

Soeren Prell (Iowa State University) Belle II Summer Workshop Virginia Tech June 23-27, 2025

What is CP Violation ?

- Discrete symmetries (C,P,T) give insight into nature of interactions
 - CPT theorem:

"any Lorentz-invariant local quantum field theory with a Hermitian Hamiltonian must have CPT symmetry"

Strong and EM interactions are invariant under C, P, and T

- *Parity violation was discovered by Wu in* β *decay in* 1957
 - Predicted by Yang and Lee to solve the $\tau \theta$ puzzle
 - Structure of weak interaction (V-A) implies PV
 - Combined symmetry of C and P (CP) still seemed to hold
- *CPV* was discovered in K_L (*CP*-odd) decays to 2π (*CP*-even) ٠ by Cronin and Fitch in 1964
 - Small effect: $BR(K_L \rightarrow 2\pi) = 0.3 \%$ (c.f. $BF(\pi l \nu + 3\pi) = 99.7\%$)
 - Not understood at the time

 $C: a \leftrightarrow \bar{a} \ (\bar{a} \text{ is the antiparticle of } a)$

- $P: \mathbf{x} \to -\mathbf{x}$ (parity inversion of spatial coordinates)
- $T: t \to -t$ (motion or "time" reversal).



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CP Violation (S. Prell)

εy 0.2

1.20

The Matter - Universe

- 13.8 × 10 ⁹ years ago big bang produced matter and antimatter in equal amounts
 - matter-antimatter symmetric EM and strong interactions



- ... but today we observe the **absence** of
 - anti-nuclei amongst cosmic rays in our galaxy
 - *intense* γ-ray *emission* due to annihilation of distant galaxies in collision with antimatter

Today, matter dominates !



"A 10⁵³lbs Gorilla in the room"





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Where did all the anti-matter go?

• Baryon to photon ratio determined from microwave anisotropy

$$\eta = \frac{N_{baryons}}{N_{photons}} = (6.5^{+0.4}_{-0.3}) \times 10^{-10}$$
WMAP

Almost all matter annihilated with anti-matter..., but not all !

• Sakharov showed that generation of a net baryon number requires:

1. Baryon number violating processes (e.g. proton decay)

Sakharov, Pisma Zh. Eksp. Teor. Fiz. 5, 32 (1967)

- 2. Non-equilibrium state during the expansion of the universe
- 3. Difference in interaction rates for particles and antiparticles (C and CP violation)

Note, SM CPV is unlikely to be sufficient to explain universe matter content. However, CPV from New Physics that played a role in the early universe might.

CPV in the Standard Model

• Kobayashi and Maskawa proposed a threegeneration **complex** quark mixing matrix between strong and weak quark eigenstates

Progress of Theoretical Physics, Vol. 49, No. 2, February 1973

CP-Violation in the Renormalizable Theory of Weak Interaction

Makoto KOBAYASHI and Toshihide MASKAWA Department of Physics, Kyoto University, Kyoto

(Received September 1, 1972)

- *CKM matrix elements modify weak charged current decay amplitude*
- Only 3 quark flavors (u, d, s) were discovered, and quark model was not widely accepted, yet
- Today, all observed CPV can be described by the CKM matrix
 - No CPV in strong interactions
 - No CPV in the lepton sector (maybe soon)

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$$V_{CKM} = \begin{pmatrix} \mathbf{d} & \mathbf{s} & \mathbf{b} \\ V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \mathbf{t}$$

Cabbibo-Kobayashi-Maskawa (CKM) matrix



g = weak Fermi coupling constant

Decay amplitude \propto (complex) V_{qp} Decay rate $\propto |V_{qp}|^2$

Complex matrix elements can lead to <u>different</u> <u>BFs for particles and antiparticles</u>

 \rightarrow CP violation

A useful parameterization ("Wolfenstein")

Unitary 3×3 matrix has only four independent parameters (λ , A, ρ , η)

$$\mathbf{V}_{CKM} = \begin{pmatrix} 1 - \lambda^{2}/2 & \lambda & A\lambda^{3}(\rho - i\eta) \\ -\lambda & 1 - \lambda^{2}/2 & A\lambda^{2} \\ A\lambda^{3}(1 - \rho - i\eta) & -A\lambda^{2} & 1 \end{pmatrix}^{Wolfenstein,} PRL 51, 1945 (1983) + O(\lambda^{4})$$

$$= \begin{pmatrix} |V_{ud}| & |V_{us}| & e^{-i\gamma} |V_{ub}| \\ -|V_{cd}| & |V_{cs}| & |V_{cb}| \\ e^{-i\beta} |V_{td}| & -|V_{ts}| & |V_{tb}| \end{pmatrix}^{Single complex phase \eta} generates all SM CPV$$
relative magnitudes
$$\begin{bmatrix} 2 & complex matrix elements: \\ V_{td} and V_{ub} \end{bmatrix}^{CKM}$$

Careful ! This is an <u>arbitrary</u> phase convention. But it allows to easily see where CPV occurs.

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SM CP violation is very <u>predictive</u>: complex CKM phase η is related to apex of UT

Can be determined from sides or angles ! Allows for consistency checks!

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3 Types of CP Violation

- 1) CPV in decay
- 2) CPV in mixing
- 3) Mixing-induced CPV, or CPV in the interference between decay with and without mixing

1) is also referred to as "direct CPV" and 2) & 3) as "indirect CPV"

CPV in decay

• Difference between magnitude of a decay amplitude and its CP-conjugate amplitude



Direct CPV:
$$|A| \neq |\overline{A}|$$
; $A_{CP} \equiv \frac{\Gamma(\overline{B} \to \overline{f}) - \Gamma(B \to f)}{\Gamma(\overline{B} \to \overline{f}) + \Gamma(B \to f)} \neq 0$

- Only type of CPV possible for charged particle decays
- *Relatively easy to measure: only BFs necessary*

Example:
$$\mathcal{B}(B^0 \to K^+\pi^-) \neq \mathcal{B}(\overline{B}{}^0 \to K^-\pi^+)$$

- Direct CPV can also show up in differential BFs (e.g. across Dalitz plot) or in individual orbital angular momentum waves for VV final states

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CPV in decay: weak and strong phases

• Easiest way to get CPV is with **2 interfering amplitudes** (e.g. tree and penguin) with different weak (CP-odd) **and** strong (CP-even) phases

$$A(B \to f) = A = A_1 + A_2$$
$$A(\bar{B} \to \bar{f}) = \bar{A} = \bar{A}_1 + \bar{A}_2$$

CP transformation: strong phase: $\delta \rightarrow \delta$ weak phase: $\phi \rightarrow -\phi$ $A_1 = |A_1|$ $A_2 = |A_2|e^{i\delta}e^{i\phi}$ $\overline{A_1} = |A_1|$ $\overline{A_2} = |A_2|e^{i\delta}e^{i\phi}$







- A_{CP} is large if the contributing amplitudes are of similar size $(r \approx 1, |A_1| \approx |A_2|)$
- Need external input on $|A_1|$, $|A_2|$, δ (usually not interesting) to measure ϕ (interesting)
- Observed in many b decays and recently in charm decays (LHCb, PRL 122, 211803 (2019))

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Neutral meson mixing ("2B or not 2B")

- Weak transitions can transform P into \overline{P} and vice versa
 - Physical neutral meson state ψ is a linear combination
- *Time-evolution given by non-hermitian Hamiltonian*
 - Diagonal terms give P mass M and width $\Gamma(1/\tau)$
 - Off-diagonal terms describe mixing (incl. CPV)
 - Eigenstates of Hamiltonian have defined $M_{H,L}$ and $\tau_{S,L}$ (if no CPV in mixing they are CP eigenstates)
 - *virtual amplitude:* $|\Delta M| = |2M_{12}|$
 - on-shell amplitude: $|\Delta\Gamma| = |2\Gamma_{12}|$
- Time-dependent $P \overline{P}$ mixing:



short-distance, virtual



$$\psi(t) = a(t) |P^0\rangle + b(t) |\bar{P}^0\rangle$$

$$\mathbf{H} = \begin{pmatrix} H_{11} & H_{12} \\ H_{21} & H_{22} \end{pmatrix} = \begin{pmatrix} M - \frac{i}{2}\Gamma & M_{12} - \frac{i}{2}\Gamma_{12} \\ M_{12}^* - \frac{i}{2}\Gamma_{12}^* & M - \frac{i}{2}\Gamma \end{pmatrix}$$

$$P(P \to \bar{P}) \propto e^{-\Gamma t} \left| \frac{q}{p} \right|^{2} \left[\cosh \frac{\Delta \Gamma t}{2} - \cos \Delta M t \right]$$

$$P(\bar{P} \to P) \propto e^{-\Gamma t} \left| \frac{p}{q} \right|^{2} \left[\cosh \frac{\Delta \Gamma t}{2} - \cos \Delta M t \right]$$

$$\frac{q}{p} = \sqrt{\frac{M_{12}^{*} - \frac{i}{2}\Gamma_{12}^{*}}{M_{12} - \frac{i}{2}\Gamma_{12}}}$$

$$V \text{ in mixing}$$

$$P(P \to \bar{P}) \neq P(\bar{P} \to P)$$

$$\left| \frac{p}{q} \right| \neq 1$$

$$M_{12} \left| |\Gamma_{12}| \sin(\theta_{M_{12}} - \theta_{\Gamma_{12}}) \neq 0$$

- *Results from interference between on-shell and virtual amplitudes*

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Neutral meson mixing comparison

Mixing asymmetry (no CPV):

$$a_{mix}(t) = \frac{P(P^0 \to P^0) - P(P^0 \to \overline{P}^0)}{P(P^0 \to P^0) + P(P^0 \to \overline{P}^0)} = \frac{\cos(x t/\tau)}{\cosh(y t/\tau)} \qquad \qquad x = \Delta M \tau$$
$$y = \Delta \Gamma \tau/2$$



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Neutral $B_{d,s}$ mixing and CPV



No CPV in B_{d/s} mixing !

Mixing-induced CP Violation



Time-dependent CP asymmetry

$$A_{CP} = \frac{\Gamma \ (\bar{B}^0 \to f_{CP}) - \Gamma \ (B^0 \to f_{CP})}{\Gamma \ (\bar{B}^0 \to f_{CP}) + \Gamma \ (B^0 \to f_{CP})}$$
$$= -C_f \cos (\Delta M t) + S_f \sin (\Delta M t)$$

single weak $|\lambda_{CP}| = 1$ amplitude: $C_f = 0$ $S_f = -\text{Im }\lambda_{CP}$

 $C_f \neq 0$ implies **direct** CPV

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Measurement of time-dependent CP Violation at Belle II



B decays sensitive to UT angles



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ϕ_1 , the B Factories' "CP or not CP"



$\sin 2\phi_1 from B^0 \rightarrow J/\psi K_{S,L} decays$

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The "gold-plated modes": $B \rightarrow J/\psi K_S^0$ and other $b \rightarrow c\bar{s}s$ transitions

- $B \to J/\psi(\to l^+l^-)K_S^0$ dominates sin $2\phi_1$ measurement
 - Relatively large BF
 - Low background
 - Small theoretical uncertainties
- LHCb error on sin 2φ₁ is now 0.014
 Dominated by statistical error
- Eventually, Belle II will be systematics limited

	No	Vertex	Leptonic
	$\operatorname{improvement}$	$\operatorname{improvement}$	categories
$S_{J/\psi K_S^0}$ (50 ab ⁻¹)			
stat.	0.0035	0.0035	0.0060
syst. reducible	0.0012	0.0012	0.0012
syst. irreducible	0.0082	0.0044	0.0040
$A_{J/\psi K_S^0}$ (50 ab ⁻¹)			
stat.	0.0025	0.0025	0.0043
syst. reducible	0.0007	0.0007	0.0007
syst. irreducible	$^{+0.043}_{-0.022}$	$^{+0.042}_{-0.011}$	0.011









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ϕ_2 , a story with some twists and turns





Original idea: time-dependent analysis of $B \rightarrow \pi^+\pi^-$

Only tree amplitudes were expected uto contribute to $B \rightarrow \pi^+\pi^ W^+$ - penguin amplitude was expected to be negligible ! $V_{ub} \propto e^{-i\gamma}$ dd*Weak phase* $2\phi_3$ *between* external tree (T) internal or color- $B \to \pi^+\pi^-$ and $\bar{B} \to \pi^+\pi^$ suppressed tree (C) combines with mixing phase $2\phi_1$ Weak phase: *Time-dependent CP asymmetries* $\lambda =$ would be - $S = \sin 2\phi_2$ and C = 0mixing decay $= e^{-i2\phi_1} e^{-i2\phi_3}$ Penguin (P) $= e^{i2\phi_2}$

The rise of the penguins



/ c²)

Events / (2 MeV

5







- Just before the start of data taking • of the B factories there was troublesome news
 - $\overline{B} \to K^+\pi^-$ and $B^+ \to \eta' K^+$ were discovered with unexpectedly large BFs
 - On the other hand $B \rightarrow \pi^+\pi^$ hadn't been seen, yet
 - Shortly after $B \to \pi^+\pi^-$ was discovered, $BF(B \rightarrow \pi^0 \pi^0)$ was found larger than expected
- *Penguins could not be neglected!* •
- But if penguins are large, ϕ_2 can't • be extracted from $C_{\pi^+\pi^-}$ and $S_{\pi^+\pi^-}$ (alone)



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ϕ_2 from $B \to \pi\pi$



 ϕ_2 can be extracted with an isospin decomposition of $B \to \pi^+\pi^-$

	isospin factors			
	α_T	α_C	α_P	
$B \to \pi^+ \pi^-$	$\sqrt{2}$	0	$\sqrt{2}$	
$B \to \pi^+ \pi^0$	1	1	0	
$B \to \pi^0 \pi^0$	0	1	-1	

$$S = \sqrt{1 - C^2} \times \sin\left(2\phi_2 - 2\Delta\phi_2^{+-}\right)$$

Gronau and London, PRL 65, 3381 (1990)



$B^0 \rightarrow \pi^+ \pi^- CPV$ and BF

PRD 87, 031103 (2013)





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Constraints on ϕ_2 from $B \rightarrow \pi\pi$ (now and with Belle II)



- Discrete ambiguities complicate constraints on ϕ_2
 - Belle II can substantially reduce uncertainties

 $S(B \rightarrow \pi^0 \pi^0)$ could resolve discrete ambiguities

- Very hard measurement
- Expected Belle II error ~0.3
- Very promising method recently (re)discovered
- $B \rightarrow \pi\pi$ was assumed to be B factories' best shot at ϕ_2
 - Other $b \rightarrow u\bar{u}d$ modes, like $B \rightarrow \pi\rho$ and $B \rightarrow \rho\rho$ would have additional problems (low π^0 efficiency, time-dep. Dalitz analysis, VV polarization amplitudes)

3 pleasant surprises in $B \rightarrow \rho \rho$

PRD 93, 032010 (2016)



Time-dependent CP asymmetry in ${\rm B^0} \to \rho^{\scriptscriptstyle +} \rho^{\scriptscriptstyle -}$

ϕ_2 World Average



 $\phi_2(PDG) = (85.2^{+4.8}_{-4.3})^{\circ}$

 ϕ_2 world average is dominated by $B \rightarrow \rho \rho$

From CKM fit (UT_{fit}, 2018): $\phi_2 = (90.1 \pm 2.2)^\circ$



Modes with $\pi^0(\rho^{\pm} \rightarrow \pi^{\pm}\pi^0)$ are much easier at Belle II than at LHCb Expected ϕ_2 error: $\Delta \phi_2 \sim 0.6^{\circ}$

The UT angle ϕ_3



Determination of ϕ_3 with $B^- \rightarrow D^{(*)}K^{(*)-}$

Measure ϕ_3 with charged B decays (interference between $\overline{D}K$ and DK intermediate states where \overline{D} and D decay to same final state)

One amplitude involves V_{ub}





Many D^0 final states investigated: CP eigenstates ($\pi\pi$, KK; K_S($\pi^0, \omega, \eta, \varphi$)) Flavor eigenstates ($K\pi$) Three-body decays (K_S $\pi\pi$, K_SKK, $\pi\pi\pi^0$)

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CP Violation (S. Prell)

Gronau & London, PLB 253, 483 (1991) Gronau & Wyler, PLB 265, 172 (1991) Atwood, Dunietz, & Soni, PRL 78, 3257 (1997), Atwood, Dunietz, & Soni, PRD 63, 036005 (2001) Giri, Grossman, Soffer, & Zupan, PRD 68, 054018 (2003) Belle, Poluektov et al., PRD 70, 072003 (2004) Bondar & Poluektov, EPJC 47,347 (2006) BPGGSZ analysis of $B \rightarrow D^{(*)} K^{(*)}$ with $D \rightarrow K_S^0 \pi^+ \pi^-$

D Dalitz plot has complicated structure of several interfering amplitudes

 Interference between D decays to Cabibbo-allowed, Cabibbo-suppressed and CP final states

 $\frac{\mathrm{d}\Gamma_{B^-}(m_+^2, m_-^2) \propto}{2r_B |A_+||A_-| (\cos \delta_D \cos(\delta_B + \phi_3) - \sin \delta_D \sin(\delta_B + \phi_3)) \,\mathrm{d}p} \qquad A_{\pm} = A_D \left(m_{\pm}^2, m_{\mp}^2 \right)$

- Model-independent method
 - Binned fit to Dalitz plot \rightarrow get δ_D for each bin from coherent $D\overline{D}$ production at CLEO-c or BESII



JHEP 2021, 169 (2021)

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ϕ_3 World Average



 $\phi_3(HFLAV) = (66.2^{+3.4}_{-3.6})^{\circ}$

From CKM fit (UT_{fit}, 2018): $\phi_3 = (65.8 \pm 2.2)^{\circ}$

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Angles of the UT show consistent picture



... also with the UT sides !



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Beyond SM CPV searches ("Hic sunt dracones")



Penguin-dominated $b \rightarrow sq\bar{q}$ decays

• Penguin loops can receive contributions from New Physics





No CPV expected in $b \rightarrow d\bar{q}q$ penguins (e.g. $B \rightarrow K_S^0 K_S^0$): Loop phase cancels mixing phase for dominant diagram

- Weak phase of $b \to s\bar{q}q$ penguin is same as for $b \to c\bar{c}s$ in SM
 - Top quark contribution dominates loop amplitude
 - *CP* asymmetries are also C = 0 and $S = -\eta_{CP} \sin 2\phi_1$
 - Contributions of sub-leading diagrams could change S and C

Penguin-dominated $b \rightarrow sq\bar{q}$ decay: $B \rightarrow \eta' K^0$



CP Woldtation (Sr Phell)

CP Asymmetries in Penguin Decays

- Theoretical predictions for $\delta S = S_{sq\bar{q}} \sin 2\phi_1$ are small and typically positive
 - Large δS could be evidence for new physics
- Most significant difference in "naïve" $S_{sq\bar{q}}$ average reached in 2004
 - Caused a lot of excitement
 - Neglects theo. uncertainties and correlation of experimental uncertainties
 - *HFLAV: "We do not advocate its use, and provide it only for academic interest.* **Use with extreme caution, if at all**."
- More precise measurements have since decreased significance below 1σ
 - Still a good place to look for New Physics
 - Note, some measurements come from complicated 3-body time-dependent Dalitz analyses



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o→ccs	World Av	rage				(0.70 ± 0.02
Ŷ	BaBar			⊷	-	0.66±0	0.17 ± 0.07
- +	Belle				*		0.90 +0.09
Ŷ	BaBar			⊢ ★-1		0.57±0	0.08 ± 0.02
َب_ ع`	Belle			н	-	0.68±0	0.07 ± 0.03
×	BaBar				*	- 0.94	+0.21 -0.24 ± 0.06
×	Belle			-		0.71 ± (0.23 ± 0.05
° ×	BaBar				•	0.55 ± 0	0.20 ± 0.03
β	Belle				-	0.67±0	0.31 ± 0.08
Ł	BaBar				0.3	5 +0.26 ± (0.06 ± 0.03
ъ.	Belle			H	 0.6	4 +0.19 ± (0.09 ± 0.10
×	BaBar				-	0.55	+0.26 -0.29 ± 0.02
3	Belle			⊢	+	- 0.9 1 ± (0.32 ± 0.05
Š	BaBar			-			0.74 +0.12
÷	Belle			•— *	-		0.63 ^{+0.16} -0.19
, K _s	BaBar		•	*	0.48	± 0.52 ± (0.06 ± 0.10
ͺ _Ϛ ϗͺ	BaBar			*	0.20	± 0.52 ± 0	0.07 ± 0.07
Ϋ́Υ	B aBar	*				-0.72 ± 0	0.71 ± 0.08
<u>ہ</u>	Belle			•	*		+0.27 -0.31 ± 0.11
∮π [®] K _s	BaBar				*		0.97 +0.03
π [⁺] π [⊺] Ϗ _s Ν	vp₿aBar				0.01	± 0.31 ± (0.05 ± 0.09
×	BaBar			⊷+	-	0.65 ± 0	0.12 ± 0.03
t t	Belle			-	*		0.76 +0.14
o→dd <u>ş</u>	Naïve av	erage		•*		(0.65 ± 0.04
2		1		`	-	1	2
-2	-	I	C)	1		2

Belle II potential for CPV in penguin decays

$B^0 \rightarrow J/\psi K_S$ (the "Golden" mode): \rightarrow constrains the UT	$A_{CP} = A \mathrm{c}$	$\mathrm{os}(\Delta M$	$\Delta t) + S$	$S\sin(\Delta N)$	$\Delta t)$		
J/ψ		WA (2017)	5 a	b^{-1}	50 e	b^{-1}
	Channel	$\sigma(S)$	$\sigma(A)$	$\sigma(S)$	$\sigma(A)$	$\sigma(S)$	$\sigma(A)$
	$J/\psi K^0$	0.022	0.021	0.012	0.011	0.0052	0.0090
\bar{B}^0	ϕK^0	0.12	0.14	0.048	0.035	0.020	0.011
	$\eta' K^0$	0.06	0.04	0.032	0.020	0.015	0.008
expected 50 ab ⁻¹ uncertainty: $\delta \phi_1 = 0.4^\circ$	ωK_S^0	0.21	0.14	0.08	0.06	0.024	0.020
(less than the current theory error of 1-2°)	$K^0_S \pi^0 \gamma$	0.20	0.12	0.10	0.07	0.031	0.021
$B^0 \rightarrow \phi K_S, \eta' K_S, \omega K_S, \pi^0 K_S$ ("penguin" modes):	$K_S^0 \pi^0$	0.17	0.10	0.09	0.06	0.028	0.018
$B^{\circ} \rightarrow \phi K_{S}, \eta' K_{S}, \omega K_{S}, \pi^{\circ} K_{S}$ ("penguin" modes):	- ~						



Belle II expected to improve precision in S from 10-20% to 2-3%

Radiative B decays: $TDCPV \text{ of } b \rightarrow s\gamma \text{ and } b \rightarrow d\gamma$



- Expect CPV to be small ($S \sim -2(m_s/m_b)sin(2\phi_l)$)
 - $-\gamma$ helicity dominantly LH for $b \rightarrow s(d)\gamma$, and RH for \overline{b}
 - $B^0 \rightarrow K^{*0}(K_S^0 \pi^0) \gamma$ behaves like effective flavor eigenstate (assuming dipole operator is dominant)
- Similar situation for $B^0 \rightarrow \rho^0 \gamma$
 - However, since weak phase from $b \rightarrow d\gamma$ decay amplitude cancels that from $B^0 \overline{B}{}^0$ mixing, CPV is suppressed further
- Observed CPV would be sign of NP amplitude emitting RH photons and with NP weak phase
- Belle II potential: $\sigma(S_{\rho\gamma}) \sim 0.07$

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All measurements consistent with no CPV, as expected in the SM





Direct CP relations: The $K\pi$ Puzzle

- First discovery of direct CPV in B decays with $B \rightarrow K^{\pm}\pi^{\mp}$ (BaBar, PRL 93 (2004) 131801; Belle, PRL 93 (2004) 131802)
 - *dCPV caused by interference between loop and tree amplitudes*



- Naive assumption that modes related by isospin of spectator quark have similar A_{CP} ... turned out to be wrong
- More accurate sum rule predicts $A_{CP}^{K^0\pi^0} = -0.138 \pm 0.025$ (with recent LHCb $K^{\pm}\pi^0$ measurement)

$$A_{CP}(K^{+}\pi^{-}) + A_{CP}(K^{0}\pi^{+}) \frac{\mathcal{B}(K^{0}\pi^{+})}{\mathcal{B}(K^{+}\pi^{-})} \frac{\tau_{0}}{\tau_{+}} = A_{CP}(K^{+}\pi^{0}) \frac{2\mathcal{B}(K^{+}\pi^{0})}{\mathcal{B}(K^{+}\pi^{-})} \frac{\tau_{0}}{\tau_{+}} + A_{CP}(K^{0}\pi^{0}) \frac{2\mathcal{B}(K^{0}\pi^{0})}{\mathcal{B}(K^{+}\pi^{-})}$$

• Expect error $\sigma(A_{CP}^{K^0\pi^0}) \sim 0.04$ with Belle II

Observables	Belle	Belle II		
	(2017)	$5 {\rm ~ab^{-1}}$	$50 \ {\rm ab}^{-1}$	
$\mathcal{A}_{\mathcal{CP}}(B \to K^0 \pi^0) [\%]$	$-0.05 \pm 0.14 \pm 0.05$	0.07	0.04	

 $A_{CP}(B^0 \to K^+\pi^-) = -0.084 \pm 0.004$ $A_{CP}(B^+ \to K^+\pi^0) = 0.040 \pm 0.021$ 5.5 standard deviations (σ)



LHCb doing π^0 final state ${\color{black} {ullet}}$

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Conclusions

- *CPV* observed in many processes, no inconsistency with SM prediction of single complex phase of CKM matrix
 - UT angles and sides are consistent
 - Still plenty of room for New Physics to hide (at few % level)
- Belle II (and LHCb & others) will tackle many open issues in CPV over next decade

- Most sensitive techniques often not even considered at start of experiment

"[CP violation] is telling us that there is a fundamental asymmetry between matter and antimatter, [...] We must continue to seek the origin of the CP symmetry violation by all means at our disposal. We know that improvements in detector technology and quality of accelerators will permit even more sensitive experiments in the coming decades. We are hopeful then, that at some epoch, perhaps distant, this cryptic message from nature will be deciphered."

James Cronin, Nobel lecture, 8 December, 1980

References

- Heavy-quark physics and CP violation, J. Richman, Les Houches Lectures (1998)
 - A little older, but still an excellent pedagogical introduction to SM CP violation
- The Physics of the B Factories, A. Bevan et al., Eur. Phys. J. C (2014) 74, 3026
 - Comprehensive description of CPV measurements at Belle and BaBar
- HFLAV website, Heavy Flavor Averaging Group,
 - Up-to-date averages of CPV measurements
- The Belle II Physics Book, E. Kou et al., Prog. Theor. Exp. Phys. (2019) 123C01
 - Best predictions for Belle II potential for CPV measurements

Back-Up Slides



CPV in Neutrinos

Neutrino mixing matrix (PMNS matrix)

$$egin{aligned} & V = egin{bmatrix} U_{e1} & U_{e2} & U_{e3} \ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \ \end{bmatrix} \ & = egin{bmatrix} 1 & 0 & 0 \ 0 & c_{23} & s_{23} \ 0 & -s_{23} & c_{23} \ \end{bmatrix} egin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \ 0 & 1 & 0 \ -s_{12} & c_{12} & 0 \ 0 & 0 & 1 \ \end{bmatrix} egin{bmatrix} c_{12} & s_{12} & 0 \ -s_{12} & c_{12} & 0 \ 0 & 0 & 1 \ \end{bmatrix} egin{bmatrix} e^{ilpha_1/2} & 0 & 0 \ 0 & e^{ilpha_2/2} & 0 \ 0 & 0 & 1 \ \end{bmatrix} \ & = egin{bmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \ -s_{12}c_{23} - c_{12}c_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}c_{23}s_{13}e^{i\delta} & s_{23}c_{13} \ s_{13}e^{i\delta} & c_{23}c_{13} \ \end{bmatrix} egin{bmatrix} e^{ilpha_1/2} & 0 & 0 \ 0 & e^{ilpha_2/2} & 0 \ 0 & 0 & 1 \ \end{bmatrix} \ & e^{ilpha_1/2} & 0 & 0 \ 0 & e^{ilpha_2/2} & 0 \ 0 & e^{ilpha_2/2} & 0 \ 0 & 0 & 1 \ \end{bmatrix},$$

CPV phase for normal (inverted) hierarchy

 $\delta_{CP} = -1.89^{+0.70}_{-0.58} (-1.38^{+0.48}_{-0.54})$

Zero CP violation ruled out at 2σ

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Removing the $90^{\circ} - \phi_1$ *ambiguity*

- Sign of $\cos 2\phi_1$ resolves $90^\circ \phi_1$ ambiguity from $\sin 2\phi_1$
- Need second interfering (strong) decay amplitude to measure $\cos 2\phi_1$
 - e.g. between CP-odd and CP-even amplitudes in 3-body or VV B decays
 - Requires time-dependent angular or Dalitz plot analysis to extract $\cos 2\phi_1$

 $\cos 2 \phi_1 < 0 \text{ excluded:} \\ \phi_1 = (22.2 \pm 0.7)^{\circ}$









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CPWidtations(SrAhell)

CP Asymmetries in Penguin Decays



CPWiddations(SrPhell)

GLW measurements for ϕ_3

• Most sensitive B final states



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ADS measurements for ϕ_3

• *Most sensitive B final states*



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0.2

5400

03

Two more variables from D decay:

Sensitivity to ϕ_3

Extract ϕ_3 from fit to $D \to K_S^0 \pi^+ \pi^-$ and $D \to K_S^0 K^+ K^-$ Dalitz-plot distributions with variables *x* and *y*







Direct CP violation, if $d = 2r_B | \sin \gamma | \neq 0$

CPV in B_s Decays

CKM matrix up to $O(\lambda^4)$



0.1

0.3

 β_s is equivalent of ϕ_1 in time-dependent B_s CPV

$$b_s^{c\bar{c}s} \approx -2\beta_s$$

World average: $\phi_s^{c\bar{c}s} = -0.050 \pm 0.019$

SM prediction: $\phi_s^{c\bar{c}s} = -0.0370^{+0.0007}_{-0.0008}$

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0.05

ATLAS 99.7 fb⁻¹

-0.3

-0.1

CP Violation (S. Prell)

 $\phi_s^{c\bar{c}s}$ [rad]

$B^{\pm} \rightarrow \pi^{\pm}\pi^{0}$ and $B \rightarrow \pi^{0}\pi^{0}$ CPV and BFs

PRD 87, 052009 (2013)	(error scaled by 1.4)	$\frac{\Gamma(\pi^0 \pi^0) / \Gamma_{\text{total}}}{\frac{VALUE \text{ (units } 10^{-6})}{1.59 \pm 0.26 \text{ OUR AVERAGE}}}$
$\begin{bmatrix} \bullet & \bullet \\ \bullet & \bullet \end{bmatrix} \xrightarrow{\bullet} B \rightarrow \pi^0 \pi^0 \qquad = \begin{bmatrix} \bullet & \bullet \\ \bullet & \bullet \end{bmatrix}$	Belle, PRD 96, 032007 (2017) BaBar, PRD 87, 052009 (2013)	$\begin{array}{c} 1.31 \!\pm\! 0.19 \!\pm\! 0.19 \\ 1.83 \!\pm\! 0.21 \!\pm\! 0.13 \end{array}$
	Belle, PRD 96, 032007 (2017) BaBar, PRD 87, 052009 (2013)	$C_{\pi^{0}\pi^{0}}(B^{0} \rightarrow \pi^{0}\pi^{0})$ <u>VALUE</u> -0.33±0.22 OUR AVERAGE -0.14±0.36±0.10 -0.43±0.26±0.05
5.2 5.22 5.24 5.26 5.28 m _{ES} (GeV/c ²)	(error scaled by 1.2)	$\frac{\Gamma(\pi^{+}\pi^{0})/\Gamma_{\text{total}}}{5.5 \pm 0.4 \text{ OUR AVERAGE}}$
	Belle, PRD 87, 031103 (2013) BaBar, PRD 76, 091102 (2007) CLEO, PRD 68, 052002 (2003)	$5.86 \pm 0.26 \pm 0.38$ $5.02 \pm 0.46 \pm 0.29$ $4.6 \begin{array}{c} +1.8 \\ -1.6 \end{array} \begin{array}{c} +0.6 \\ -0.7 \end{array}$
	verage (no scale factor)	$5.48^{+0.35}_{-0.34}$ HFLAV a $A_{CP}(B^+ \rightarrow \pi^+ \pi^0)$
	elle, PRD 87, 031103 (2013) aBar, PRD 76, 091102 (2007)	VALUE 0.03 ±0.04 OUR AVERAGE 0.025±0.043±0.007 Be 0.03 ±0.08 ±0.01 Ba
-0.2 -0.1 0 0.1 0.1 $\Delta E (GeV)$	/ average	0.026 ± 0.039 HFLAV

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0.2

ϕ_2 from $B \rightarrow \rho \rho$ with Belle II

- Most measurements will still be statistically limited with 50/ab
 - Expected error $\Delta \phi_2 \sim 1^{\circ}$ without $S_{\rho^0 \rho^0}$ constraint, and $\Delta \phi_2 \sim 0.6^{\circ}$ with $S_{\rho^0 \rho^0}$ constraint
- The $B^0 \to (\pi \rho)^0$ Dalitz analysis
 - Done by both Belle and BaBar, but limited by spurious ambiguities due to small data samples
 - Analysis should be repeated with a few ab^{-1} , which will allow to estimate sensitivity with 50 ab^{-1}

	Value	$0.8 \ {\rm ab}^{-1}$	50 ab^{-1}
$f_{L, ho^+ ho^-}$	0.988	$\pm 0.012 \pm 0.023$ [725]	$\pm 0.002 \pm 0.003$
$f_{L, ho^{\mathrm{o}} ho^{\mathrm{o}}}$	0.21	$\pm 0.20 \pm 0.15$ [729]	$\pm 0.03 \pm 0.02$
${\cal B}_{ ho^+ ho^-}$ [10 ⁻⁶]	28.3	$\pm 1.5 \pm 1.5$ [725]	$\pm 0.19 \pm 0.4$
$\mathcal{B}_{ ho^0 ho^0}$ [10 ⁻⁶]	1.02	$\pm 0.30 \pm 0.15$ [729]	$\pm 0.04 \pm 0.02$
$A_{\rho^+\rho^-}$	0.00	$\pm 0.10 \pm 0.06$ [725]	$\pm 0.01 \pm 0.01$
$S_{\rho^+\rho^-}$	-0.13	$\pm 0.15 \pm 0.05$ [725]	$\pm 0.02 \pm 0.01$
	Value	0.08 ab^{-1}	$50 {\rm ~ab^{-1}}$
$f_{L, ho^+ ho^0}$	0.95	$\pm 0.11 \pm 0.02$ [716]	$\pm 0.004 \pm 0.003$
$\mathcal{B}_{ ho^+ ho^0}$ [10 ⁻⁶]	31.7	$\pm 7.1 \pm 5.3$ [716]	$\pm 0.3 \pm 0.5$
	Value	$0.5 { m ~ab^{-1}}$	$50 { m ~ab^{-1}}$
$A_{ ho^{\mathrm{o}} ho^{\mathrm{o}}}$	-0.2	$\pm 0.8 \pm 0.3$ [715]	$\pm 0.08 \pm 0.01$
$S_{ ho^{0} ho^{0}}$	0.3	$\pm 0.7 \pm 0.2$ [715]	$\pm 0.07 \pm 0.01$