

# THE UNIVERSITY OF **MELBOURNE**

# Particle identification at *Belle II* with the Electromagnetic Calorimeter (ECL)

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- 1. Overview of particle identification at *Belle II*.
- 2. Review of particle interactions with matter  $\rightarrow$  cf. Christian Wessel's lecture.
- 3. The Electromagnetic Calorimeter (ECL).

3.1 Pulse Shape Discrimination.

- 4. Overview of ECL clustering algorithm.
- 5. Neutral particle identification with the ECL.
  - 5.1 Photons

5.2 Ku

- 6. Charged particle identification with the ECL.
  - 6.1 Standard proxy: *E/p*
  - 6.2 Novel improvements at low momentum  $\rightarrow$  machine learning

#### Outline

1. Particle identification at Belle II

### 1. Particle identification at Belle II

• Particle Identification (PID): identify "long lived" particles passing through the detector by means of their interaction with matter.

- Often one of the most crucial factors determining sensitivity/precision of a physics measurement.
- PID algorithm works by encoding measurements from different sub-detectors into a likelihood ratio  $\rightarrow$  cf. Umberto Tamponi's lecture. This lecture focuses on the PID reach of the Belle II Electromagnetic Calorimeter (ECL).



(In Belle II) "standard charged": { $e^{\pm}$ ,  $\mu^{\pm}$ ,  $\pi^{\pm}$ ,  $K^{\pm}$ ,  $p^{\pm}$ ,  $d^{\pm}$ }, "standard neutral": { $\gamma$ ,  $K^{0}_{L}$ }

### The Belle II detector

"Barrel"

B = 1.5

#### "Backward"

EM Calorimeter:

CsI(TI), waveform sampling (barrel) Pure CsI + waveform sampling (end-caps)

electron (7GeV)

Beryllium beam pipe 2cm diameter

Vertex Detector

2 layers DEPFET + 4 layers DSSD

Central Drift Chamber He(50%):C<sub>2</sub>H<sub>6</sub>(50%), Small cells, long lever arm, fast electronics KL and muon detector: Resistive Plate Counter (barrel) Scintillator + WLSF + MPPC (end-caps)

#### "Forward"

Particle Identification Time-of-Propagation counter (barrel) Prox. focusing Aerogel RICH (fwd)



# 2. Particle interactions with matter

#### 2. (Very brief) review of particle interactions with matter

Muons:

• In  $p \sim [10^{-1}-10^3]$  GeV/c, mostly lose energy by *ionisation*  $\rightarrow$ 

minimum ionising (M.I.P) up to several GeV,  $dE/dx \sim O(2 \cdot \rho \text{ MeV/cm})$ .

> As a result, muons unlikely to be stopped by detector material, will escape through. ightarrow m<sub>µ</sub> = 105 MeV/c<sup>2</sup>  $\approx$  m<sub>π</sub> = 139 MeV/c<sup>2</sup> implies measuring dE/dx has low power to discriminate the two species.



Bethe-Bloch (approx.) formula of average *E* loss per unit distance:

$$\frac{dE}{dx} \approx \rho (2 \text{MeVcm}^2/\text{g}) \frac{Z^2}{\beta^2} \qquad \left( p = m\beta\gamma c, \quad \gamma = \right)$$

CDC-dE/dx distribution and predictions







# 2. (Very brief) review of particle interactions with matter

#### Electrons:

• almost always  $\beta \approx 1$ ,  $dE/dx \sim O(2 \text{ MeV/cm})$  in light absorbers  $\rightarrow$  not stopped by VXD and CDC

(unless tracks curl in B field:  $p_T < 200 \text{ MeV/c}$ ).

• For  $E \sim [100-1000]$  MeV  $\rightarrow$  loss dominated by *bremsstrahlung*:

$$\frac{dE}{dx} = -\frac{E}{X_0} \qquad (X_0: r)$$

radiation length  $\rightarrow$  property of material)

(NB: for muons, brems loss suppressed by a factor  $(m_e/m_\mu)^2$  up to  $E \approx O(\text{TeV})$ .

Photons in same range lose E by analogous mechanism: pair production  $(\gamma \rightarrow ee)$ 

 $\blacktriangleright$  EM shower progresses until critical energy reached:  $E_c$ : brems = ionisation

> Max *longitudinal* shower depth  $t_{max}$  depends only on  $\ln(E_0)$ .

▶ 95% of *lateral* width  $< 2R_M$ , independently of  $E_0$ .

Moliere radius





# 2. (Very brief) review of particle interactions with matter

#### Hadrons $(\pi, K, p)$ :

• As other charged massive particles, in  $E \sim [10^{-1}-10^3]$  GeV, show M.I.P behaviour.

• Strong interactions with material atoms also lead to inelastic scattering loss.

► For  $Z \ge 6$ ,  $\lambda_I \gg X_0 \rightarrow$  hadrons are likely to punch-through the EM calorimeter.

► Inelastic interactions lead to fuzzier shapes of hadronic showers vs. EM showers.

► Modelling of simulated hadronic interactions w/ detector material not a trivial task:

cross-section energy dependence, different particle type responses...



 $\lambda_I = \frac{1}{N\sigma} \approx \frac{A^{1/3}}{\rho} 35 \text{ g/cm}^2$ 

 $(\lambda_{l}: interaction \ length \rightarrow hadronic \ mean \ free \ path)$ 





3. The Electromagnetic Calorimeter

### 3. The *Belle II* Electromagnetic Calorimeter (ECL)

Main tasks of the ECL:

Made up of 8736 laterally segmented CsI(TI) crystals.

Designed to longitudinally contain ~any EM cascade.





### 3. The *Belle II* Electromagnetic Calorimeter (ECL)

We can clearly observe resonances decaying to photons in the early *Belle II* data, with good mass resolution.

 $\eta \rightarrow \gamma \gamma$ 





 $\pi^0 \rightarrow \gamma \gamma$ 



Upgraded ECL readout electronics in *Belle II* (waveform sampling) allows offline analysis of the shape of the CsI(TI) crystal signal waveform  $\rightarrow$  pulse shape discrimination.

 $\blacktriangleright$  Exploit different hadronic ( $\pi$ ,K,p) vs.E.M. ( $\gamma$ , e) scintillation response as a powerful handle for particle identification.



# 3.1 Pulse Shape Discrimination (PSD)

3. Overview of ECL clustering algorithm

### 4. Overview of ECL clustering (very simplistic)



# 4. Overview of ECL clustering

Simulation of single particles' energy deposition in a 15x15 ECL crystal array (pre-clustering):

#### γ

- Radially symmetric shape
- Usually contained in
  5x5 cells

#### **e**

- Similar shape of  $\gamma$
- Less symmetric

(*B* field bend, brems  $\gamma$ 

emitted before the ECL)



#### π

Ionisation loss contained

in 1-2 cells.

• Asymmetric lateral spread

#### due to hadronic interactions

#### μ

- Pure MIP behaviour.
- $\langle E_{cluster} \rangle \sim 200 \text{ MeV}$



# 5. Neutral PID with the ECL

#### 5. Neutral PID with the ECL

Photons can be mimicked by:

- Neutral hadrons
- Charged hadrons  $\rightarrow$  "secondary" clusters

w/o matching track due to hadronic splitoffs.

► Identification mostly relies on variables describing the lateral shower shape development  $\rightarrow$  E1/E9, Zernike moments.



Standard "candle" to test ID performance in data:  $ee \rightarrow \mu\mu\gamma$ 







# 6. Charged PID with the ECL

#### 6.1 Standard proxy for *charged* ECL PID: *E/p*

Ratio of  $E_{cluster}$  over  $p_{track}$ : "standard" variable, mostly designed for electron identification ( $\rightarrow$  peaks at 1!)

- *PDFs* defined in MC simulation from one-dimensional fits of *E/p* templates (single particle samples).
- Used to calculate likelihood of i-th particle type:

e



$$\mathcal{L}_i^{ECL} = \mathcal{L}(x|i) = \mathcal{L}\left((E/p)_{obs}|(E/p)_i^{MC}\right)$$

Events / (0.01035) 1701 (0.01035) 1001 (0.01035) Simulation Belle I  $cluster \le 2.24$ 0.56 < 0.60 $\square$  PDF for particle class:  $\pi$ 80 60 40 20 Pulls 0.2 0.4 0.6 0.8 20 E/p

π

# 6.1 Standard proxy for *charged* ECL PID: *E/p*





### 6.2 The low momentum PID challenge

- material, averaged over  $\phi$ 1.2 -**B2TIP** SVD 1.0 -ARICH 0.8 XX radiation length, - 9.0 • Separation power partially recovered by CDC. 0.2 - $0.2 \le p < 0.6$  [GeV/c], ECL Barrel  $0.2 \le p < 0.6$  [GeV/c], ECL Barrel 4.5 4.0
- At low momenta, E/p by itself becomes sub-optimal for PID purposes: • Electrons: larger bremsstrahlung losses before the ECL + stronger track bending. • Muons: if  $p_T < 600$  GeV/c, they fall outside of the KLM acceptance.





### 6.2 Novel improvements to (low p) ECL PID - BDT



### 6.2 Novel improvements to (low p) ECL PID - Convolutional NN

Feed raw energy information (instead of higher level variables like shower shapes) in a convolutional NN architecture. • Train algorithm on (pre-processed) calorimeter cells images for  $\mu$  and  $\pi$ .

- Idea is by-pass clustering algorithm shortcomings when dealing with "atypical" energy deposition patterns.







# Questions?

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