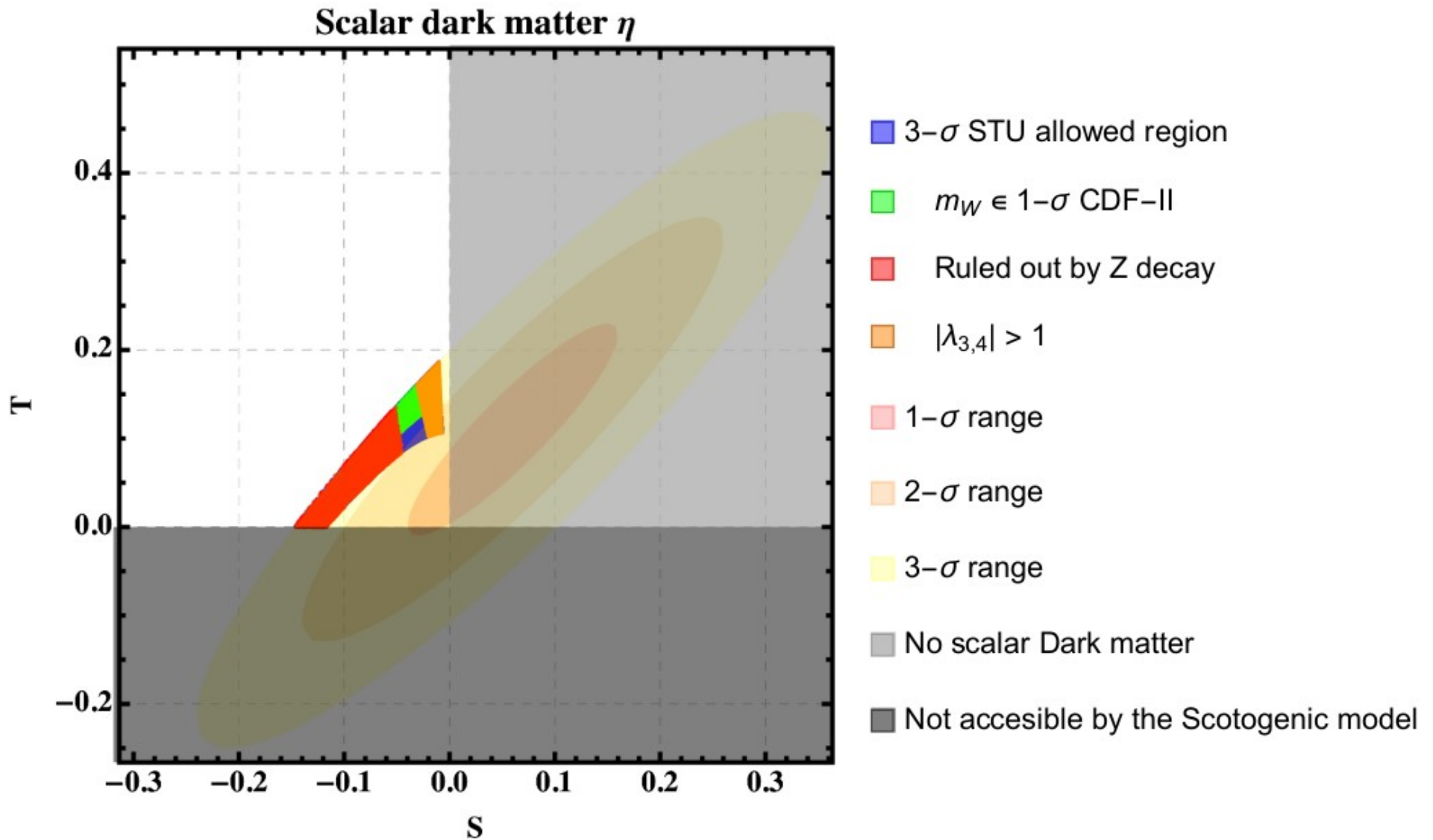


$$\mathcal{M}_{\alpha\beta}^{\nu} = \sum_i \frac{Y_{\alpha i}^N Y_{\beta i}^N}{32\pi^2} m_{N_i} \left[ \frac{m_{\eta_R}^2}{m_{\eta_R}^2 - m_{N_i}^2} \log \left( \frac{m_{\eta_R}^2}{m_{N_i}^2} \right) - \frac{m_{\eta_I}^2}{m_{\eta_I}^2 - m_{N_i}^2} \log \left( \frac{m_{\eta_I}^2}{m_{N_i}^2} \right) \right]$$

# Doublet Scalar: $W$ mass

- Doublet scalars lead to loop corrections

RS. et. al. '22, '22, '22

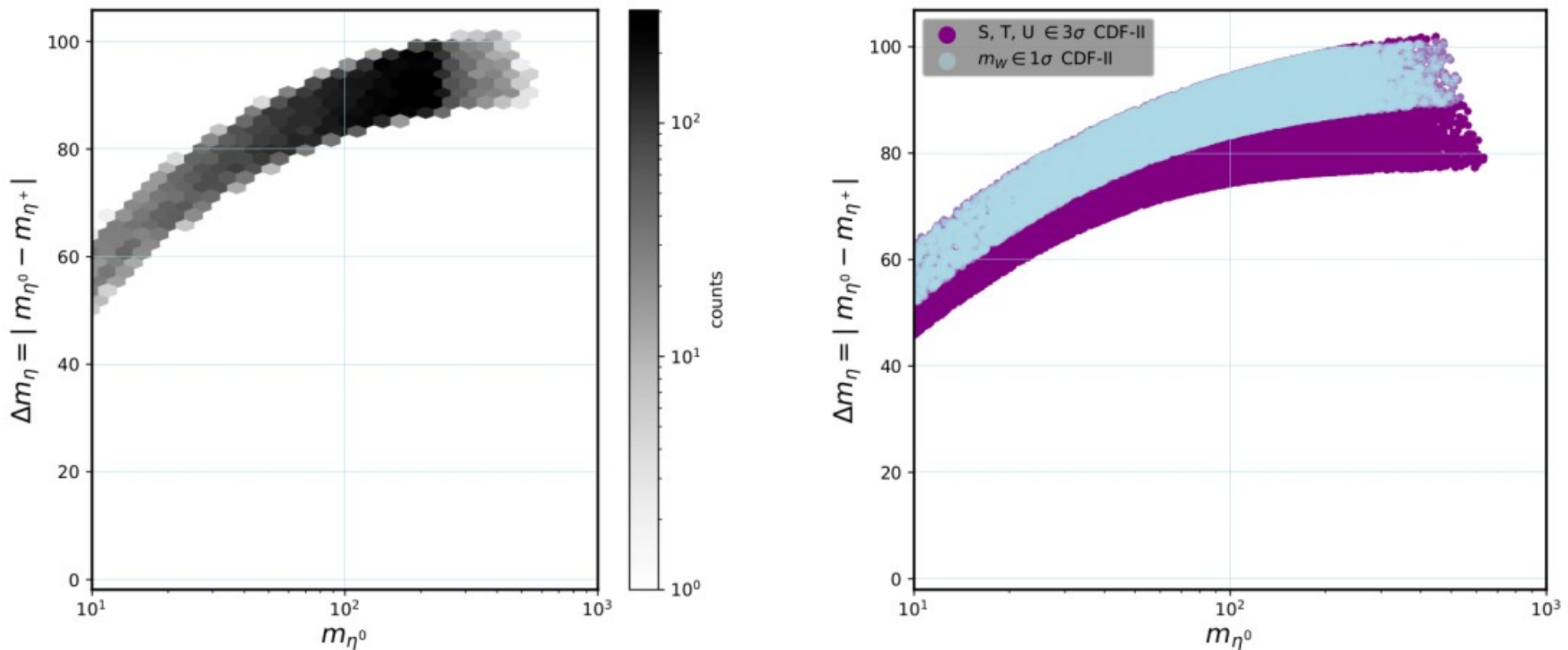


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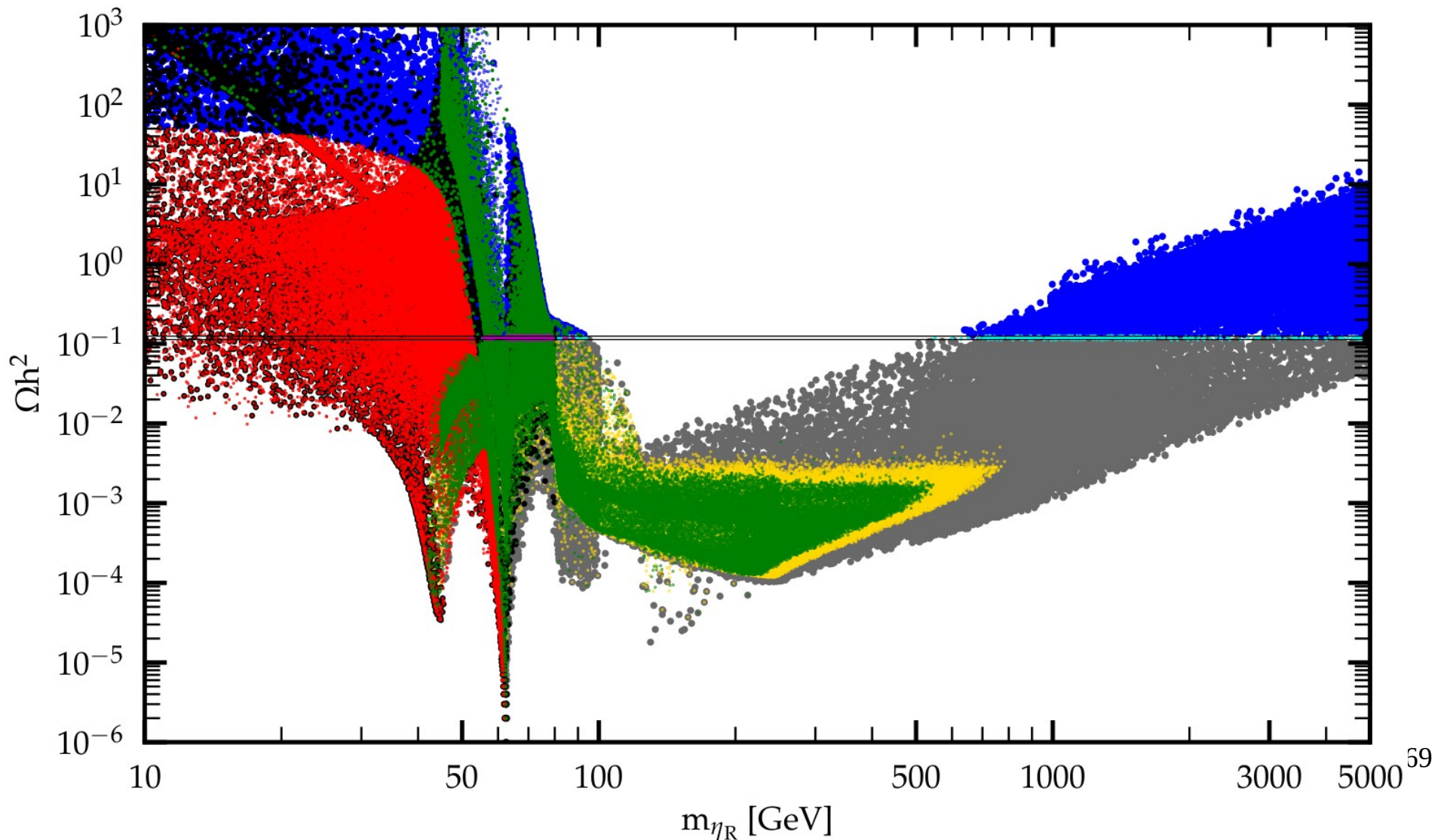
Scalar dark matter  $\eta$



# Doublet Scalar Dark Matter

- Doublet scalars can simultaneously be dark matter and satisfy the CDF-II measurement

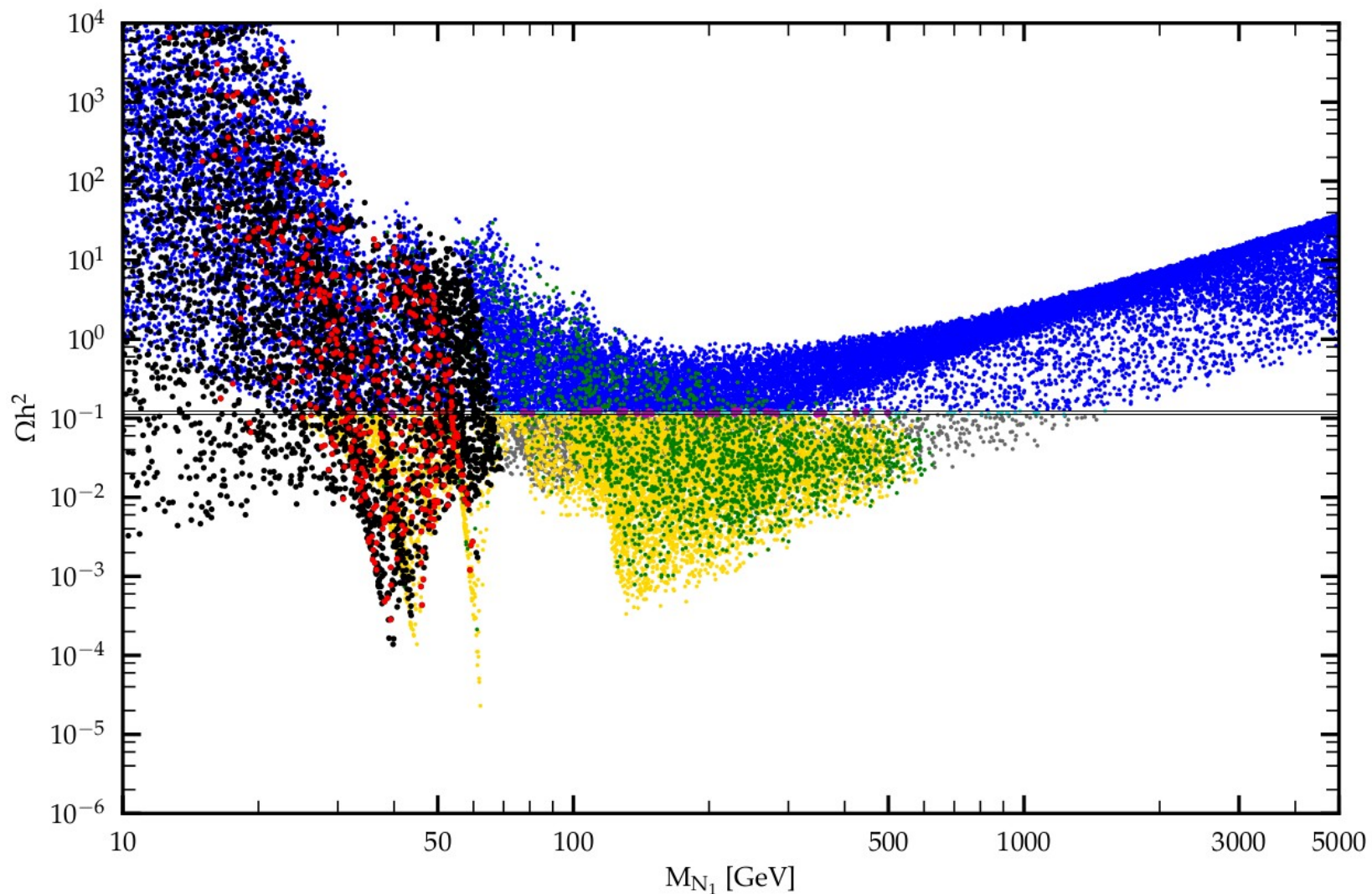
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# Fermionic Dark Matter

- Doublet scalars provide loop correction to  $W$  mass while the dark fermion is dark matter

RS. et. al. '22, '22, '22



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- **Presence of a Z' originating from say a new U(1) gauge symmetry can change the Z mass at tree level**

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E. Ma, RS. '14

Fields	( $SU(3)_C \otimes SU(2)_L \otimes U(1)_Y \otimes U(1)_{B-L}$ )
$L_L$	(1, 2, -1/2, -1)
$Q_L$	(3, 2, 1/6, 1/3)
$e_R$	(1, 1, -1, -1)
$u_R$	(3, 1, 2/3, 1/3)
$d_R$	(3, 1, -1/3, 1/3)
$\nu_R^1$	(1, 1, 0, 5)
$\nu_R^{2,3}$	(1, 1, 0, -4)
$\Phi$	(1, 2, 1/2, 0)
$\varphi$	(1, 2, 1/2, -3)
$\sigma$	(1, 1, 0, 3)
$\chi_d$	(1, 1, 0, 1/2)

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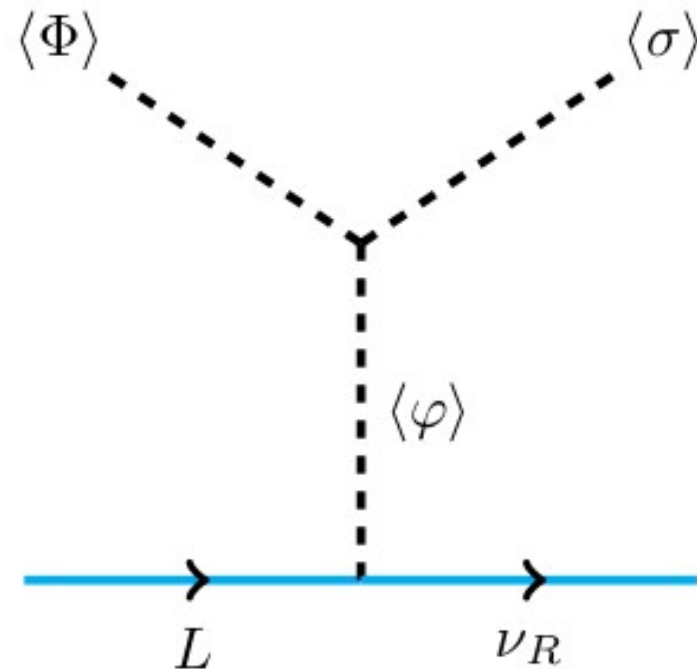
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# Gauged B-L Symmetry: Z mass

- Due to presence of new Z' the neutral gauge boson mass matrix becomes

Mandal, Prajapati, RS '22

$$\mathcal{M}_V^2 = \frac{v^2}{4} \begin{bmatrix} g'^2 & -gg' & -6u^2g'g_x \\ -gg' & g^2 & 6u^2gg_x \\ -6u^2g'g_x & 6u^2gg_x & 36b^2g_x^2 \end{bmatrix}, \text{ where } u = \frac{v_\varphi}{v}, \text{ and } b^2 = u^2 + \frac{v_\sigma^2}{v^2}$$

- This leads to

$$\begin{bmatrix} A^\mu \\ Z^\mu \\ Z'^\mu \end{bmatrix} = \begin{bmatrix} \cos \theta_w & \sin \theta_w & 0 \\ -\cos \alpha' \cos \theta_w & \cos \alpha' \cos \theta_w & -\sin \alpha' \\ -\sin \alpha' \sin \theta_w & \sin \alpha' \cos \theta_w & \cos \alpha' \end{bmatrix} \begin{bmatrix} B^\mu \\ W_3^\mu \\ X^\mu \end{bmatrix}$$

- Changing the mass of Z as

$$M_A = 0, \quad M_Z^2 = \frac{v^2}{8} \left( A' - \sqrt{B'^2 + C'^2} \right) \quad \text{and} \quad M_{Z'}^2 = \frac{v^2}{8} \left( A' + \sqrt{B'^2 + C'^2} \right),$$

where,  $A' = 36b^2g_x^2 + (g^2 + g'^2)$ ,  $B' = 36b^2g_x^2 - (g^2 + g'^2)$  and  $C' = 12g_xu^2\sqrt{g^2 + g'^2}$ .

# Gauged B-L Symmetry: Z mass

- **The modified Z mass in turn feeds into a erroneous SM expectation for W mass** Mandal, Prajapati, RS '22

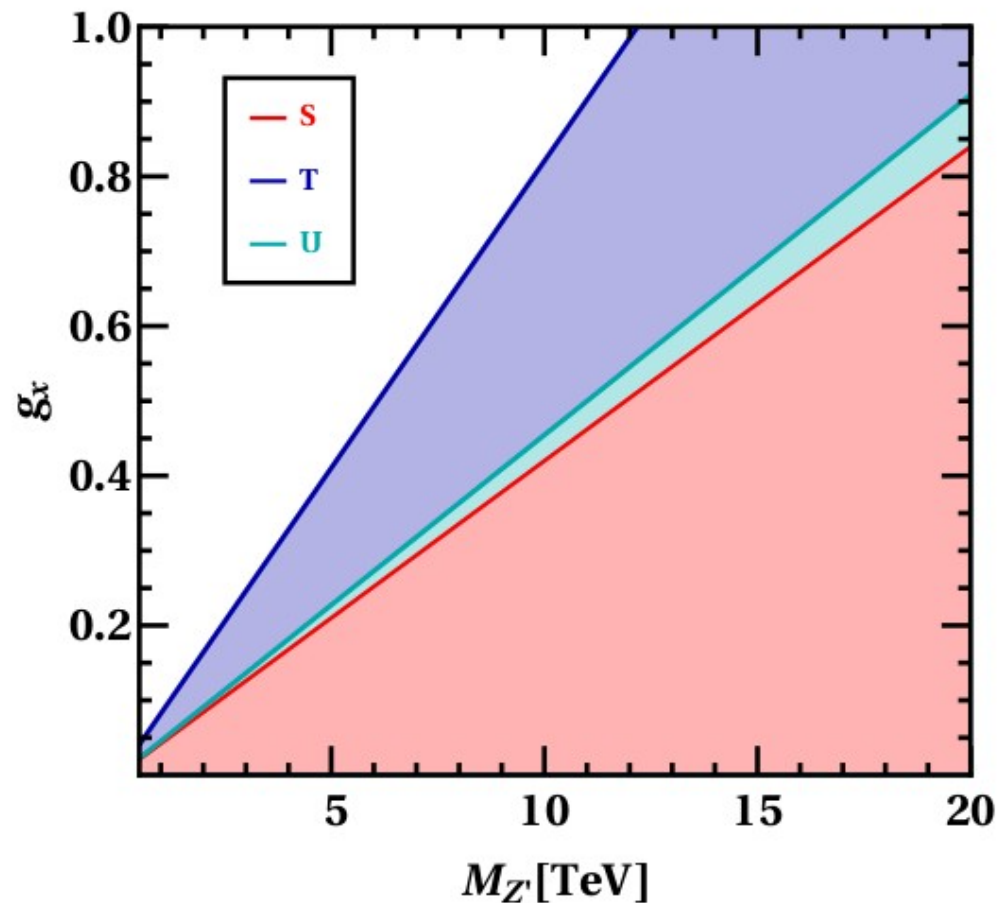
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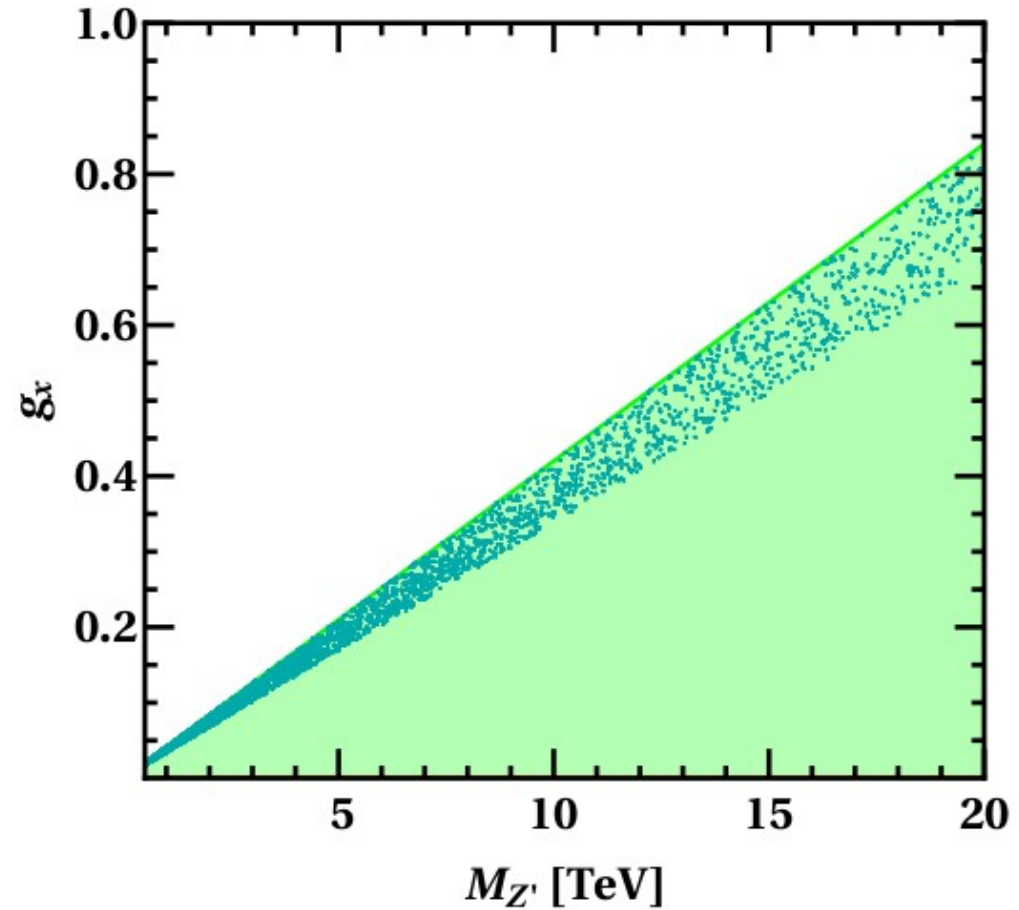
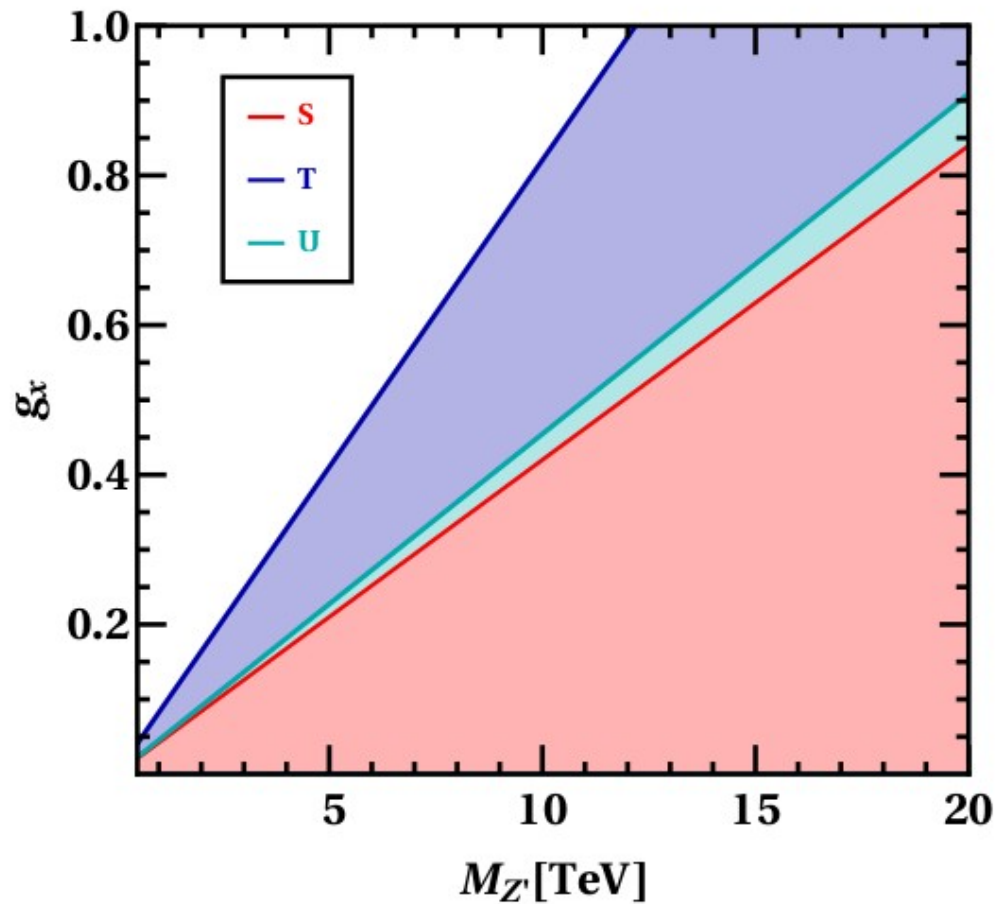
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# Conclusions

- The CDF-II  $W$  mass measurement disagrees with both
  - Other experimental measurements
  - Current SM expectation of  $W$  mass
- In my opinion it is too early to say if this is conclusive proof of new physics
- Still, if this is new physics it can be accommodated in several ways. I discussed few such options
  - Triplet scalars modifying  $W$  mass at tree level
  - Doublet scalars leading to loop corrections
  - Modification in  $Z$  mass leading to  $W$  mass anomaly
- Updated ATLAS results and CMS measurements will hopefully clarify the situation

**Thank You**