

Outlook for theoretical calculations

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Points for discussion

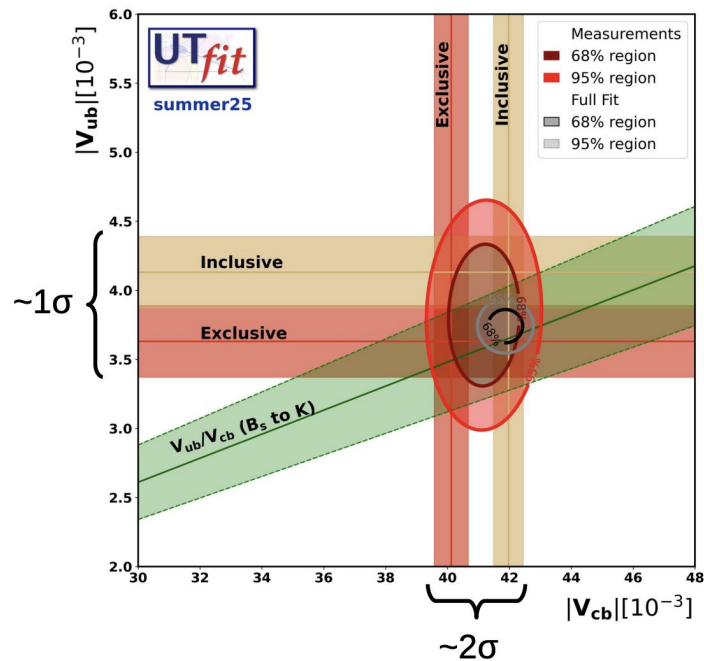
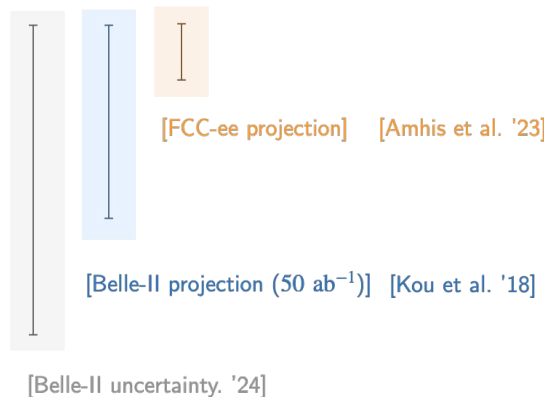
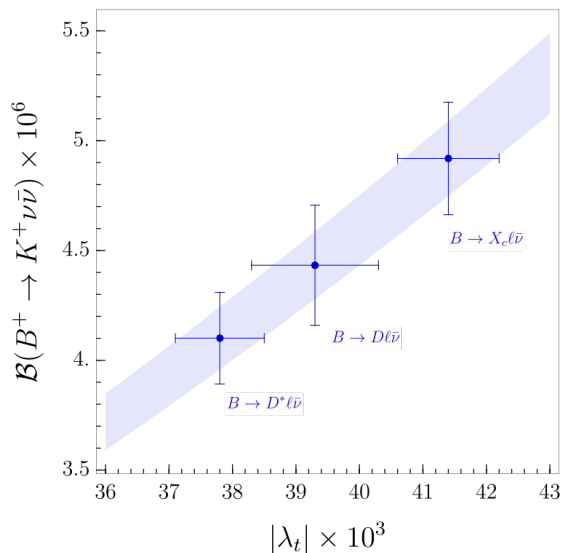
- Improving SM parametric inputs:
 - Exclusive V_{cb} and $B \rightarrow D^{(*)}$ form factors from LQCD
 - Comparing $B \rightarrow D$ with $B_s \rightarrow D_s^{(*)}$ form-factors
 - Improving kinematic parameterization with unitarity constraints
 - V_{cb} inclusive: Inclusive decays from the lattice
- Hadronic matrix elements for rare decays:
 - Charm loops: lattice vs dispersive bound vs hadronic models
 - QCD-unstable final states, e.g., $B \rightarrow K^*$
- Inputs for NP calculations:
 - Improvements in inclusive calculations
 - Need f_T for $B \rightarrow D$, tensor FFs have been calculated for $B \rightarrow D^*$
 - New observables to reduce theoretical uncertainties

Which CKM value to use?

$$\frac{|V_{ts}^* V_{tb}|}{|V_{cb}|} = 0.965(1)$$

$$|\lambda_t| \times 10^3 = \begin{cases} 41.4 \pm 0.8, & (B \rightarrow X_c l \bar{\nu}) \\ 39.3 \pm 1.0, & (B \rightarrow D l \bar{\nu}) \\ 37.8 \pm 0.7, & (B \rightarrow D^* l \bar{\nu}) \end{cases}$$

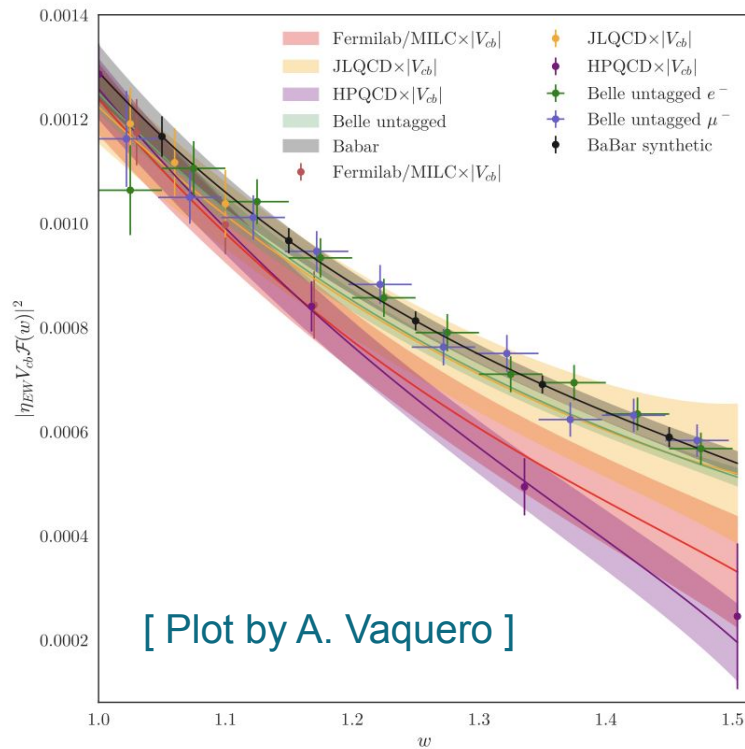
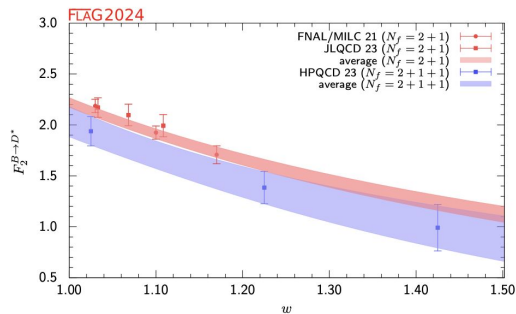
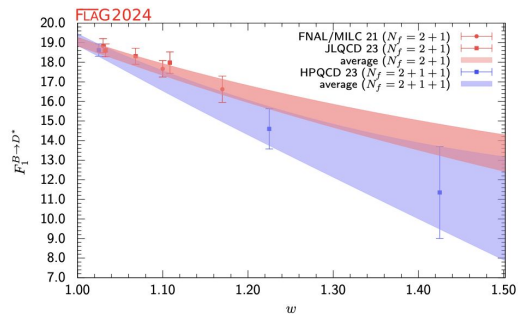
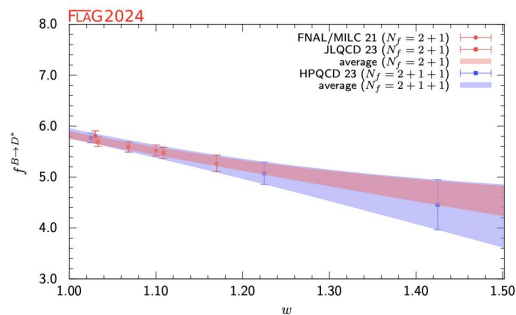
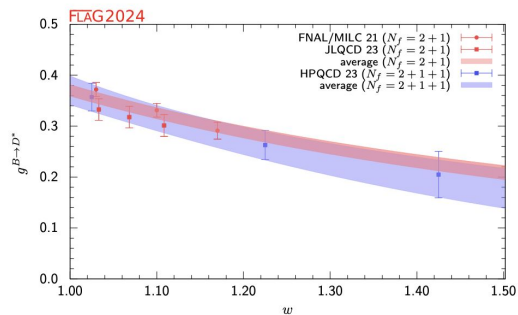
[HFLAV, '22] [FLAG, '21] [HFLAV, '22]



B → D*: Lattice vs. Exp. data

$$\frac{d\mathcal{B}}{dq^2}(B \rightarrow D^* \ell \nu) \propto |V_{cb}|^2 |\mathcal{F}(w)|^2$$

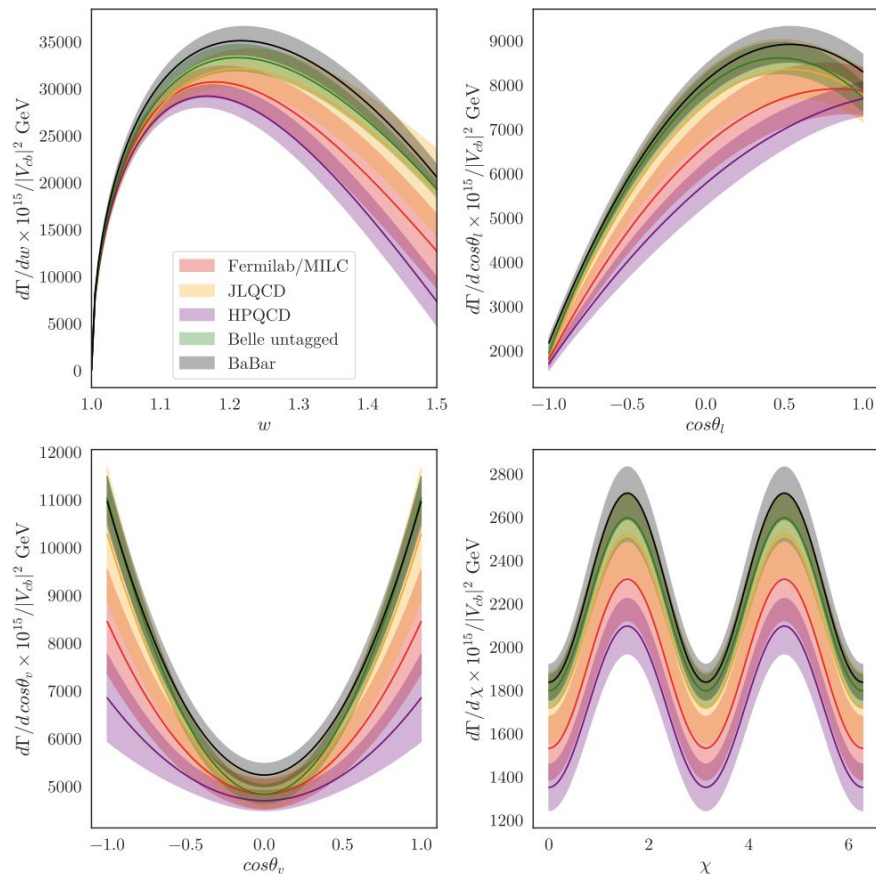
$$\left[w = \frac{m_B^2 + m_{D^*}^2 - q^2}{2m_B m_{D^*}} \right]$$



$$|V_{cb}|^{\text{HPQCD}} = 39.31(74) \times 10^{-3}$$

$$|V_{cb}|^{\text{FerMILC}} = 38.17(85) \times 10^{-3}$$

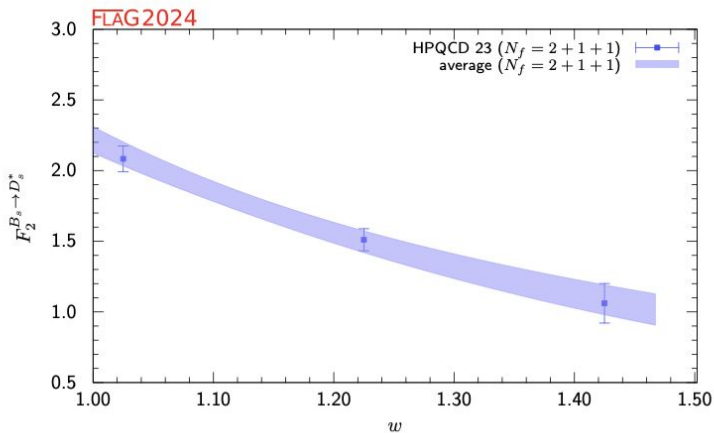
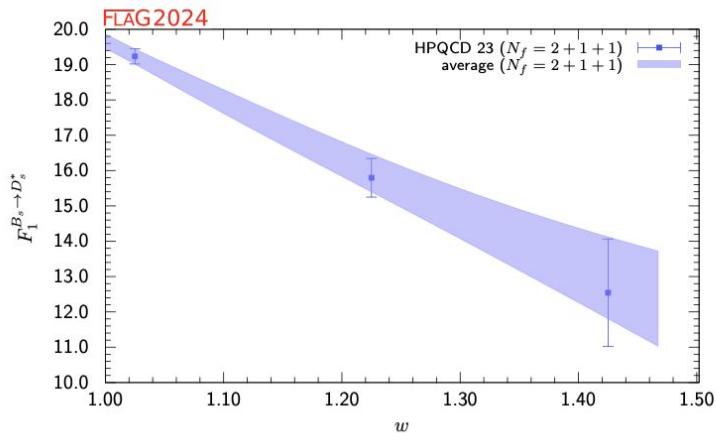
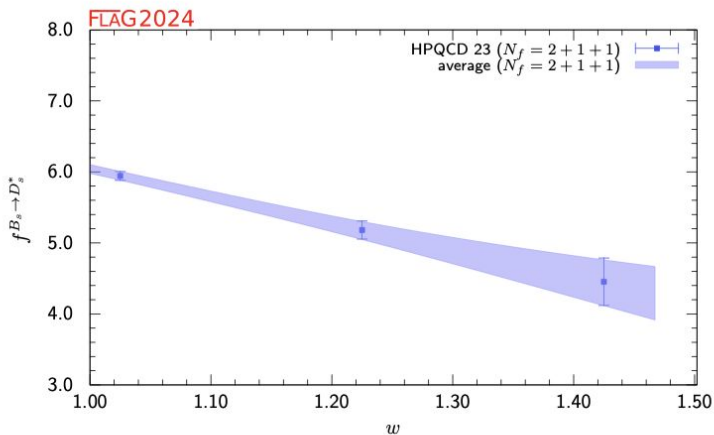
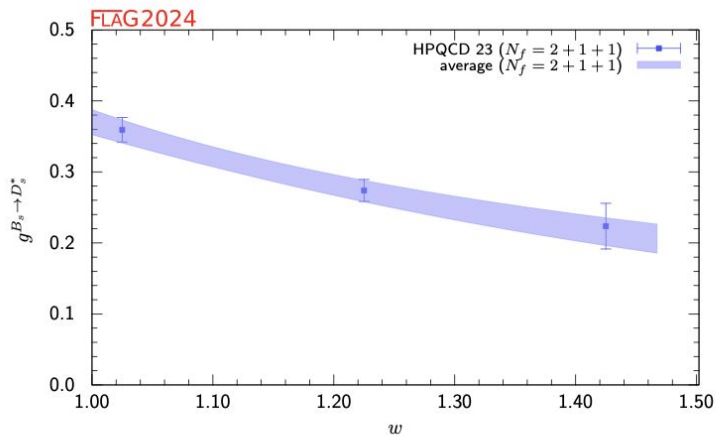
B $\rightarrow$ D*: Lattice vs. Exp. data



- Belle-II data and more/improved lattice inputs will be fundamental to solve this discrepancy.
- HPQCD has also computed $B_s \rightarrow D_s^*$.

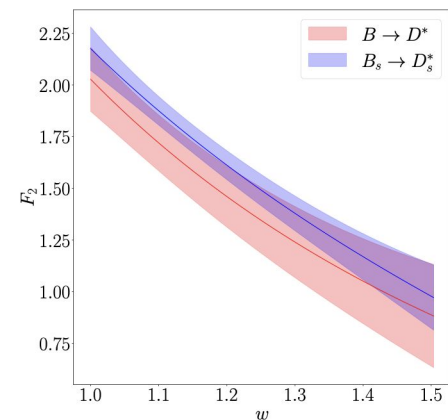
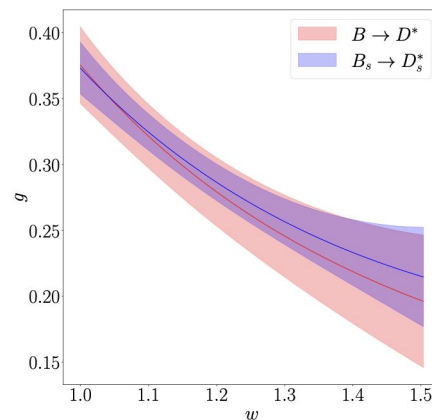
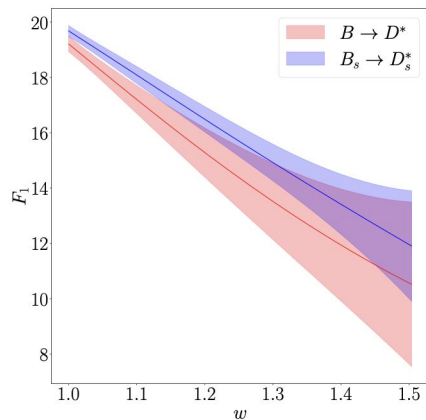
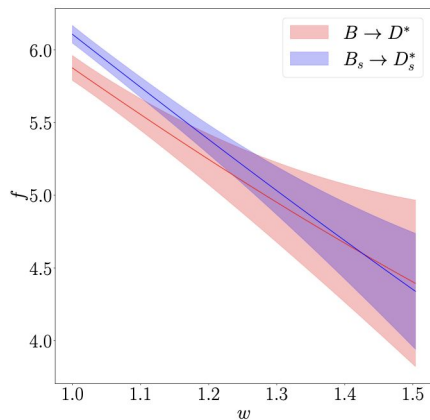
[Plots by A. Vaquero]

Bs $\rightarrow$ Ds*



- Bs $\rightarrow$ Ds* easier on lattice than B $\rightarrow$ D*
- Should be similar to B $\rightarrow$ D* ; offers a sanity check
- Experimentally determined shape would be useful

Comparing $B \rightarrow D^*$ with $B_s \rightarrow D_s^*$



[(HPQCD) Harrison and Davies, [2304.03137](#)]

Improving kinematic parameterizations

- Dispersive matrix methods - from continuum results at large q^2 , provides shape for all q^2 without explicit parameterization or associated truncation error

[Martinelli et al. [2105.08674](#) , ...]

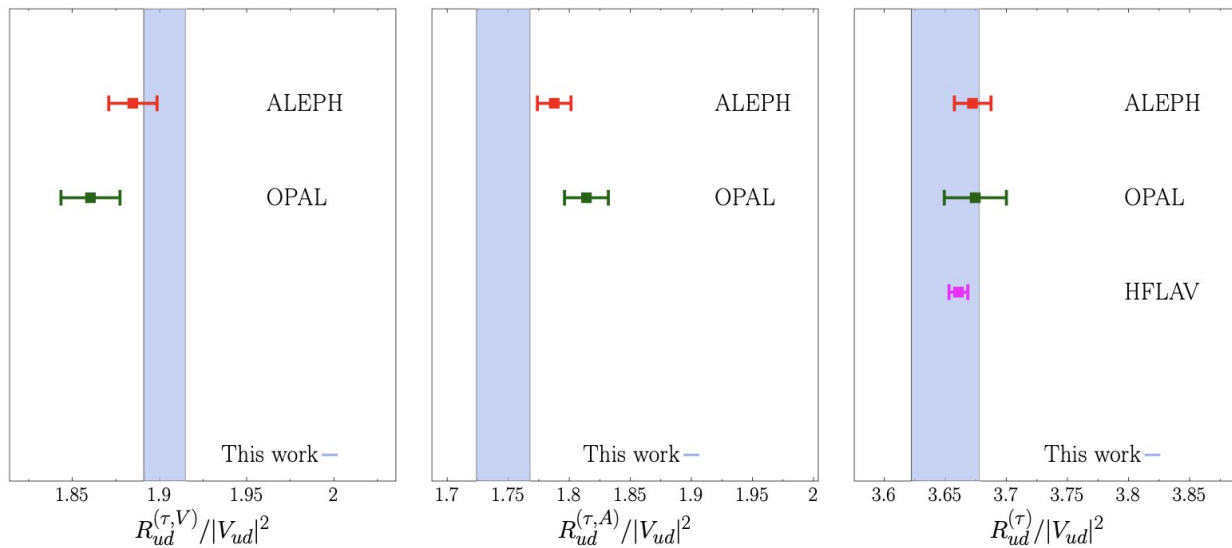
- Improved parameterization - build unitarity constraints into z-parameterization
- Evaluate truncation error

[(RBC/UKQCD) Flynn et al. [2303.11280](#)]

[Danny's talk]

Lattice for inclusive decays: Tau

$$\Gamma(\tau \rightarrow X \nu_\tau) \propto \text{Im} \left(\text{Diagram} \right)$$



$$R_{ud}^{(\tau)} = \frac{\Gamma(\tau \rightarrow X_{ud} \nu_\tau)}{\Gamma(\tau \rightarrow e \bar{\nu}_e \nu_\tau)}$$

$$\Delta_{V-A}^{(\tau)} \equiv \frac{R_{ud}^{(\tau, V)} - R_{ud}^{(\tau, A)}}{R_{ud}^{(\tau)}}$$

$$\Delta_{V-A}^{(\tau)} = 0.042 \quad (5)$$

vs experiment

$$\Delta_{V-A}^{(\tau)}(\text{ALEPH}) = 0.026 \quad (7)$$

$$\Delta_{V-A}^{(\tau)}(\text{OPAL}) = 0.013 \quad (7)$$

[(ETMC) Evangelista et al. [2308.03125](#)]

Lattice for inclusive decays: Charm

Ds -> X lv

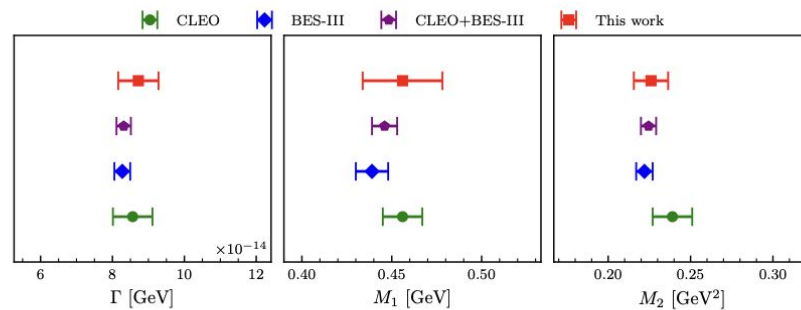
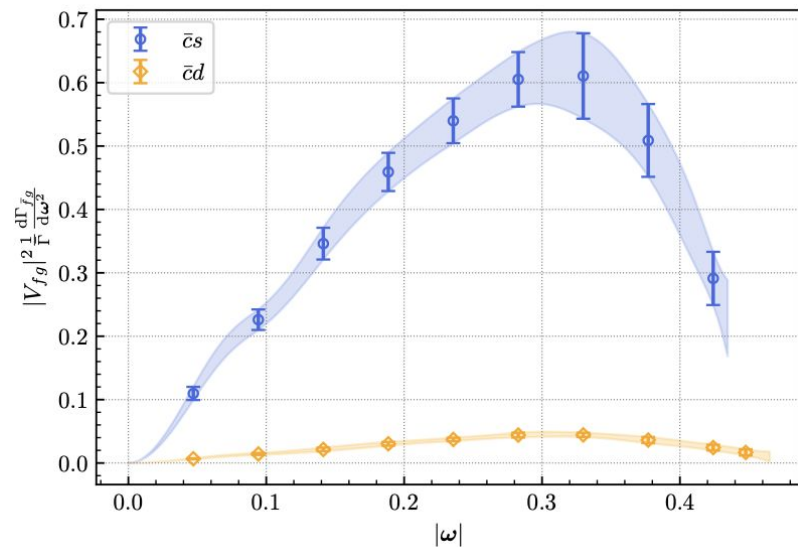
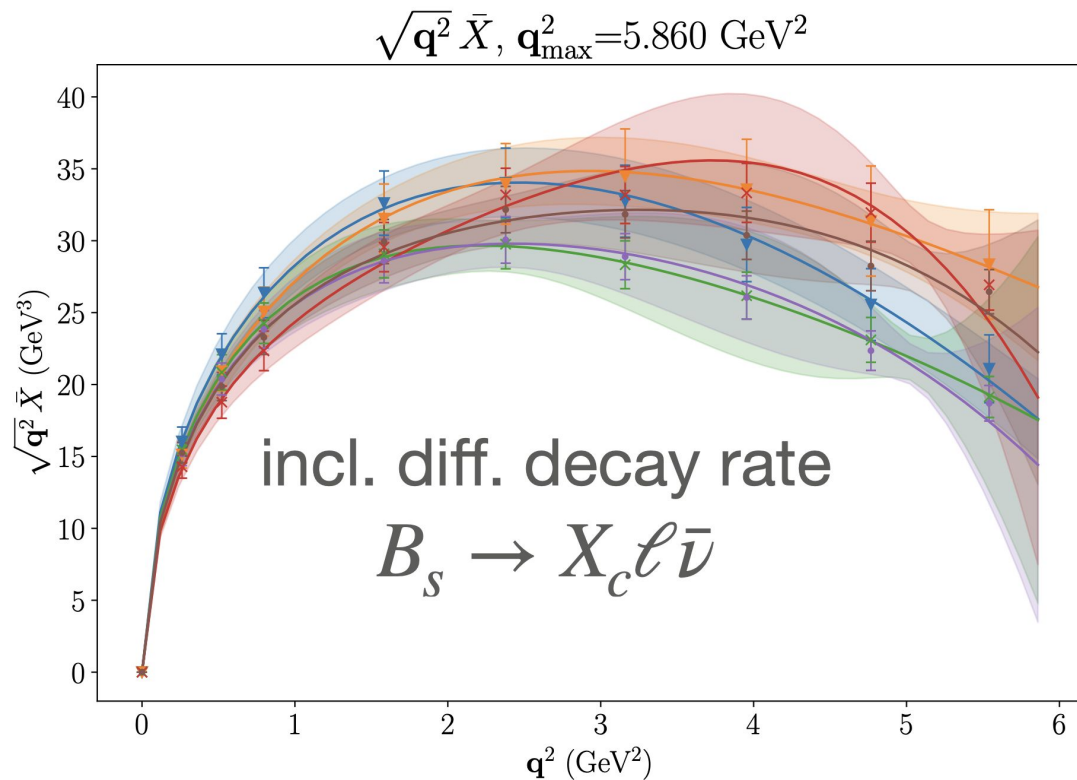


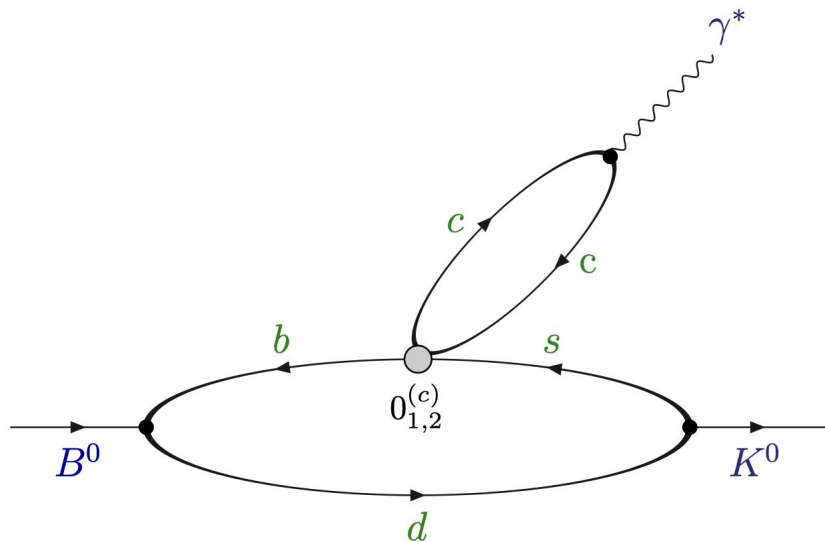
FIG. 4. Comparison between the experimental results from Refs. [44–46, 56] and our theoretical prediction (red points), for the decay rate (left-panel) and for the first (middle-panel) and second (right-panel) lepton moment.

Lattice for inclusive decays: Bs

- Lattice inclusive calculations are making rapid progress, working toward the B
- Barone et al. calculated inclusive differential decay rate for $B_s \rightarrow X_c \ell \bar{\nu}$
[Barone et al. [2305.14092](#)]
- Important to understand systematic effects involved
[Kellermann et al. [2504.03358](#)]



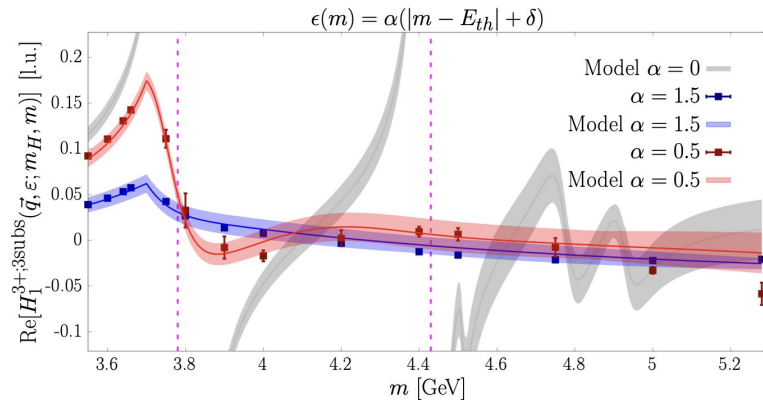
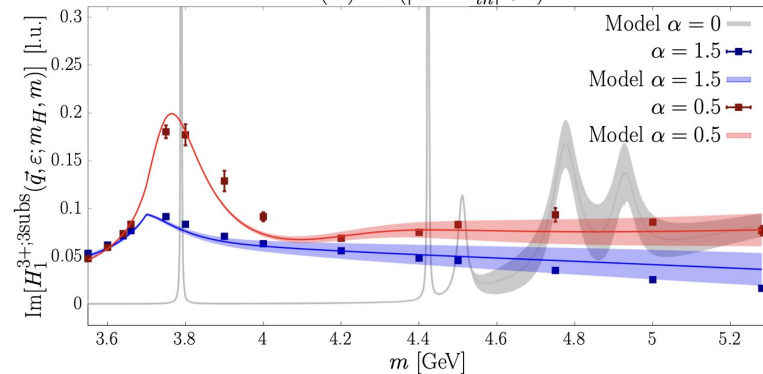
Charm loops



- As with inclusive decays, 4 pt correlators give access to the relevant physics, ie. non-local form factors
- Free from assumptions about factorization

$$B_s \rightarrow \eta_s \ell^+ \ell^-$$

$$\epsilon(m) = \alpha(|m - E_{th}| + \delta)$$

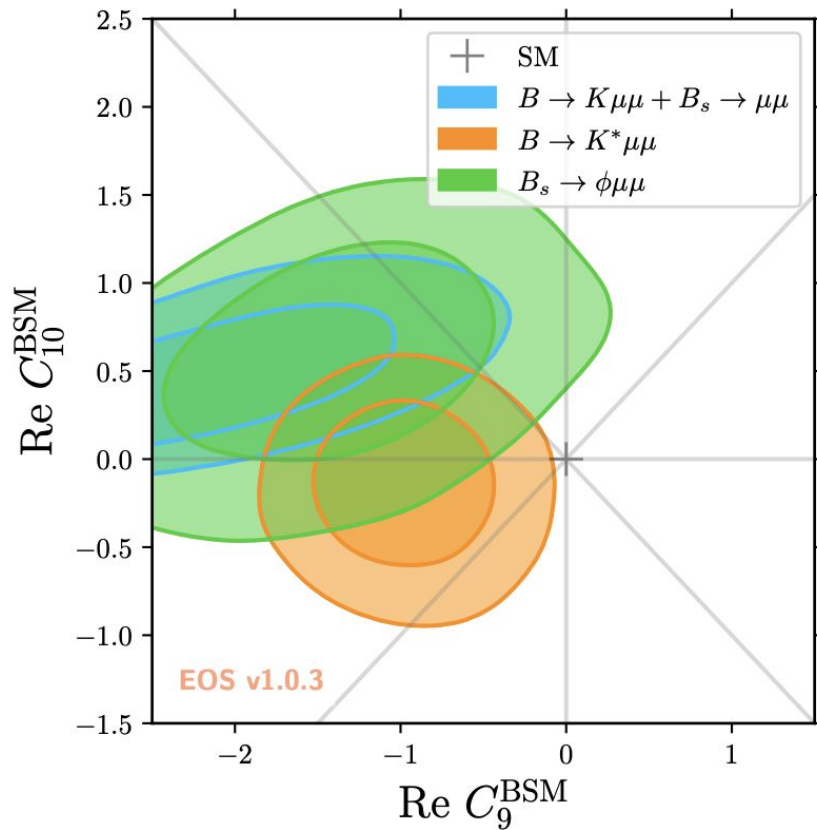


[Frezotti et al. [2508.03655](#)]

Charm loops

[Gubernari et al. [2206.03797](#)]

- Instead of calculating directly, Gubernari et al. use dispersive bounds to constrain the contribution from non-local form factors
- Find local form factors responsible for majority of theory uncertainty
- Tension between SM and experiment remains



New Physics in $B \rightarrow Xc l \nu$

[Carvunis et al. [2507.22123](#)]

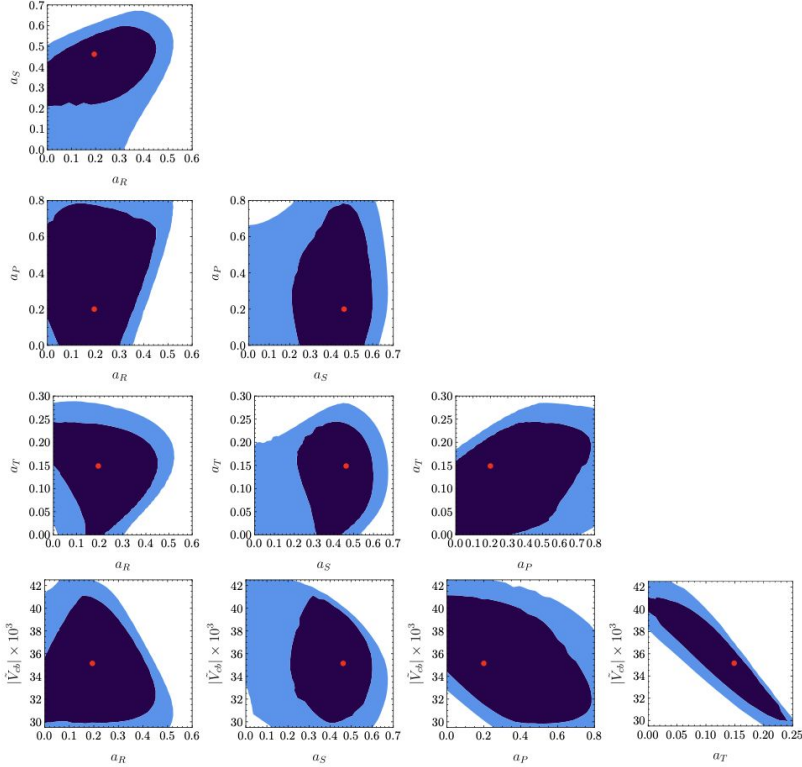


Figure 3. 2D contour plots for pairs of fit parameters in the global fit with $\Delta\chi^2 \leq 2.30$ (dark blue) and $\Delta\chi^2 \leq 5.99$ (light blue) (plots for the phases are not shown for simplicity). The red dot indicates the best fit point.

m_b [GeV]	m_c [GeV]	μ_π^2 [GeV ²]	μ_G^2 [GeV ²]	ρ_D^3 [GeV ³]	ρ_{LS}^3 [GeV ³]	$\text{BR}_{cl\nu}$ [%]	$10^3 V_{cb} $
4.574	1.090	0.435	0.278	0.164	-0.090	10.62	41.64
0.012	0.010	0.040	0.048	0.018	0.089	0.15	0.47
1	0.390	-0.229	0.560	-0.022	-0.181	-0.064	-0.421
	1	0.015	-0.238	-0.028	0.084	0.034	0.071
		1	-0.097	0.535	0.266	0.144	0.346
			1	-0.261	0.004	0.001	-0.271
				1	-0.014	0.025	0.172
					1	-0.010	0.056
						1	0.694
							1

Table 3. Results for the global fit with SM only. Central values on the first line, uncertainties on the second line and correlation matrix below.

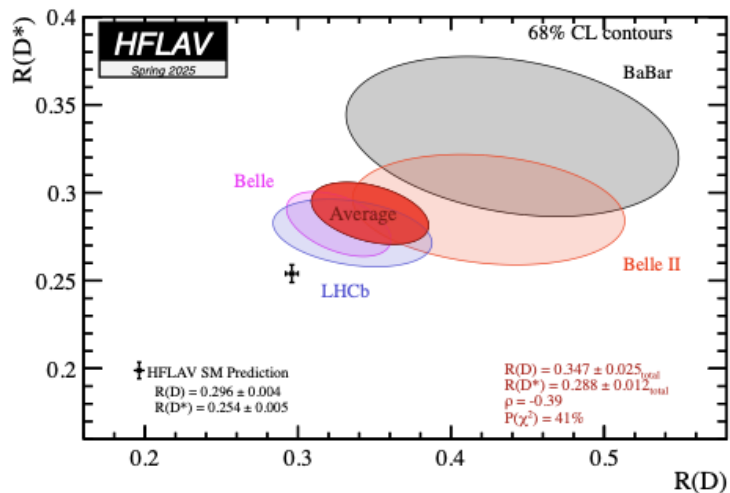
m_b [GeV]	m_c [GeV]	μ_π^2 [GeV ²]	μ_G^2 [GeV ²]	ρ_D^3 [GeV ³]	ρ_{LS}^3 [GeV ³]	$\text{BR}_{cl\nu}$ [%]
4.569	1.092	0.454	0.342	0.195	-0.126	10.72
+0.017	+0.011	+0.049	+0.065	+0.027	+0.096	+0.16
-0.016	-0.011	-0.050	-0.067	-0.029	-0.104	-0.16
a_R	a_S	a_P	a_T	$\cos(\delta_R)$	$\cos(\delta_{ST})$	$\cos(\delta_{PT})$
0.19	0.46	0.20	0.15	0.99	-0.94	0.91
+0.21	+0.10	+0.43	+0.07	+0.01	+0.92	+0.09
-0.16	-0.14	-0.20	-0.11	-0.81	-0.06	-1.91

Table 4. Results for the global fit with full set of NP operators. Central values on the first line, 68.3% confidence level intervals on the second line. The value of $|\tilde{V}_{cb}|$ extracted from the fit parameters is shown in (3.12).

Optimized observables

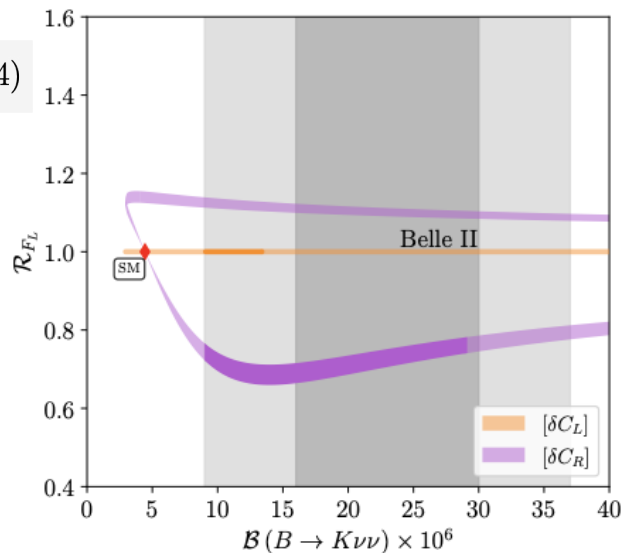
$$R_{H_c} = \frac{\mathcal{B}(B \rightarrow H_c \tau \nu)}{\mathcal{B}(B \rightarrow H_c \ell \nu)}$$

$$F_L(B \rightarrow K^* \nu \nu) = \frac{d\Gamma_L}{dq^2} \bigg/ \frac{d\Gamma}{dq^2}$$



$$F_L(B \rightarrow K^* \nu \bar{\nu})^{\text{SM}} = 0.49(4)$$

$$\mathcal{R}_{F_L} \equiv \frac{F_L}{F_L^{\text{SM}}}$$



- Can we devise other observables less sensitive to hadronic uncertainties (while remaining sensitive to New Physics effects)?

Final comments/questions

- Opportunity for progress for parametric (V_{cb}) and form-factor uncertainties.
- Many exploratory directions on the lattice (e.g., $B \rightarrow K^*$ form factors, inclusive decays, charm loops, ...).
- The public $B \rightarrow K_{VV}$ likelihood is very useful for theory interpretations (both in the SM and beyond). Can this be extended to other channels (e.g., $b \rightarrow c$ tau nu modes)?
- Interplay between theory and experiment is fundamental to verify precision theory calculations.
- Anything else?