

# The Flavor Path to New Physics

$b \rightarrow s\nu\bar{\nu}$  & related modes

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# Key questions

**Questions** I'd like to discuss:

- What does current **data** tell us about the **flavor** of **physics beyond the SM**?  
Connection to the SM flavor puzzle?
- Which role does  $b \rightarrow s\nu\bar{\nu}$  play in shaping this picture? What's its **interplay** with other processes, e.g. with other flavor changing transitions ( $s \rightarrow d\nu\bar{\nu}$ ,  $b \rightarrow s\tau\tau\dots$ )?
- How will this change in the **future**, with existing and future facilities?

## **Caveats:**

- Some overlap with other theory talks
- Not covered here:
  - light new physics scenarios
  - SM prediction challenges in  $b \rightarrow s\nu\bar{\nu}$

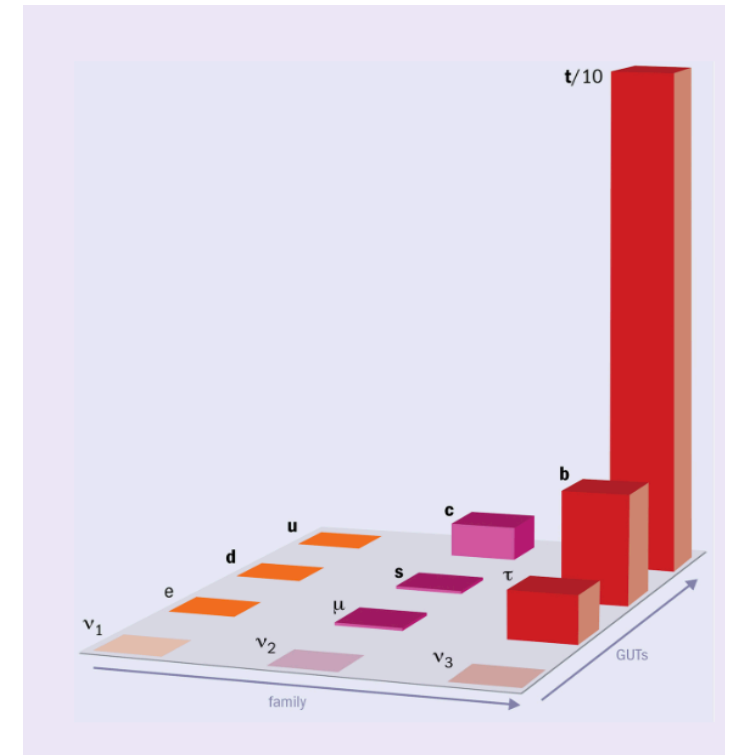
# Flavor in the Standard Model: puzzling aspects

**SM flavor puzzle = a series of puzzling observations**

3 copies x species, identical from the point of view of gauge interactions, yet seen very differently by the Higgs

[and by BSM giving mass to neutrinos]

- 12 orders of magnitude from neutrinos to the top mass
- pronounced mass hierarchies for charged fermions
- mixing looks very different in lepton vs quark sector



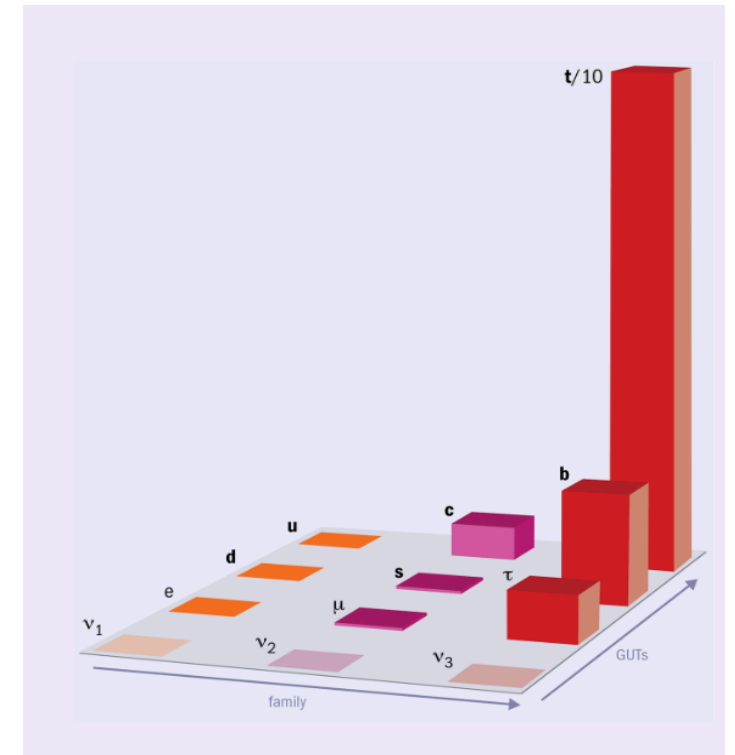
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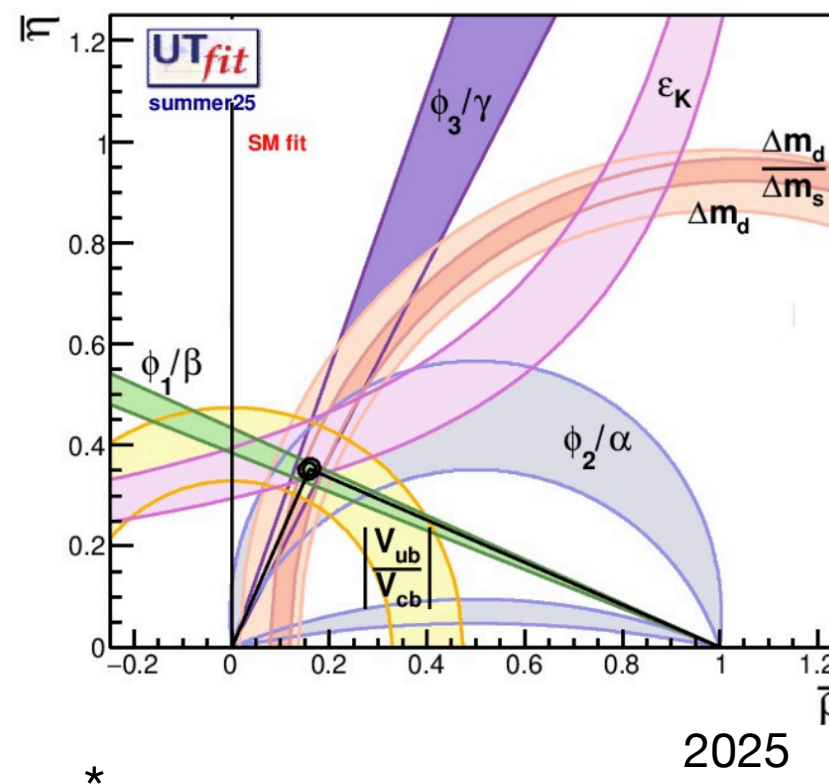
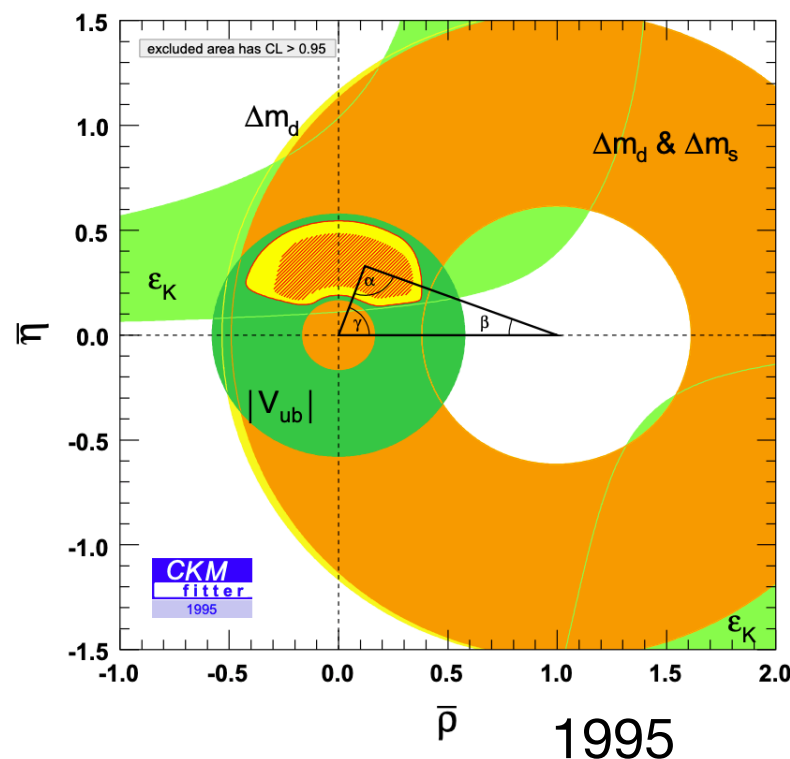
No explanation in the SM: just free parameters “fixed” via measurements. *Technically natural*, yet **suggestive of an organising principle** (necessarily) *beyond* the SM.

- Many ideas: Froggatt-Nielsen, Randall-Sundrum, GUTs, Flavor deconstruction...
- Not pointing to specific scales unless linked to other “problems” (hierarchy/gauge coupling unification)

# Flavor in the Standard Model: strong predictions

Despite its mysterious origin, the flavor structure of the SM leads to a set of remarkably **successful predictions**:

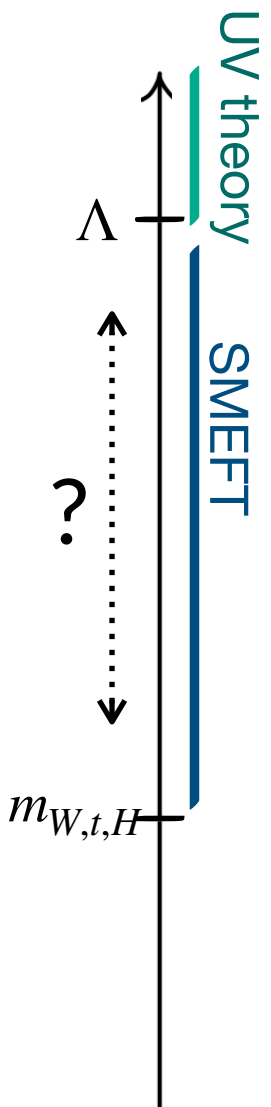
- absence of charged lepton flavor violation [up to small mv effects]
- lepton flavor universality [up to ml effects]
- suppression of flavor changing neutral currents [GIM + loop]
- unitarity of the CKM



...all expression of the **SM matter content** and the resulting **accidental symmetries**.

# The scale of New Physics, from theory

We have many more reasons beyond the flavor puzzle to know that BSM exists — yet the absence of direct hints suggests a mass gap. How large? **What is the scale of NP?**



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## Top-down considerations

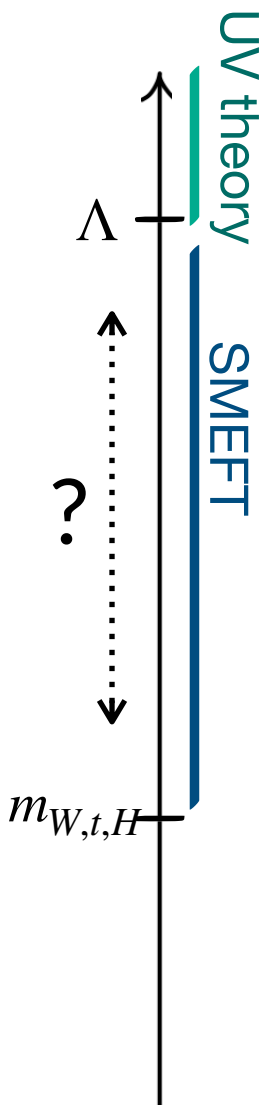
- the **main hint for a “low”  $\Lambda$**  is a “natural” solution to the **hierarchy problem**

Any heavy BSM coupled to the Higgs destabilises  $m_h$

⇒ Need some NP coupled to Higgs & top at  $\Lambda \sim \text{TeV}$  to stabilise it.

In general, these solutions modify the Yukawa sector: necessarily connected to flavor.

- Other challenges (dark matter, dark energy, inflation...) are **more difficult to link to a  $\Lambda$  accessible by colliders** — *Exception*: the WIMP miracle

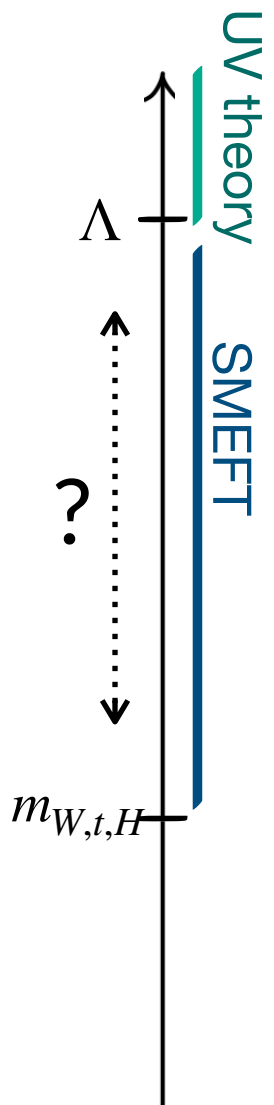


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## Bottom-up considerations

- **Indirect searches** at  $E \ll \Lambda$  can pinpoint  $\Lambda$  far beyond directly accessible scales  
many historical precedents:  $m_c$  from K mixing,  $m_t$  from EWPOs
- best performed with **processes** that we can **predict and measure precisely**
  - null tests proton decay ( $U(1)_B$ ),  $0\nu\beta\beta$  ( $U(1)_L$ ),  $\mu \rightarrow e\gamma$  ( $U(1)_{Li}$ ), LFUV
  - flavor-changing transitions
  - EWPOs
- likely where we'll see the **biggest experimental progress** in the next 50 yrs
  - @ existing & planned facilities: (HL-)LHC, KEK, PSI, JPARC...
  - @ a future Tera Z factory



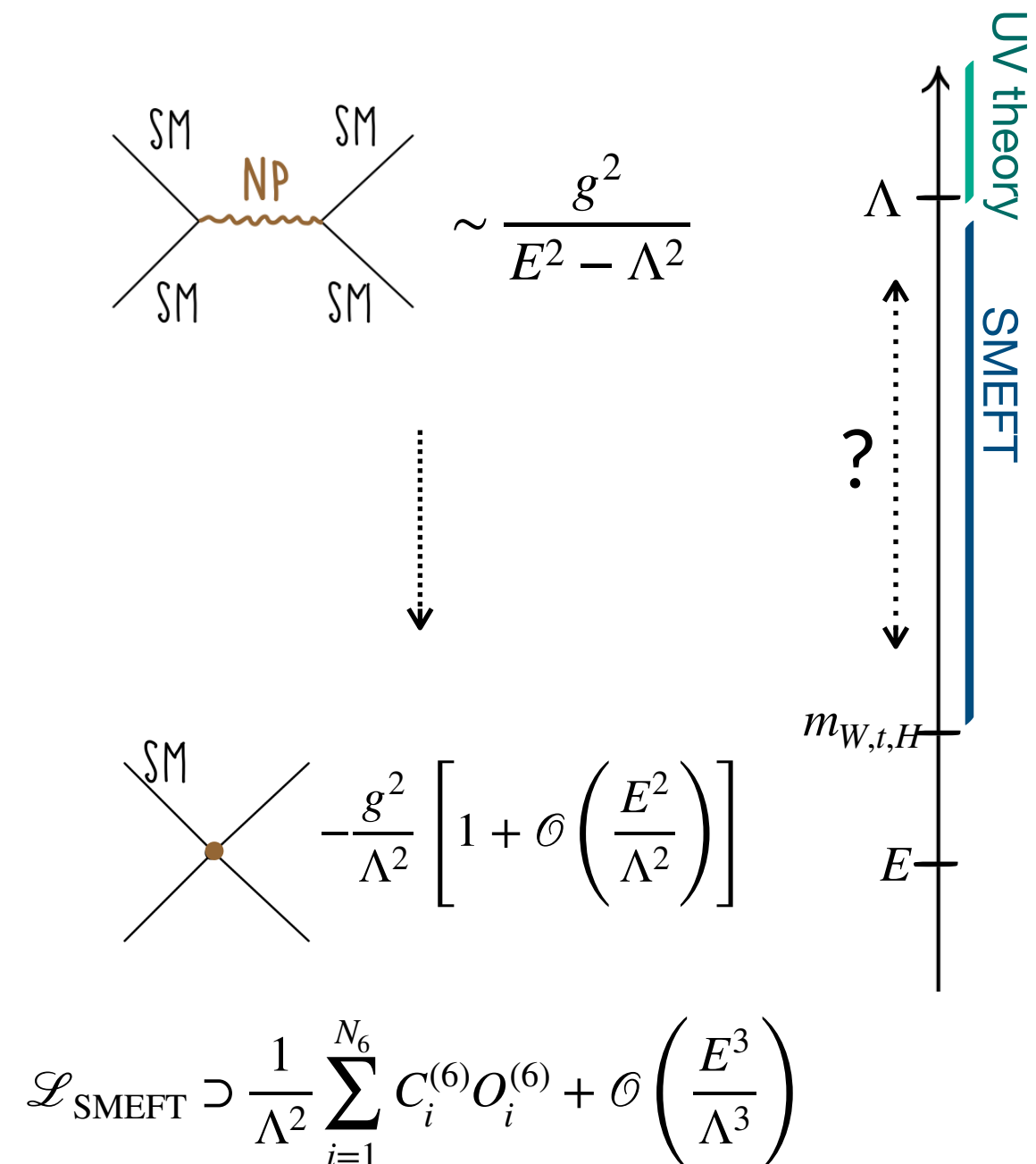


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- If NP is much heavier than the scale we're probing, we can describe its effects “model independently” with an **effective field theory**



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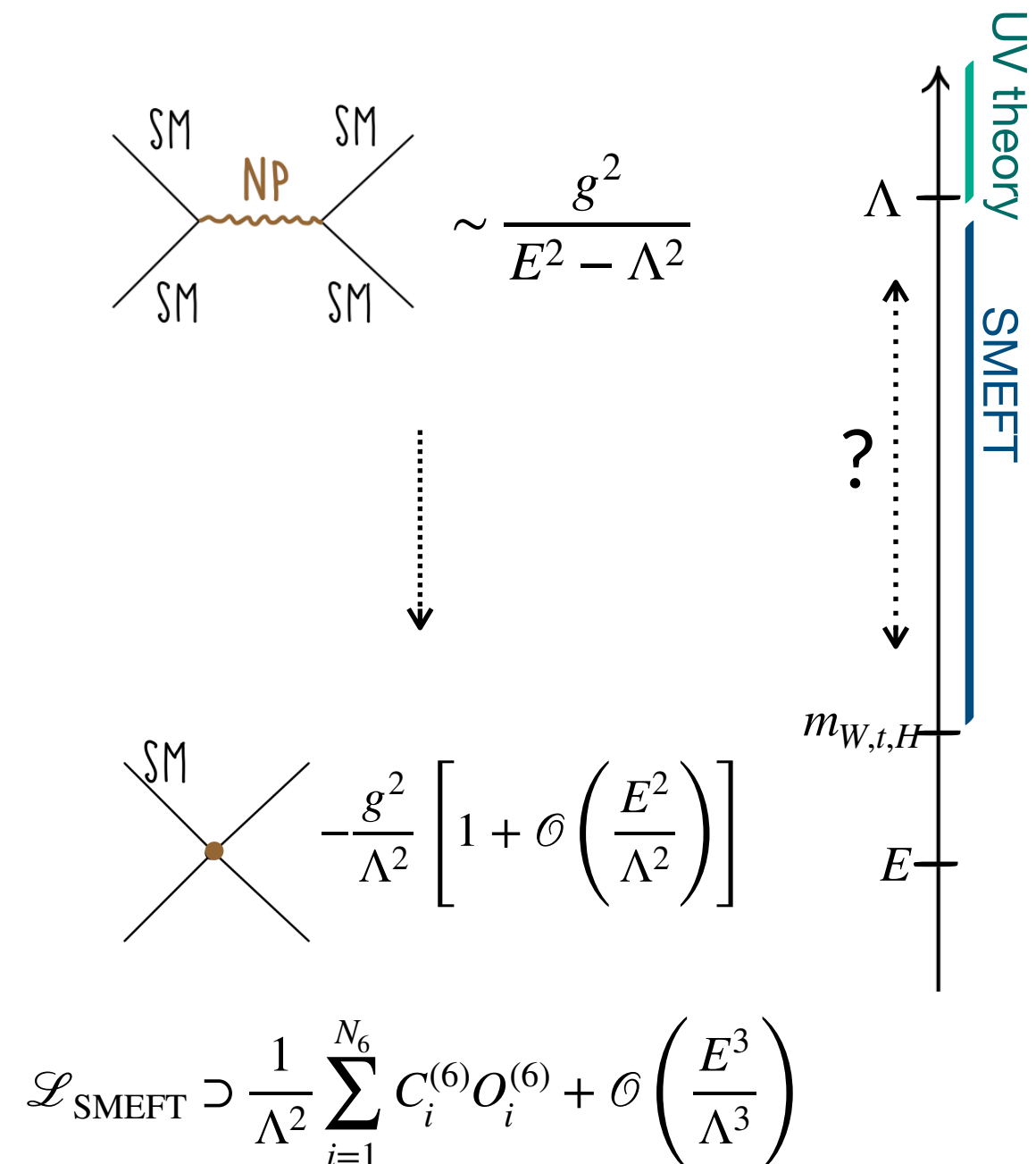
- If NP is much heavier than the scale we're probing, we can describe its effects “model independently” with an **effective field theory**

- EFTs translate **data** → **info on heavy NP**:

- use data (EWPOs, flavor, collider) to constrain WCs
- constraints interpreted as bounds on an *effective* NP scale

$$\Lambda_{\text{eff}}^i = \frac{\Lambda}{\sqrt{C_i}} \sim \frac{M}{g}$$

- *absolute* scale  $\Lambda$  depends on assumptions on NP couplings!



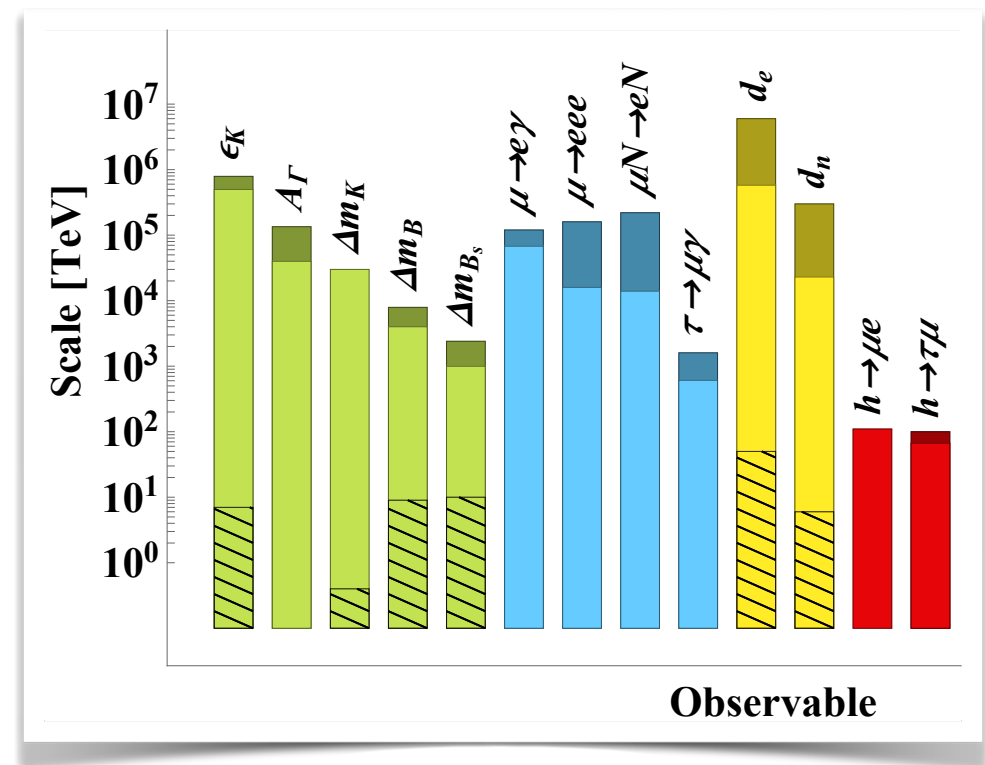
# On the importance of flavor assumptions

When looking at flavor data through this EFT lens, we encounter the well-known

**new physics flavor puzzle**

= nothing forbids bad violations of the SM accidental symmetries, yet we don't see any

$$\Lambda_{\text{eff}}^{\text{FV}} \gtrsim 10^{4-6} \text{ TeV}$$



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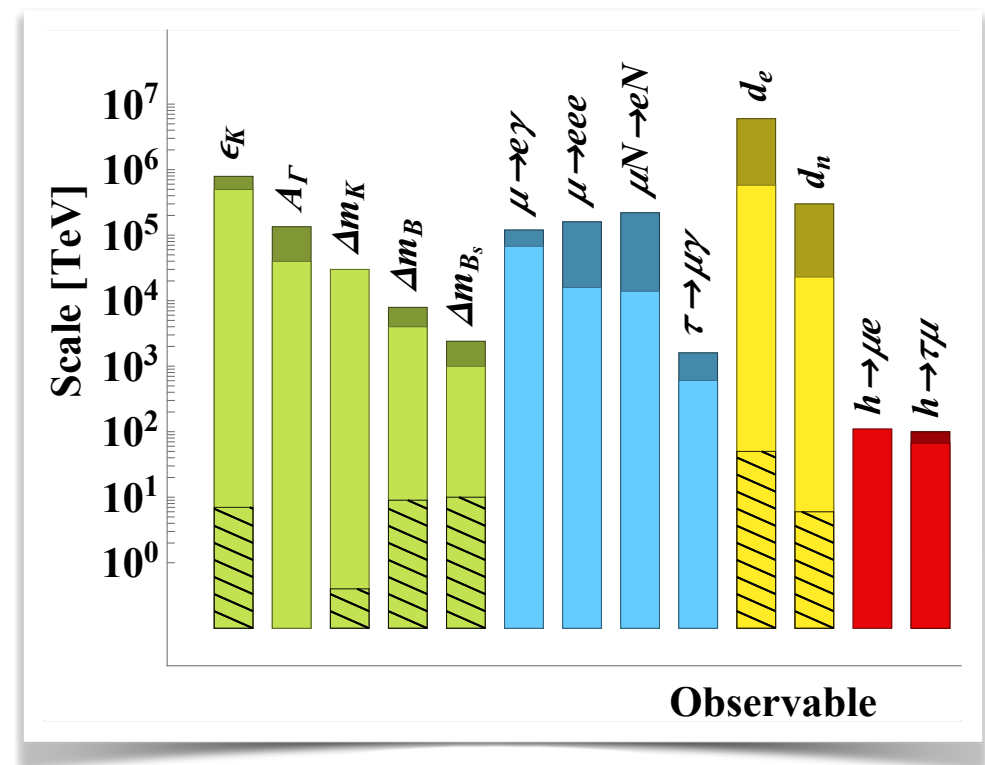
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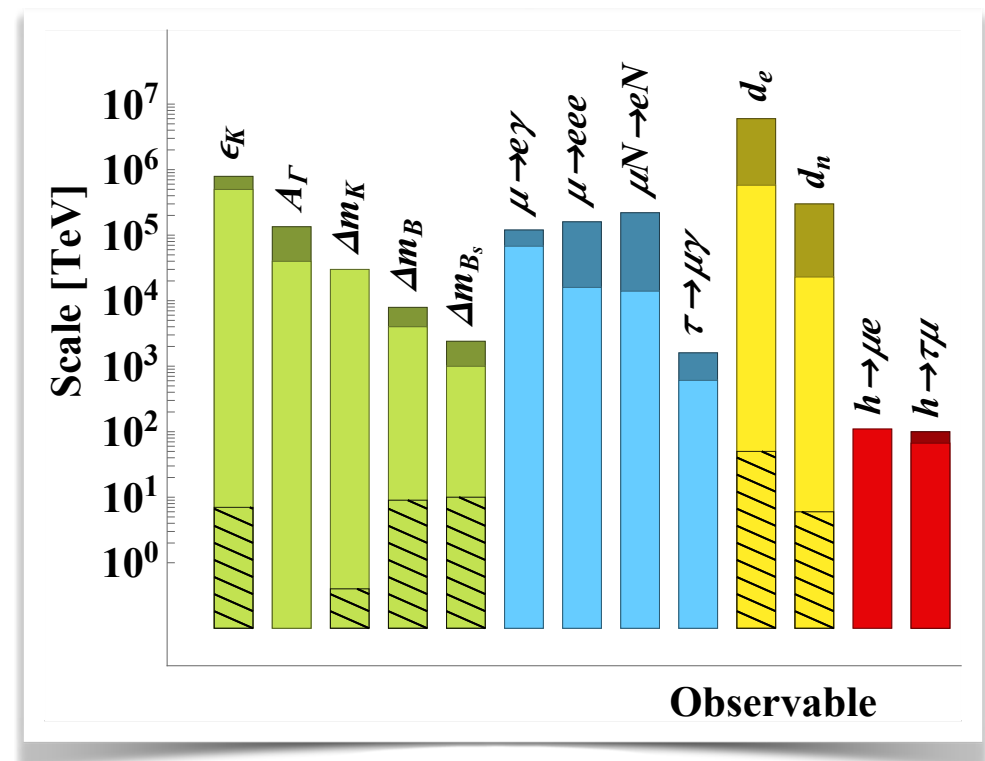
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Two options:

New physics is **anarchic**:

$$C_{\text{FC}} \sim \mathcal{O}(1) \Rightarrow \Lambda = \Lambda_{\text{eff}}$$

- Very heavy, hence **untestable**
- Higgs stabilised in some other way [relaxion, landscape..?]
- unclear how flavor patterns could arise



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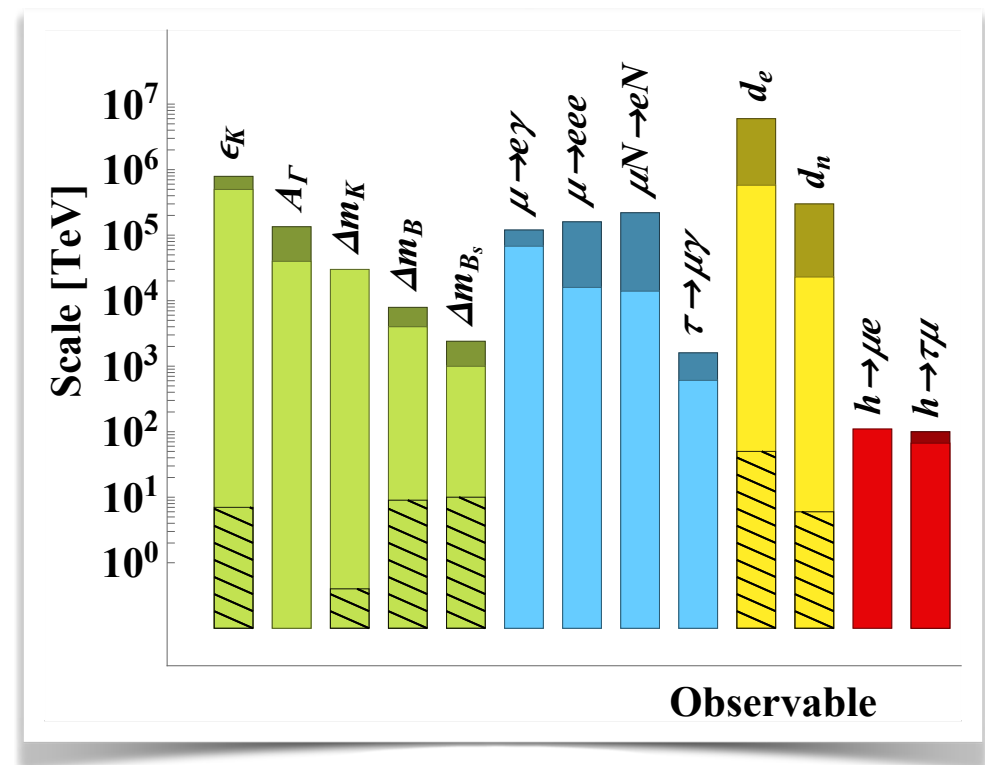
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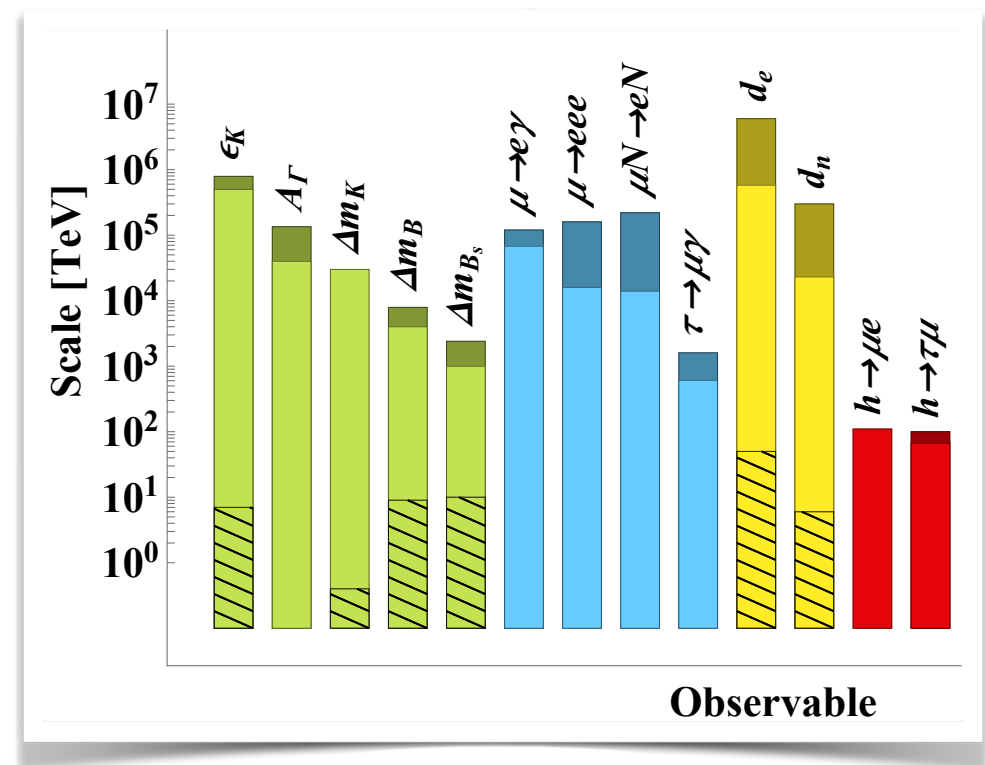
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**Which one?**

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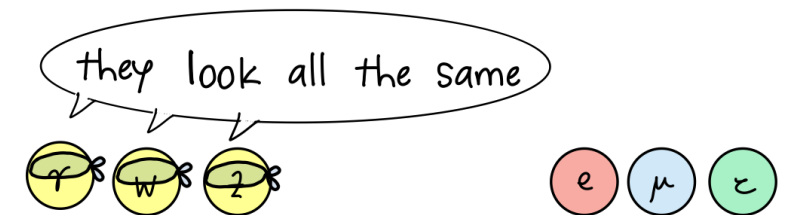
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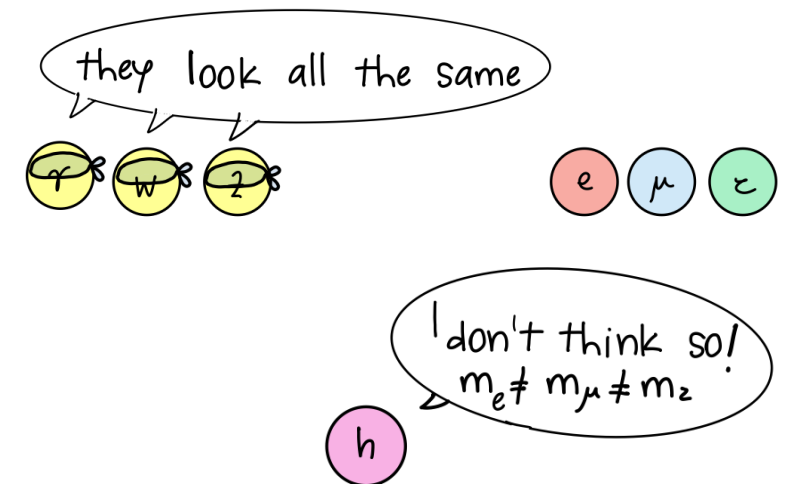
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....broken by Yukawa interactions to an *approximate*  $U(2)^5$ .

— *approximate*, because it needs to be broken to reproduce light quark masses and CKM.

$U(2)$  exact

$$Y_{u,d} = \begin{bmatrix} & & \\ & & \\ \cdots & & \\ & & \blacksquare \end{bmatrix}$$

only  $m_{b,t} \neq 0$ , CKM = 1

"Minimal" breaking

$$Y_{u,d} = \begin{bmatrix} \text{gray box} & \\ & \text{gray box } V_q \\ \cdots & \\ & \blacksquare \end{bmatrix}$$

[non minimal options also possible]

$$V_q \sim 2_q \sim \epsilon_F \begin{pmatrix} V_{td} \\ V_{ts} \end{pmatrix}$$

$$\Delta_{u,d} \sim (\bar{2}_{d,u}, 2_q) \sim y_{c,s}$$

only two "free" params:  $\epsilon_F, \phi_q$

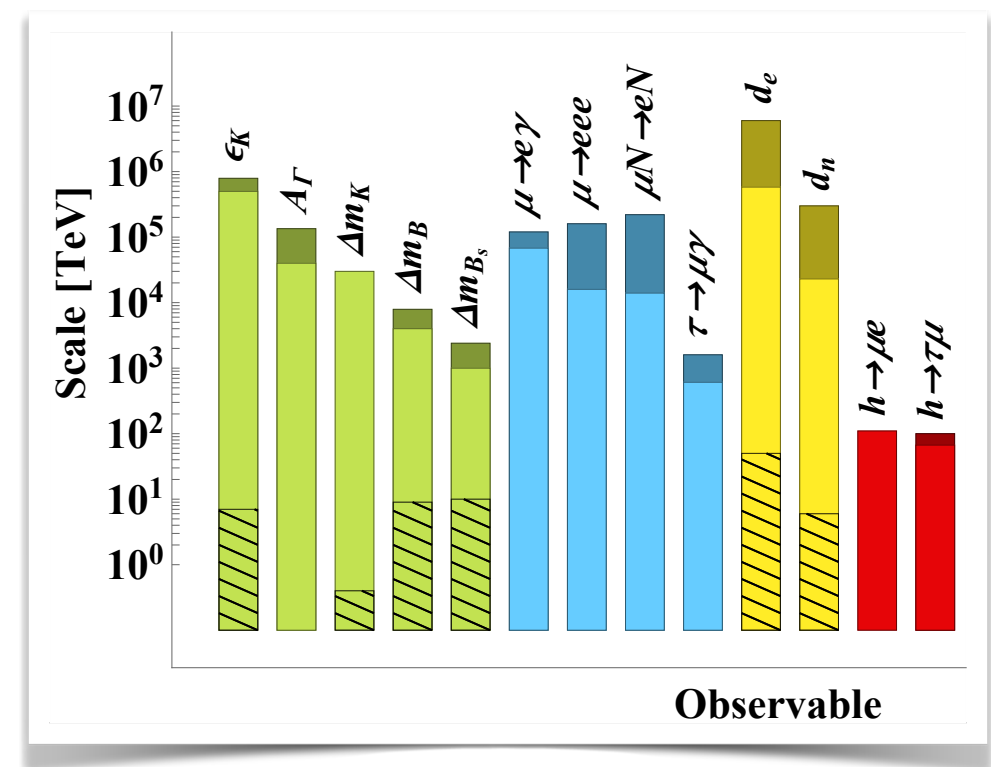
# Using the SM accidental symmetries as a guidance

This symmetry seems to hold **also beyond the SM** — at least **if NP is not too far**.

In other words, the only unambiguous message of flavor bounds is that

if there is NP “close by” ( $< 10^4$  TeV),  
it must respect a U(2)-like structure

...which in turn is suggestive of the idea that  
this symmetry might have a more  
“fundamental” origin.



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  - Yukawas *only* sources of flavor violation.
  - describes approximately flavor-universal NP
- CKM-like suppression of FCNCs, but unsuppressed valence-quark couplings  
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## Non-universal (3rd family) New Physics

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- $U(2)^n$  can host **NP** coupling dominantly to the **3rd family**
- theoretical motivation: possible **link to flavor puzzle** (and hierarchy problem)  
[as of today, natural solutions to the hierarchy problem require flavor-non-universal NP]
- **valence-quark couplings suppressed**  $\Rightarrow$  high pT bounds relax to  $\Lambda \gtrsim O(1)$  TeV



# Adding flavor to the SMEFT: $U(2)^5$

To study these scenarios” agnostically, **impose  $U(2)^5$  in the SMEFT.**

It's an efficient organising principle

- the SMEFT has 1350 (w/o CPV) + 1149 (w CPV) = 2499 WCs at  $d=6$ .
- In the *exact*  $U(2)^5$  limit, this reduces to 124 + 23 = 147 WCs
- breaking terms can be incorporated systematically order by order

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*Example*  $[Q_{Hq}^{(1)}]_{ij} = (H i \overleftrightarrow{D}_\mu H)(\bar{q}_i \gamma^\mu q_j)$

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**Caveat: need to define flavor basis**

- $q_3$  is in general neither down- nor up-aligned
- misalignment described by  $V_q$

$$q_t = \begin{pmatrix} t_L \\ V_{td}d_L + V_{ts}s_L + V_{tb}b_L \end{pmatrix}$$

$$q_b = \begin{pmatrix} V_{ub}^*u_L + V_{cb}^*c_L + V_{tb}^*t_L \\ b_L \end{pmatrix}$$

# Adding flavor to the SMEFT: some comments

**Why using flavor assumptions in the SMEFT?** Don't we lose agnosticity? Yes, but...

- the **SMEFT** is *not a model*, and fully general  $\neq$  informative
- realistic models populate only some directions  $\rightarrow$  new correlations & stronger bounds
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**Including flavor data in SMEFT analyses is challenging**

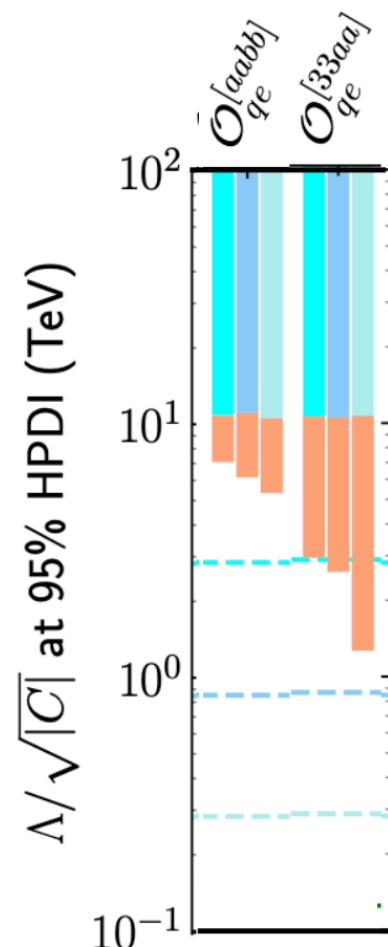
- For **EW and collider** data, **exact** flavor **symmetries** [hence CKM = 1] are **enough**
- **Flavor** data requires the inclusion of **breaking terms**:
  - “mandatory” breakings to reproduce masses + CKM
  - additional non-standard sources of breaking, to be constrained from data
- Ongoing, non-trivial effort [Aoude, Hurth, Renner, Shepherd 2003.05432 Bruggisser, van Dyk, Westhoff 2212.02532 Grundwald, Hiller, Kröninger, Nollen 2304.12837 Allwicher, CC, Isidori, Stefanek 2311.00020 Bartocci, Biekötter, Hurth 2311.04963...]  
First global SMEFT fit with  $U(2)^5$  and real CKM [de Blas, Goncalves, Miralles, Reina, Silvestrini, Valli 2507.06191]

# The role of flavor data: complementarity

Flavour measurements play a crucial role in constraining flavor-conserving new interactions\*, **complementary** to EW and high pT

[\* = **overall scale** & suppression of light fam. couplings]

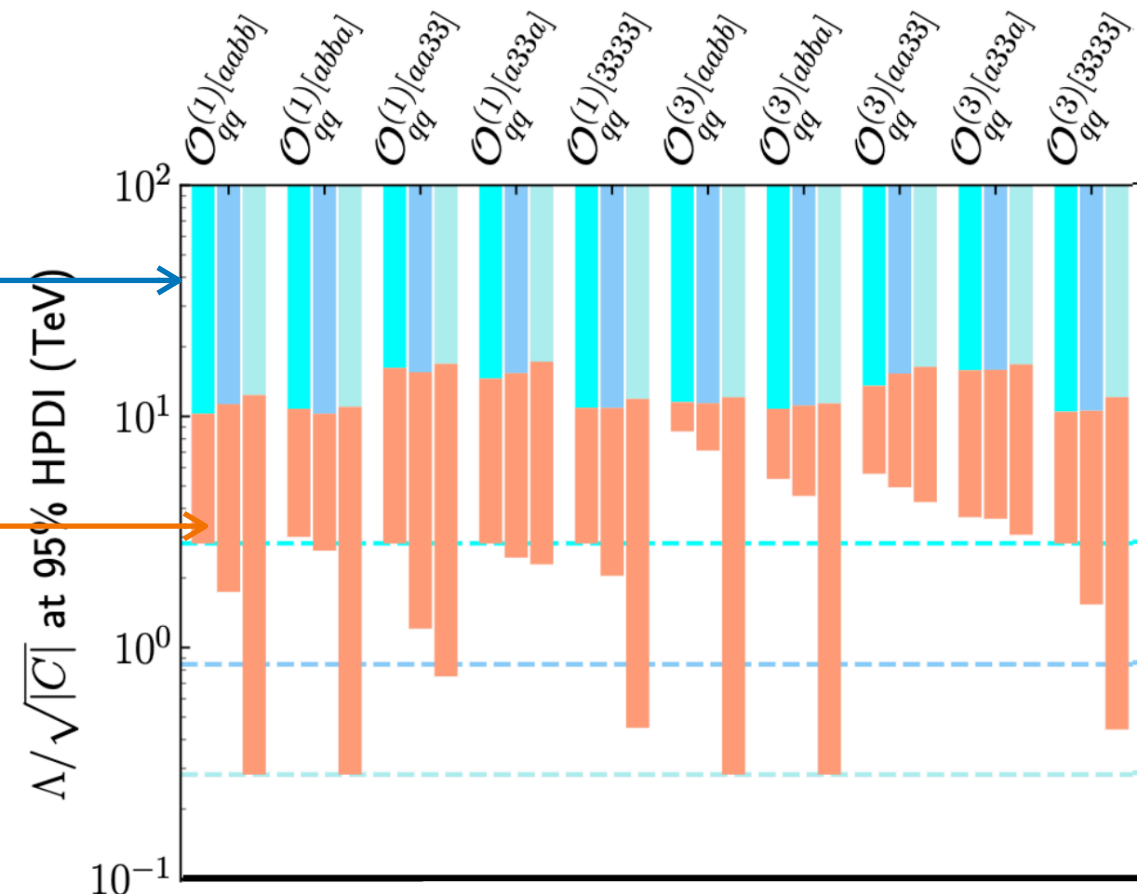
Examples:  $B_s \rightarrow \mu^+ \mu^-$  for lepton-quark ops



← Full fit

← W/o flavor observables

meson mixing for 4-quark ops.



[de Blas, Goncalves, Miralles, Reina, Silvestrini, Valli 2507.06191]

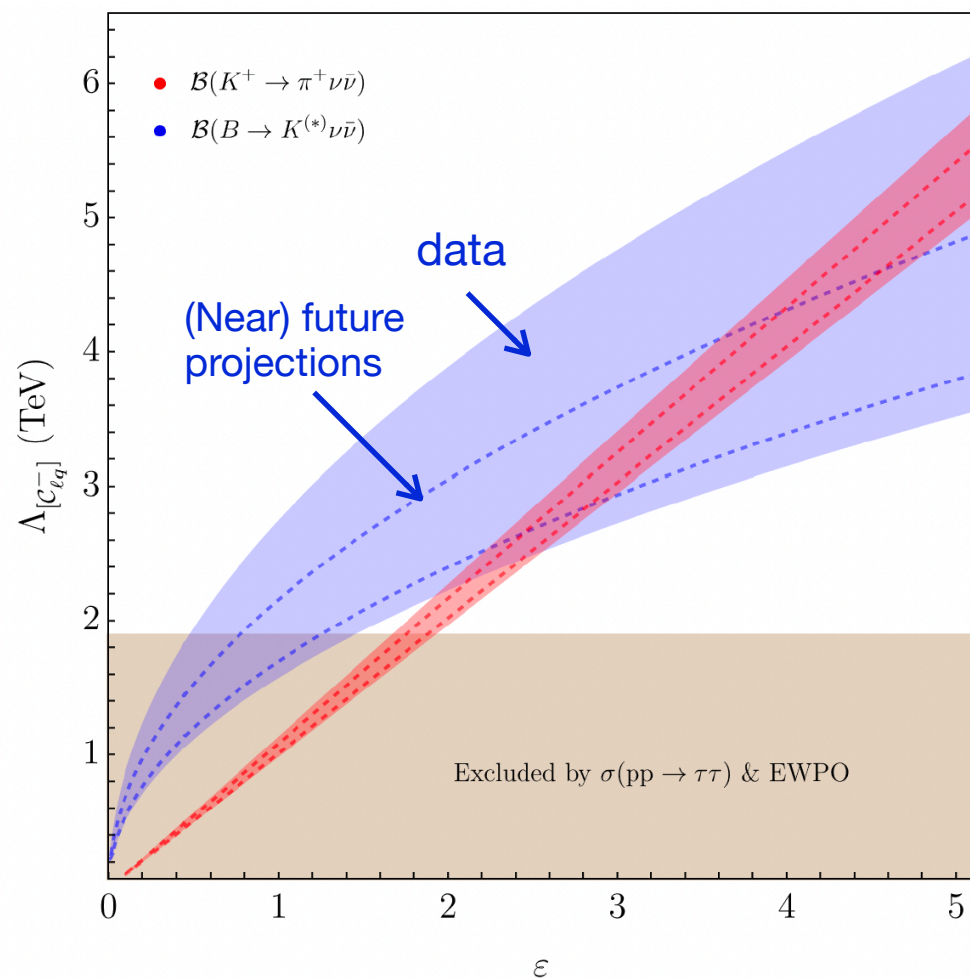


# The role of flavor data: unicity

but flavor also provides **unique** info on the NP **flavor structure** by constraining genuinely flavor violating interactions

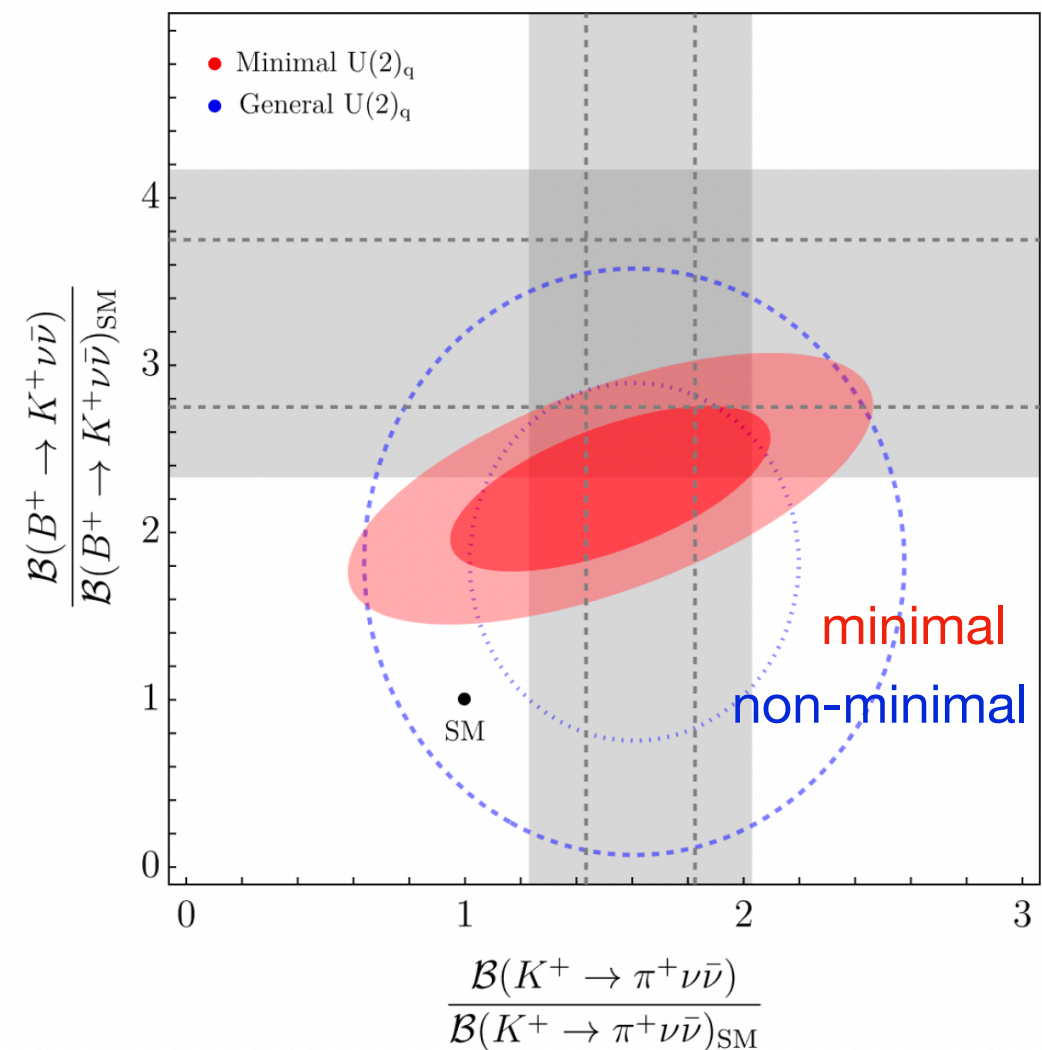
determine size of breaking terms

[those not already fixed by masses & CKM]



e.g. leading  $U(2)_q$  breaking can be determined by comparing  $b \rightarrow s$  vs  $s \rightarrow d$

discriminate between  $\neq$  breaking patterns via characteristic correlations



[Allwicher, Bordone, Isidori, Piazza, Stanzione, 2410.21444]



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As an example, two **predictions of the “minimal” breaking ansatz** for down-type FCNCs:

**NP in  $b \rightarrow s$  fixes NP in  $b \rightarrow d$**

**NP in  $s \rightarrow d$  is related to NP in  $b \rightarrow s, d$**

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- $\frac{b \rightarrow sX}{b \rightarrow dX} = \frac{b \rightarrow sX}{b \rightarrow dX} \Big|_{\text{SM}} = \frac{V_{td}}{V_{ts}}, \quad X = \ell^- \ell^+, \nu \bar{\nu}$

- possible clean tests:  $\frac{\mathcal{B}(B_s \rightarrow \mu\mu)}{\mathcal{B}(B_d \rightarrow \mu\mu)}, \quad \frac{\mathcal{B}(B \rightarrow K \bar{\nu}\nu)}{\mathcal{B}(B \rightarrow \pi \bar{\nu}\nu)}$

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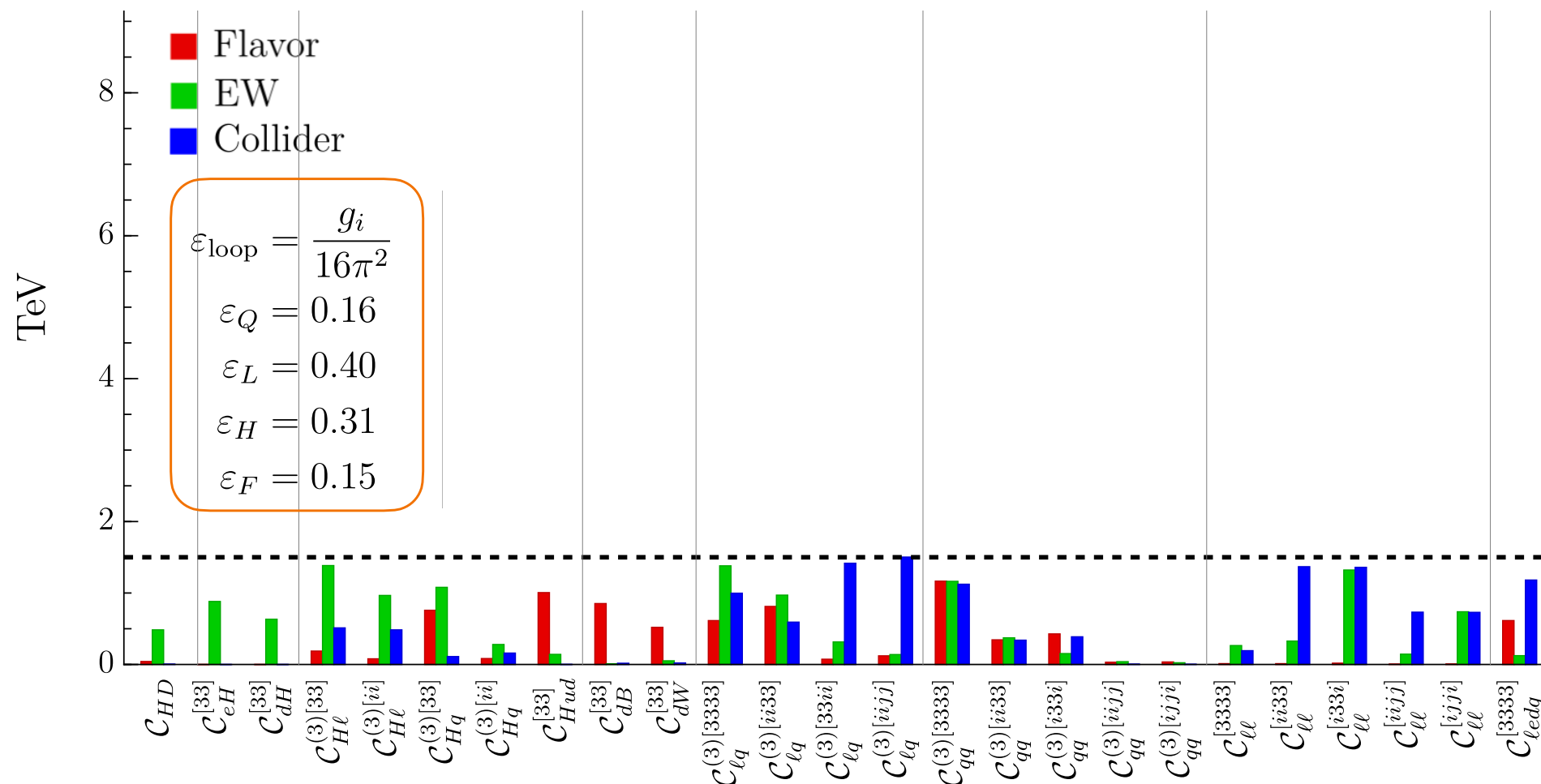
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- $s \rightarrow d$  feels subleading breaking terms — the connection becomes 1:1 only with additional hypotheses [e.g. “rank 1”: NP coupled only to 3rd family]
- only clean test is comparing dineutrino modes:  $\mathcal{B}(K \rightarrow \pi\bar{\nu}\nu)$  vs  $\mathcal{B}(B \rightarrow K\bar{\nu}\nu)$   
[both charged & neutral modes]

# SMEFT Bounds on Third Family New Physics

Single operator analysis in **SMEFT + minimally broken  $U(2)^5$**

[ = no sources of quark FV apart from the minimal ones required for the CKM,  $V_q$ ,  $\Delta_{u,d}$ ]



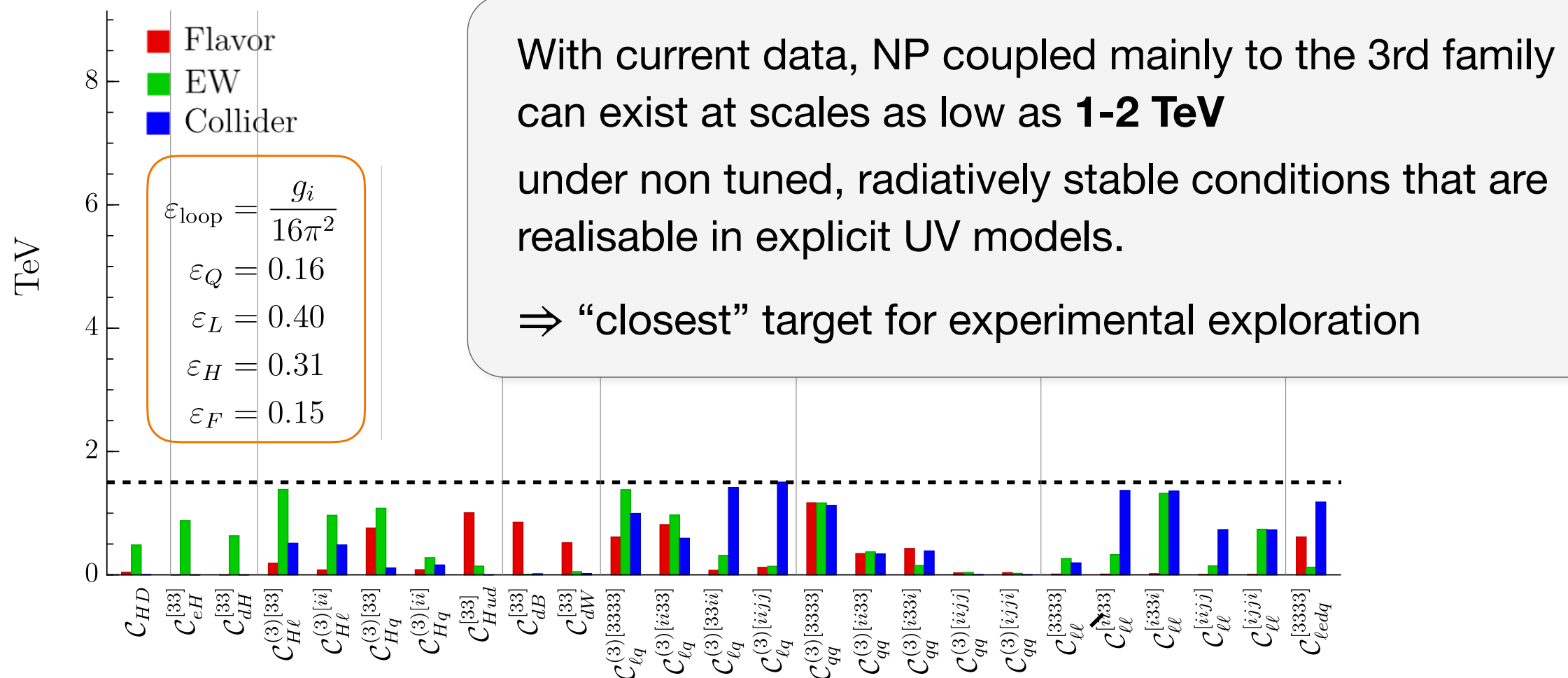
[Allwicher, CC, Isidori, Stefaneke, 2311.00020]

(\*) Minimal suppression of NP couplings to light families, Higgs & orientation in flavor space fixed by high-PT, EW and flavor data, respectively.

# SMEFT Bounds on Third Family New Physics

Single operator analysis in **SMEFT + minimally broken  $U(2)^5$**

[ = no sources of quark FV apart from the minimal ones required for the CKM,  $V_q$ ,  $\Delta_{u,d}$ ]

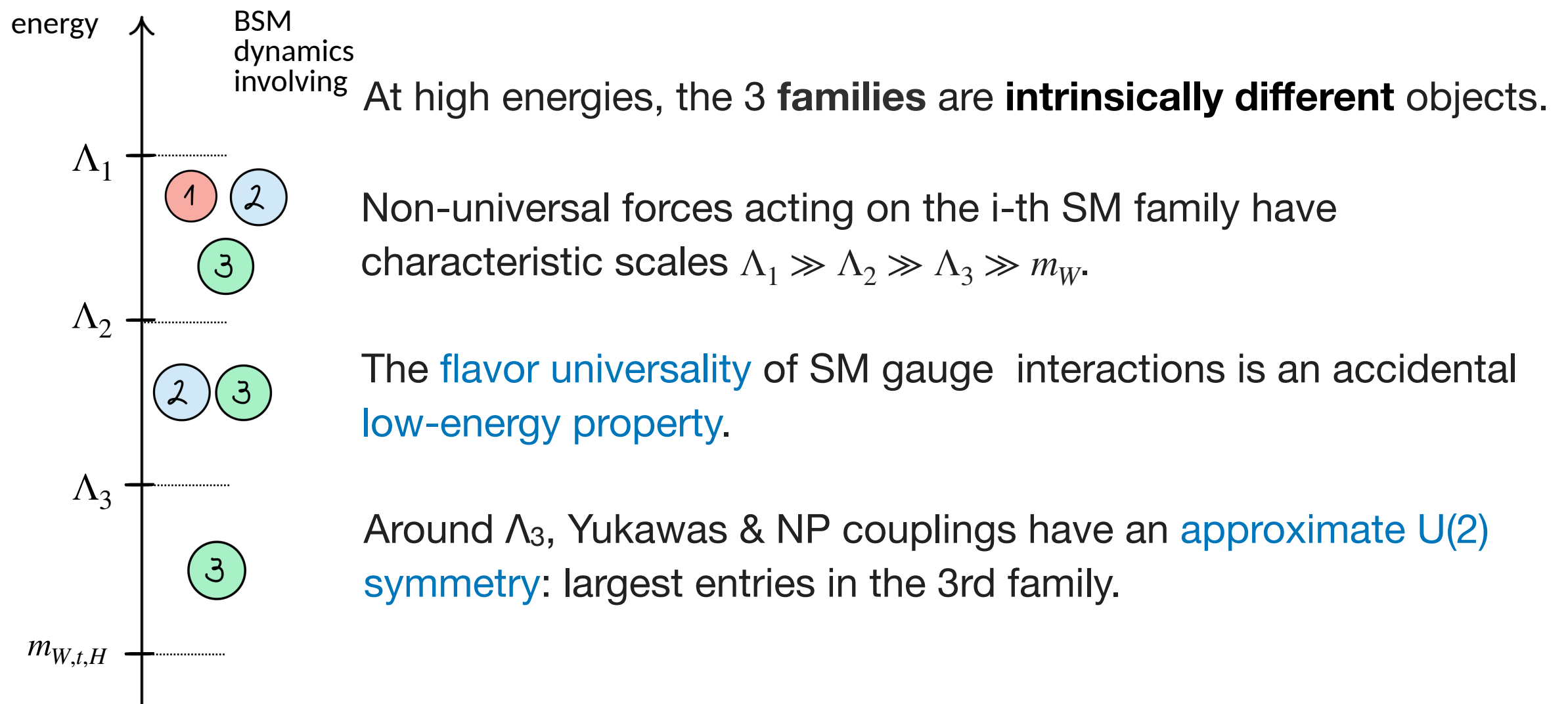


[Allwicher, CC, Isidori, Stefaneke, 2311.00020]

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# Aside: Model building for Third Family New Physics

*Key idea:* The U(2) symmetry in the Yukawas and in the new physics couplings has the same dynamical origin & is a remnant of a fundamental difference



# Aside: Model building for Third Family New Physics

*Key idea:* The  $U(2)$  symmetry in the Yukawas and in the new physics couplings has the same dynamical origin & is a remnant of a fundamental difference

- Explicit realisation via **flavor deconstruction** of the SM gauge group:

$$G = \boxed{G_{3,\text{SM}}} \times \boxed{G_{12,\text{SM}}} \xrightarrow{\Lambda_3} G_{\text{SM}}$$

acts on 3rd fam. & Higgs      acts on light families

Many examples for  $G$  have been studied:

$$SU(4)_3 \times SU(3)_{12} \times SU(2)_L \times U(1)'_Y$$

$$SU(3) \times SU(2)_{L,3} \times SU(2)_{L,12} \times U(1)_Y$$

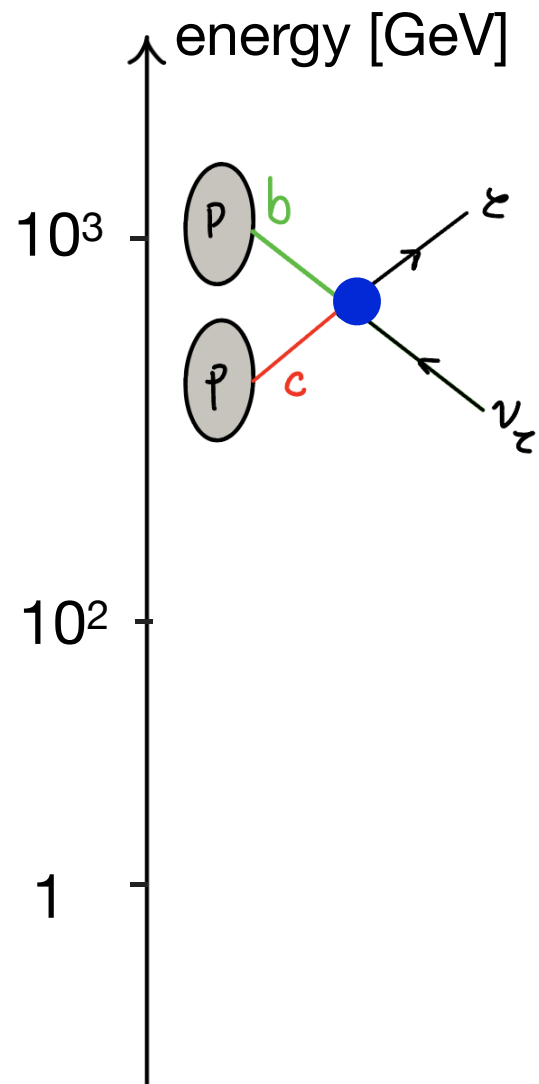
$$SU(3) \times SU(2)_L \times U(1)_{Y,3} \times U(1)_{Y,12}$$

- built-in  $U(2)^5$  in the gauge sector; only  $y_3 \neq 0$
- SSB to SM generates **new gauge bosons** with  $M \sim \mathcal{O}(\Lambda_3)$  **coupled mostly to the 3rd family** — rich phenomenology for  $\Lambda_3 \sim \text{TeV}$
- same breaking of  $U(2)^5$  generates light Yukawas and couplings of the new gauge bosons to light families



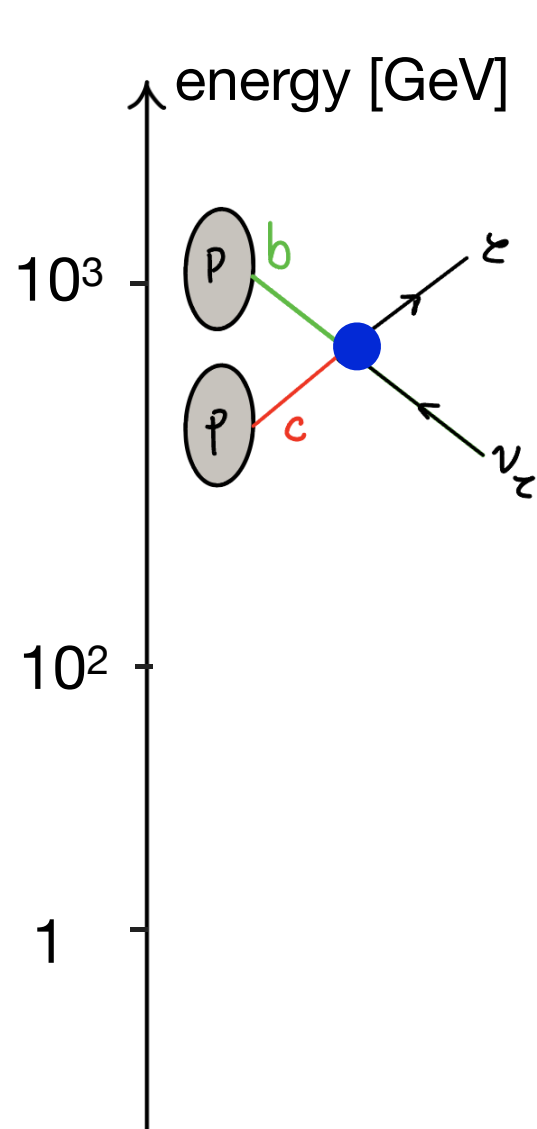
# High-energy signatures of 3rd family new physics

**LHC searches** [  &  ] [See talks by T.Vazquez & M. Martinez]



- largest effects in 3rd-family processes:
  - lepton sector:  $pp \rightarrow t\bar{t}, pp \rightarrow b\bar{b} \dots$
  - quark sector:  $pp \rightarrow \tau\tau, pp \rightarrow \tau\nu$
  - also LFU, e.g. comparing  $pp \rightarrow \tau\tau$  to  $pp \rightarrow \mu\mu$
- energy enhancement in **tails** helps overcome pdf suppression of heavy flavours in the proton

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## LHC searches



&



[See talks by T.Vazquez & M. Martinez]

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- energy enhancement in **tails** helps overcome pdf suppression of heavy flavours in the proton

## Status and prospects

- **currently**, LHC probes scales  $\sim 1 \text{ TeV}$
- **HL-LHC**: improvement in WCs bounds range **from 20% to 4 x** for semileptonic operators (factor 2x in the scale)

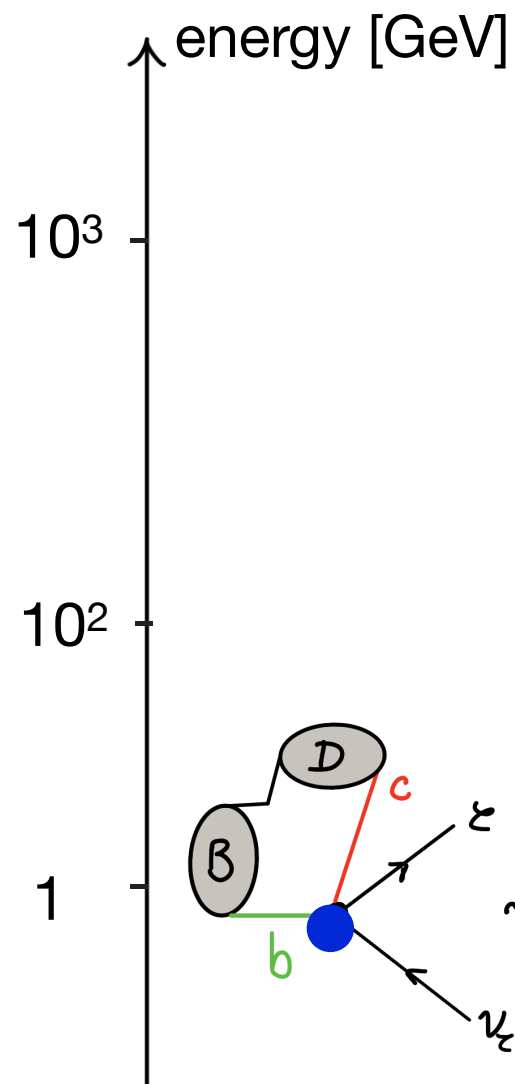
# Low-energy signatures of 3rd family new physics

## Flavor-changing low-energy probes

- Leading effects in  $3 \rightarrow \text{light}$  transitions: B physics

e.g. semileptonic  $3 \rightarrow 2$  transitions:  $b \rightarrow s(d)\ell\ell^{(\prime)}, b \rightarrow s(d)\nu\nu$   
 largest effects expected for  $\tau, \nu_\tau$ :  $b \rightarrow c(u)\ell\nu$

- Subleading, correlated effects in  $2 \rightarrow 1$  transitions: K physics



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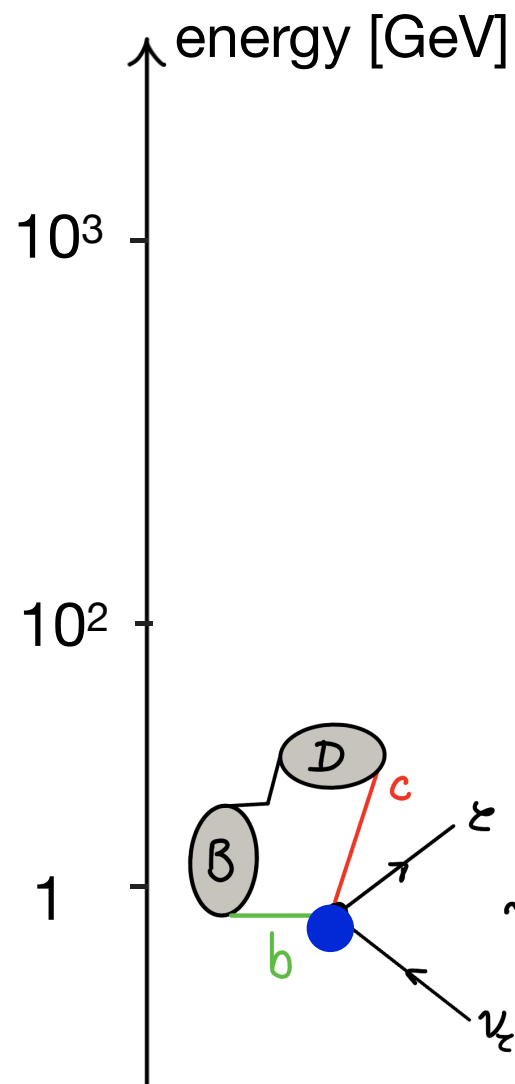
## Status & prospects in B physics

- In the **next 15 years**, LHCb & Belle II should collect  **$\sim 100\times$**  the **B** mesons they have now.

This means:

- CKM matrix elements  $< 1\%$
- LFU ratios in SL decays to  $O(1\%)$  level
- observe CPV in  $B_s$
- measure CPV in charm precisely

Important progress:  $B_s$  and D mixing  
 are already leading constraints on flavored heavy NP!



# FCNCs with Taus

Probing  $b \rightarrow s\tau\tau$  directly is experimentally very challenging.

Recently, several **remarkable results**:

- $B^+ \rightarrow K^+\tau^+\tau^-$

[★ =new world best]

[CKM2025] Belle [711 fb-1] incl. tagging + Belle II [365 fb-1]:  $\text{BR} < 8.7 \times 10^{-4}$  at 90% C.L [★]

- $B^0 \rightarrow K^{*0}\tau^+\tau^-$

[CKM2025] Belle II [365 fb-1]:  $\text{BR} < 1.8 \times 10^{-3}$  at 90% CL

- other:

[2024] bound on  $C_{9\tau}$  from  $\tau^+\tau^-$  rescattering in  $B^0 \rightarrow K^{*0}\mu^+\mu^-$  competitive with direct

[CKM2025] LHCb [5.4 fb-1]: searches for  $B^0 \rightarrow K^+\pi^-\tau^+\tau^-$  &  $B^0 \rightarrow K^+K^-\tau^+\tau^-$  translate in  $\text{BR} < \mathcal{O}(10^{-4})$  on  $B^0 \rightarrow K^{*0}\tau^+\tau^-$  and  $B_s \rightarrow \phi\tau^+\tau^-$  [★]

**Limit:** even with full LHCb and Belle II datasets, bounds likely to exceed SM ( $\sim 10^{-7}$ ) by  $10^{2-3}$ .

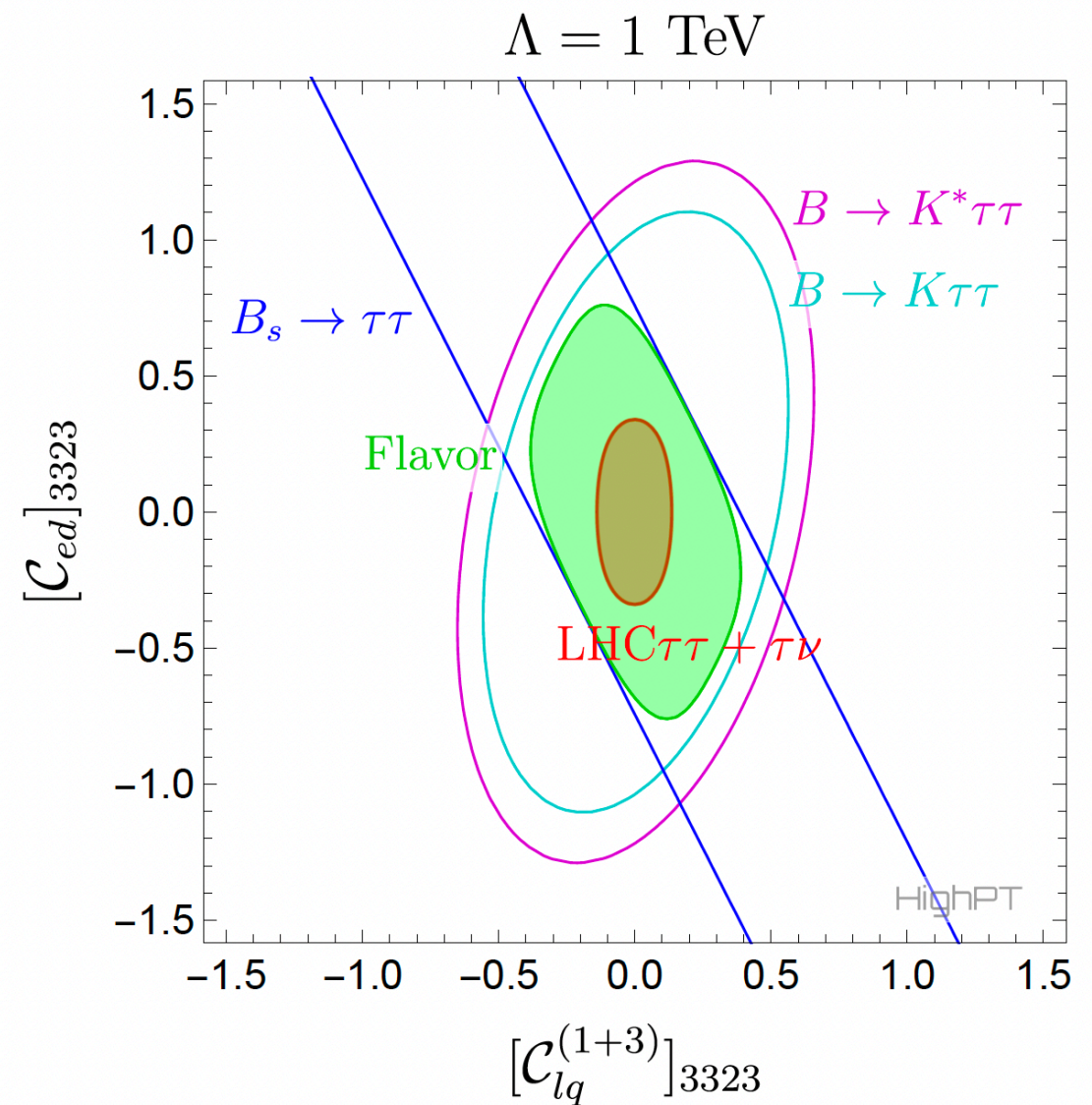
⇒ Will need Tera-Z to go beyond!

# FCNCs with Taus

Still, we start being **sensitive** to scenarios with **large NP couplings** to 3rd family

Complementary with **dilepton tails** at the LHC! [though not as powerful just yet]

[see e.g. Faroughy, Greljio, Kamenik, 1609.07138  
Greljo, Marzocca. 1704.09015  
Allwicher, Faroughy, Jaffredo, Sumensari, Wilsch 2207.10714 ...]



# FCNCs with Neutrinos

Currently the only measured FCNCs sensitive to NP interacting with 3rd family leptons are the **dineutrino** modes  $b \rightarrow s\nu\bar{\nu}$  and  $s \rightarrow d\nu\bar{\nu}$ .

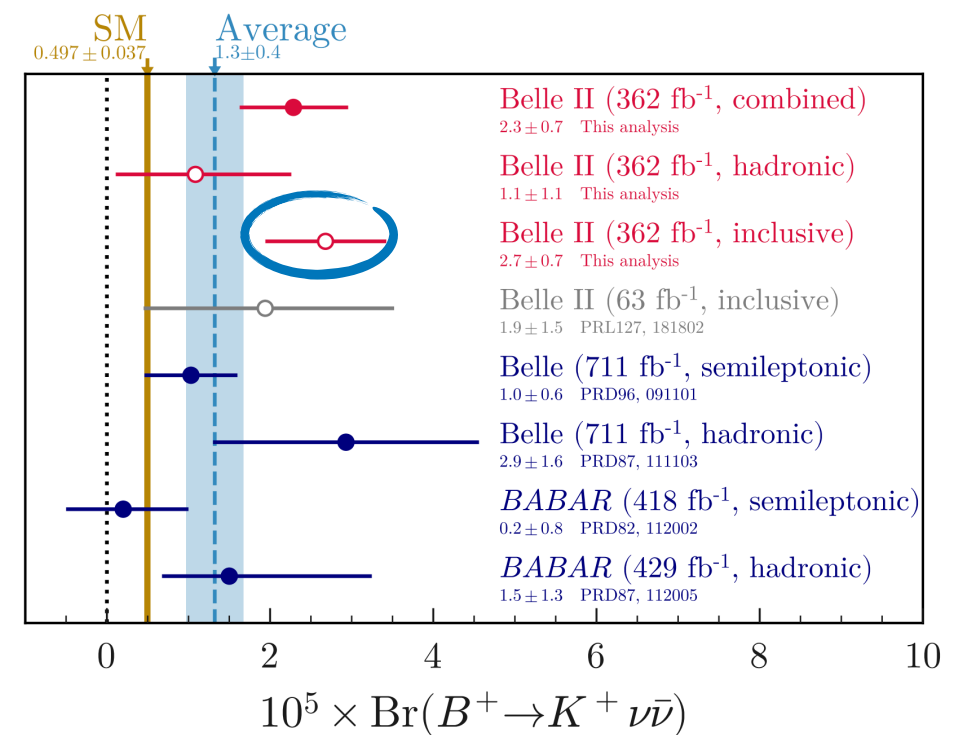
- **Very precise SM prediction:**
  - advantage wrt to dilepton modes: **no “charm loop”** effects ( $\nu$  don't couple to  $\gamma$ )
  - **theory uncertainty** dominated by  $V_{cb}$  ( & form factors for  $B \rightarrow K/\pi$ ).
  - not yet a showstopper, but solving the  $V_{cb}$  puzzle + lattice improvements will be important to exploit future exp. precision
- **Powerful tests of NP flavor structure:**
  - all sensitive to **leading  $U(2)_q$  breaking**
  - $s \rightarrow d$  sensitive also to **subleading  $U(2)_q$ ,  $U(2)_d$  spurions**

# FCNCs with Neutrinos

Snapshot of the experimental situation [much more in the dedicated talks!]

$$b \rightarrow s\nu\bar{\nu}$$

- Belle II [2023]: first evidence for  $B^+ \rightarrow K^+ \nu\bar{\nu}$ ;
- combined result  $\sim 2\sigma$  above SM ( $\Lambda_{\text{eff}} \sim 6 \text{ TeV}$ )
- Target: 10% precision @ Belle II
- work ongoing on  $K^{*0,+}$  and  $K_S$



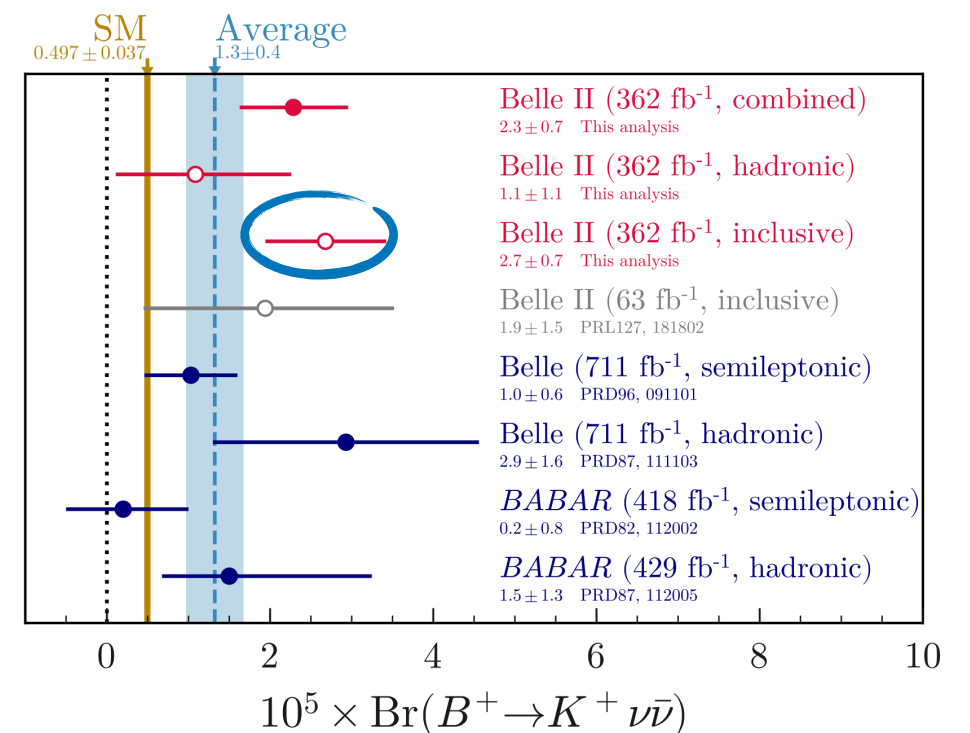


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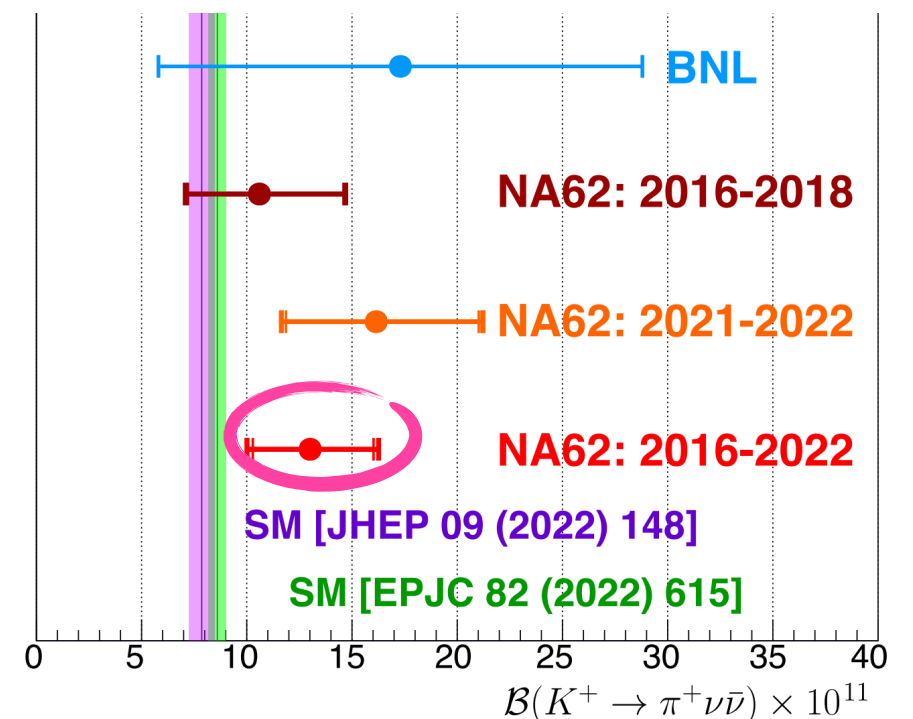
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## $s \rightarrow d \nu \bar{\nu}$

- Na62 [2024]: first evidence for  $K^+ \rightarrow \pi^+ \nu \bar{\nu}$
- combined result  $\sim 2\sigma$  above SM ( $\Lambda_{\text{eff}} \sim 80 \text{ TeV}$ )
- Target: 15% precision @Na62 (5% @HIKE<sup>+</sup>)
- $K_L \rightarrow \pi^0 \nu \bar{\nu}$ :  $\text{BR}_{\text{SM}} \sim \text{O}(10^{-11})$ ,  $\text{BR}_{\text{exp}} < \text{O}(10^{-9})$  [KOTO], atm not competitive with charged mode for NP searches



# LFUV in $b \rightarrow c \ell \bar{\nu}$

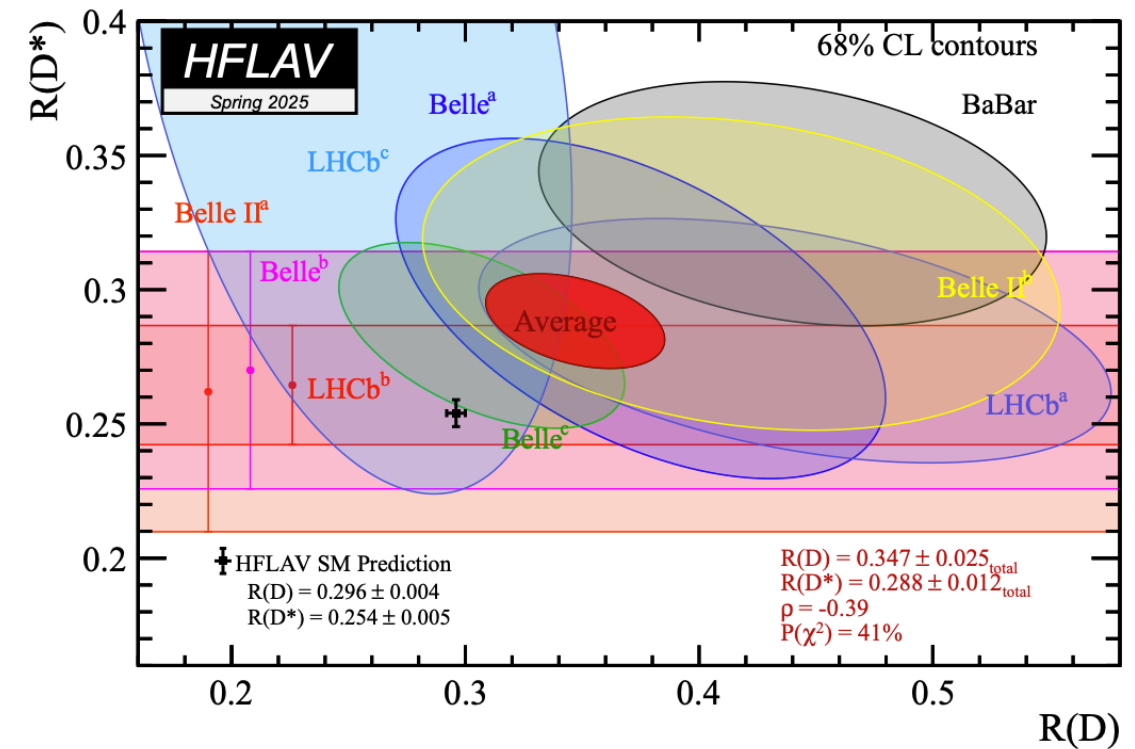
[\*w/o the Belle II measurement presented @CKM2025 by I. Tsaklidis, 1.3 $\sigma$  from the HFLAV average displayed here]

15-year old tension far from being settled:

$$R_{D^{(*)}} = \frac{\mathcal{B}(B \rightarrow D^{(*)} \tau \bar{\nu})}{\mathcal{B}(B \rightarrow D^{(*)} \ell \bar{\nu})}$$

~ 10% enhancement due to excess in  $\tau$  mode combined\* ~3.8 $\sigma$  above SM ( $\Lambda_{\text{eff}} \sim \text{O}(1) \text{ TeV}$ )

- SM well under control
- Two recent results [2025]: Belle II semil. & hadronic tag
- target for 2040: 1% @ Belle II [50ab<sup>-1</sup>], 3% @ LHCb[300fb<sup>-1</sup>]



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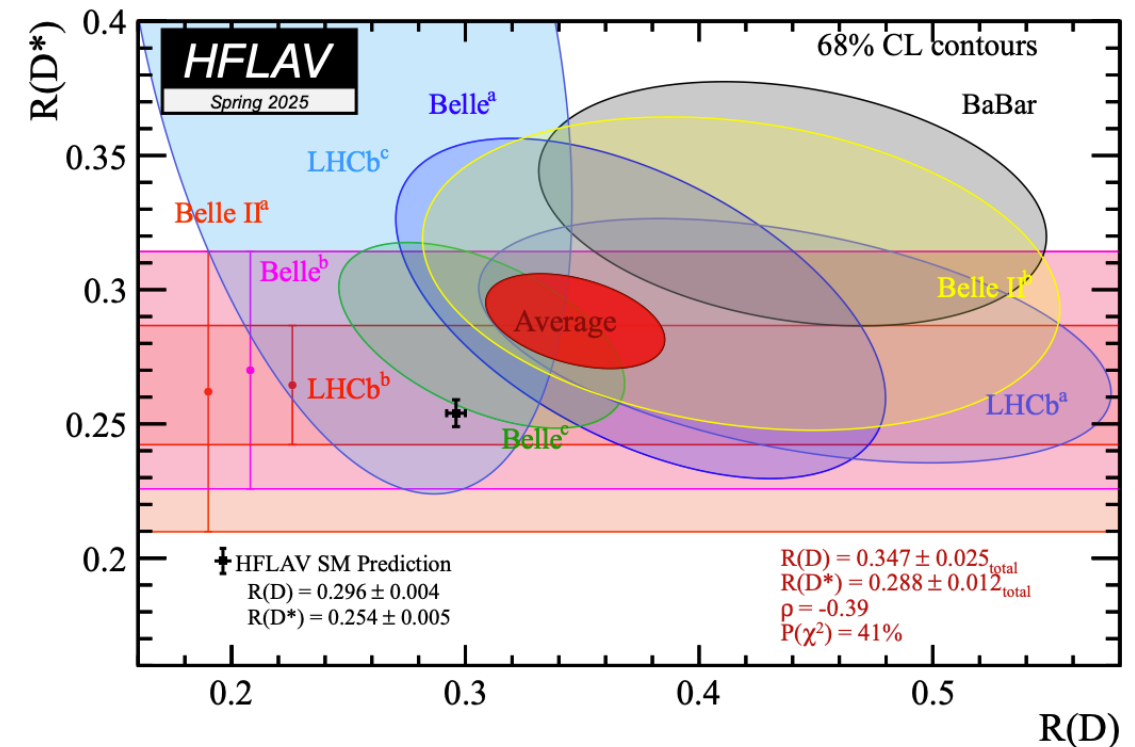
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If due to new physics, expect correlated excesses at low and high energies:

- $B \rightarrow K\tau\tau, B_s \rightarrow \tau\tau \sim \text{O}(100) \times \text{SM}$
- distortion of tails in  $pp \rightarrow \tau\tau, pp \rightarrow \tau + E_{\text{miss}}$   
 ATLAS/CMS already constrain a relevant portion of parameter space
- ...plus many more more “model dependent” signatures, incl. enhancement of  $b \rightarrow s\nu\bar{\nu}$

# Implications of a NP interpretation of $b \rightarrow s\nu\bar{\nu}$ data

**IF the  $b \rightarrow s\nu\bar{\nu}$  excess is due to [heavy] new physics**

- $\Lambda_{\text{eff}} \sim 6\text{-}7 \text{ TeV}$
- **should be  $\nu_\tau$**  to avoid constraints from  $B_s \rightarrow \mu^+\mu^-$  &  $R_{K^{(*)}}$   
 $\Rightarrow$  by  $\text{SU}(2)_L$  invariance, unavoidable contribution\* to  $b \rightarrow s\tau\tau$  (&  $b \rightarrow s\mu\mu, ee$  via RGE)  
[\*exact contribution depends on the Lorentz structure at work: vector, scalar, tensor...]

[Allwicher, Becirevic, Piazza, Rosauero-Alcaraz, Sumensari, 2309.02246; Bause, Gisbert Hiller 2309.00075; Allwicher, CC, Isidori, Stefanek, 2311.00020; Marzocca, Nardecchia, Stanzione, Toni 2404.06533; Allwicher, Bordone, Isidori, Piazza, Stanzione, 2410.21444]

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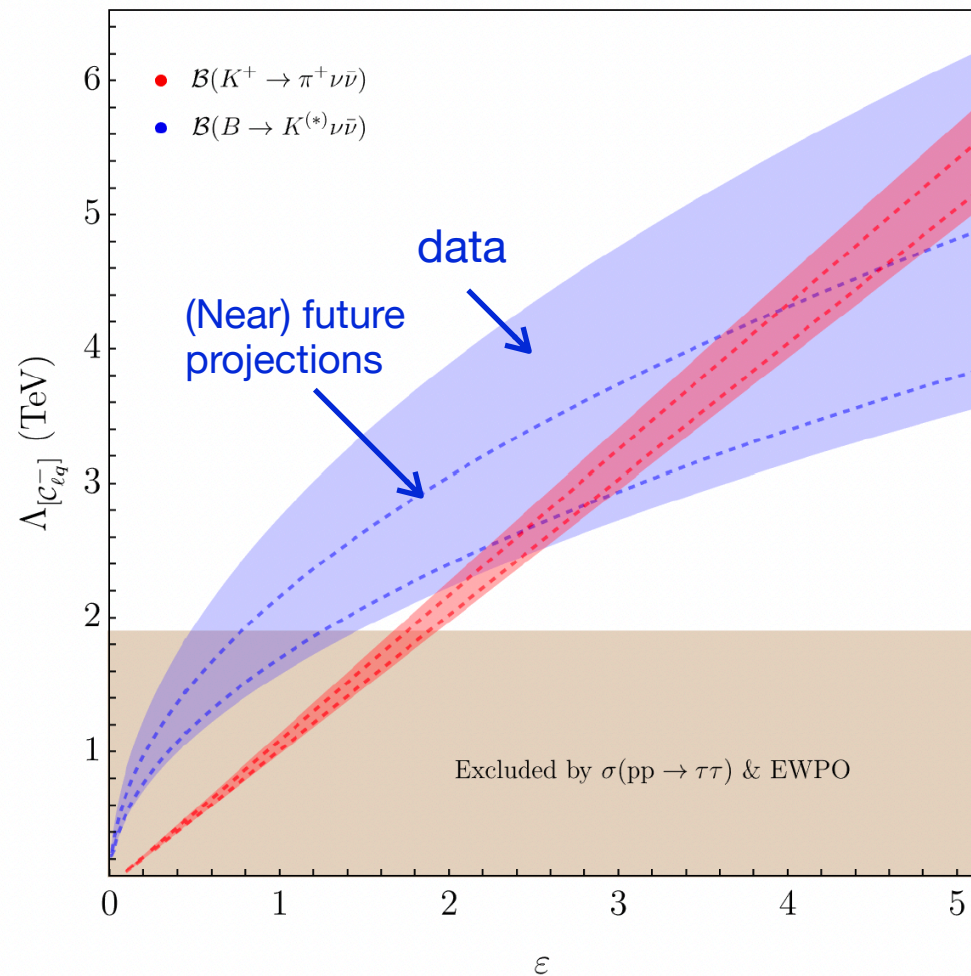
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 $\Rightarrow$  **leptoquarks** favoured  
 $\Rightarrow$  same NP would contribute to  $R_{D^{(*)}}^*$
- If the underlying flavor structure is **U(2)-like**  
 $\Rightarrow$  expect correlated effects in  $s \rightarrow d\nu\nu, b \rightarrow d\nu\nu$  (+  $\tau\tau$  modes)

Currently, excesses in  $b \rightarrow s\nu\bar{\nu}, s \rightarrow d\nu\nu$  &  $R_{D^{(*)}}$  are reconcilable with U(2)-like structure and absolute scale  $\Lambda \sim \mathcal{O}(1 \text{ TeV})$  [works at the EFT level, less trivial in explicit modes]



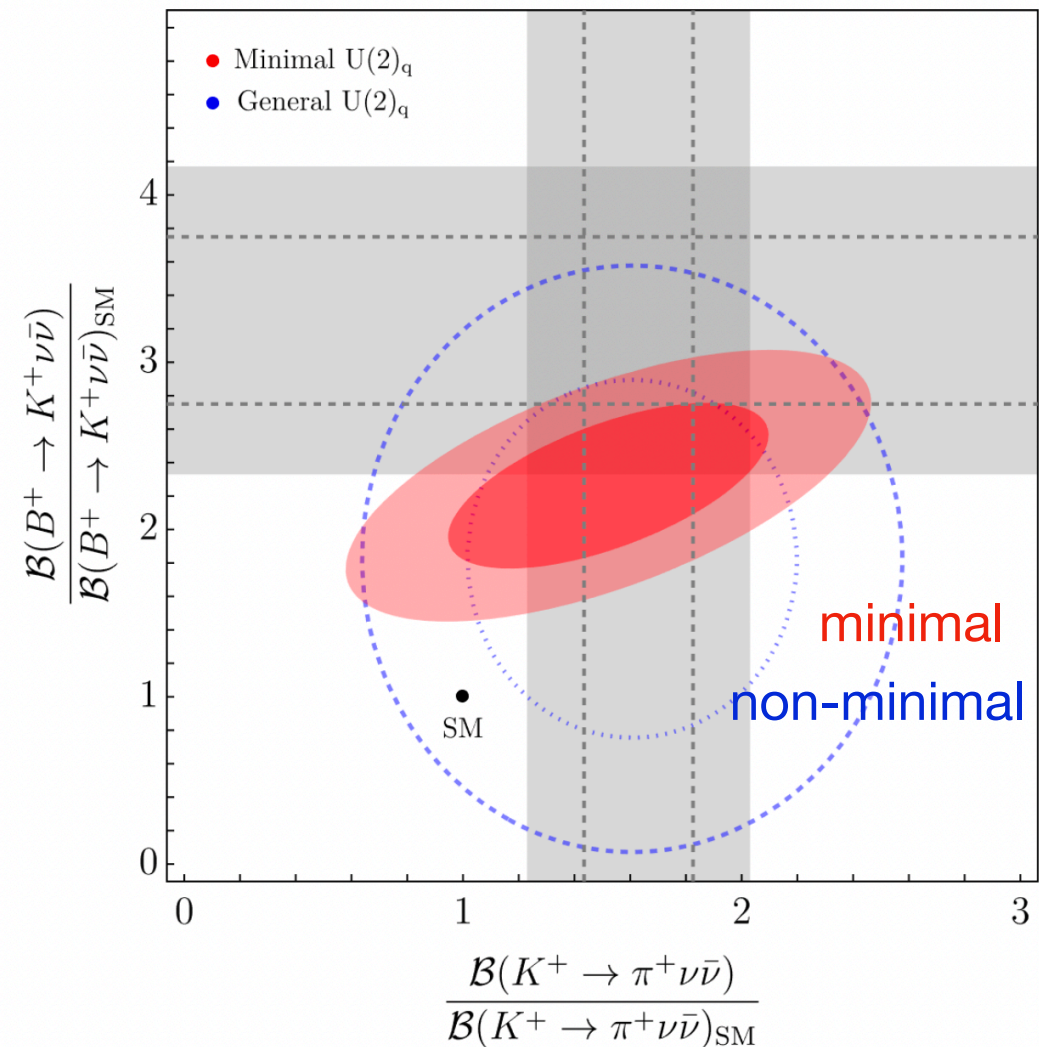
# Implications of a NP interpretation of $b \rightarrow s\nu\bar{\nu}$ data

determine size of breaking terms  
[those not already fixed by masses & CKM]



e.g. leading  $U(2)_q$  breaking can be determined by comparing  $b \rightarrow s$  vs  $s \rightarrow d$

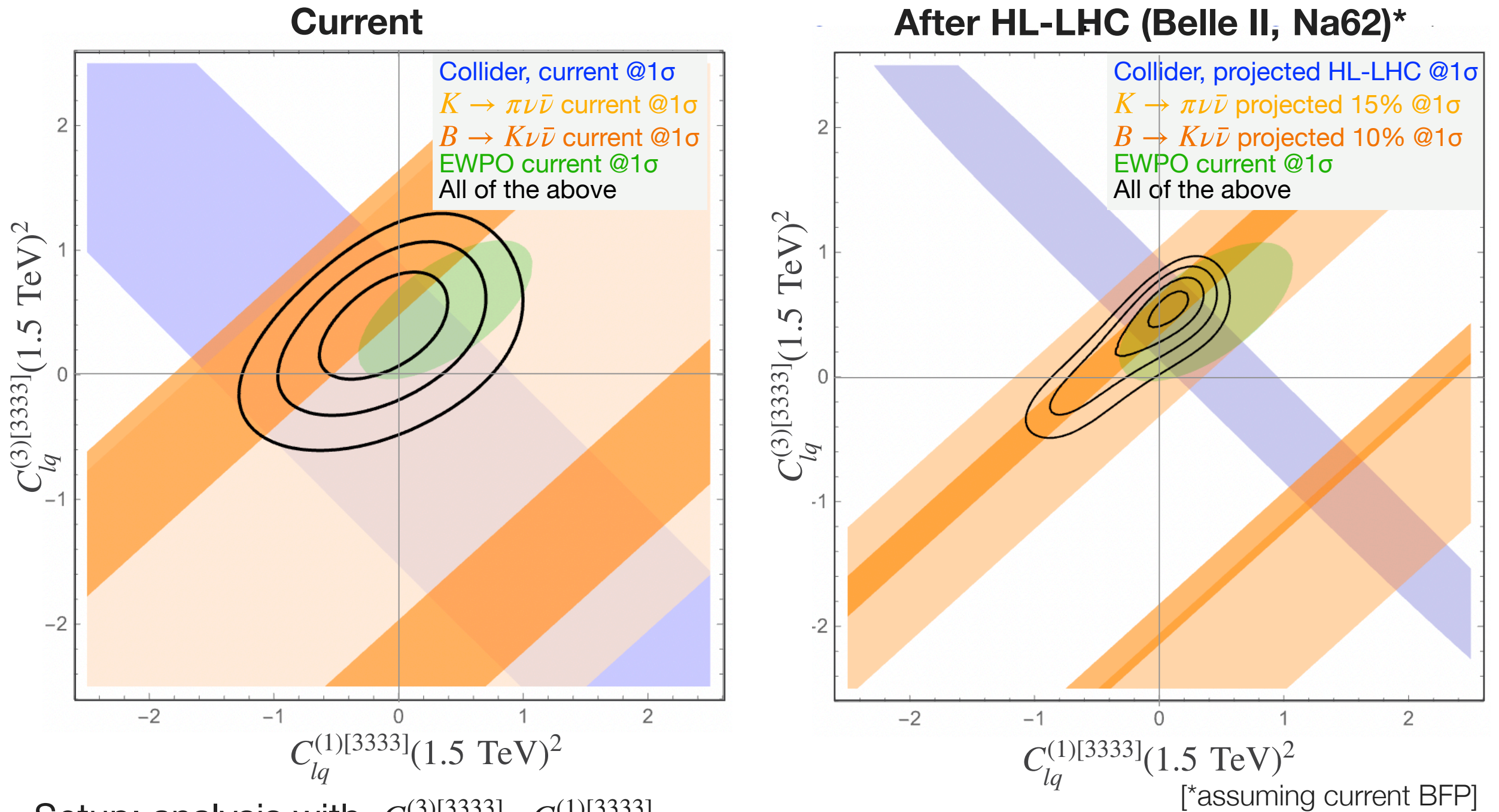
discriminate between  $\neq$  breaking patterns  
via characteristic correlations



[Allwicher, Bordone, Isidori, Piazza, Stanzione, 2410.21444]

# Implications of a NP interpretation of $b \rightarrow s\nu\bar{\nu}$ data

Combining  $\nu\bar{\nu}$  data with EW precision and collider:



Setup: analysis with  $C_{\ell q}^{(3)[3333]}$ ,  $C_{\ell q}^{(1)[3333]}$ ,  $\epsilon_F$

[Plots from Allwicher, CC, Isidori, Stefaneek, 2311.00020 updated with more recent data]



# Why Tera-Z can help probe Flavored New Physics

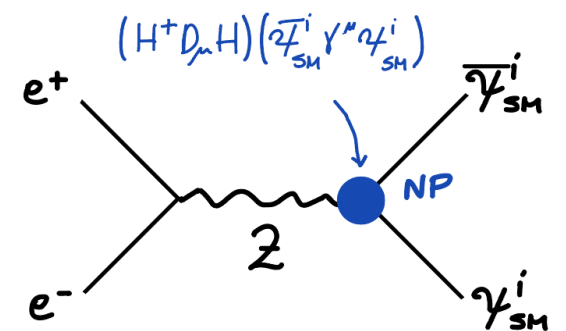
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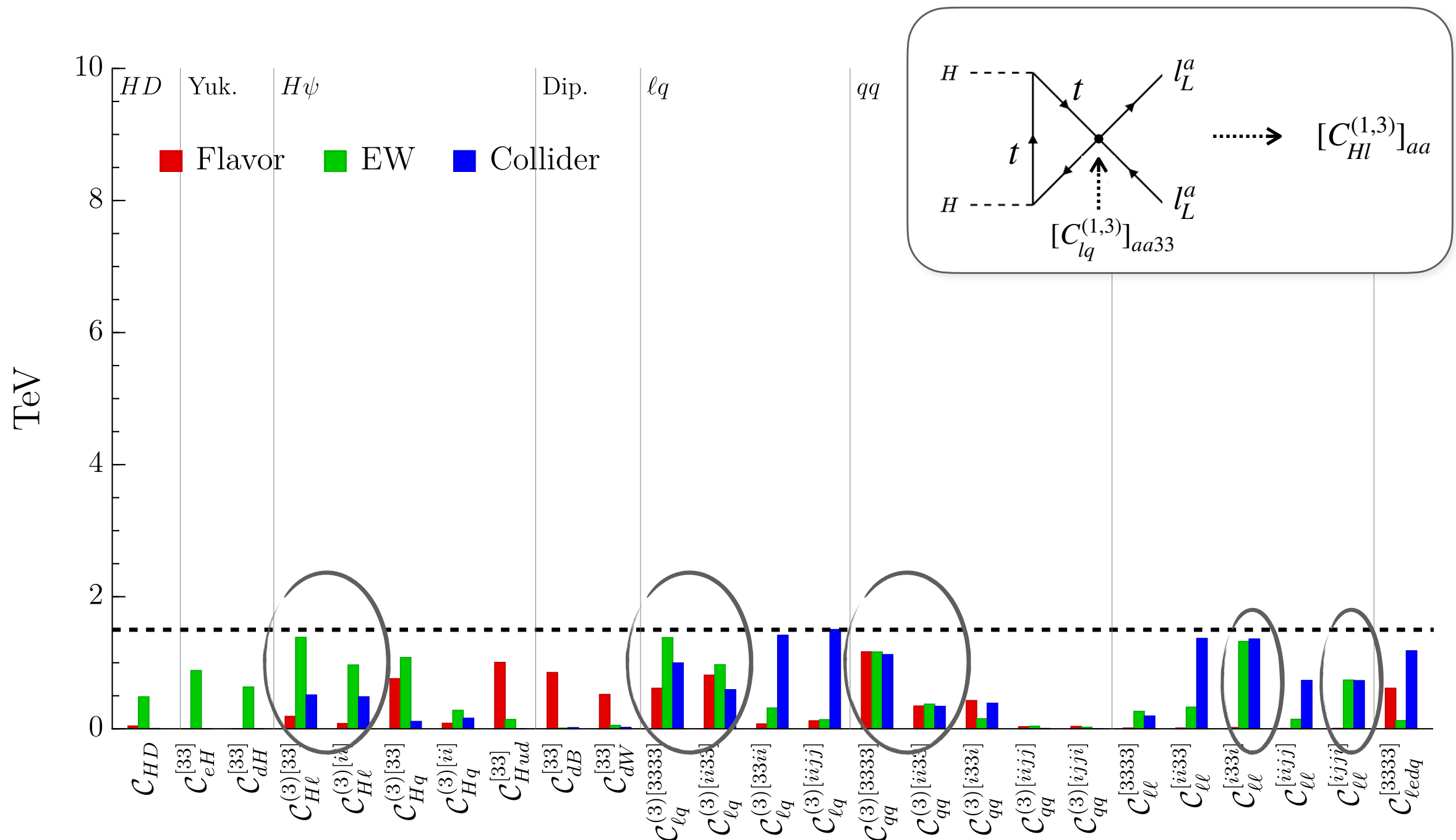
## EW precision

- 3rd fam. NP is “protected” against direct searches at the LHC & flavor, but *not* protected vs EWPT - “*everything runs into EW*”
- At a Z factory, we can use the flavor blindness of SM gauge interactions to probe NP coupled to any generation via EWPT  
⇒ **EWPT** are powerful probes of flavor **non-universal NP**



# Why Tera-Z can help probe Flavored New Physics

Even now LEP bounds have a strength comparable to current direct searches for operators involving mostly the 3rd generation:

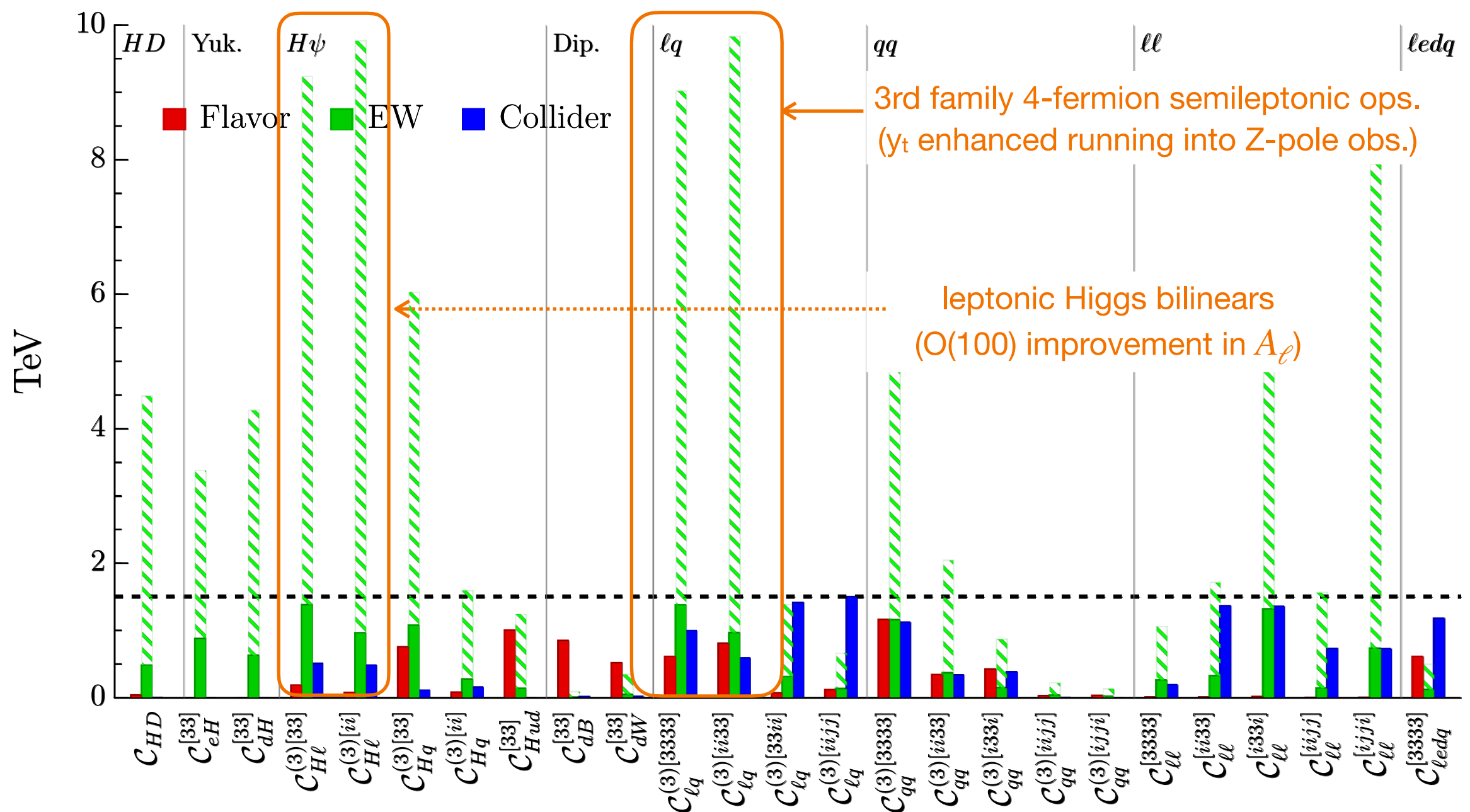


# Why Tera-Z can help probe Flavored New Physics

....with  $\approx 10^5$  more Z bosons than LEP,

A tera-Z machine in its Z-pole run could probe 3rd-family NP up to  $\sim 10$  TeV!

[And flavor universal NP up to much higher scales, O(30-50 TeV)]



# Why Tera-Z can help probe Flavored New Physics

Looking into the future, a Tera-Z facility can help test flavoured NP via the *combination* of its EW precision & flavor program. How?

## Flavor

- combines the best features of pp colliders and B factories  
high statistics, “closed” kinematics, high boost of  $b$  and  $\tau$ , access to all b hadrons
- precise measurements of  $b \rightarrow s\tau\tau$  &  $b \rightarrow s\nu\nu$ , incl.  $b \rightarrow d$  counterpart
- test LFU in  $\tau$  decays @  $O(10^{-4})$
- dedicated studies with detector simulation (IDEA baseline) + background modelling available for a few channels, many more under development

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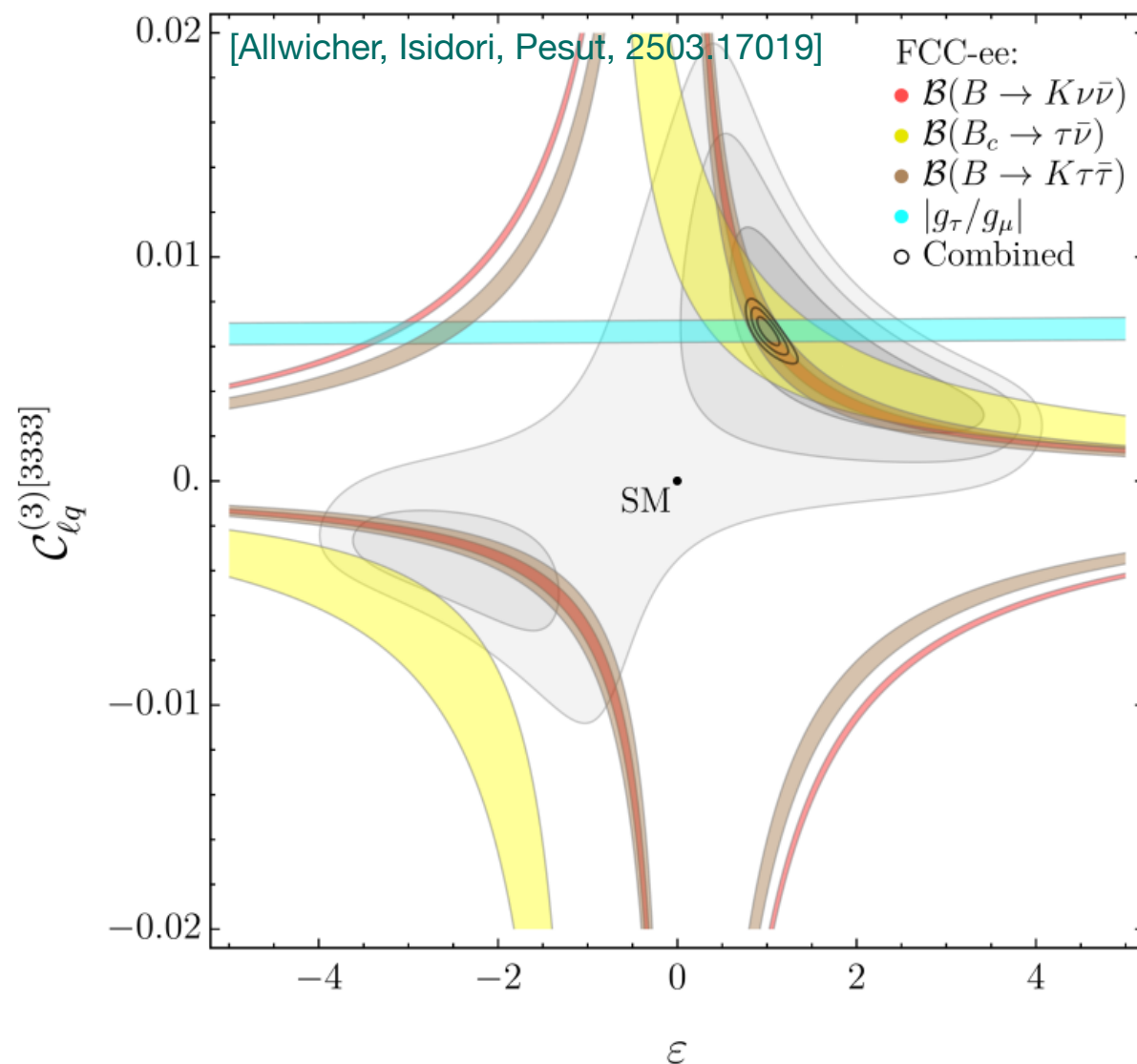
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## Flavor

**Example: determination of alignment of NP in flavor space** [ $\epsilon_F$  = leading U(2)<sub>q</sub> breaking]



- current data (gray):  
main sensitivity to  $\epsilon_F$  from  $B \rightarrow K \nu \bar{\nu}$  &  $K \rightarrow \pi \nu \bar{\nu}$ , subleading  $R_{D^{(*)}}$
- with FCC-ee (flavor), two “new” hyperbolae:  
 $B \rightarrow K \tau \bar{\tau}$  &  $B_c \rightarrow \tau \nu$

[setup as before: analysis with  $C_{\ell q}^{(3)[3333]}$ ,  $C_{\ell q}^{(1)[3333]}$ ,  $\epsilon_F$ ]



# Conclusions

NP with **flavor protection** can exist at the **TeV** if coupled dominantly to the 3rd family.

**Many signatures** to look for at **existing experiments**:

- direct 3rd family searches
- precision measurements in B & K decays

These are the best path to discovery until the next collider.

In this context, **dineutrino modes** provide privileged test of the **flavor structure of NP** (alignment in flavor space, info on non minimal breaking patterns).

Should the mild excess solidify, it would imply a plethora of signatures.

**Looking forward**, a **tera-Z** machine is ideal in testing “flavored” NP scenarios

- unprecedentedly precise EWPT that cannot be bypassed by flavor symmetries
- major advancements incl. access to new channels