

Physics

... personal selection (with more focus on Belle II)



Aglia tau, the tau emperor, is a moth of the family Saturniidae



Ami Rostomyan
Jeniffer Summer School
26 July 2021



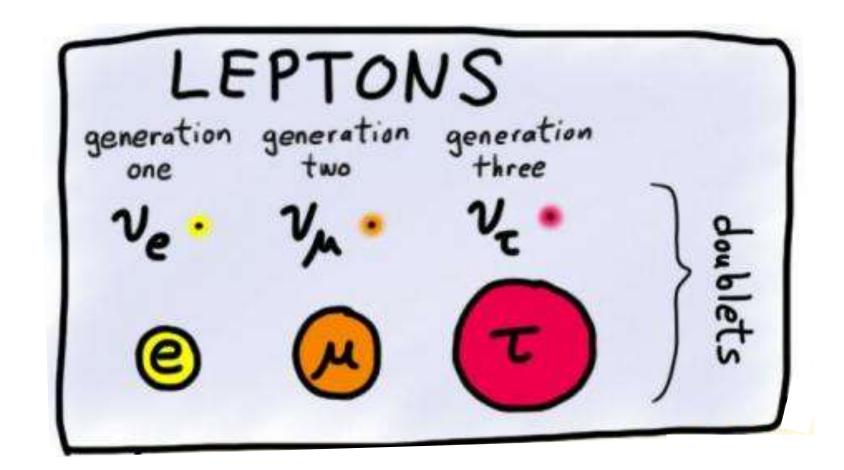
The third charged lepton - τοίτον

First observed in 1974 at SPEAR in SLAC as an anomalous decay in e⁺e⁻ collisions

We have discovered 64 events of the form

 $e^{+} + e^{-} \rightarrow e^{\pm} + \mu^{\mp} + at least two undetected particles$

for which we have no conventional explanation.



SLAC-PUB-1626 LBL-4228 August 1975 (T/E)

Proposal:

these events were the production and subsequent decay of a new particle pair:

$$e^{+} + e^{-} \rightarrow \tau^{+} + \tau^{-} \rightarrow e^{\pm} + \mu^{\mp} + 4\nu$$

The mass and spin of the τ was established at DESY with the Double Arm Spectrometer (DASP), and at SLAC-Stanford with the SPEAR Direct Electron Counter (DELCO).

EVIDENCE FOR ANOMALOUS LEPTON PRODUCTION IN e - e ANNIHILATION*

ABSTRACT

We have found events of the form $e^+ + e^- \rightarrow e^- + \mu^+ + \text{missing}$ energy, in which no other charged particles or photons are detected. Most of these events are detected at or above a center-of-mass energy of 4 GeV. The missing energy and missing momentum spectra require that at least two additional particles be produced in each event. We have no conventional explanation for these events.

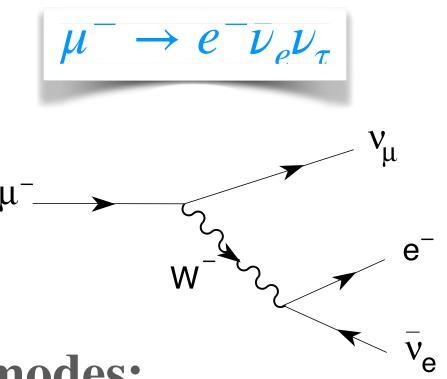
The t lepton

Very massive:

→ ~3500 times more massive than electron

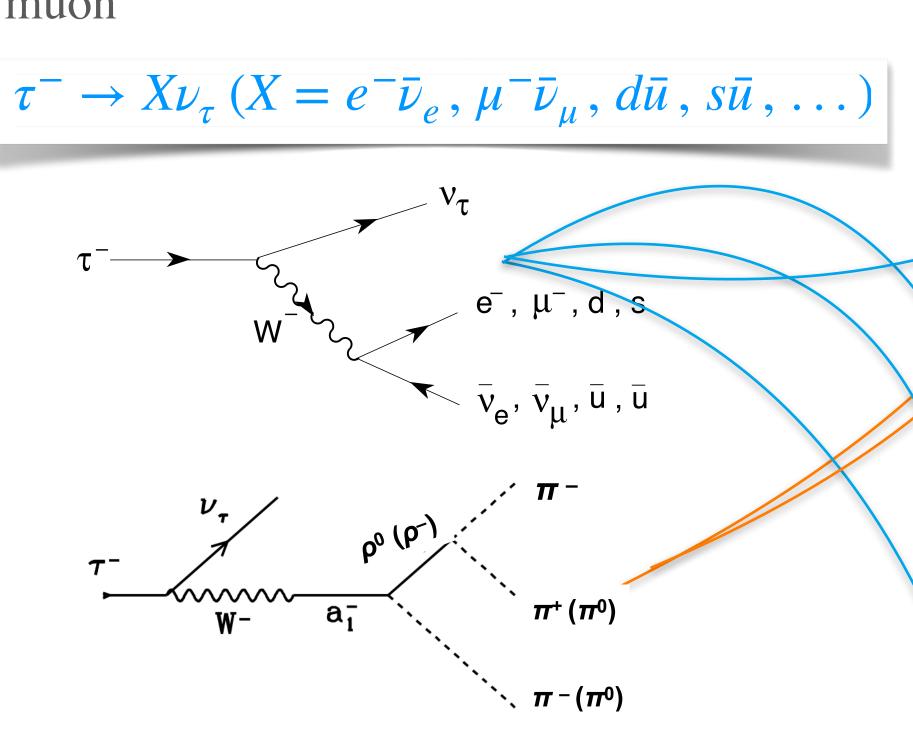
Short lifetime:

→ ~1.3*10-7 times smaller lifetime than muon



Decay modes:

- → ~ 250 decay modes in PDG
- $\rightarrow \tau \rightarrow e \text{ or } \mu \text{ (~35\%)}$
- heavy enough to decay to hadrons



 $K^{-3}\pi^{0}\nu_{\tau}$

K ην_τ

 $K^{-2}\pi^{0}\nu_{\tau}$

 $K^-\pi^+\pi^-\nu_{\tau}$

 $K^-K^+\pi^-\pi^0\nu_{\tau}$

 $K_S K_I \pi^- \nu_{\tau}$

 $K^-K^0\pi^0\nu_{\tau}$

 $2h^{-}h^{+}3\pi^{0}\nu_{\tau}$

 $3h^{-2}h^{+}\pi^{0}\nu_{\tau}$

 $2h^{-}h^{+}2\pi^{0}\nu_{\tau}$

 $h^-4\pi^0\nu_{\tau}$

 $3h^{-}2h^{+}v_{\tau}$

 $\pi^-\pi^0\eta\nu_{\tau}$

 $K^-K^0v_{\tau}$

 $\bar{K}^0\pi^-\pi^0\nu_{\tau}$

 $\bar{K}^0\pi^-\nu_{ au}$

 $0.84 \pm 0.04 \%$

 $K^-\pi^0\nu_{\tau}$

 K^-v_{τ}

0.700 ± 0.010 %

 $h^-\omega\pi^0\nu_{\tau}$

 $\pi^{-3}\pi^{0}\nu_{\tau}$

 $1.05 \pm 0.07 \%$

 $\pi^{+}\pi^{0}2\pi^{-}\nu_{\tau}$

 $2.70 \pm 0.08 \%$

 $h^-\omega v_{\tau}$

 $2.00 \pm 0.08 \%$

 $K^{-}\pi^{+}\pi^{-}\pi^{0}\nu_{\tau}$

100

- 80

×10

(%)

40

- 20

25 modes

10.12 %

 $\pi^{-}2\pi^{0}\nu_{\tau}$

 $9.30 \pm 0.11 \%$

 $\pi^+2\pi^-\nu_{\tau}$

 $8.99 \pm 0.06 \%$

 $\pi^-\pi^0\nu_{\tau}$

 $10.83 \pm 0.06 \%$

 $\mu^- \overline{\nu}_{\mu} \nu_{\tau}$

 $17.41 \pm 0.04 \%$

 $e^{-\overline{\mathsf{v}}}e^{\mathsf{v}}_{\tau}$

 $17.83 \pm 0.04 \%$

 25.52 ± 0.09

Laboratory to test the structure of the weak currents, the universality of their coupling to the gauge bosons and the low-energy aspects of strong interactions.



B factories

Not just B-Factories but also τ factories!

- Collision energy at Y(nS)
 - \rightarrow e.g. $E_{CMS} = 10.58 \text{ GeV}$
 - \rightarrow BR(Y(4S) \rightarrow BB) > 96%

→ Asymmetric beam energies

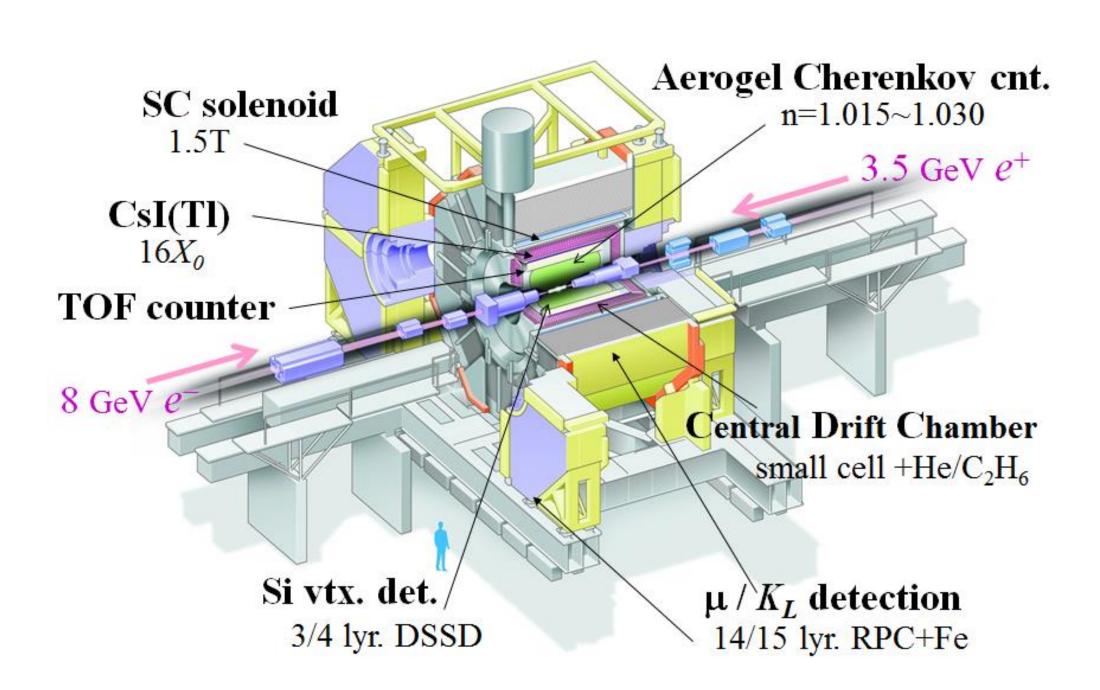
Belle: 8.0 GeV (e⁻) / 3.5 GeV (e⁺)

BaBar: 9.1 GeV (e⁻) / 3.0 GeV (e⁺)

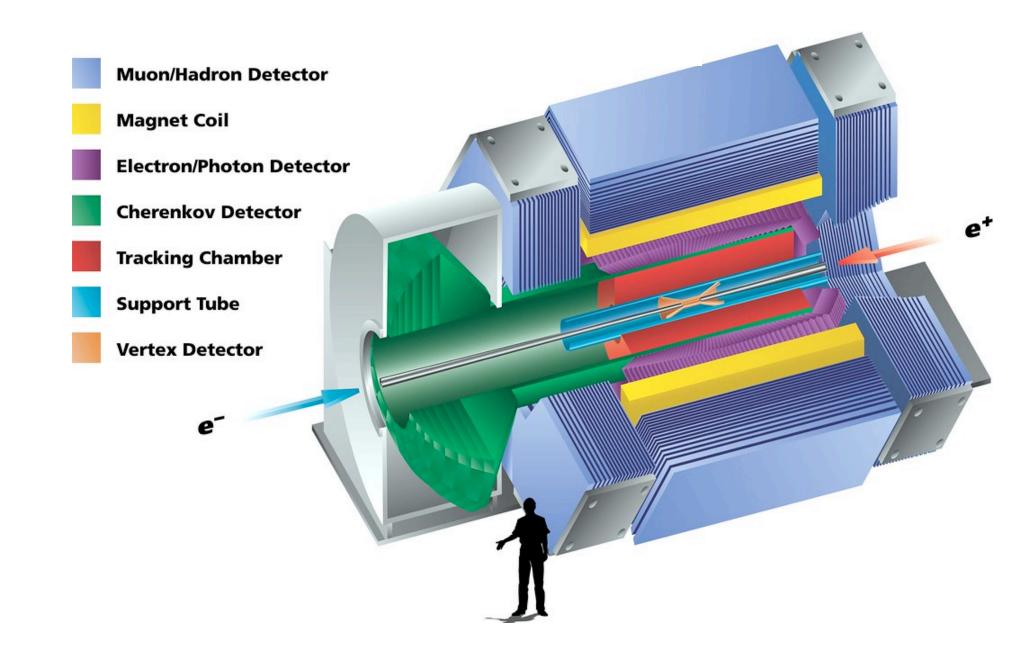
→ Boosted BB pairs

- → High luminosities
 - ~Belle: 710 fb-1 @ Y(4S)
 - ~ BaBar: 424 fb-1 @ Y(4S)

Belle@KEKB

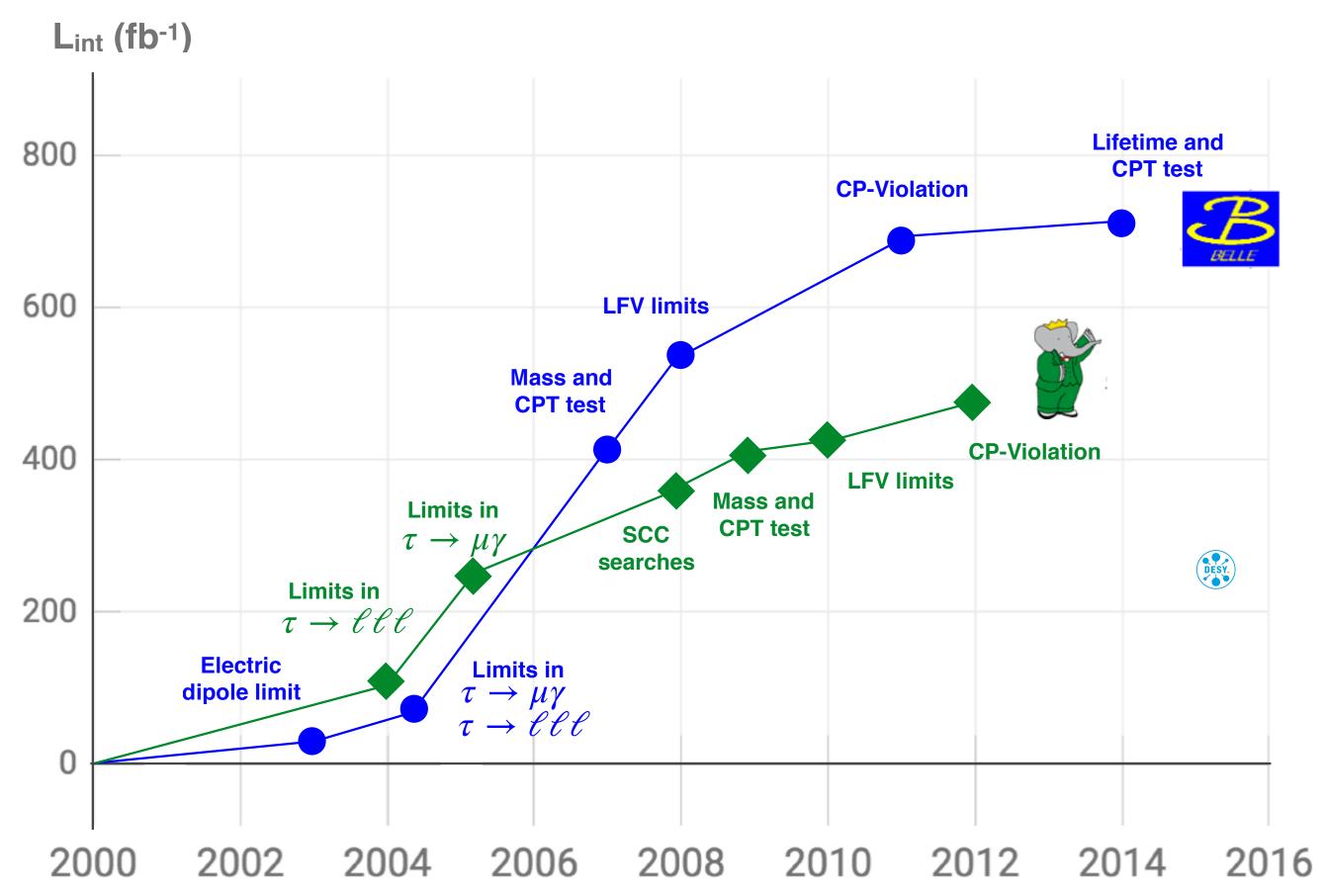


BaBar@PEP-II

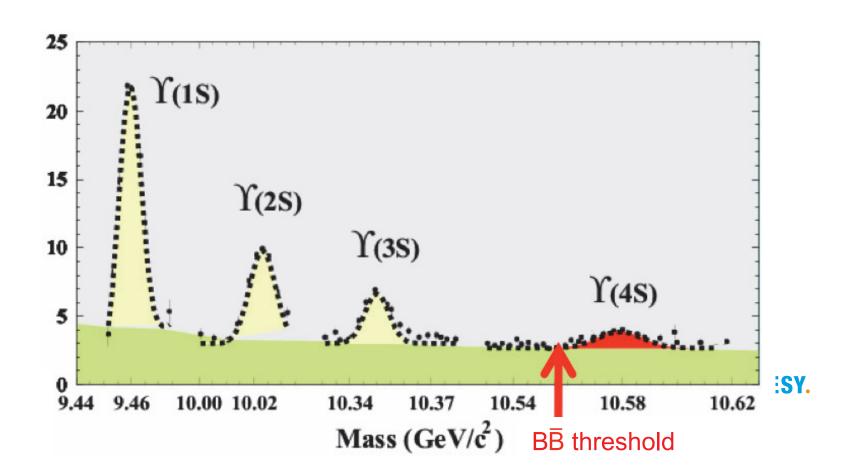


τ physics program @ B factories

Historically B-factories provided a variety of very interesting results in the last two decades.



B-factories: Belle@KEKB and BaBar@PEP-II



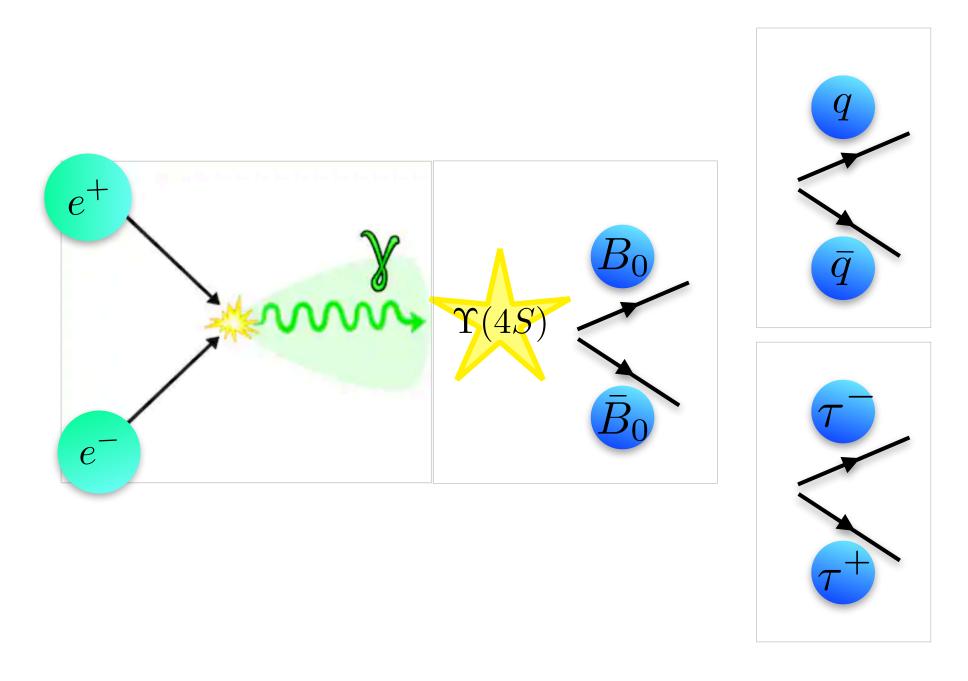
Wide range of observables to confront theory

- → Precision SM measurement
- → CP asymmetries
- → Rare decays, forbidden decays



B-Factories

Not just B-Factories but also τ factories!



$$\sigma(e^{+}e^{-} \to \Upsilon(4S)) = 1.05 \text{ [nb]}$$
 $\sigma(e^{+}e^{-} \to q\bar{q}) = 3.69 \text{ [nb]}$
 $\sigma(e^{+}e^{-} \to \tau^{+}\tau^{-}) = 0.919 \text{ [nb]}$

Clean environment

- → the kinematics of the initial state is precisely known
- the neutrino energy can be determined precisely

Hermetic detectors with

- → high track reconstruction efficiency
- good kinematic and vertex resolution
- excellent PID capability
- \rightarrow excellent γ and π^0 reconstruction

Wide physics program

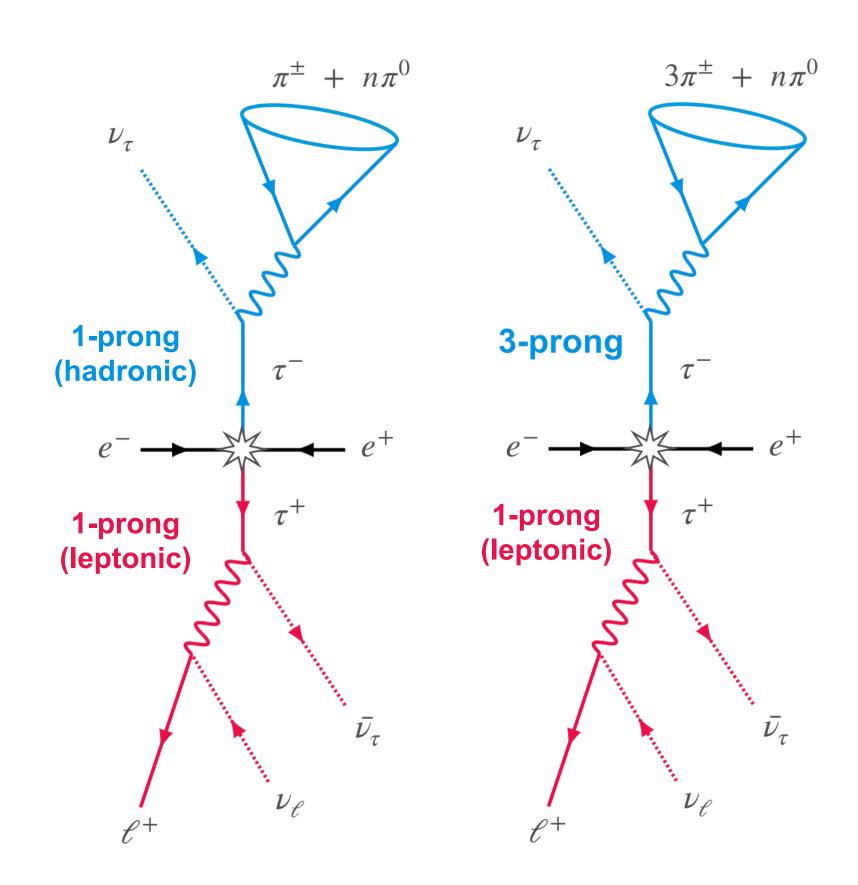
- precision measurements of time-dependent CPV and CKM parameters
- searches for lepton flavor/universality/number violations
- **→**

6

$e^+e^- \rightarrow \tau^+\tau^-$

e+e- annihilation data is ideal for missing energy channels

the neutrino energy can be determined precisely



* "prong" is the number of charged tracks in the decay

Reconstruct 1-prong*, 3-prong \tau decays in various channels

wide variety of signatures involving tracks (e, μ, π, K) and neutrals

Precise tests of SM & NP searches

- T mass and lifetime measurements
- → Lepton universality tests
- → V_{us} measurement
- → Electric Dipole Moment (CP/T violation)
- ightharpoonup CP violation $au
 ightharpoonup K_S \pi \nu$
- → LFV and LNV decays

•

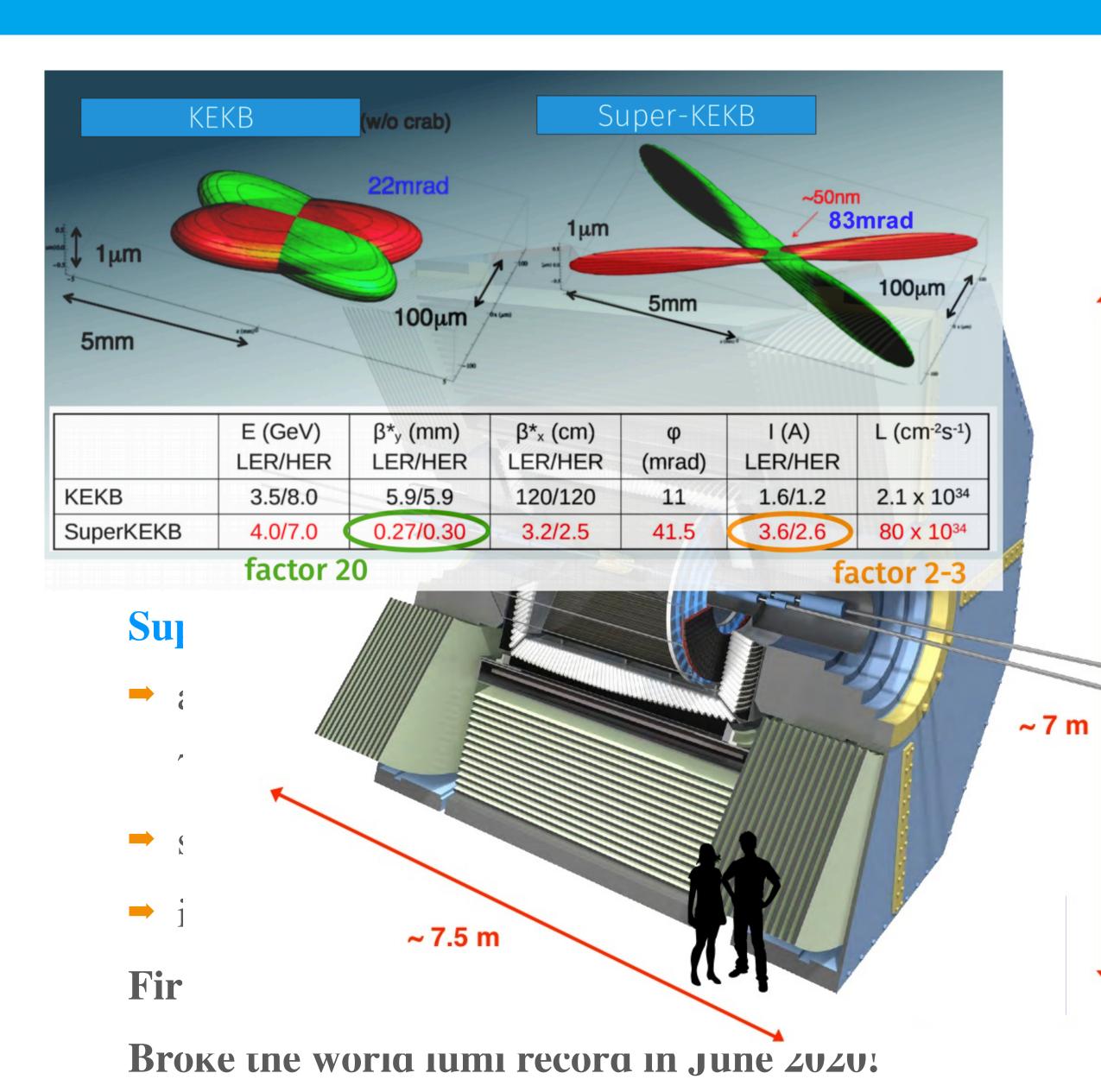
Belle II @ SuperKEKB

Next generation B-factory



Unprecedented design luminosity of ~6×10³⁵ cm⁻²s⁻¹

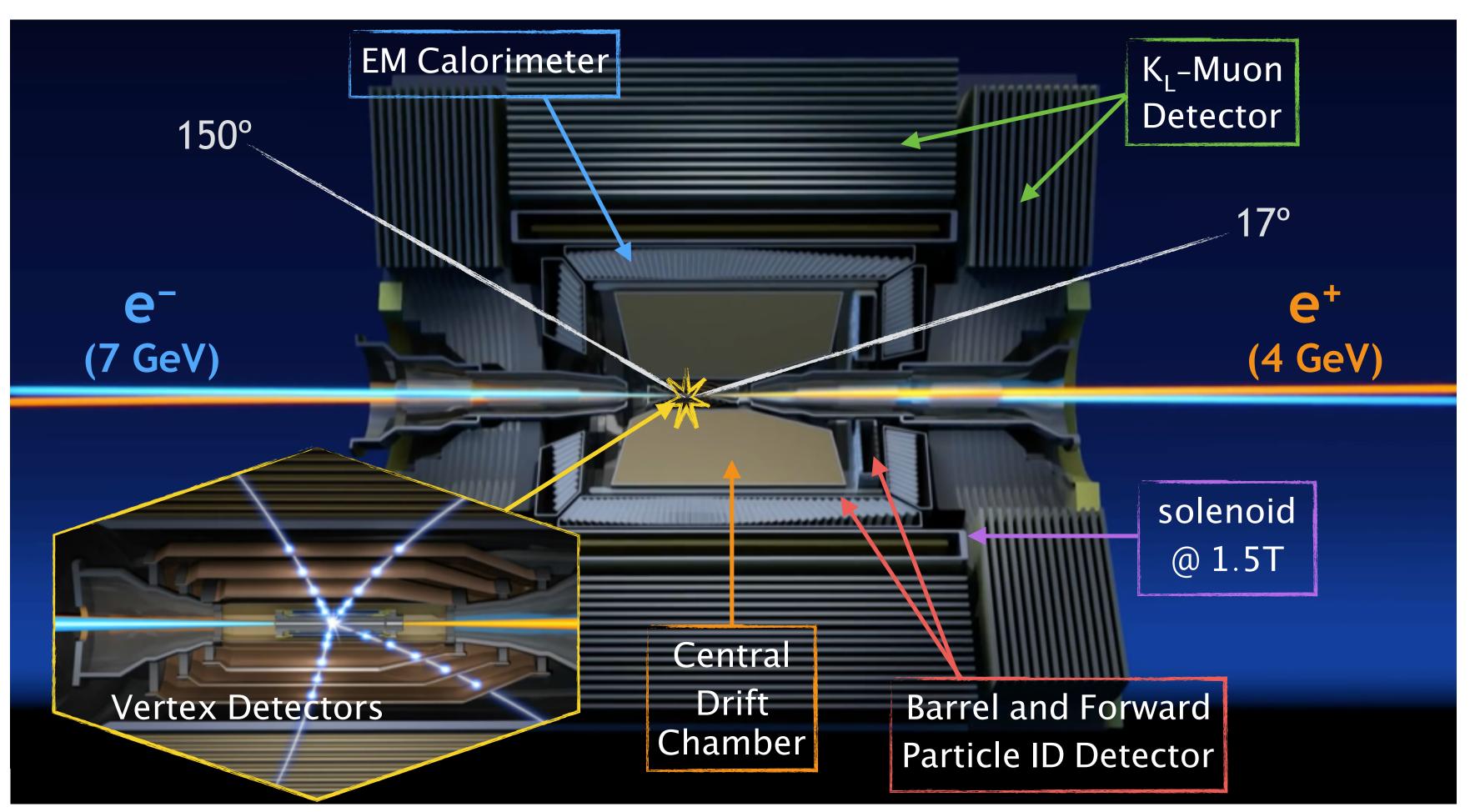




8

Belle II @ SuperKEKB

Belle II detector – upgraded Belle detector



- improved tracking efficiency
- improved particle identification
- → smarter software
- more precise algorithms
- rolled in April 2017

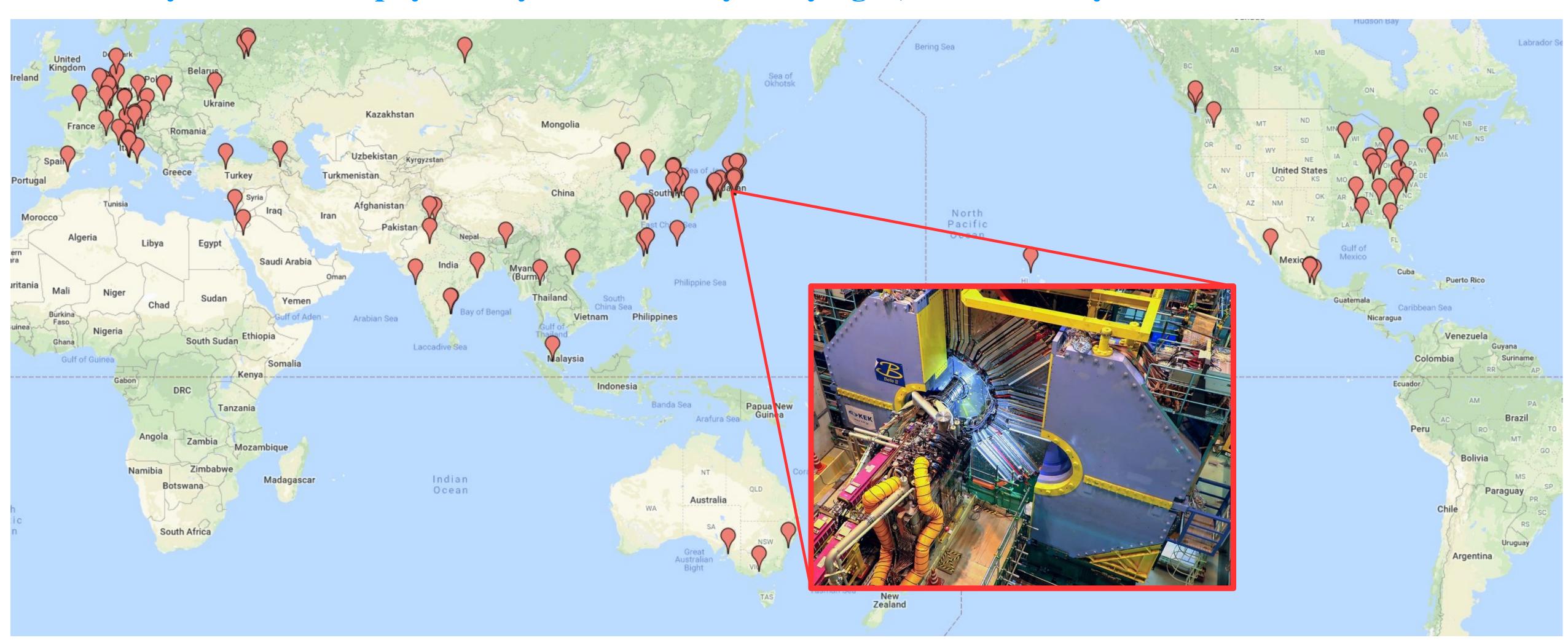
First recorded events in April 2018

Important for τ analysis: discriminate between e, μ , π , K; reconstruct neutrals!



The Belle II Collaboration

New facility to search for physics beyond the SM by studying B, D and τ decays



The world largest set of $e^+e^- \rightarrow \tau^+\tau^-$ for τ physics with high precision.

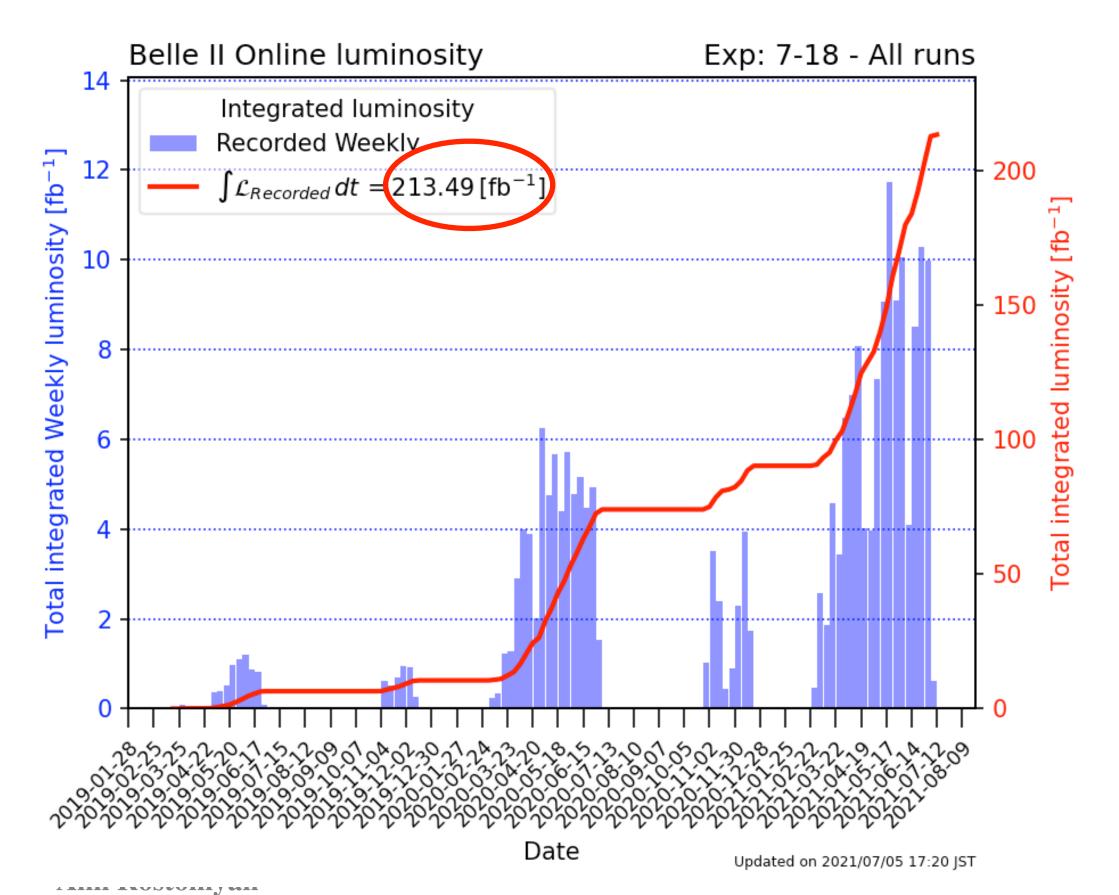


Luminosity status and plans

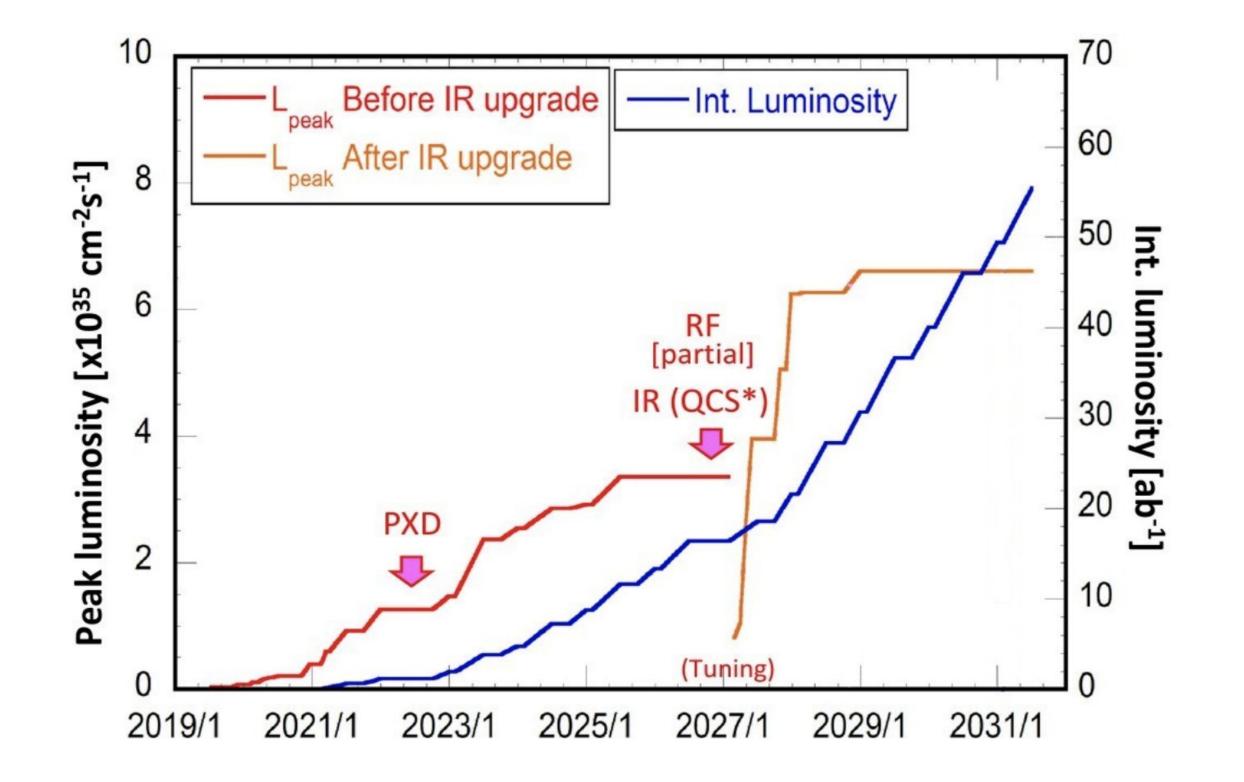
~ 200/fb of data already collected since 2019

Data taking is

- challenging in the pandemic
- however efficiency largely unaffected



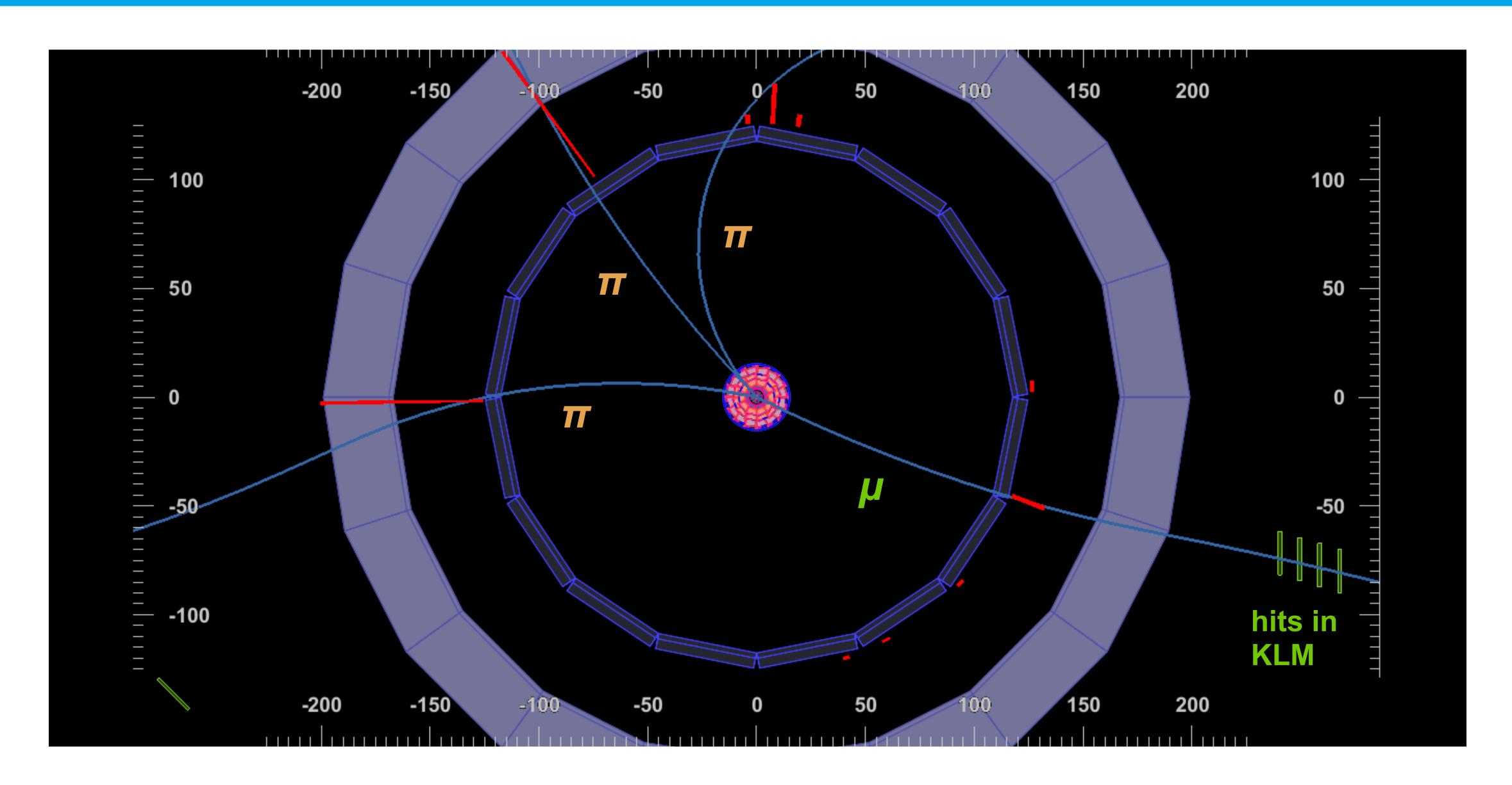
→ similar data sample as at 1st generation B-factories by 2022





Jeniffer Summer School 2021

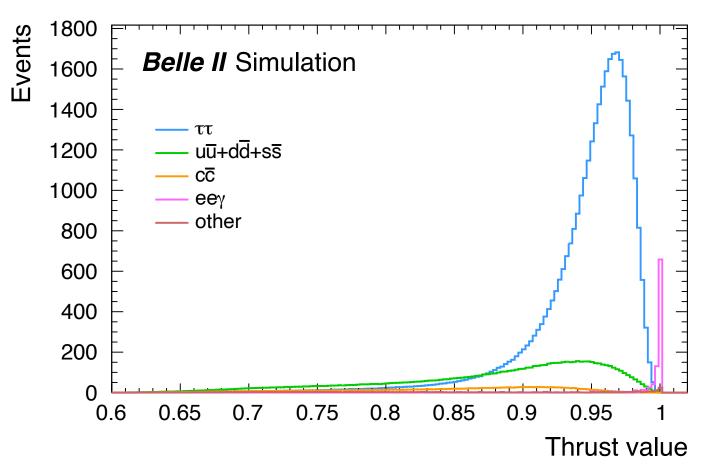
τ leptons @ Belle II



$e^+e^- \rightarrow \tau^+\tau^-$: An example of 3x1 topology

Event topology and kinematics crucial to observe τ leptons

 \rightarrow relatively mild deviation of the τ decay particles from the primary trajectory

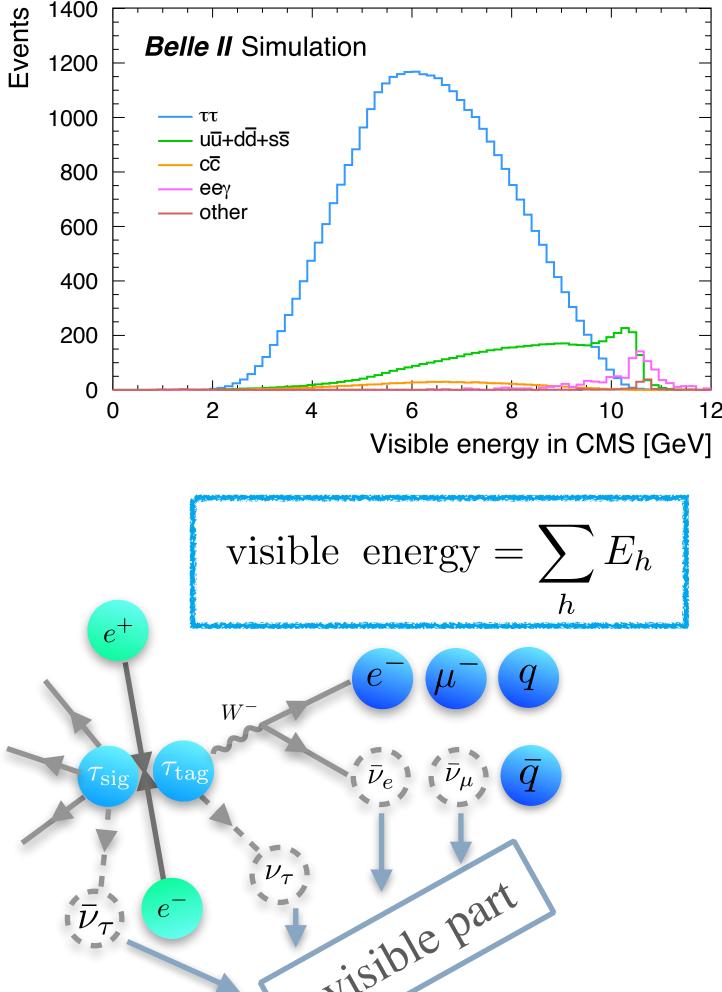


Thrust axis (T) is maximising the event shape variable

thrust value
$$\sum_{h}^{T} \frac{\vec{P}_h \cdot \hat{n}_{,h} \cdot \hat{T}}{|p_h|}$$
soft particles

thrust value $\sum_{h}^{T} \frac{\vec{P}_h \cdot \hat{n}_{,h} \cdot \hat{T}}{|p_h|}$
thrust axis hemisphere-a

 \rightarrow undetected neutrinos in τ events

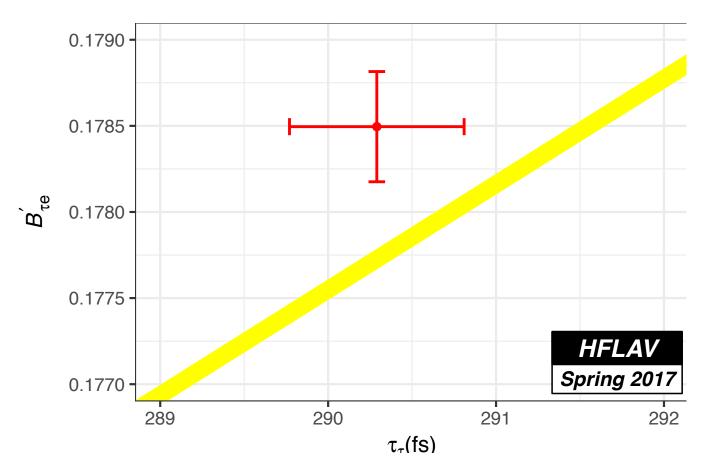


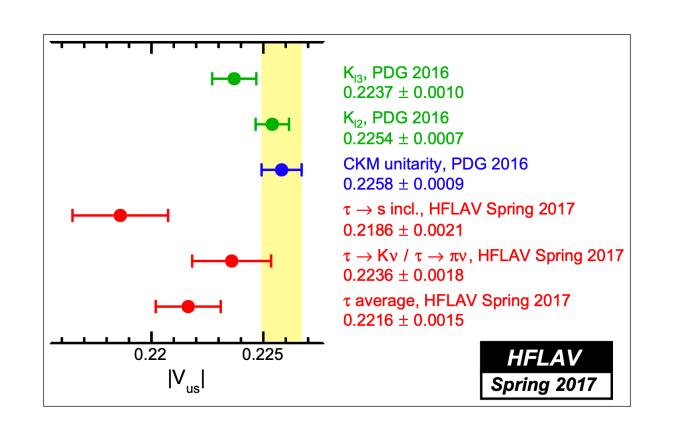
Outline of my lecture: SM & BSM

Does NP couple to 3rd generation strongly?

Precision measurements or indirect search of BSM

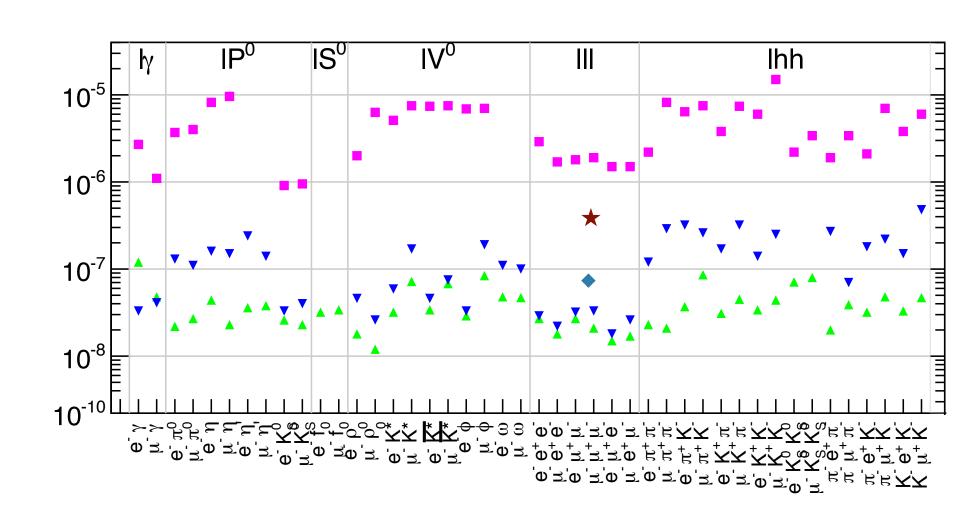
- → elementary particles *talk* to each other through loops and boxes
 - → learn about *new particles* by making precision measurements of the "talking" particle's properties
- 1-prong τ decays probe the SM predictions of charged current lepton universality and the unitarity relation of the first row of the CKM quark mixing matrix.
- → *significant deviations from SM* are unambiguous signatures of NP





Direct search of BSM

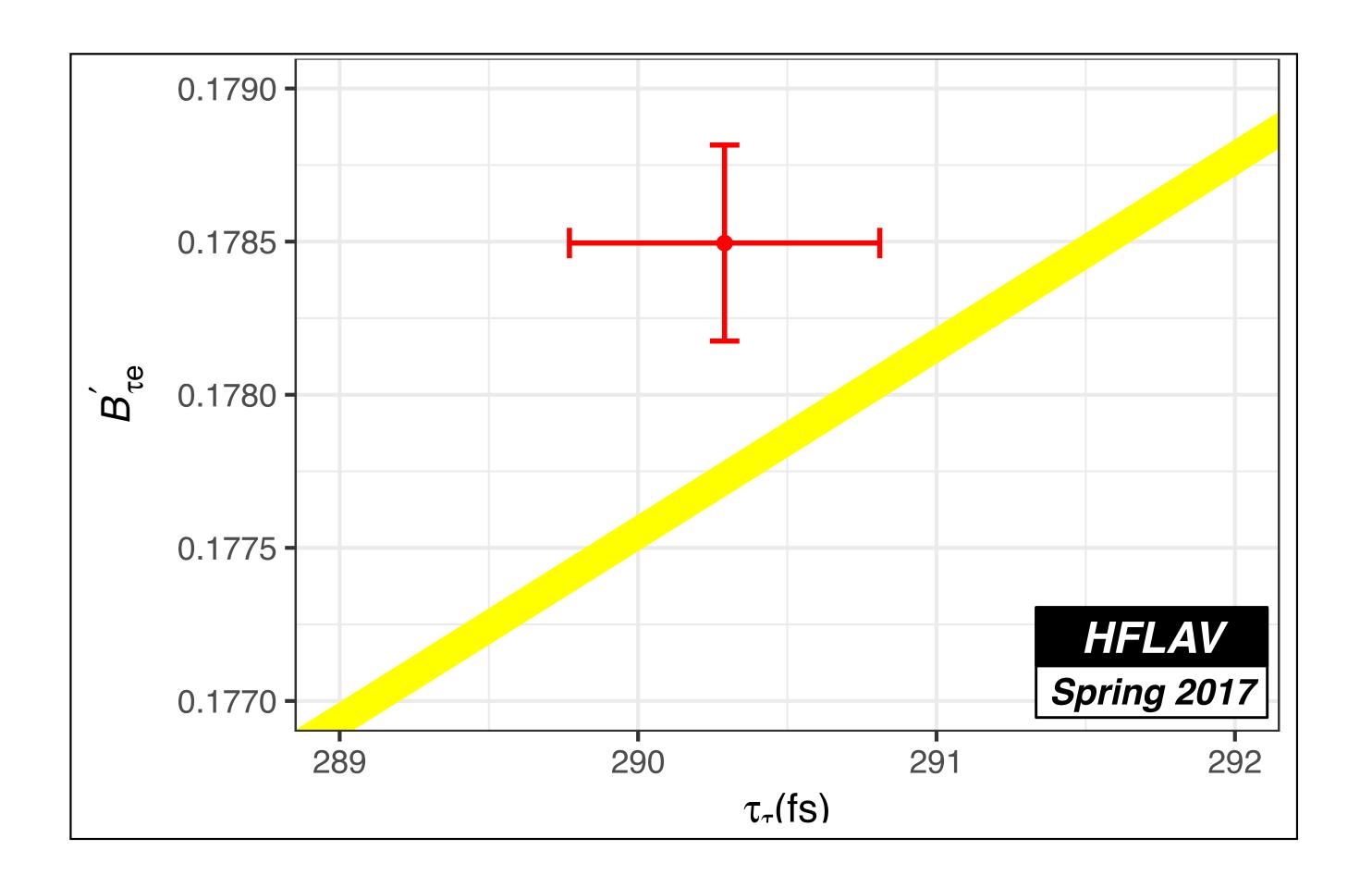
- search of forbidden decays
 - → LFV and LNV decays
- → any signal is unambiguous signature of NP





Jeniffer Summer School 2021

The mass, lifetime and leptonic decays of τ



The mass, lifetime and leptonic decays of τ

Lepton masses and lifetimes are fundamental parameters of SM!

- → A precise tau mass and lifetime measurements are crucial for lepton universality tests of SM
- Possibility to test CPT conservation measuring τ^- and τ^+ lifetimes and masses separately.

Lepton Masses (MeV):

 \rightarrow m_e= 0.5109989461 ± 0.0000000031

 $\delta m/m \sim 6*10^{-9}$

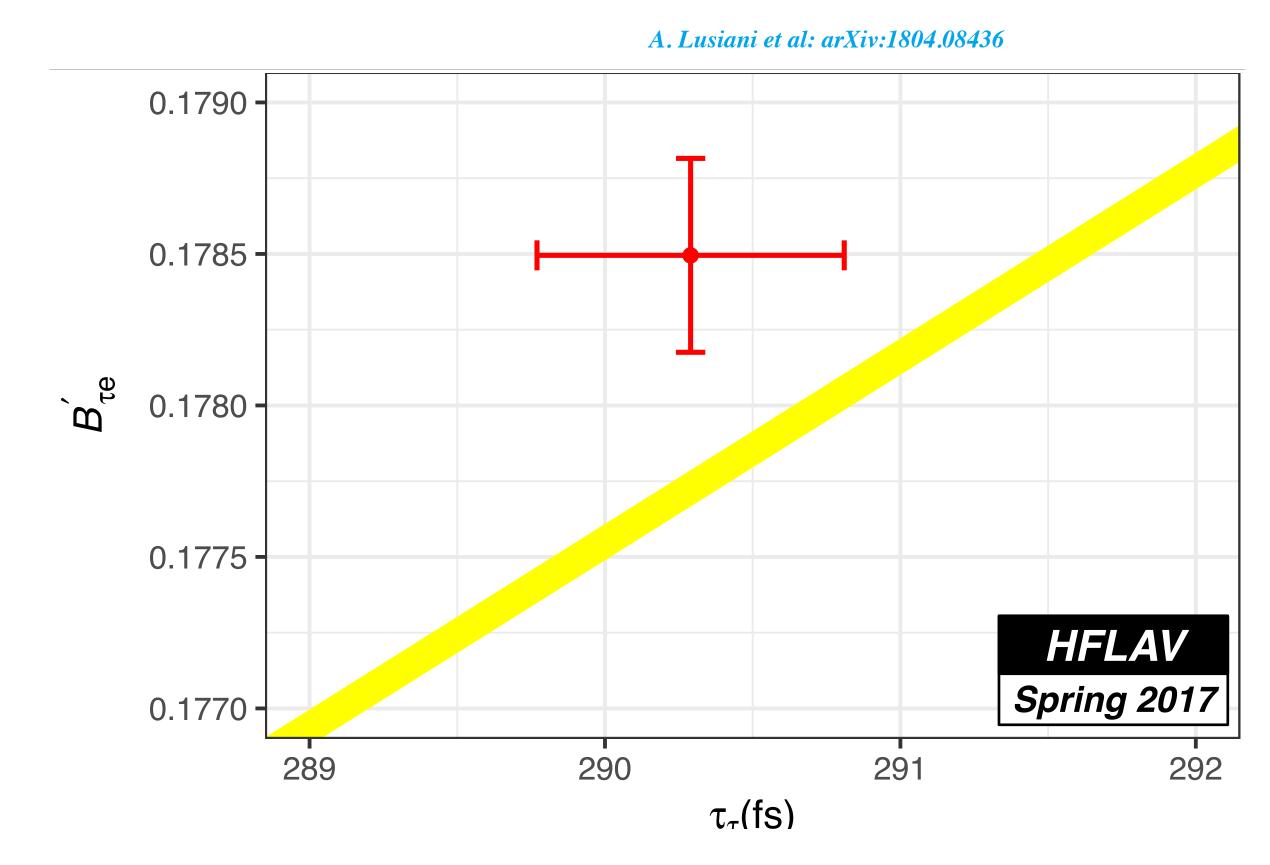
 \rightarrow m_{μ}= 105.6583745 ± 0.0000024

 $\delta m/m \sim 2*10^{-8}$

 \rightarrow m_{τ}= 1776.86 ± 0.12

 $\delta m/m \sim 7*10^{-5}$

Similar situation for lifetime

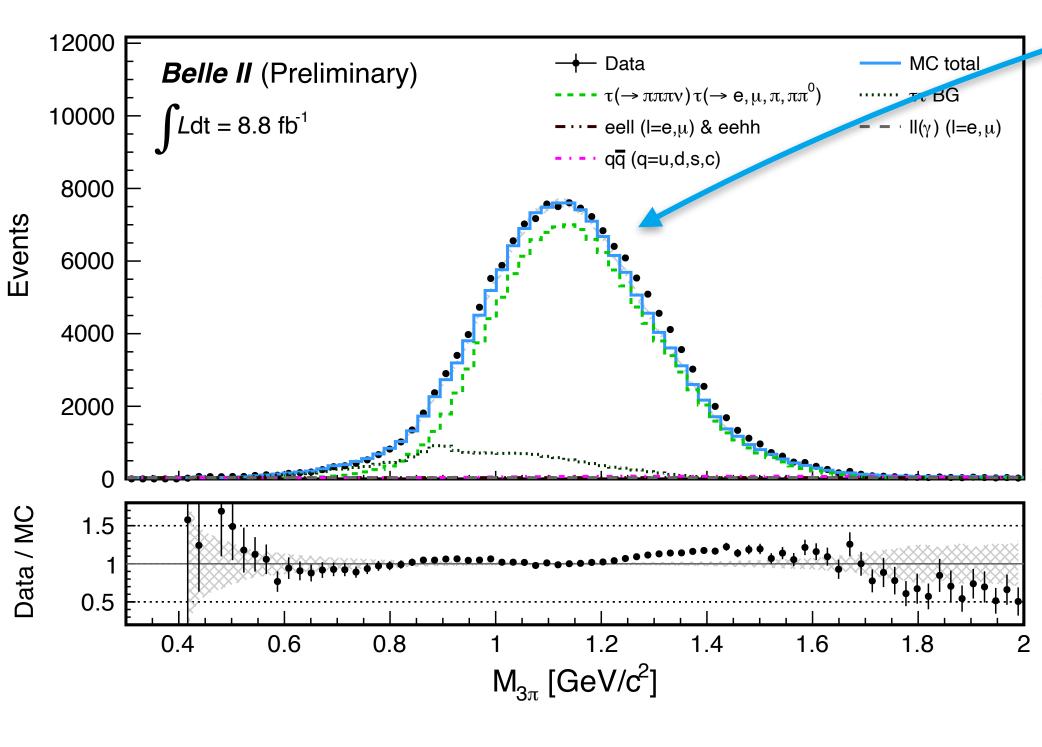


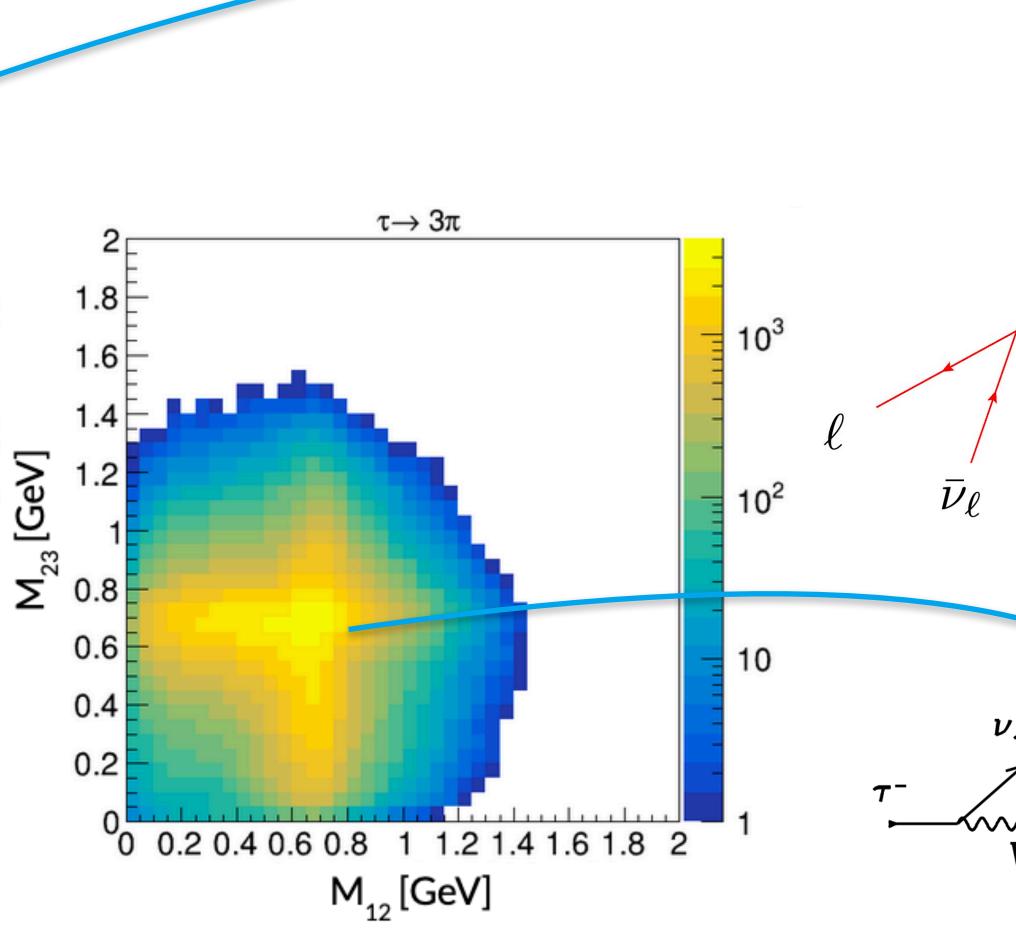
$$B_{ au l} \propto B_{\mu e} rac{ au_{ au}}{ au_{\mu}} rac{m_{ au}^5}{m_{\mu}^5}$$

t mass measurement

The τ mass cannot be measured directly

neutrinos in the final state







Jeniffer Summer School 2021

hadrons

 \hat{n}_{thrust}

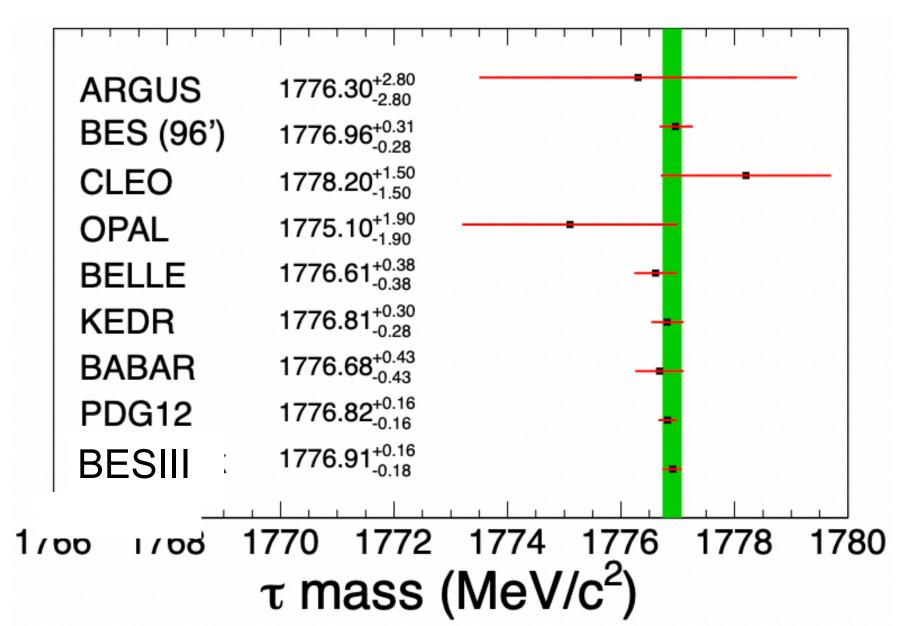
Why important to have a precise measurement of τ mass?

Threshold scan method

- → DELCO, BES, KEDR, BESIII
- exploits the threshold behaviour of the τ-pair production cross section in e+e- collisions

Pseudomass technique

- → ARGUS, OPAL, BELLE, BABAR
 - relies on the reconstruction of the invariant mass and energy of the hadronic system in the hadronic τ decay



→ SM prediction for the relationship between the τ lifetime, mass, B($\tau \rightarrow e\nu\nu$) and weak coupling constant

$$\frac{B(\tau \to e\nu\nu)}{\tau_{\tau}} = \frac{g_{\tau}^2 m_{\tau}^5}{192\pi^3}$$

violated before the first precise mass measurement by BES

BES - PRL V69 (1992) 3021 -

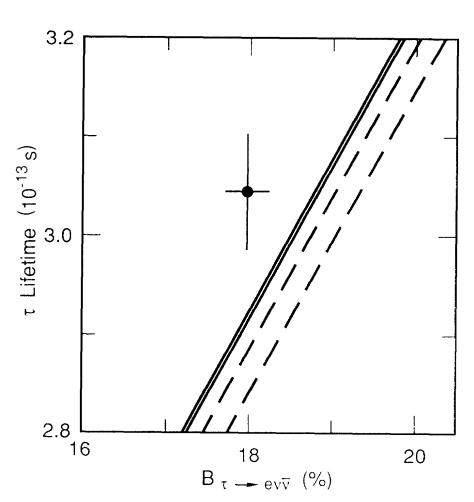


FIG. 3. The variation of τ_{τ} with B_{τ}^{e} , given by Eq. (1) under the assumption of lepton universality; the $\pm 1\sigma$ bands obtained using m_{τ} from this experiment (solid lines) and using the PDG value (dashed lines) are shown in comparison to the point corresponding to the PDG values (1σ error bars).

Threshold scan method

Determination of the beam energy and the beam energy spread precisely is of great importance

measure the electron and positron beam energies with one laser using Compton backscattering principle

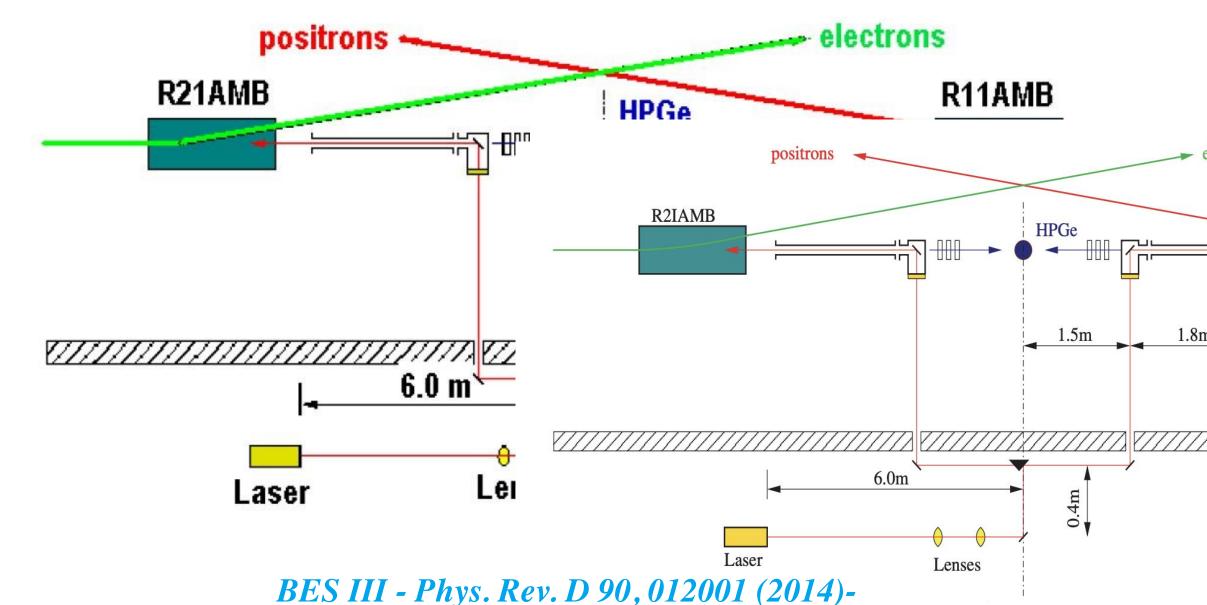
$$E_e = \frac{E_\gamma}{2} \left[1 + \sqrt{1 + \frac{m_e^2}{\epsilon_\gamma E_\gamma}} \right]$$
 beam energy of the laser beam max energy of scattered photon

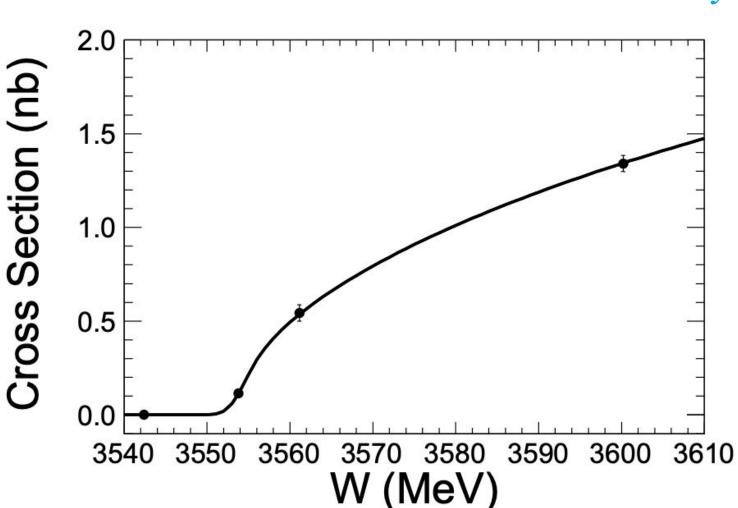
- \rightarrow E_e accuracy at the level of 10-5.
- data at 4 scan points near τ pair production threshold
 - the first τ scan point is below the mass of τ pair while the other three are above
- \rightarrow Likelihood function to estimate the τ mass

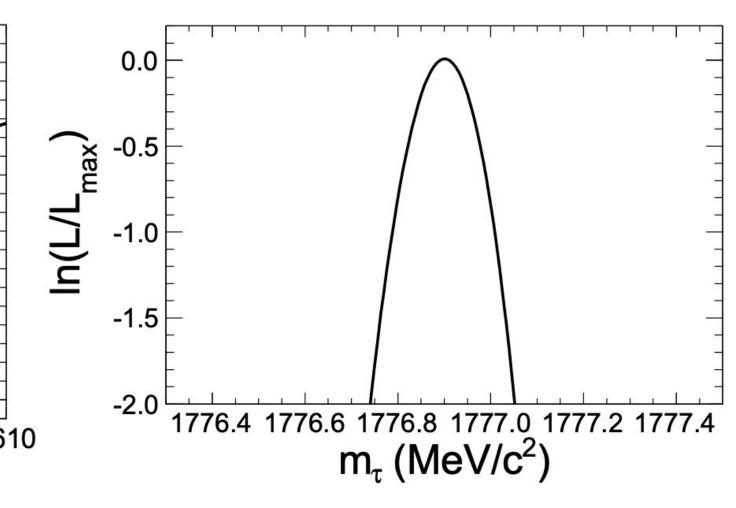
$$L(m_{\tau}, \mathcal{R}_{Data/MC}, \sigma_B) = \prod_{i=1}^{4} \frac{\mu_i^{N_i} e^{-\mu_i}}{N_i!}$$

$$\mu_i = [\mathcal{R}_{Data/MC} \times \epsilon_i \times \sigma(E_{CM}^i, m_\tau) + \sigma_B] \times \mathcal{L}_i.$$

Beam Energy Measurement System







Pseudomass technique

Use conservation of momentum and energy:

$$\mathcal{P}_{\tau}^{2} = (\mathcal{P}_{\nu} + \mathcal{P}_{3\pi})^{2}$$

$$\Rightarrow m_{\tau}^{2} = m_{\nu}^{2} + m_{3\pi}^{2} + 2(E_{\nu} E_{3\pi} - \vec{p_{\nu}} \cdot \vec{p_{3\pi}})$$

$$= m_{\nu}^{2} + m_{3\pi}^{2} + 2(E_{\nu} E_{3\pi} - p_{\nu} p_{3\pi} \cos \theta)$$
(1)

Use:

$$E_{\nu} = E_{\tau} - E_{3\pi}, and$$

$$p_{\nu} = \sqrt{E_{\nu}^2 - m_{\nu}^2} = E_{\nu} = E_{\tau} - E_{3\pi}$$
(2)

To get:

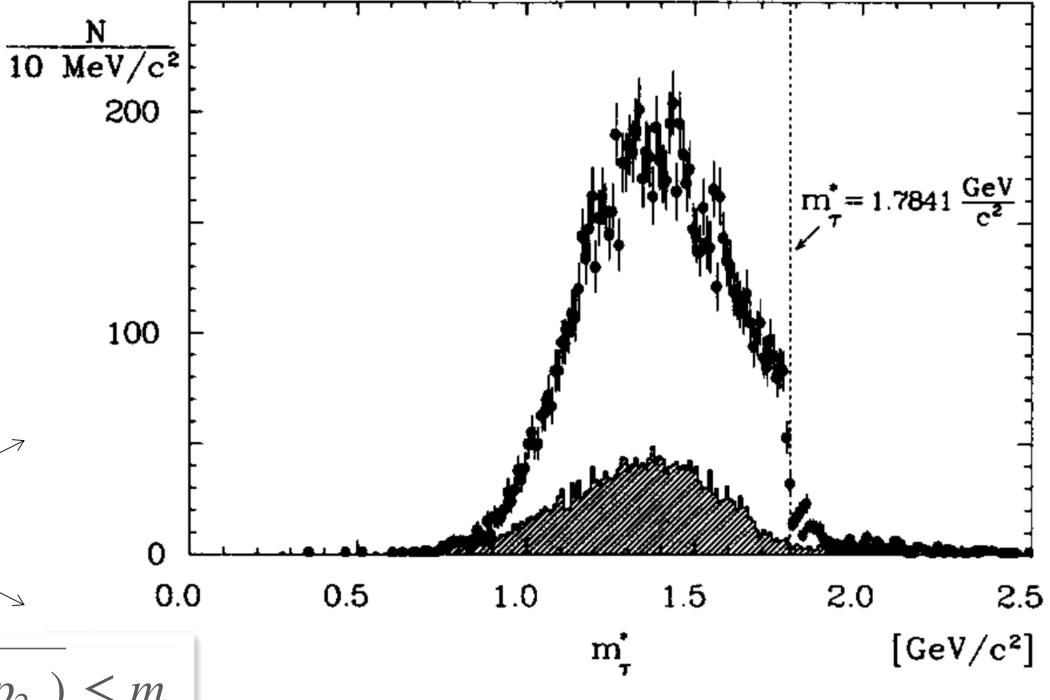
$$m_{\tau}^{2} = m_{3\pi}^{2} + 2\left((E_{\tau} - E_{3\pi})E_{3\pi} - (E_{\tau} - E_{3\pi})p_{3\pi}\cos\theta_{\nu,3\pi}\right)$$

$$= m_{3\pi}^{2} + 2(E_{\tau} - E_{3\pi})(E_{3\pi} - p_{3\pi}\cos\theta_{\nu,3\pi})$$
(3)

- \rightarrow in the centre of mass $E_{\tau} = E_{beam} = \sqrt{s/2}$
- \rightarrow the equation has a minimum when $\cos\theta_{\nu,3\pi}=1$

The distribution has a kinematic edge around the τ mass

- \rightarrow a sharp threshold behaviour in the region close to the nominal value of the τ mass
- irst used by ARGUS in 1992, later by Opal, BELLE and now by BELLEII

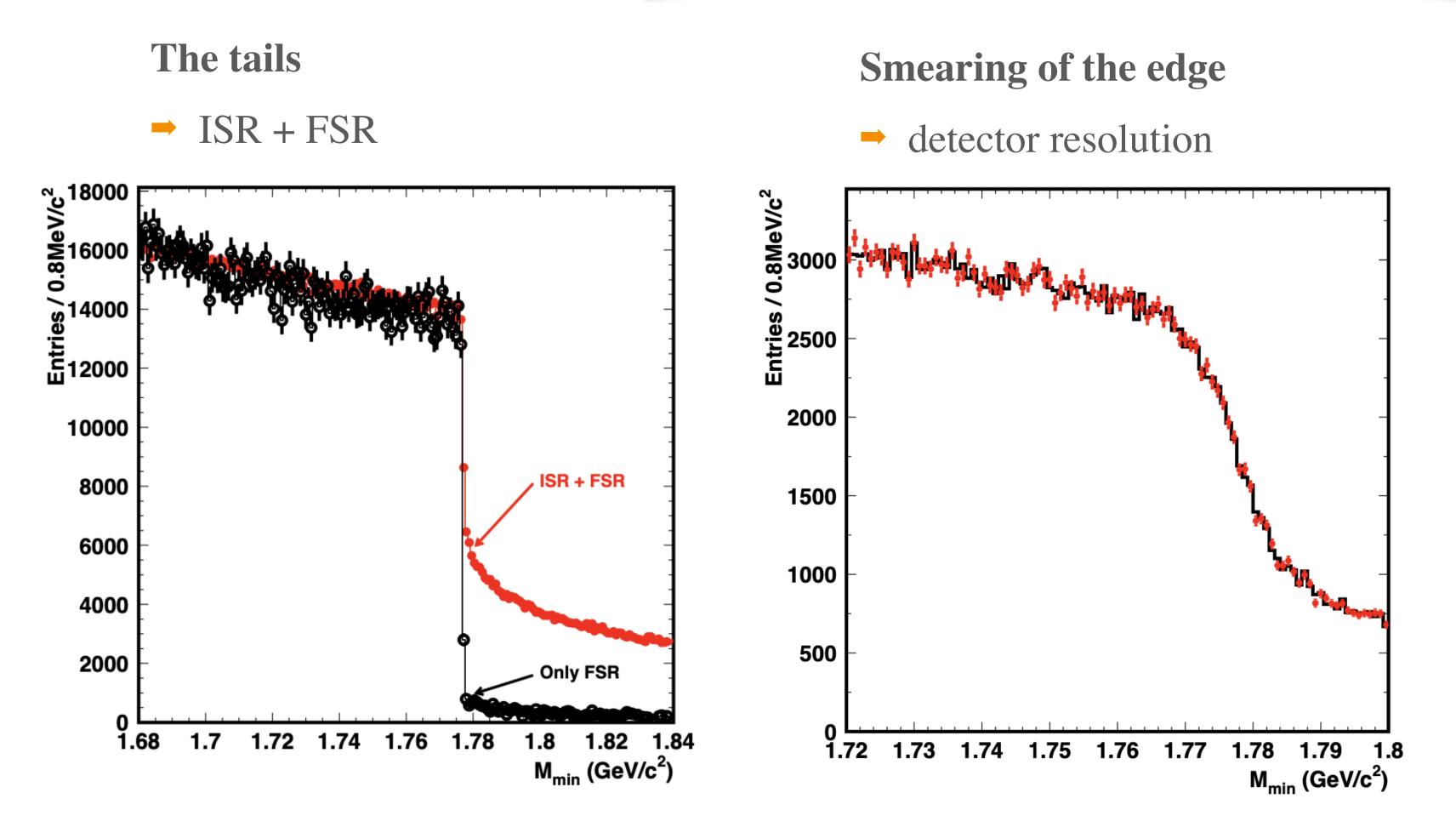


This is called pseudomass

$$m_{\min} = \sqrt{m_{3\pi}^2 + 2(E_{\text{beam}} - E_{3\pi})(E_{3\pi} - p_{3\pi})} \le m_{\tau}$$

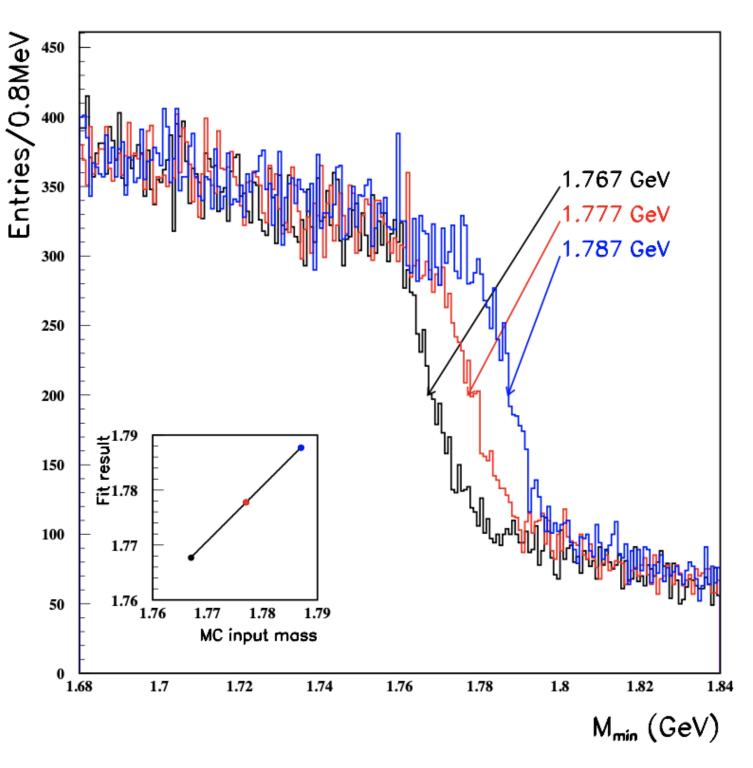
The threshold behaviour

$$m_{\min} = \sqrt{m_{3\pi}^2 + 2(E_{\text{beam}} - E_{3\pi})(E_{3\pi} - p_{3\pi})} \le m_{\tau}$$



Position of the edge

depends on mass



→ the edge position can be exploited to extract the mass

The t lepton mass

High signal purity

- the remaining continuum backgrounds are flat
 - don't impact the shape of the distribution

Mass extraction using ML fit

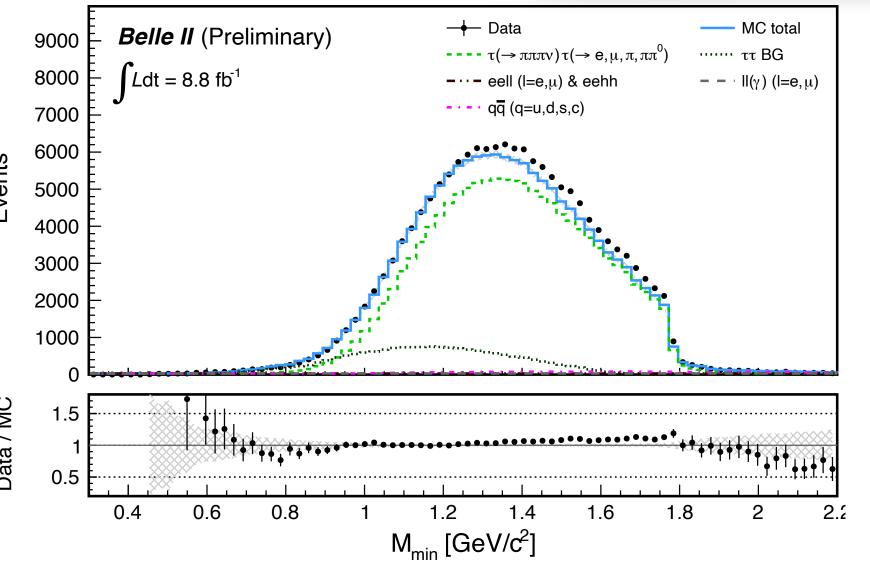
- \rightarrow P1 is an estimator for the τ mass
- with multi/additive components to describe the tails

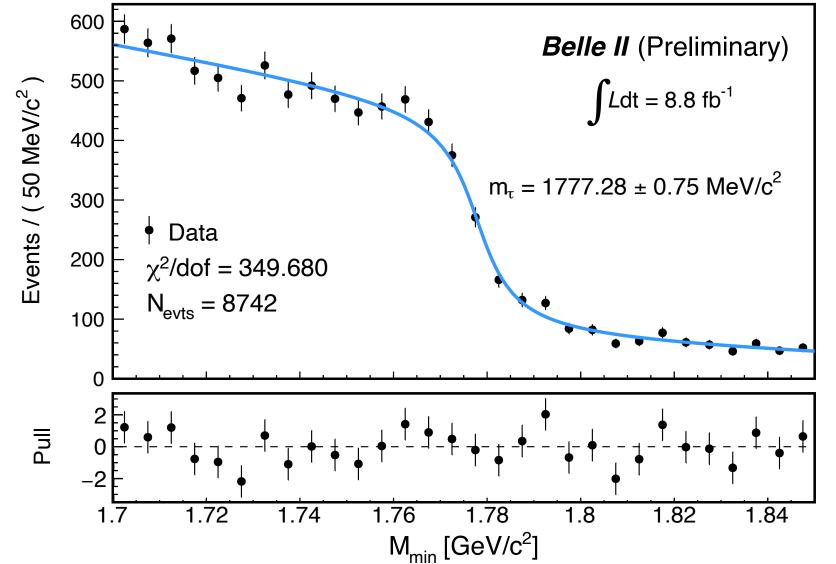
Systematics

- Compatible precision with previous B factory results
- dominated by uncertainty on the track momentum scale
- expected to improve

$$F(M_{min} | \overrightarrow{P}) = (P_3 + P_4 \cdot M_{min}) \cdot \tan^{-1}[(M_{min} - P_1)/P_2] + P_5 \cdot M_{min} + 1$$

BELLE2-NOTE-PH-2020-001



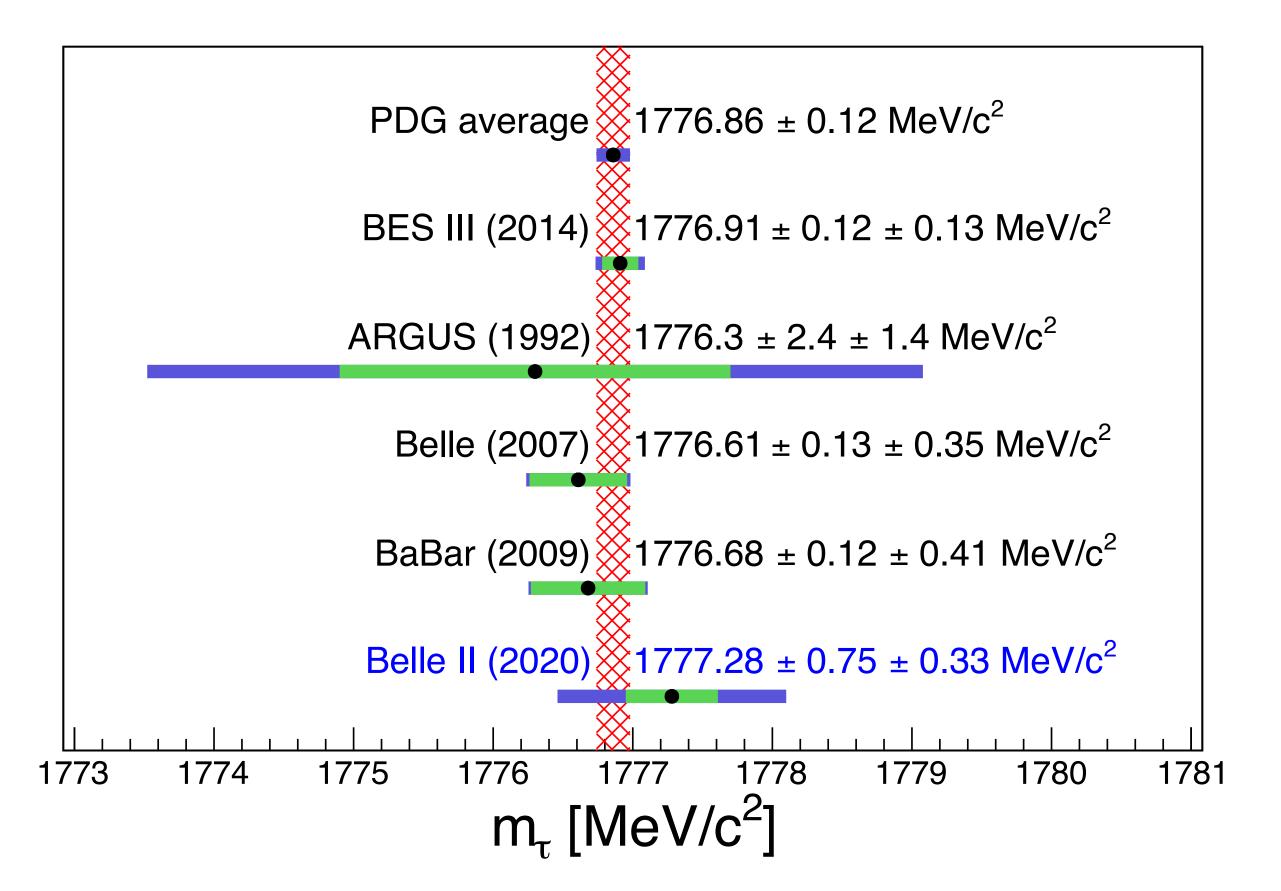


Systematic uncertainty	MeV/c^2
Momentum shift due to the B-field map	0.29
Estimator bias	0.12
Choice of p.d.f.	0.08
Fit window	0.04
Beam energy shifts	0.03
Mass dependence of bias	0.02
Trigger efficiency	≤ 0.01
Initial parameters	≤ 0.01
Background processes	≤ 0.01
Tracking efficiency	≤ 0.01

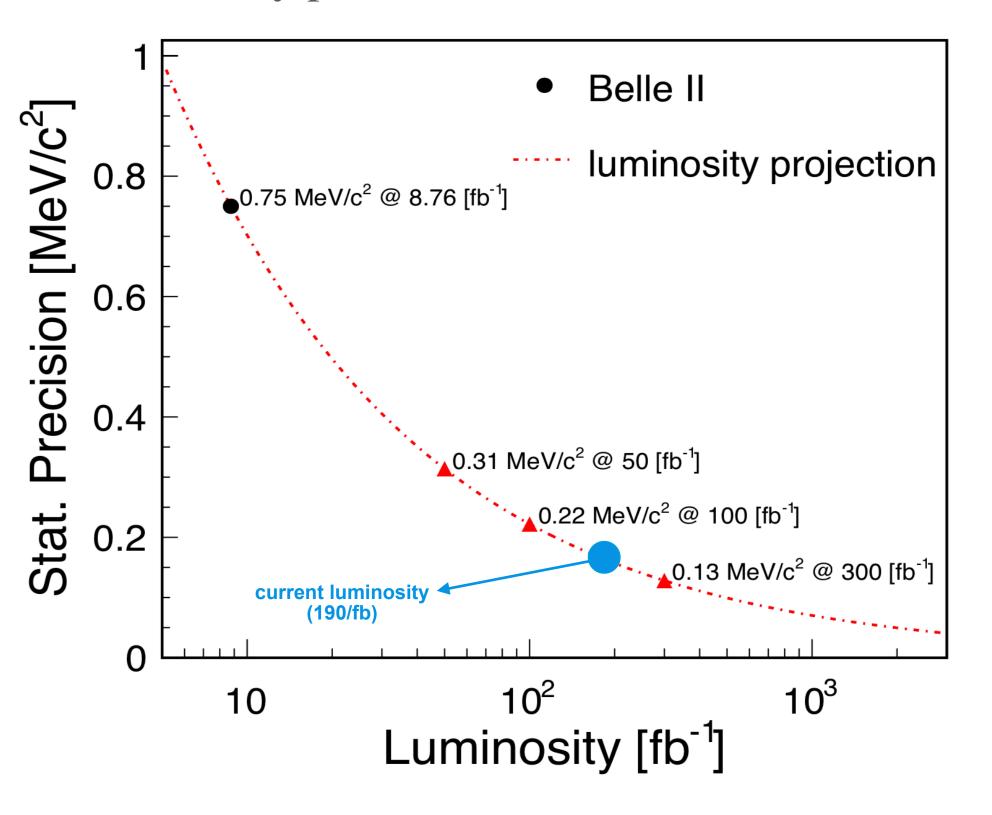
The t leptons mass

Goal: achieve best precision among pseudomass measurements

- best result from BES III from pair production at threshold energy
- best measurement from pseudomass technique by Belle



- → expect to match statistical precision of Belle/BABAR with ~300 fb⁻¹
- → future improvements of systematic uncertainty
- eventually perform CPV test as well



Jeniffer Summer School 2021

The t lifetime

The lifetimes of elementary particles are statistical in nature.

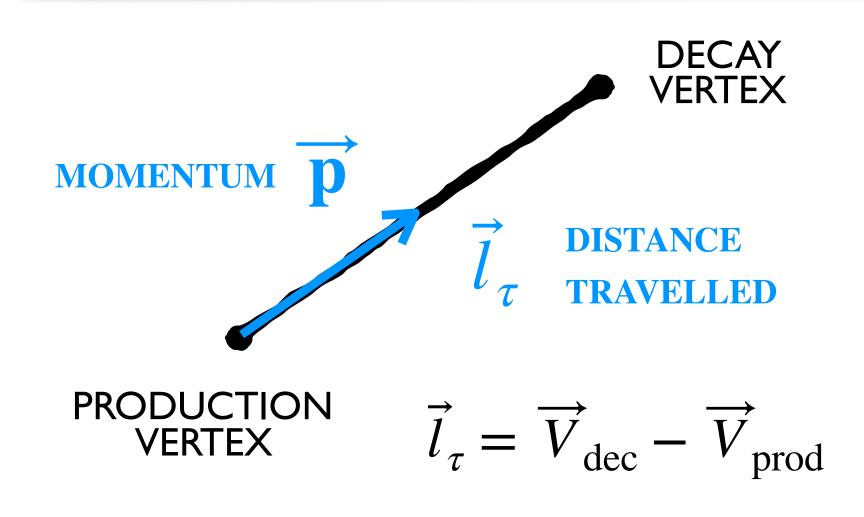
the lifetime is the time it takes for a sample to decay so that 1/e (~30%) of the kample is left

$$t = t_{decay} - t_{prod}$$

$$N(t) = N(0) e^{-t/\tau}$$

- $t = t_{decay} t_{prod}$ in an ideal experiment the lifetime τ is extracted with an exponential fit to the proper time distribution $t = t_{dec} t_{prod}$
- the finite resolution of a real experiment has a to be taken into account

$$p(t;\tau) = \frac{1}{\tau}e^{-t/\tau} \times \mathcal{R}(t)$$
 Proper time resolution



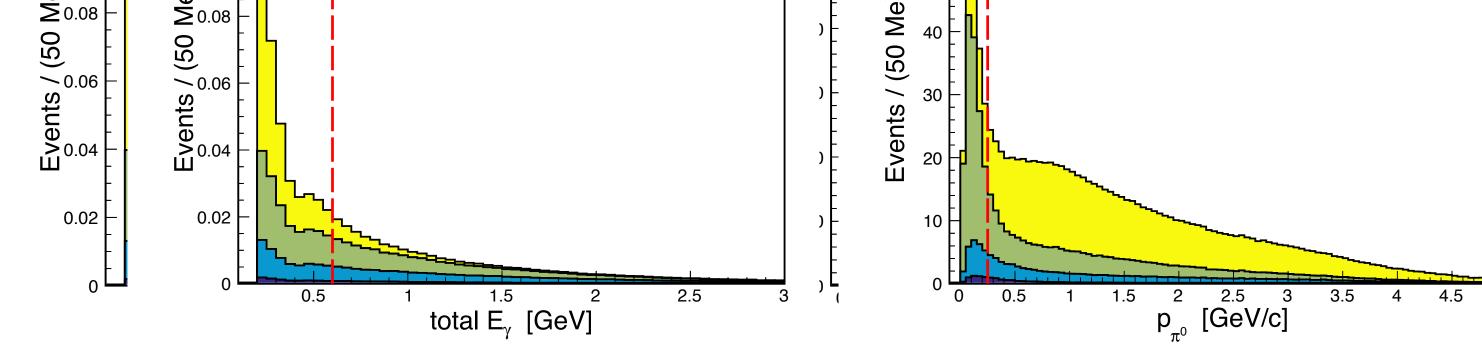
The proper time is measured using the particle's

- \rightarrow flight distance in the lab frame \vec{l}_{τ}
- \rightarrow momentum $\overrightarrow{\mathbf{p}}$ in the lab frame

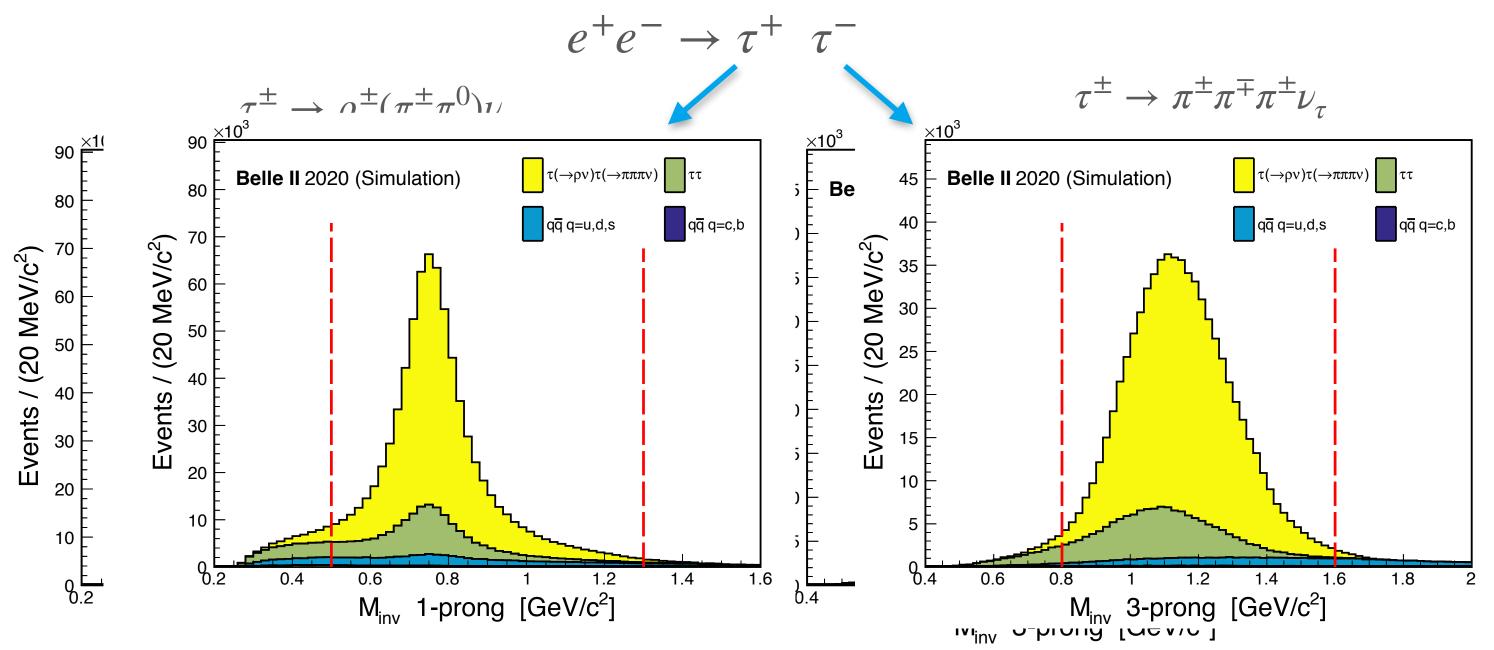
$$t_{lab} = l_{\tau} / \beta c$$

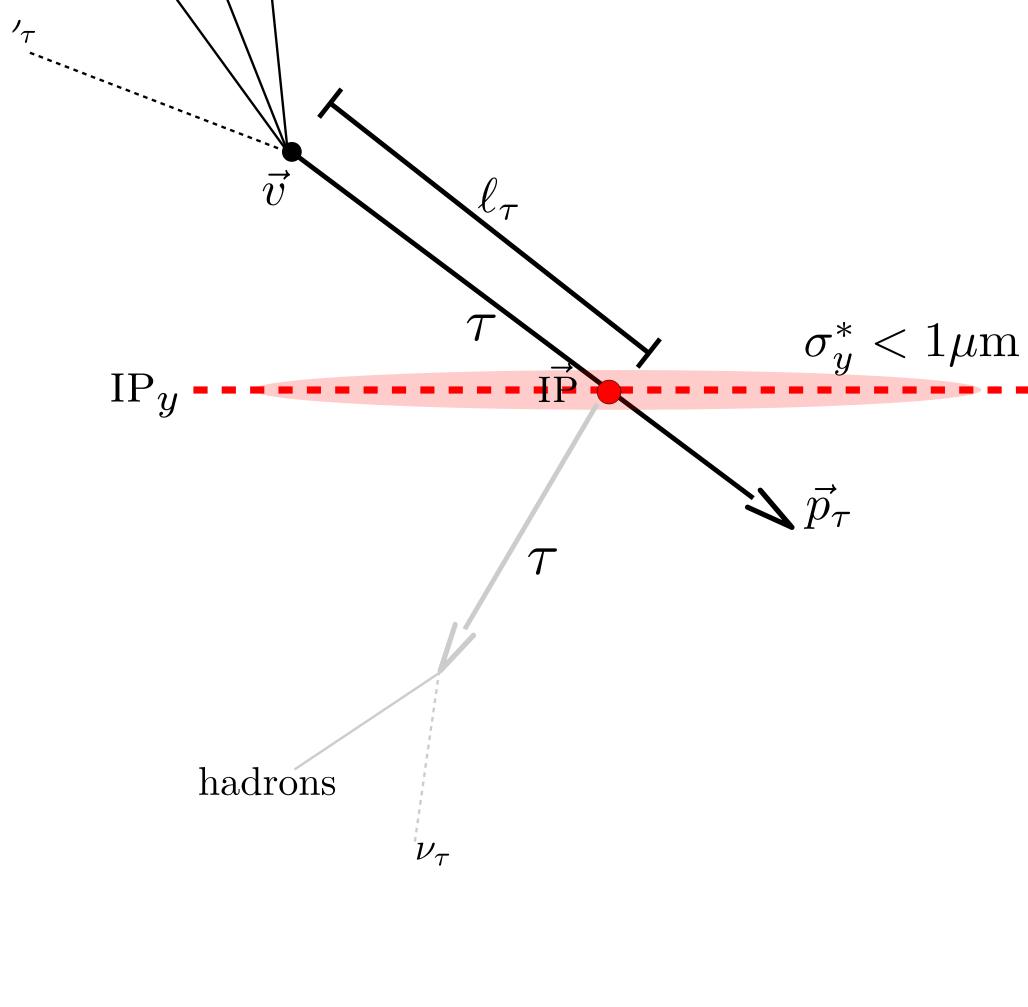
$$t_{true} = \frac{l_{\tau}}{\beta \gamma c} = m \frac{l_{\tau}}{p}$$
measure these!

24



- \rightarrow (2) estimate τ momentum using τ decay products p_{τ}
 - increase the statistical precision by using 3x1 topology





3-prong

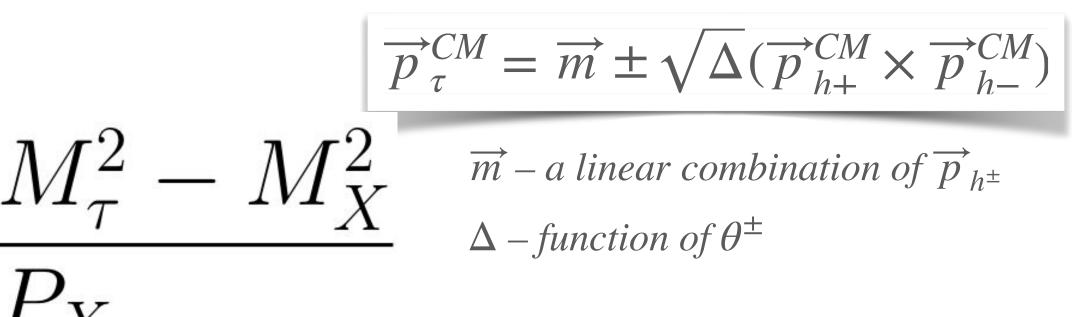
- \rightarrow (3) production vertex that is intersection of momentum direction with plane $y = IP_y$
 - exploits the tiny beam spot size at the interaction

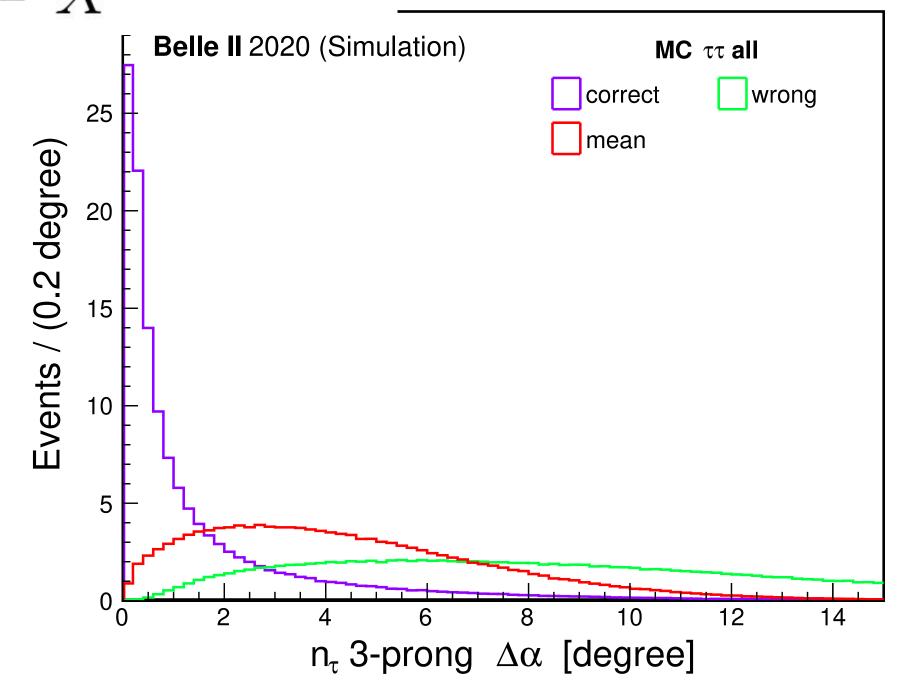


Momentum of n=

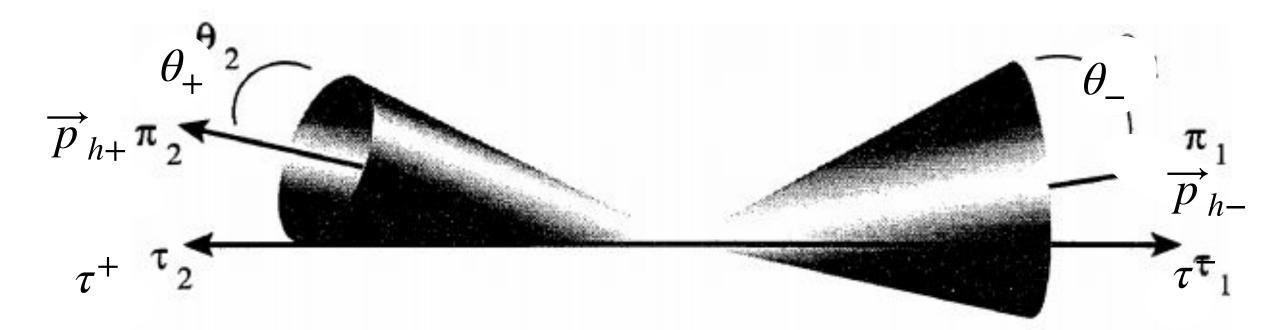
In the CMS the τ^+ (τ^-) direction lies on the surface of a cone with axis \overrightarrow{p}_{h+} (\overrightarrow{p}_{h-}) and opening angle θ_+ (θ_-).

 \rightarrow the τ momentum is given by the intersection of the two cones

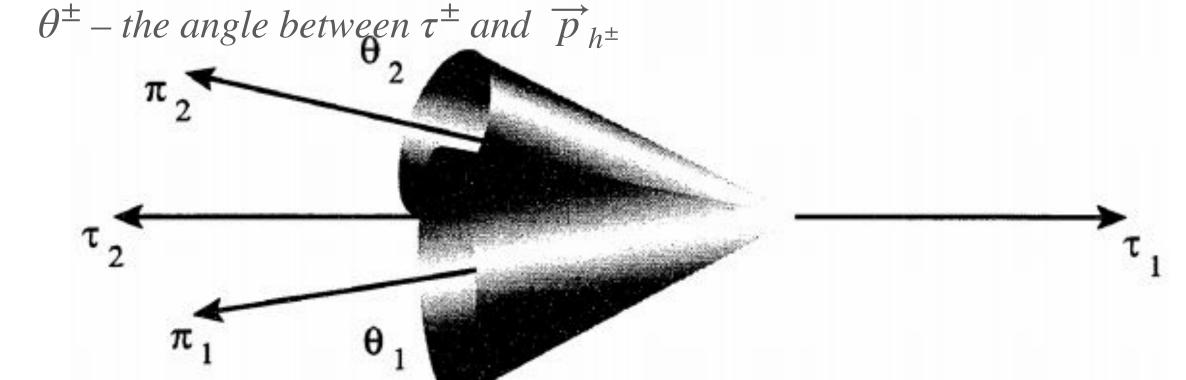




- $\Delta > 0$: the
- $\Delta = 0$: then
- $\Delta < 0$: then



 $\overrightarrow{p}_{h^{\pm}}$ – the total momentum of visible daughters

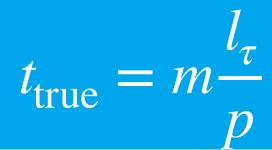


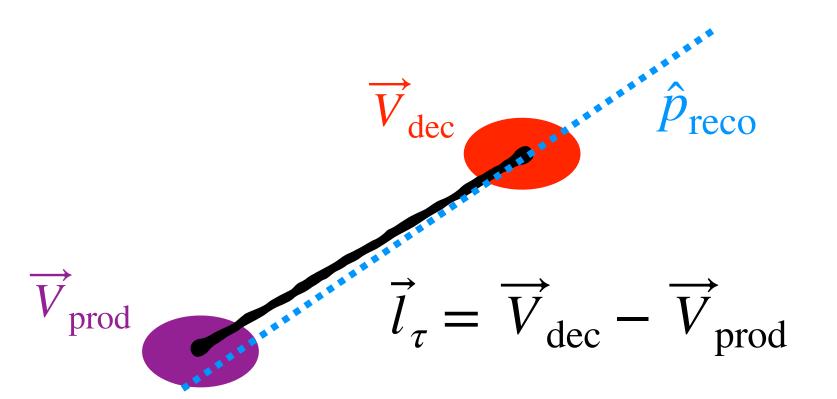
Consider only events having $\Delta \ge 0$











The flight distance is determined using the positions of decay and production vertices and the direction of the momentum

the direction is better measured by the momentum

$$ec{l}_{ au}\cdot\hat{p}_{reco}$$

$$\overrightarrow{V}_{\text{dec}} \qquad \overrightarrow{l}_{\tau} \cdot \widehat{p}_{reco}$$

$$\widehat{p}_{reco} = \overrightarrow{p}_{\tau} / |\overrightarrow{p}_{\tau}|$$

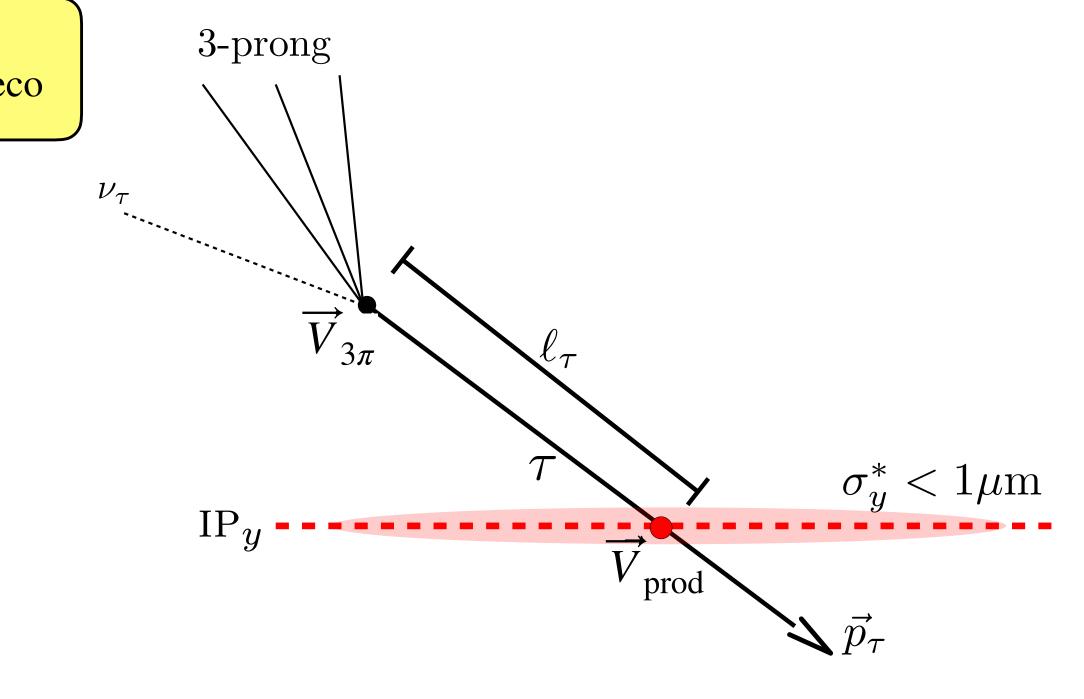
The decay length l_{τ} is extracted minimising a function $l = \overrightarrow{d}_{\text{reco}} \cdot \hat{p}_{\text{reco}}$

- beam-spot constraint
- \rightarrow the relation between τ displacement and production/decay vertices

$$\overrightarrow{V}_{\text{prod}} + l_{\tau} \cdot \hat{p}_{\text{reco}} - \overrightarrow{V}_{3\pi} = 0$$

strong constraint on IP_v

3-prong vertex



Proper decay time

$$t = m \frac{l_{\tau}}{p}$$

[µm]

computed from the reconstructed decay length l_{τ} and the estimated momentum p

Proper decay time resolution

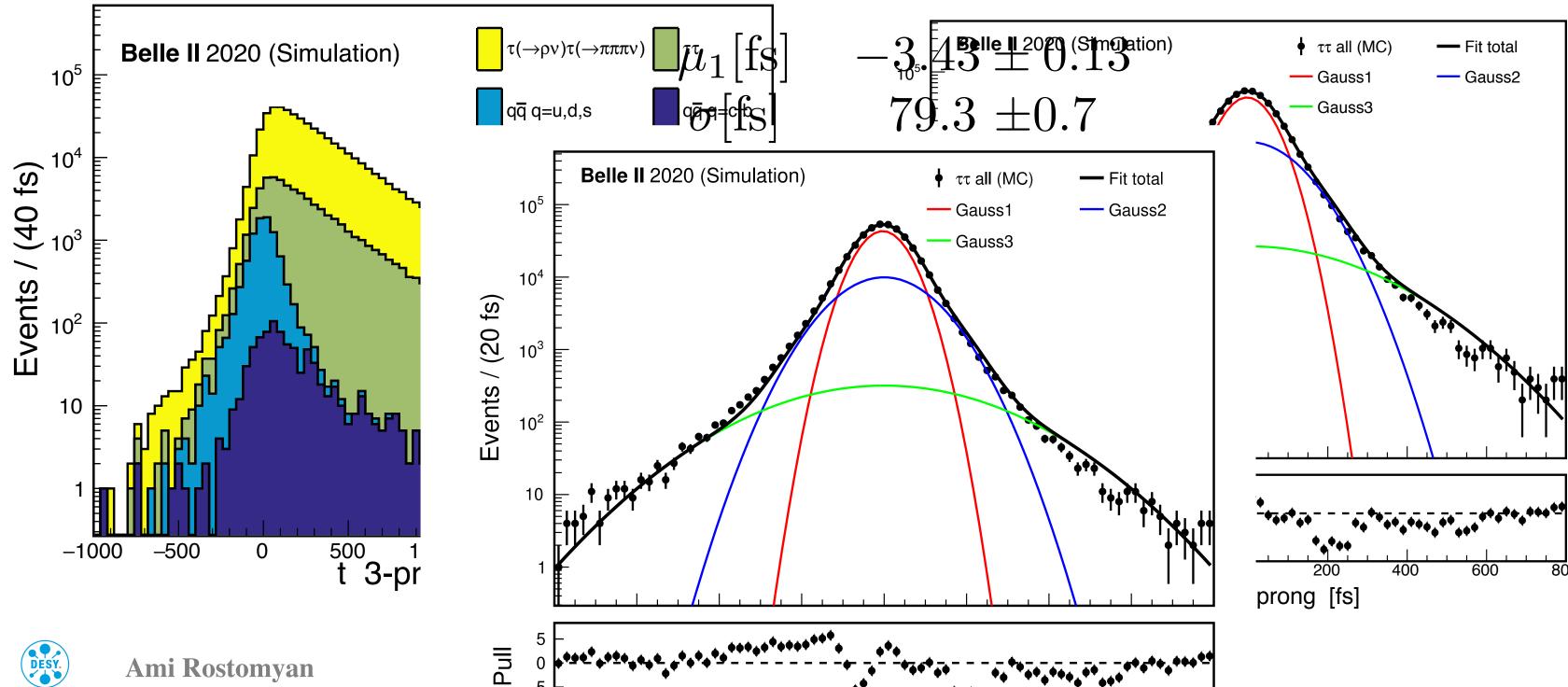
$$\Delta t = t^{rec} - t^{gen}$$

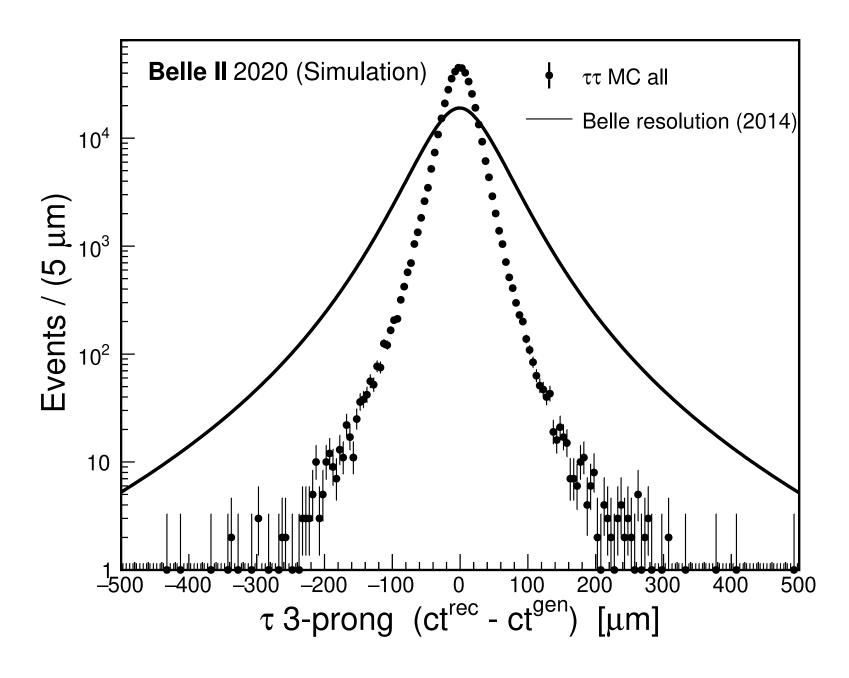
binned ML fit with 3 Gaussians

$$\mu[fs] = -3.43 \pm 0.13$$

 $\sigma[fs] = -79.3 \pm 0.7$

→ the resolution @Belle II is nearly x2 narrower than @Belle



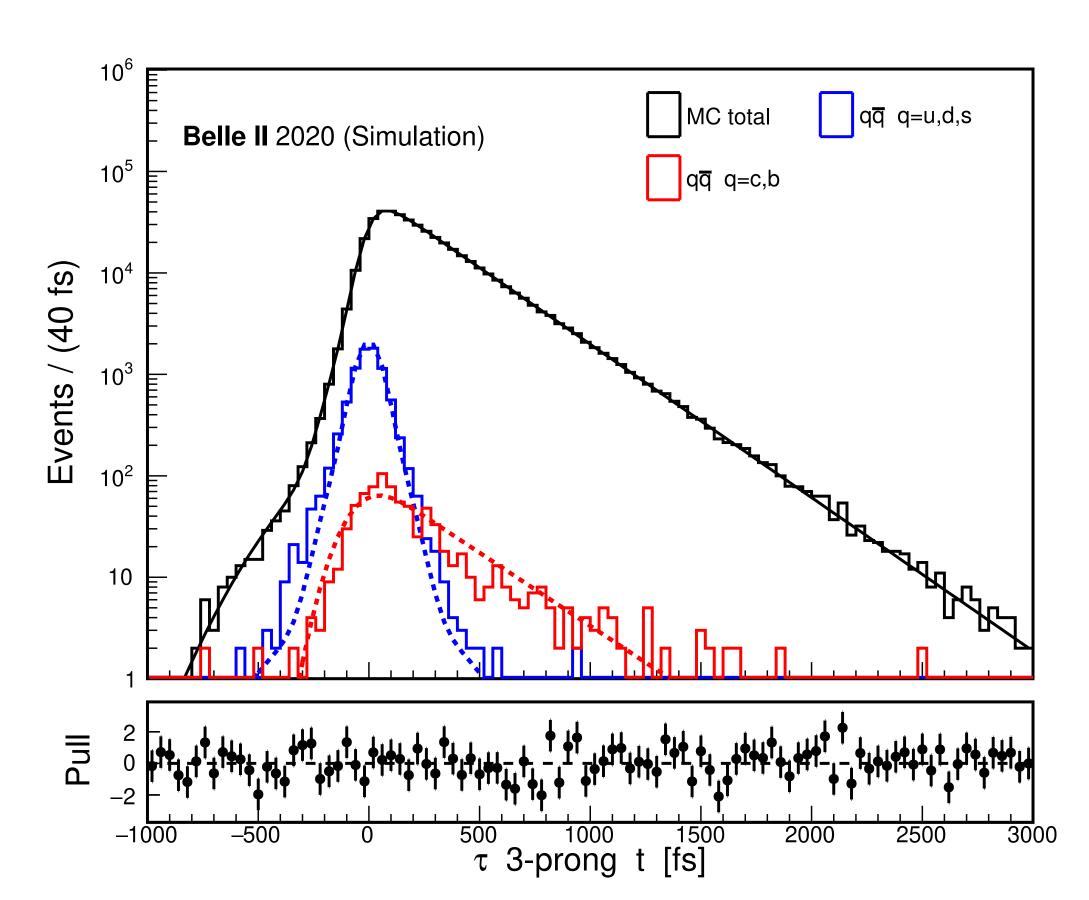


The t lepton lifetime

Fit the proper time distribution with a convolution of an exponential distribution and resolution function

$$p(t;\tau) = \frac{1}{\tau}e^{-t/\tau} \times \mathcal{R}(t) \longrightarrow \mathcal{T} = 287.2 \pm 0.5 \text{ fs}$$

BELLE2-NOTE-PH-2020-076



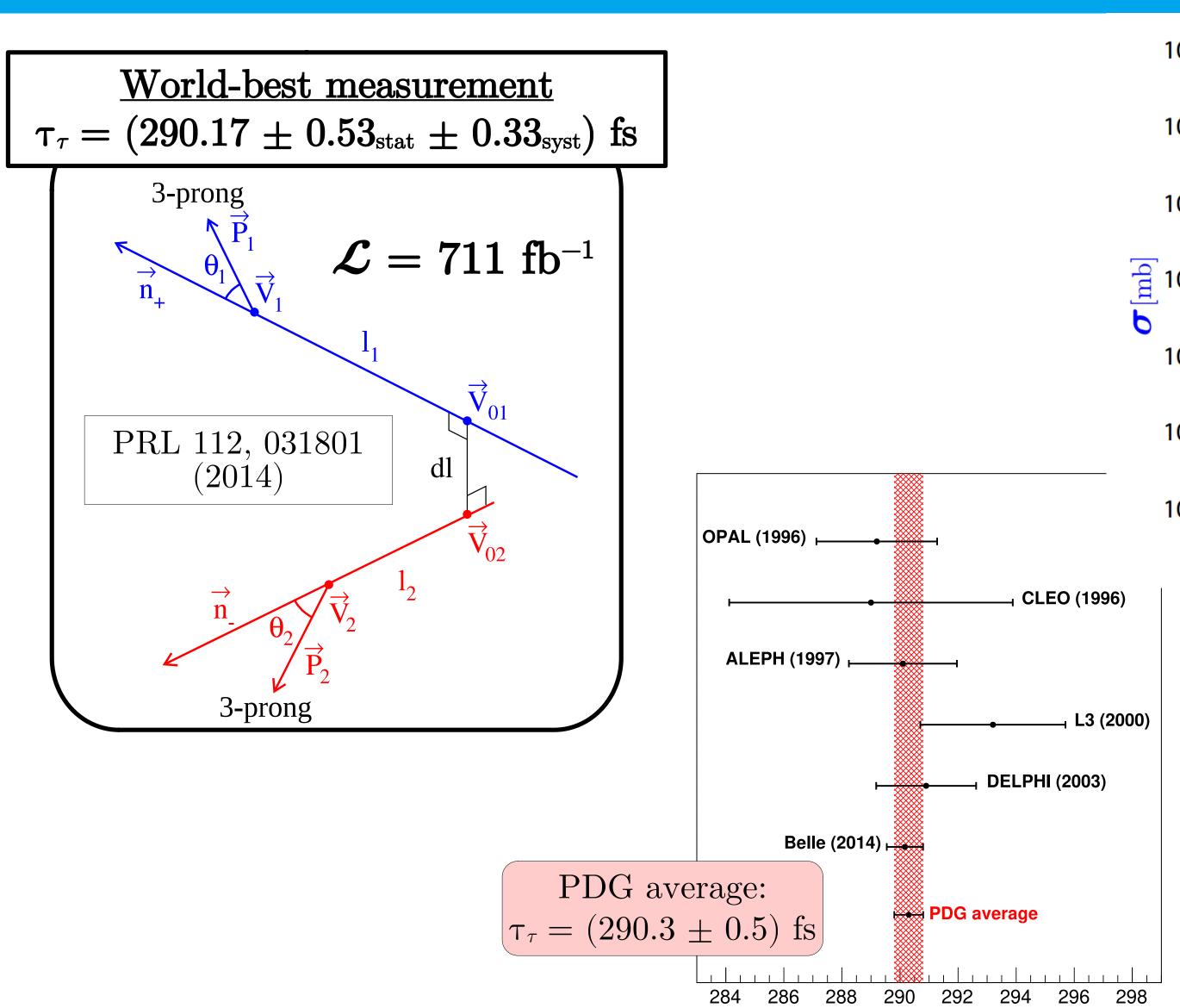
Generated lifetime $\tau = 290.57$ fs

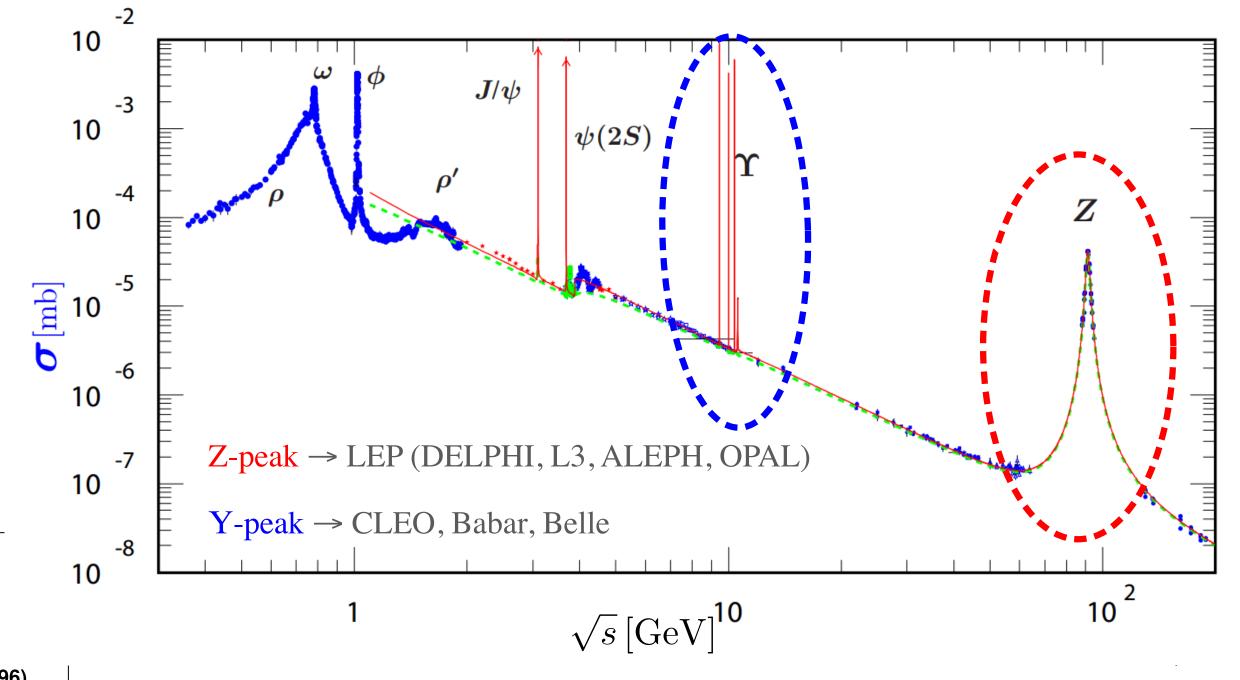
- → ~3 fs bias in the measurement
 - → ISR/FSR losses
 - overestimation of p_{τ}^{τ} results in underestimation of proper time $p_{\tau}c$
 - intrinsic bias of the measurement
 - estimate the bias from MC and correct the measurement $\mu_1[\mathrm{fs}] \quad -0.80 \pm 0.20 \quad \tau[\mathrm{fs}]$

DESY.

Jeniffer Summer School 2021

Previous measurements of τ lifetime in $e^+e^- \to \tau^+\tau^-$





With respect to Belle:

- exploit the tiny beam spot size at the IP
- increase the statistical precision by a factor of 5 using 3x1 topology
- → competitive statistical precision can already be reached with 200/fb

 $au_{ au}$ [fs]

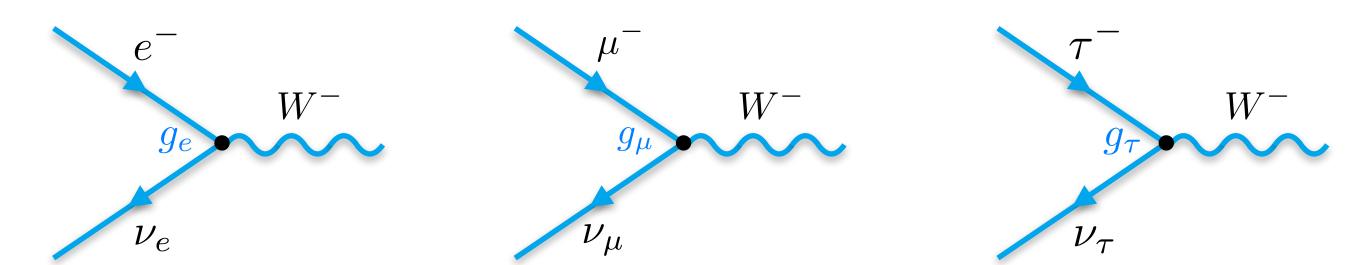
e-μ-τ universality

e, μ and τ differ only by

The coupling of leptons to W bosons is flavour-independent:

 $g_e = g_\mu = g_\tau$

- → the mass
- → different and separately conserved lepton numbers



Anomalies in quark sector

- → $R(D)-R(D^*)$ plane (~3.1 σ)
- ⇒ R(K) (3.1 σ), also P₅' in B→K* $\mu\mu$ (~3.4 σ)
- → and more...

$$R(D^*) = \frac{\mathcal{BR}(B \to D^* \tau \nu)}{\mathcal{BR}(B \to D^* \ell \nu)} \text{ with } \ell = e, \mu$$

Significant tensions in lepton sector

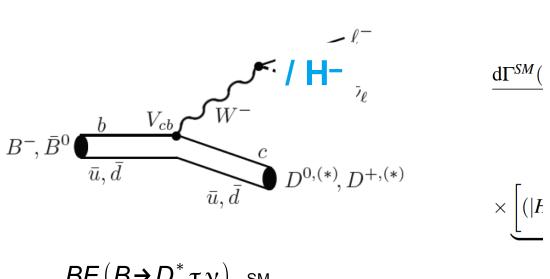
 \rightarrow anomalous magnetic moment of μ (4.2 σ) and e (~2.5 σ)

Imperial College

Semileptonic decly son interaction rates involving e, μ or τ

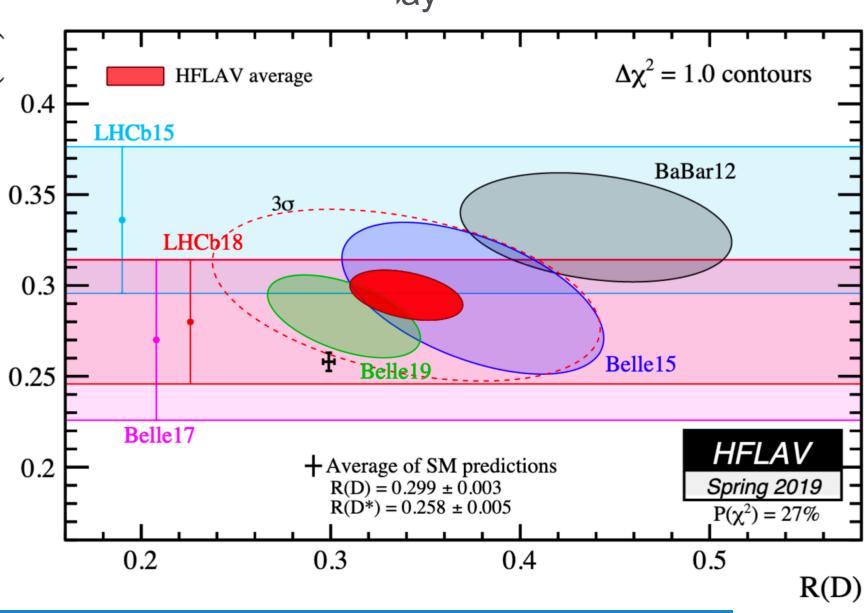
London In the SM the decay $R^0 \rightarrow D^{*-}I^+\nu$ proceed through a tree level decay

Semileptonic d $\stackrel{\frown}{\mathbb{Z}}$ | London In the SM, the decay $B^0 \rightarrow D^{*-}I^+v$ proceed 1 0.4



$$R(D^*) = \frac{Br(B \rightarrow D^* \psi \nu)}{BF(B \rightarrow D^* \psi \nu)} \stackrel{\text{SM}}{=} 0.252 \pm 0.00$$

11 July 2017 Ulrik Egede



Are these hints of a new fundamental interaction that violates LFU?

e-u-t universality

If so, then we could also see hints in the tau sector

Test of e-µ universality

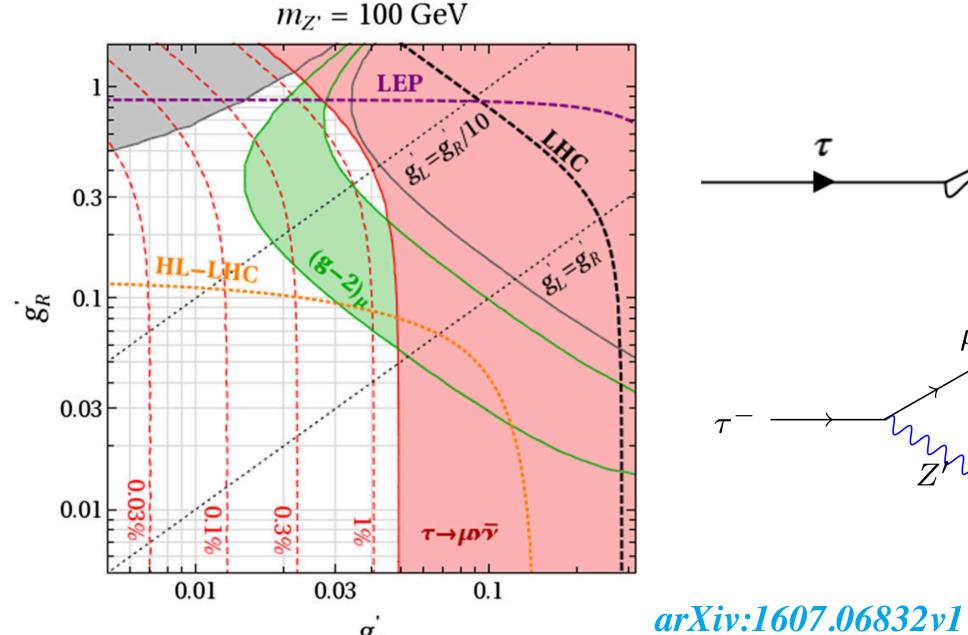
$$\left(\frac{g_{\mu}}{g_{e}}\right)_{\tau}^{2} \propto \frac{B(\tau^{-} \to \mu^{-}\bar{\nu}_{\mu}\nu_{\tau})}{B(\tau^{-} \to e^{-}\bar{\nu}_{e}\nu_{\tau})} \frac{f(m_{e}^{2}/m_{\tau}^{2})}{f(m_{\mu}^{2}/m_{\tau}^{2})}$$

$$f(x) = 1 - 8x + 8x^3 - x^4 - 12x^2 \log x$$

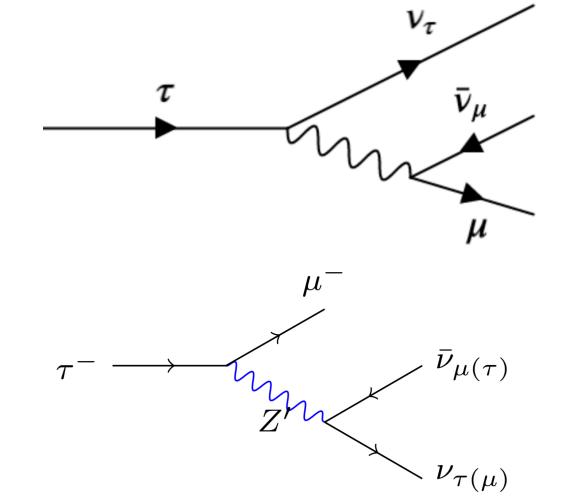
Test of τ–μ universality

$$\left(\frac{g_{\tau}}{g_{\mu}}\right)_{h}^{2} \propto \frac{B(\tau \to h\nu_{\tau})}{B(h \to \mu\nu_{\mu})} \frac{2m_{h}m_{\mu}^{2}\tau_{h}}{(1 + \delta_{h})m_{\tau}^{3}\tau_{\tau}} \left(\frac{1 - m_{\mu}^{2}/m_{h}^{2}}{1 - m_{h}^{2}/m_{\tau}^{2}}\right)^{2}$$

Can put strong constraints on lepton flavour violating Z' models



 g_L



W Decays (0.997±0.010) K₁₃ Decays (1.0021 ± 0.0025) K₁₂ Decays (1.004 ± 0.007) π Decays (1.0023±0.0016) τ Decays (1.0000±0.0020) World Average (2008) (1.0015±0.0011) This Work (τ Decays) (1.0036 ± 0.0020) **New World Average** (1.0020±0.0010) 0.98 0.985 0.99 0.995 1.005 1.01 g_{μ}/g_{e}



Jeniffer Summer School 2021

arXiv:1607.06832v1

e-μ-τ universality

Most precise measurement from BaBar

$$\left(\frac{g_{\mu}}{g_{e}}\right)_{\tau} = 1.0036 \pm 0.0020$$
 $\left(\frac{g_{\tau}}{g_{\mu}}\right)_{h} = 0.9850 \pm 0.0054$

in agreement with SM

→ 2.8 σ below the SM expectation

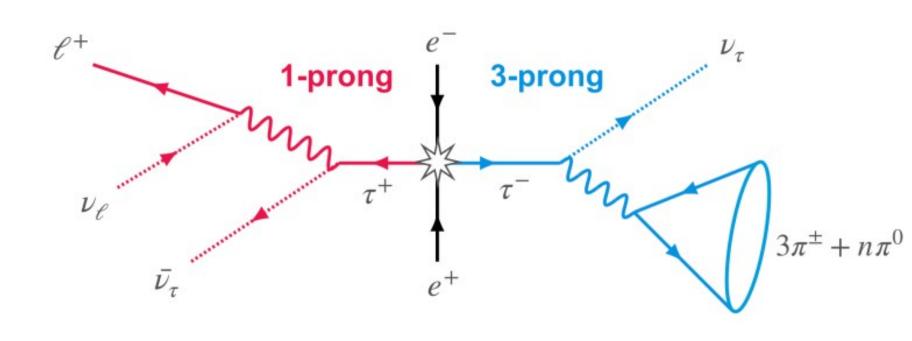
- → The BR measurements dominated by systematic uncertainty
 - → µ: PID due limited size of data and MC samples
 - → h: additional contribution to systematics detector modelling and associated BGs

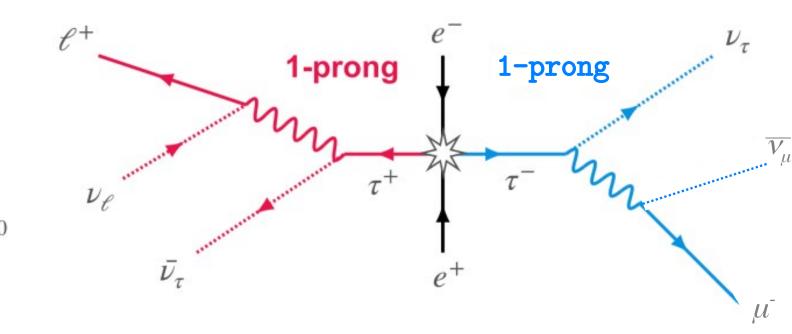
- Phys. Rev. Lett. 105 051602 -

	μ	π	K
\mathbf{N}^{D}	731102	369091	25123
Purity	97.3%	78.7%	76.6%
Total Efficiency	0.485%	0.324%	0.330%
Particle ID Efficiency	74.5%	74.6%	84.6%
Systematic uncertainties:			
Particle ID	0.32	0.51	0.94
Detector response	0.08	0.64	0.54
Backgrounds	0.08	0.44	0.85
Trigger	0.10	0.10	0.10
$\pi^-\pi^-\pi^+$ modelling	0.01	0.07	0.27
Radiation	0.04	0.10	0.04
$\mathcal{B}(au^- o \pi^- \pi^- \pi^+ u_ au)$	0.05	0.15	0.40
$\mathcal{L}\sigma_{e^+e^- o au^+ au^-}$	0.02	0.39	0.20
Total [%]	0.36	1.0	1.5

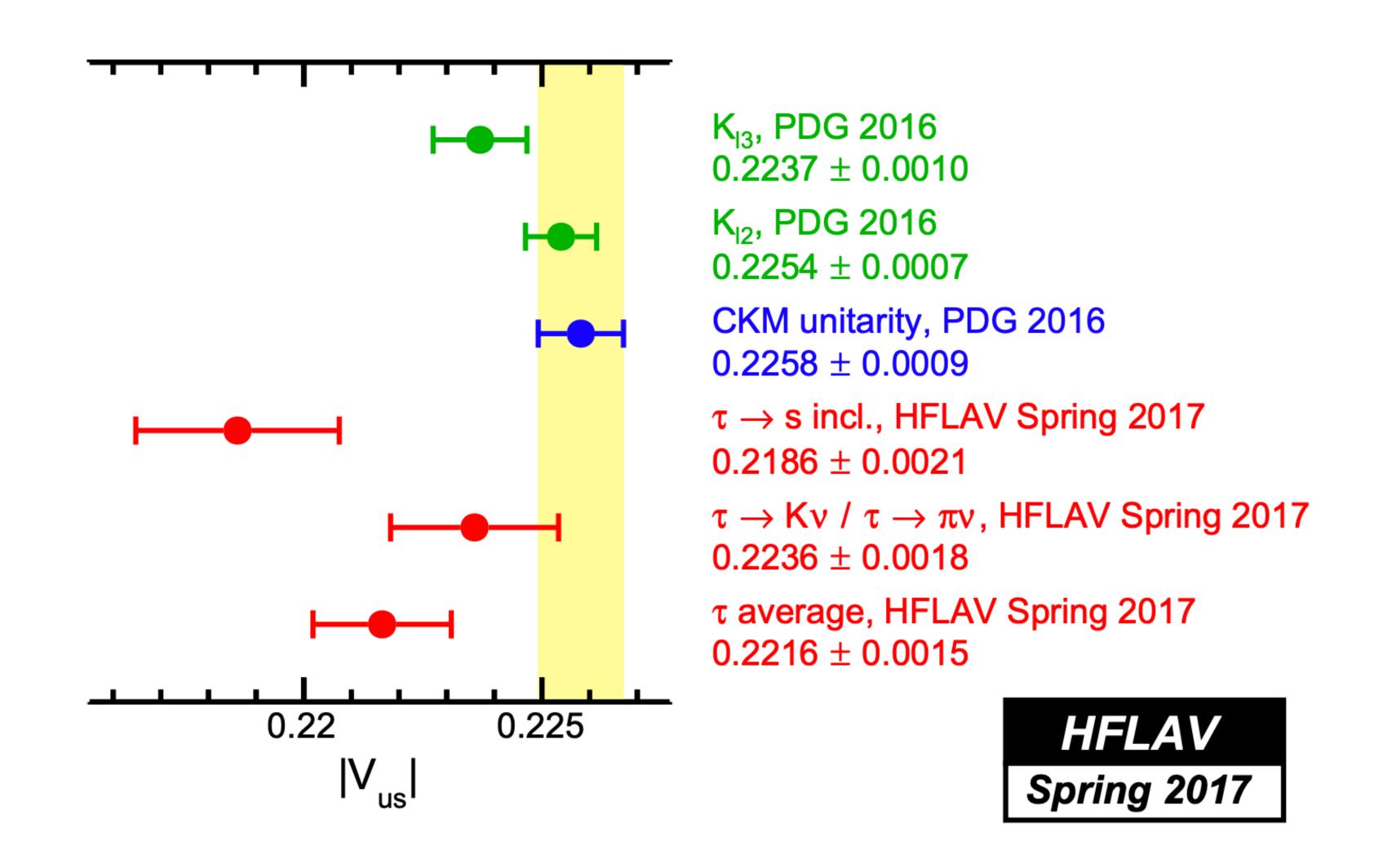
Can Belle II improve this?

- → use 3x1 and 1x1 (not used @BaBar) to improve the statistical precision
- work hard to improve the systematics





Vus from t decays





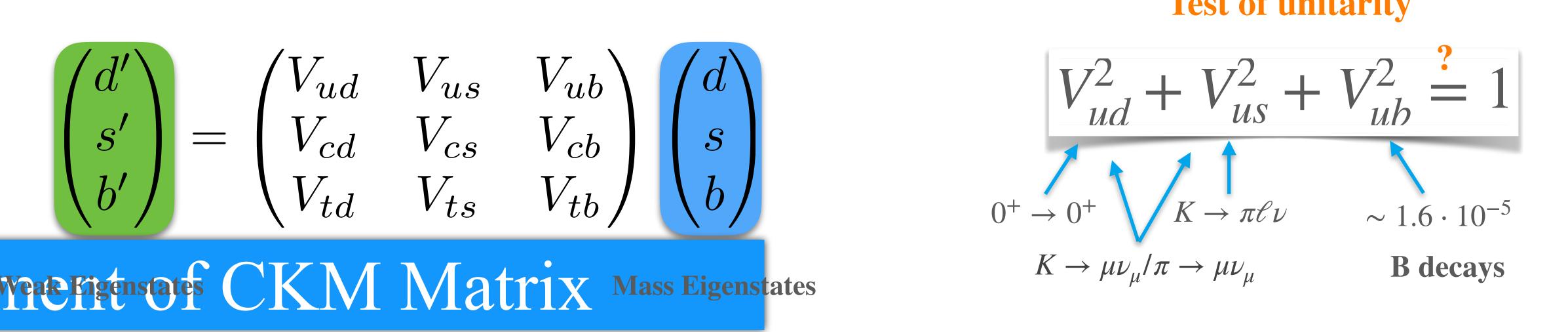
Jeniffer Summer School 2021

Test of unitarity

Unique opportunity for probing the coupling strength of the weak current to the first and second generation of quarks to a very high precision

Test of unitarity

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$



le Weak Figenstates CKM Matrix Mass Eigenstates

ween Weak and Mass Eigenstates From kaon, pion, baryon and nuclear decays

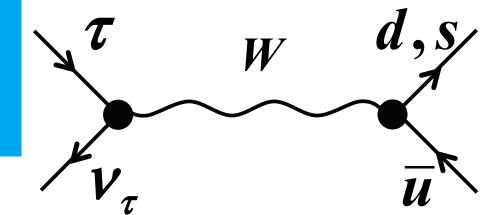
	1/ -	\ /		. V.
us cs	V_{ud}	$0^{+} \rightarrow 0^{+}$ $\pi \rightarrow \pi e \nu_{e}$	$n \rightarrow pe\nu_e$	$\pi o \ell u_{\ell}$
ts	V_{us}	$K o \pi \ell \nu$	$\Lambda \to pe\nu_e$	$K \to \ell \nu_{\ell}$

\rightarrow From τ decays

V_{ud}	$ au o \pi \pi^0 u_{ au}$	$ au ightarrow \pi u_{ au}$	$ au ightarrow h_{NS} u_{ au}$
V_{us}	$ au o K\pi u_{ au}$	$ au o K u_{ au}$	$ au ightarrow h_S u_{ au}$

Jeniffer Summer School 2021

Two methods of V_{us} from τ decays



Exclusive: compare the BR of $\tau \to \pi \nu$ and $\tau \to K \nu$

- BaBar, Phys. Rev. Lett. 105 051602 -

Fermi constant

$$B(\tau^{-} \to K^{-}\nu_{\tau}) = \frac{G_F^2 f_K^2 |V_{us}|^2 m_{\tau}^3 \tau_{\tau}}{16\pi} \left(1 - \frac{m_K^2}{m_{\tau}^2}\right)^2 S_{\text{EW}}$$

$$\frac{B(\tau^{-} \to K^{-}\nu_{\tau})}{B(\tau^{-} \to \pi^{-}\nu_{\tau})} = \frac{f_{K}^{2} |V_{us}|^{2}}{|f_{\pi}^{2}| |V_{ud}|^{2}} \frac{(m_{\tau}^{2} - m_{K}^{2})^{2}}{(m_{\tau}^{2} - m_{\pi}^{2})^{2}} (1 + \delta_{LD})$$

electroweak corrections

electroweak corrections

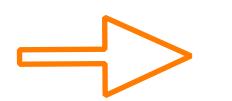
$$V_{us} = 0.2193 \pm 0.0032$$

$$V_{us} = 0.2255 \pm 0.0024$$

- \rightarrow within 2 σ of the value predicted by the CKM unitarity
- consistent with CKM unitarity

Inclusive: compare the BR of $\tau \to (\bar{u}d) \nu$ and $\tau \to (\bar{u}s) \nu$

 $BR \text{ w.r.t. } BR(\tau \rightarrow e\nu\nu)$



fundamental parameters of SM

$$(\alpha_s, |V_{us}|, m_s)$$

hadrons with S=0

$$\Delta R_{\text{SU(3) breaking}} = \frac{R_{\text{NS}}}{|V_{ud}|^2} - \frac{R_{\text{S}}}{|V_{us}|^2}$$

decay constant

hadrons with S=1

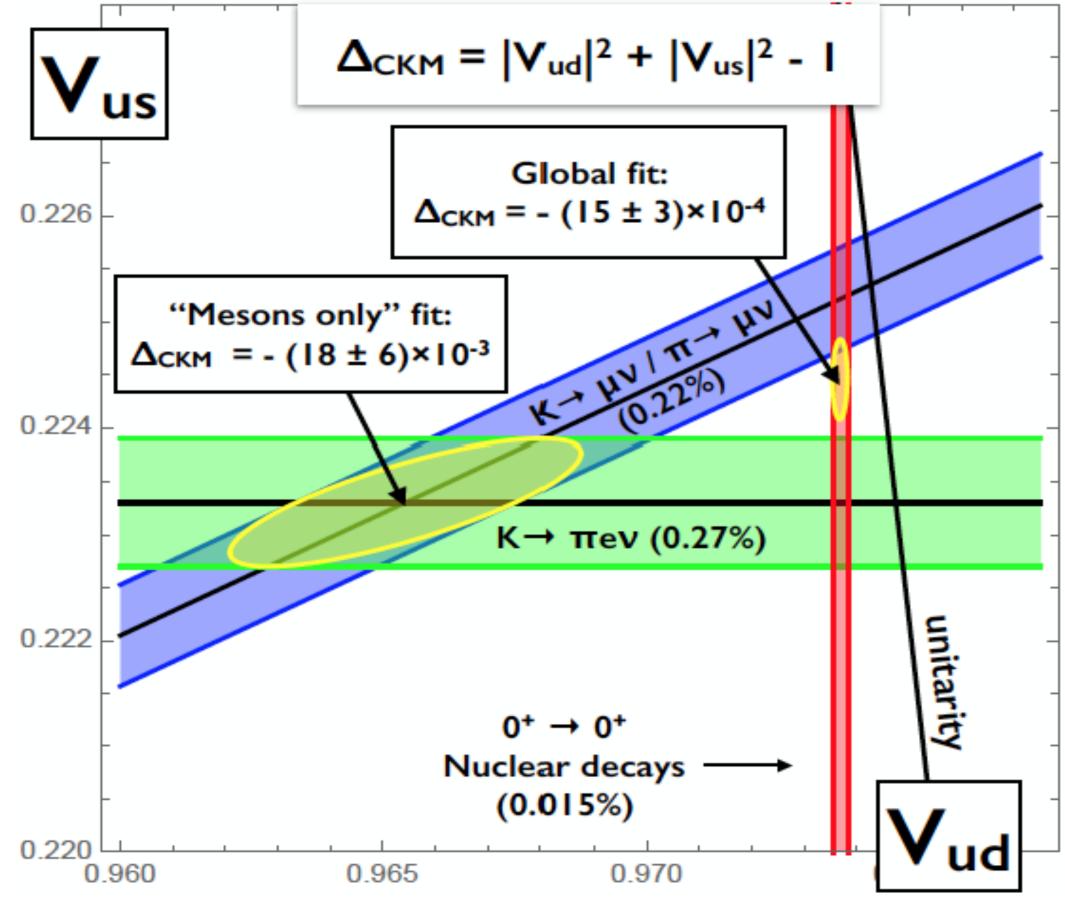
$$V_{\rm us} = 0.2186 \pm 0.0021$$

 \rightarrow within 3.1 σ of the value predicted by the CKM unitarity

Vus from t decays @Belle II

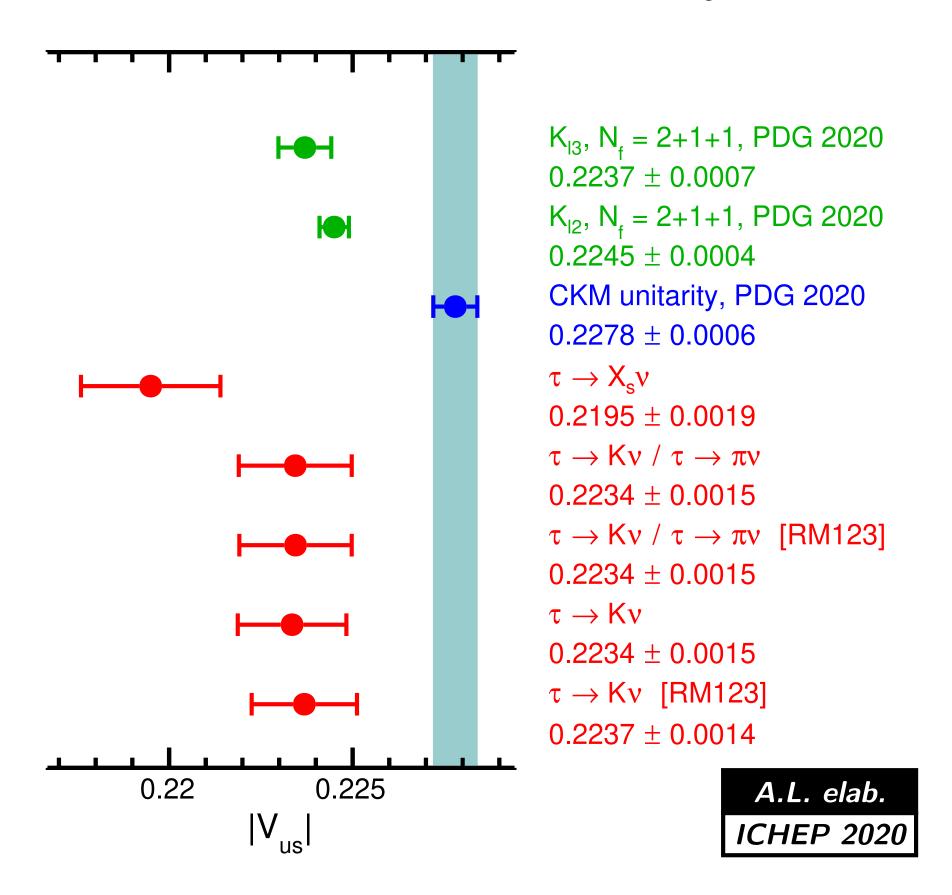
New results with improved theoretical input

- \rightarrow precise determination of V_{us} from kaon and nuclear decays
- \rightarrow discrepancy with CKM unitarity at 4.8 σ



Can t physics help?

- \rightarrow currently less precise determination of V_{us}
- → large PID systematic uncertainties @BaBar
- inclusive measurement not truly inclusive



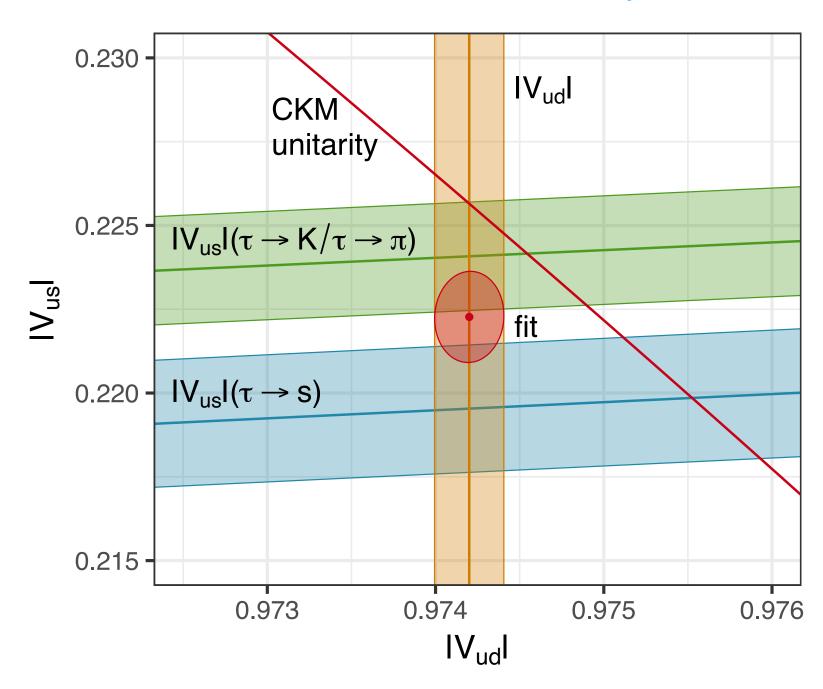


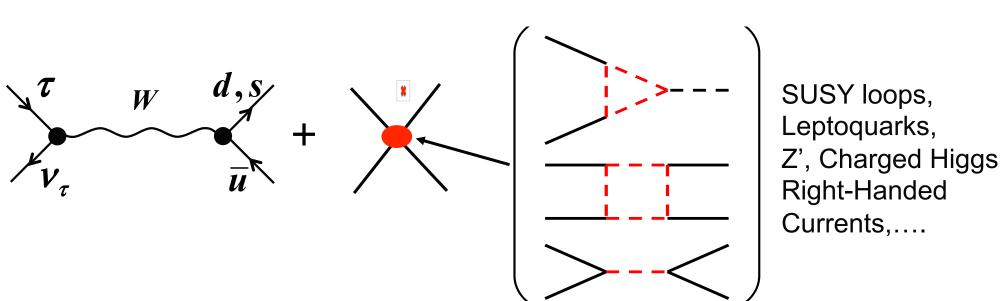
Jeniffer Summer School 2021

us

3σ tension between $|V_{us}|$ from the CKM matrix unitarity and $\tau \rightarrow s$.

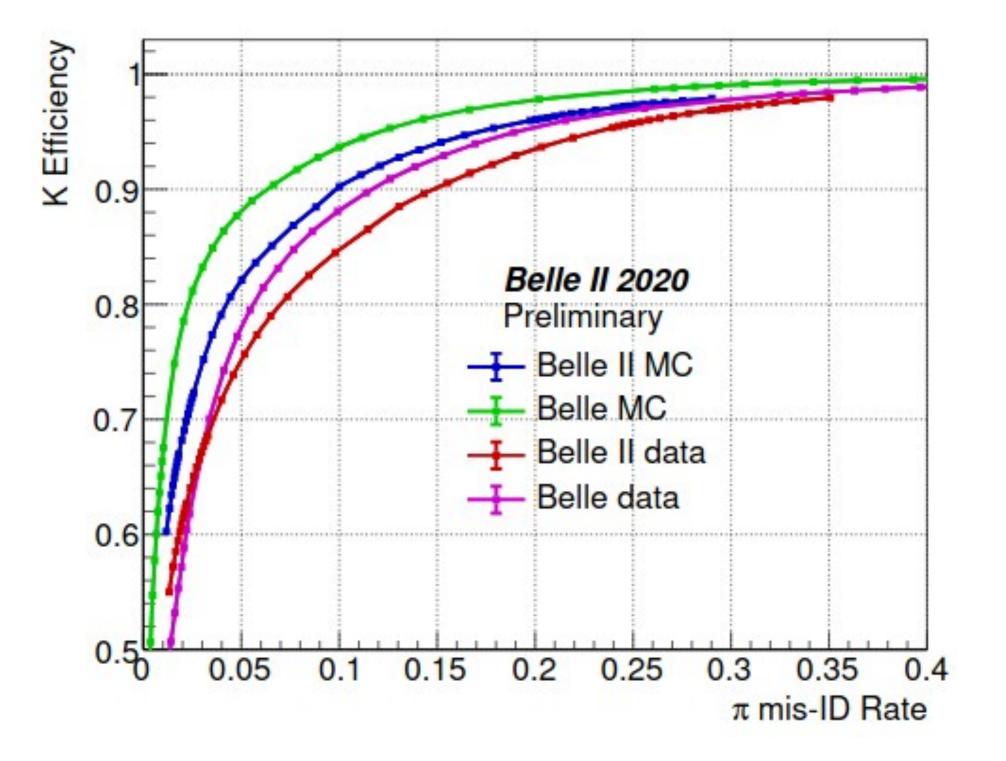
SciPost Phys. Proc. 1, 001 (2019)





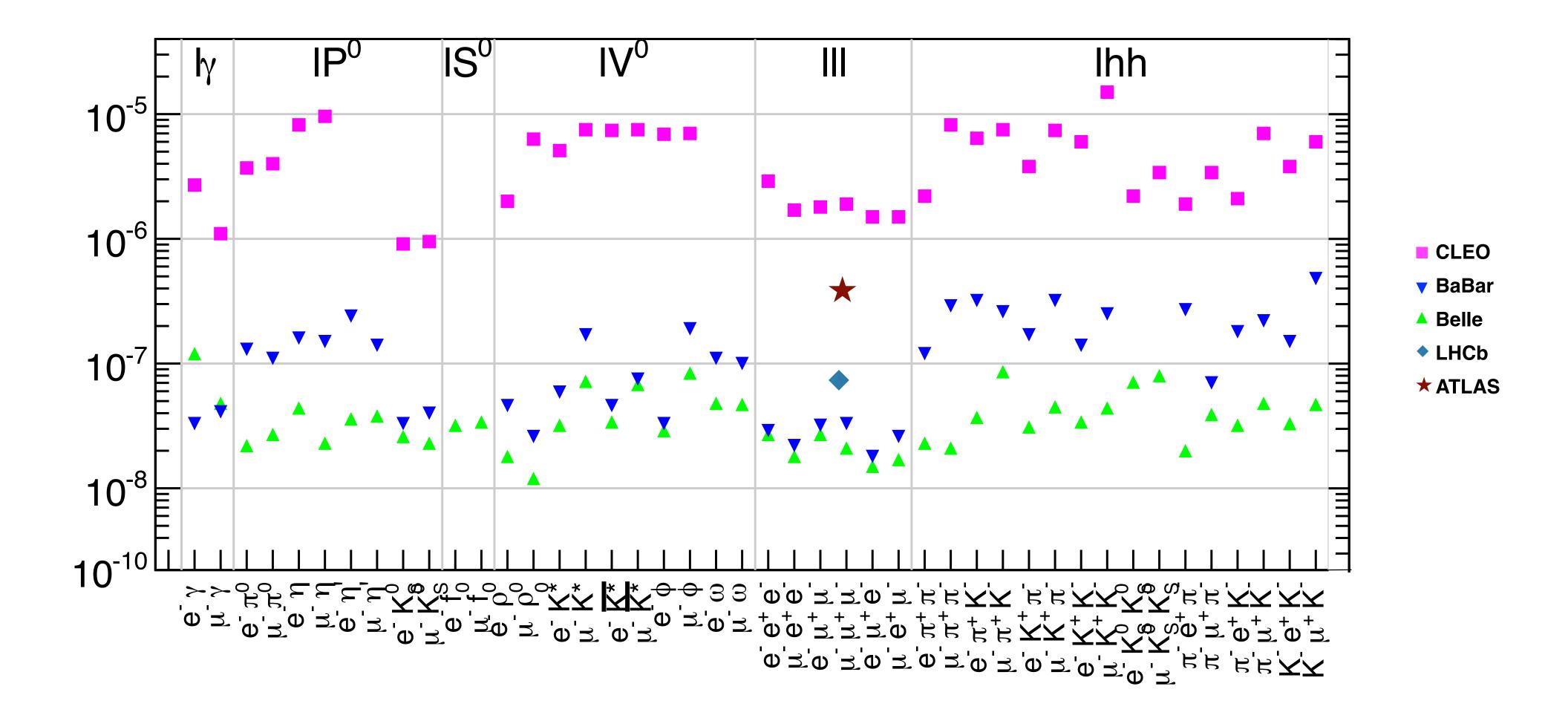
What can we do @BelleII?

- → larger data sample will be available
- → similar to LFU analysis use 3x1 and 1x1 topologies
- improve the understanding of the detector (PID, trigger, ...)



Belle II will improve the experimental precision!

LVF & LNV



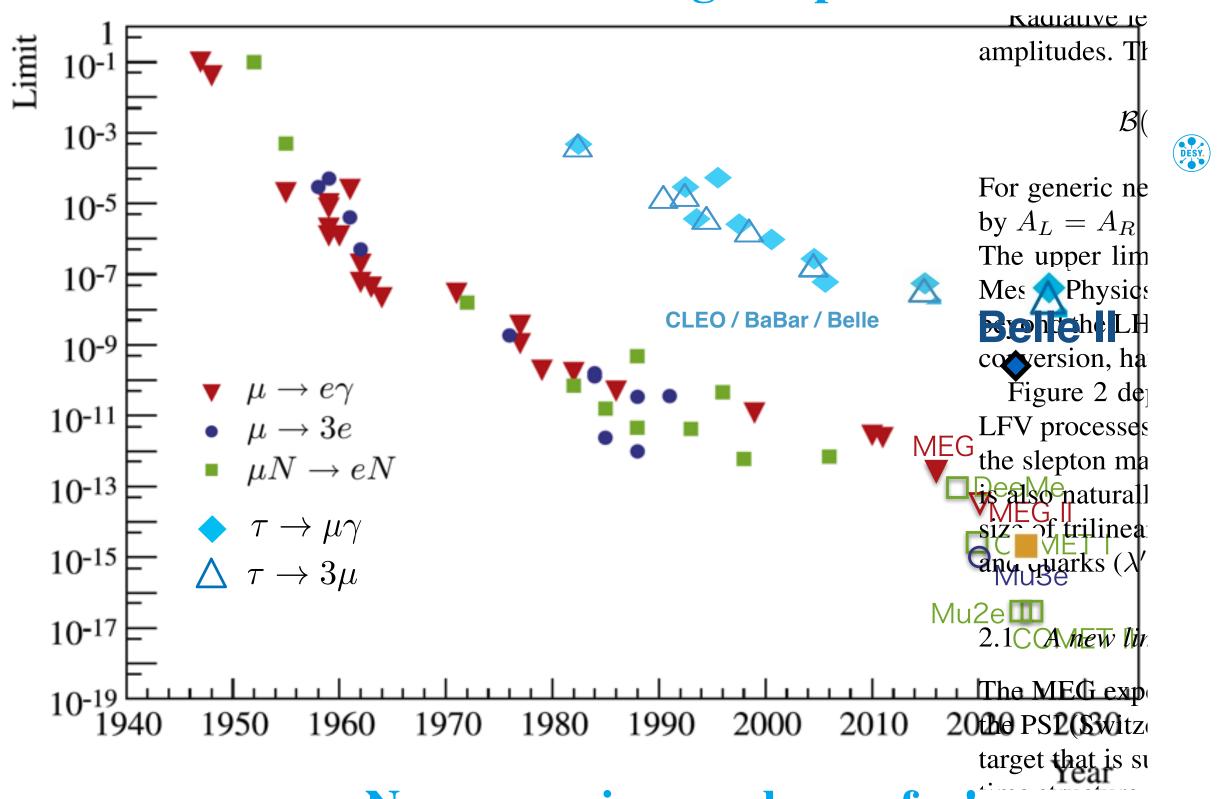
 $\mathbf{M}\;(m_{\mathrm{v}}=\mathbf{0})$

Ve Vµ

Data taking shifts

Tuning of PYTHIA8; charged track and neutral meson reconstruction

What about the charged leptons?



No success in searches so far!



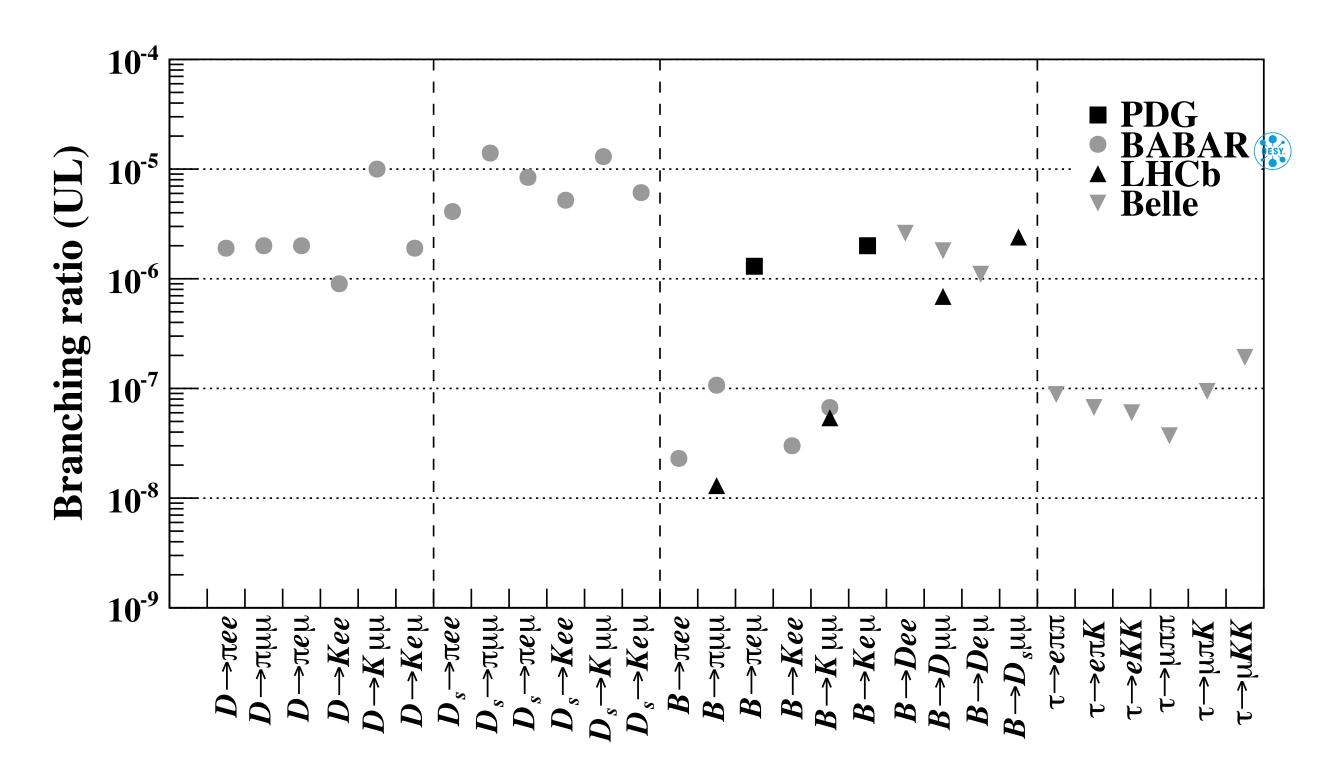
Lepton number conservation

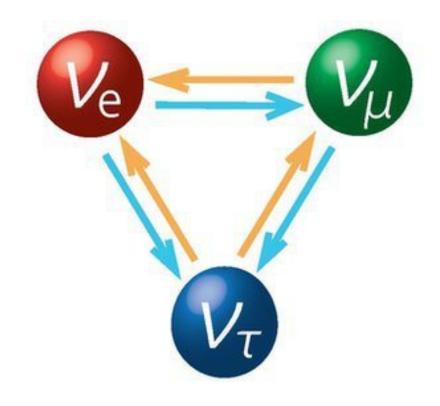
Conservation of the individual lepton-flavour and the total lepton numbers within the SM ($m_v = 0$)

$$G_{SM}^{global} = U(1)_B \times U(1)_{L_e} \times U(1)_{L_{\mu}} \times U(1)_{L_{\tau}}$$

→ The observation of neutrino oscillations as a first sign of LFV beyond the SM!

Are neutrinos Dirac ($|\Delta L| = 0$) or Majorana ($|\Delta L| = 2$) particles?





ell as to improve the for searches of rare field theory approach unique combination of

ent as compared with a

ioNorcompellingoevidence for new particles mediating LFV processes

nation, the sensitivity is

 \rightarrow Strong experimental constraints on the scale Λ for new degrees of freedom

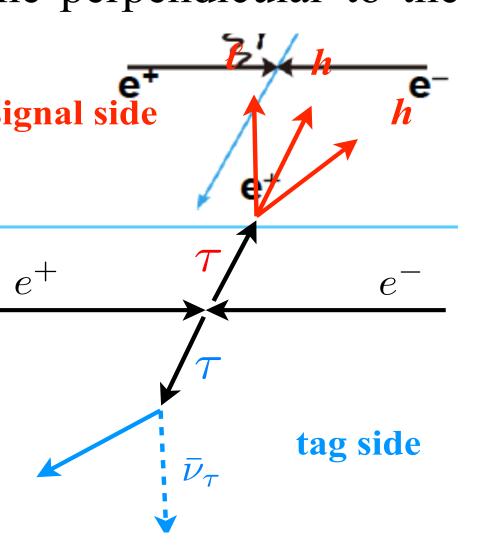
I data Parameterise the LFV τ decays via the effective field theory (EFT)

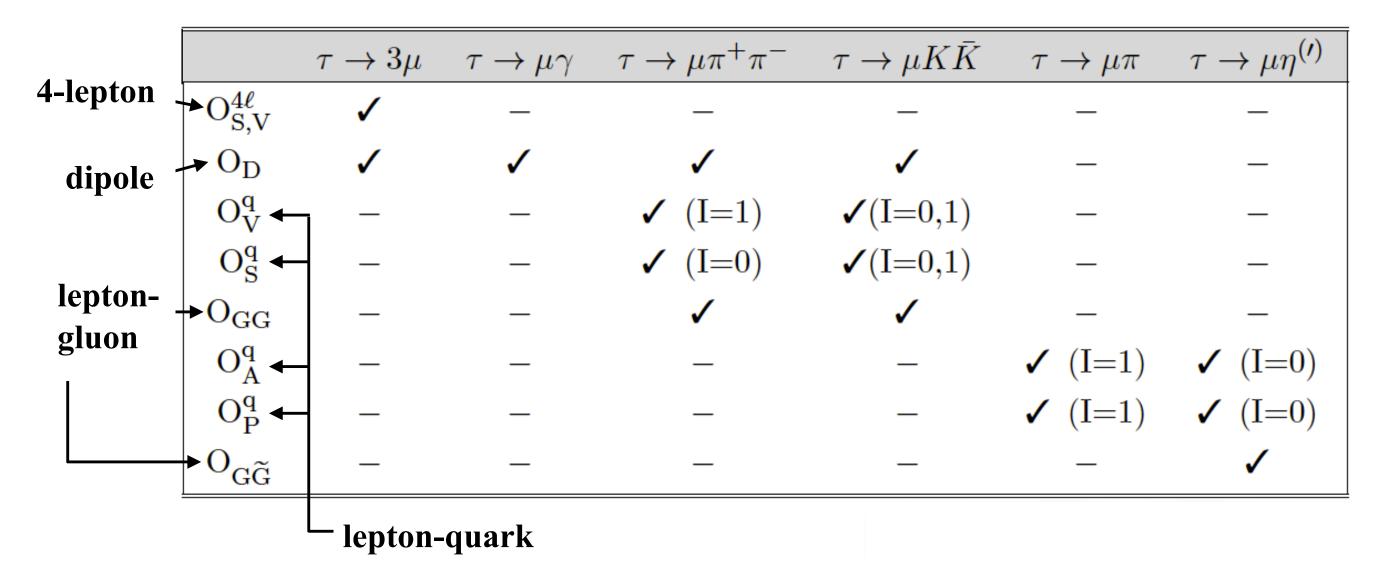
 $L = L_{SM} - C_i^{(5)} + \sum_i \frac{c_i^{(6)}}{\Lambda^2} O_i^{(6)} + \dots$

be made. Their effect will show up at low energies as a series of non-renormalisable operators:

le topology – one (τ_{tag})

as minimal Repode been erates a specific pattern of operators are boosted and their $\mathcal{L} = \mathcal{L}_{SM} + \frac{C^{(5)}}{O^{(5)}} + \sum_{i=1}^{C} \frac{C^{(6)}}{O^{(6)}} + \dots$ ne perpendicular to the hadronic final states, the semi-lepton $\mathbf{L} = \mathbf{L}_{SM} + \frac{C^{(6)}}{O^{(6)}} + \mathbf{L}_{SM} + \frac{C^{(6)}}{O^{(6)}} + \dots$





- Celis, Cirigliano, Passemar (2014) -

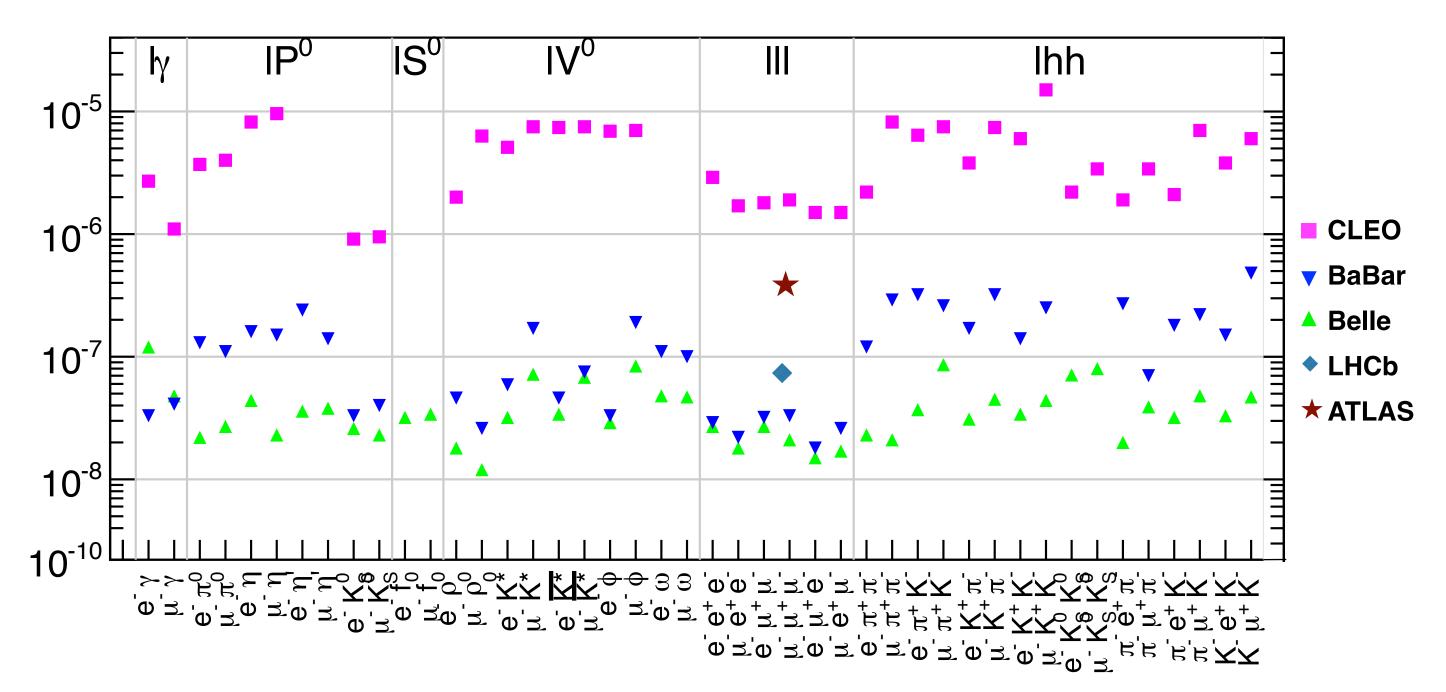
butions from Bhabha, $\mu_{\text{hierogeometric probession}}$ aminations are channel Francesco Tenchini to probe the underlying NP responsible for the LFV.

annel-by-channel. This

42

The progress of t LFV and LNV searches

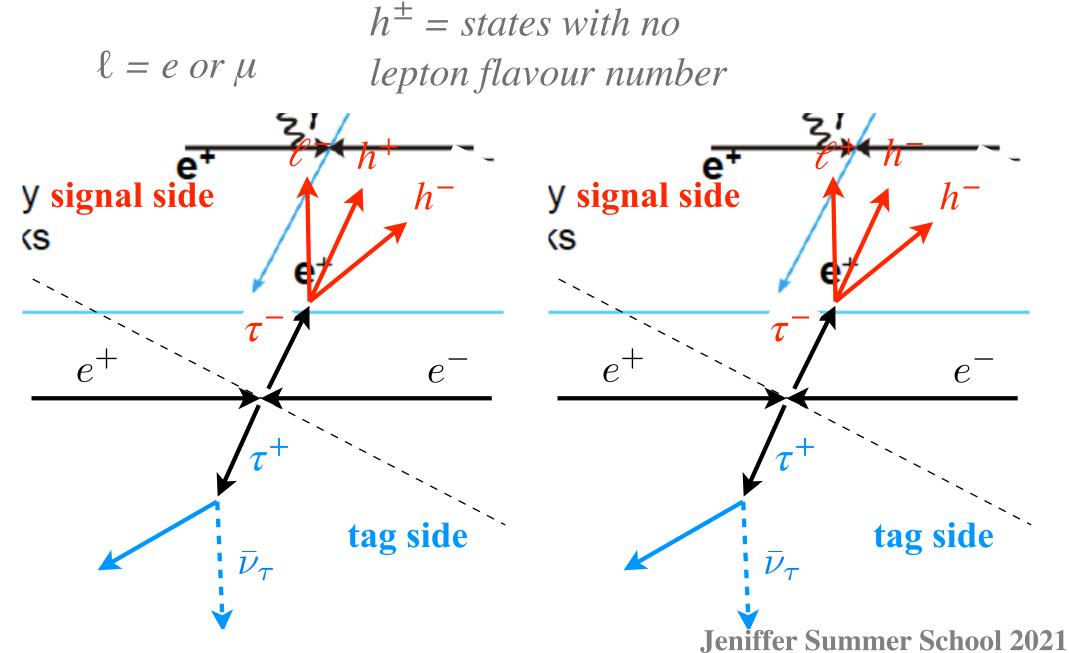
... mostly occurred at the B-factories



The upper limits reached for τ decays approached the regions sensitive to NP.

Test the SM in a variety of ways

- \rightarrow radiative $(\tau \rightarrow \ell \gamma)$
- \rightarrow leptonic decays $(\tau \rightarrow \ell \ell \ell)$
- → a large variety of LFV and LNV semi-leptonic decays
- \rightarrow $\tau \rightarrow \mu$ and $\tau \rightarrow$ e: test of the lepton flavour structure

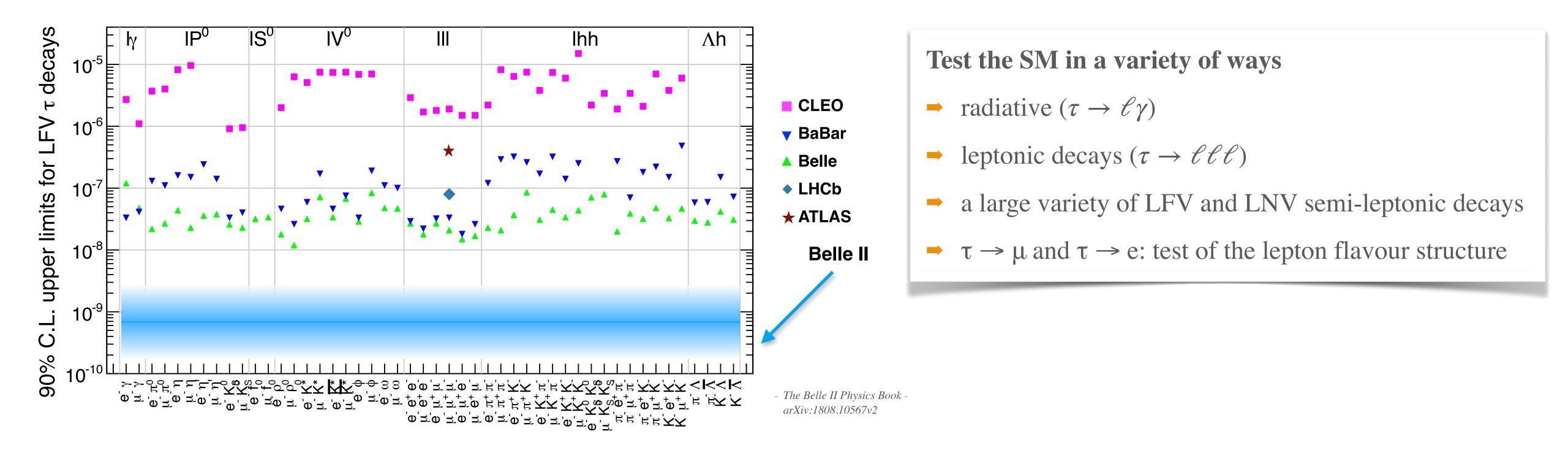




Ami Rostomyan

Perspectives at Belle II

... mostly occurred at the B-factories



- One of the factors pushing up the sensitivity of probes is the increase of the luminosity
- → Equally important is the increase of the signal detection efficiency
 - high trigger efficiencies; improvements in the vertex reconstruction, charged track and neutral-meson reconstructions, particle identification, refinements in the analysis techniques...

The searches at Belle II will push the current bounds further by more than one order of magnitude



$au o \mu \mu \mu$

Signal-background discrimination using kinematics of the event

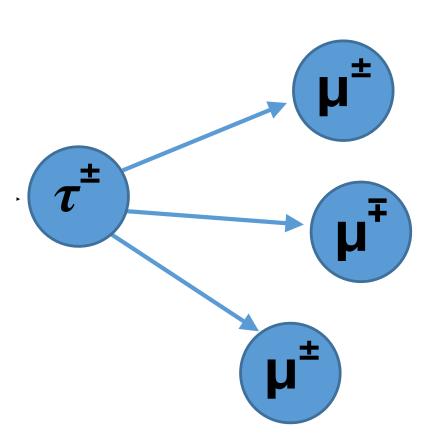
μID - the most powerful discriminating variable

Signal-background discrimination depends on the tag-side track pendent optimisation of the muID requirement

Momentum dependent optimisation of the muID requirement

- $\rightarrow P_{\mu} < 0.7 \text{ GeV}$
 - Leptonic tag

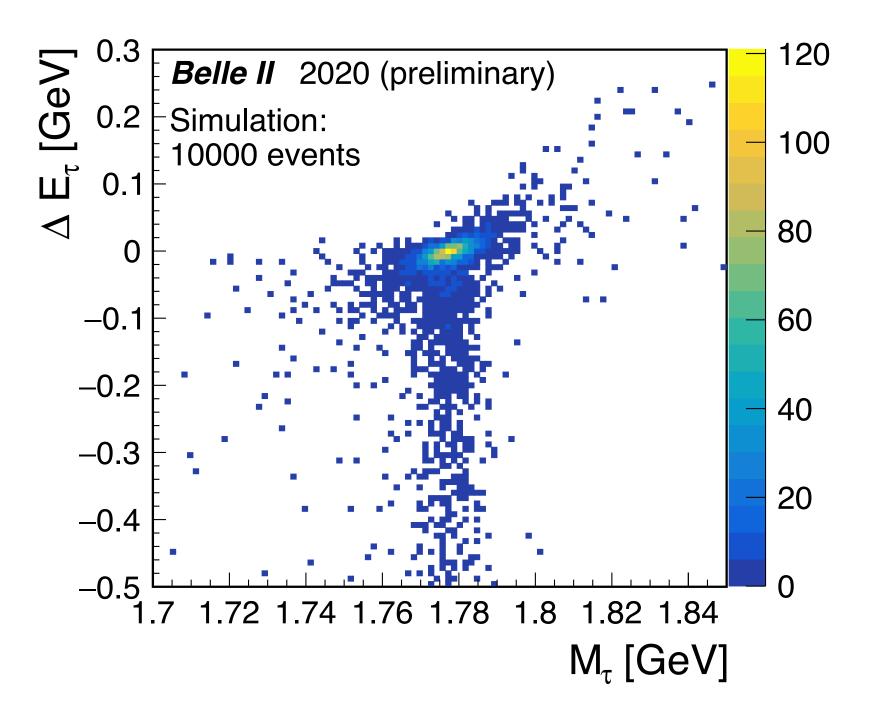
 μ do not reachphon detector (KLM)
- → $0.7 < P_{\mu} < 1 \text{ GeV}$
 - → μ reach KLM but not many layers are crossed
- $\rightarrow P_u > 1 \text{ GeV}$
 - \rightarrow μ reach KLM and many layers are crossed Tag side ν_{τ}



Other requirement used @Belle but not @Belle II:

- μ veto on tag track
- \rightarrow P_{μ}> 0.6 GeV

Higher efficiency is foreseen @Belle II than @Belle or @BaBar



Two independent variables:

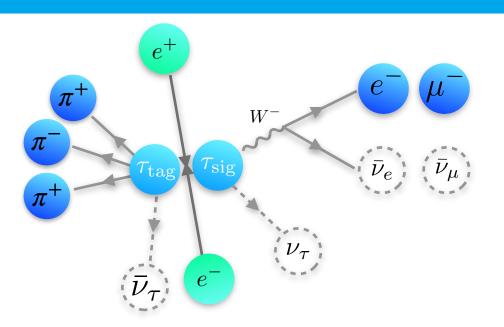
$$M_{\tau} = \sqrt{E_{\mu\mu\mu}^2 - P_{\mu\mu\mu}^2}$$

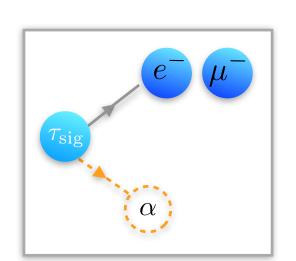
$$\Delta E = E_{\mu\mu\mu}^{CMS} - E_{\text{beam}}^{CMS}$$

 \rightarrow For signal \rightarrow ΔE close to 0 and $M_{\mu\gamma}$ close to τ mass

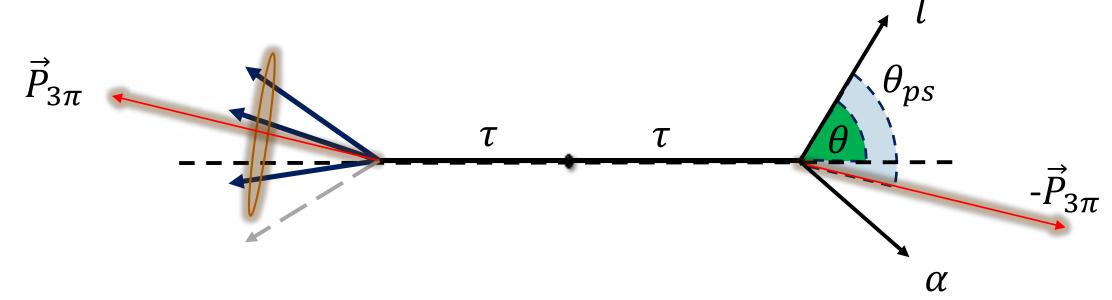
Search for LFV $\tau \rightarrow \ell \alpha$ ($\alpha \rightarrow$ invisible)

Probe the existence of a new boson α





- → previous studied at Mark III (9.4 pb⁻¹) and ARGUS (476 pb⁻¹)
- search for a two body decay spectrum
- \rightarrow signal will manifest itself as a peak in the τ rest frame

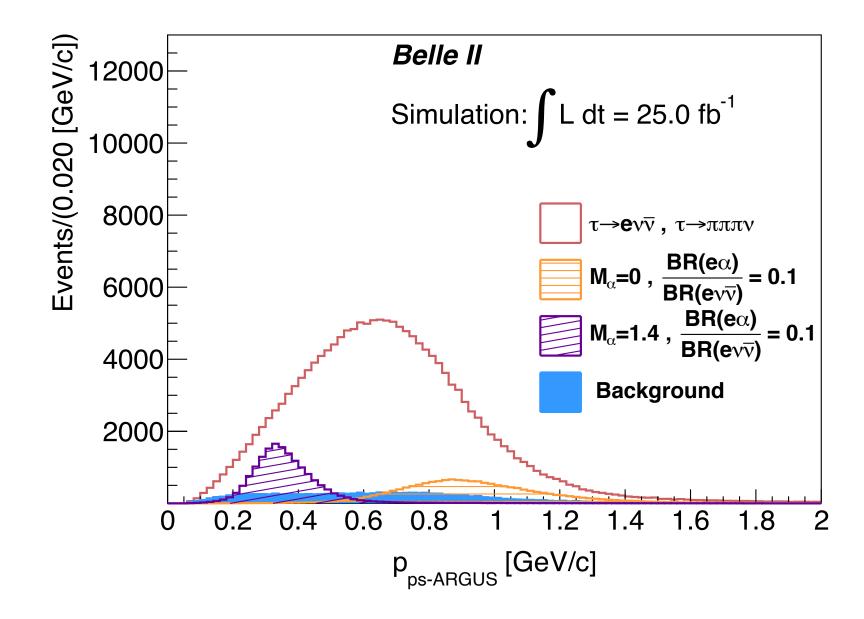


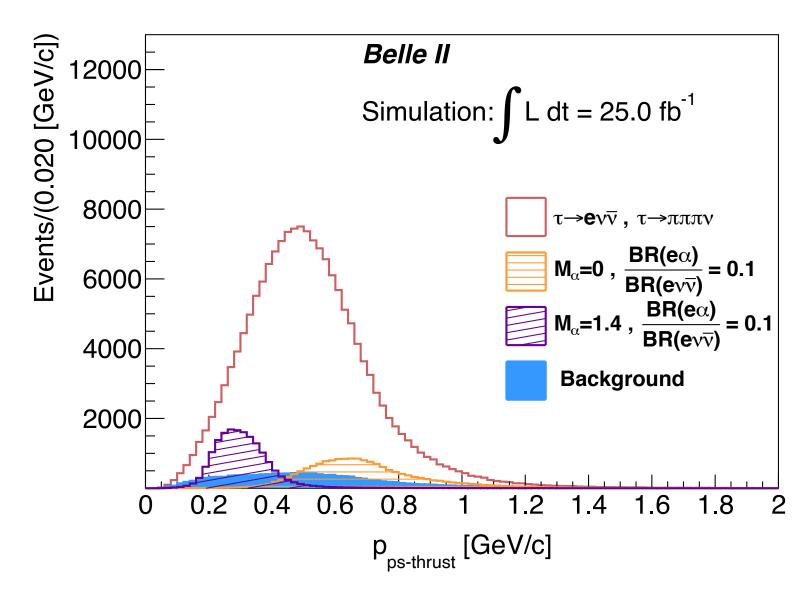
- \rightarrow cannot access the τ rest frame directly due to the missing neutrino
- approximate with the following assumptions:

$$E_{\tau} = \sqrt{s/2}$$

ARGUS method: $\hat{p}_{\tau} \approx -\hat{p}_{3\pi}$

Thrust method: $\hat{p}_{\tau} \approx \hat{T}$





Search for LFV $\tau \rightarrow \ell \alpha$ ($\alpha \rightarrow$ invisible)

UL is provided for the ratio $Br(\tau \to e\alpha)/Br(\tau \to e\nu\nu)$

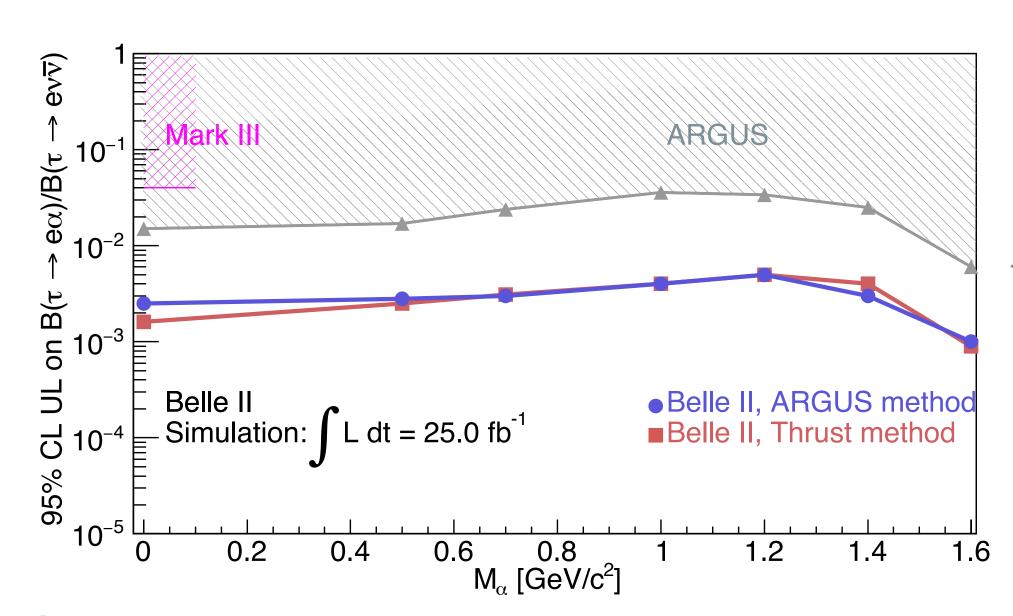
BELLE2-NOTE-PH-2019-009

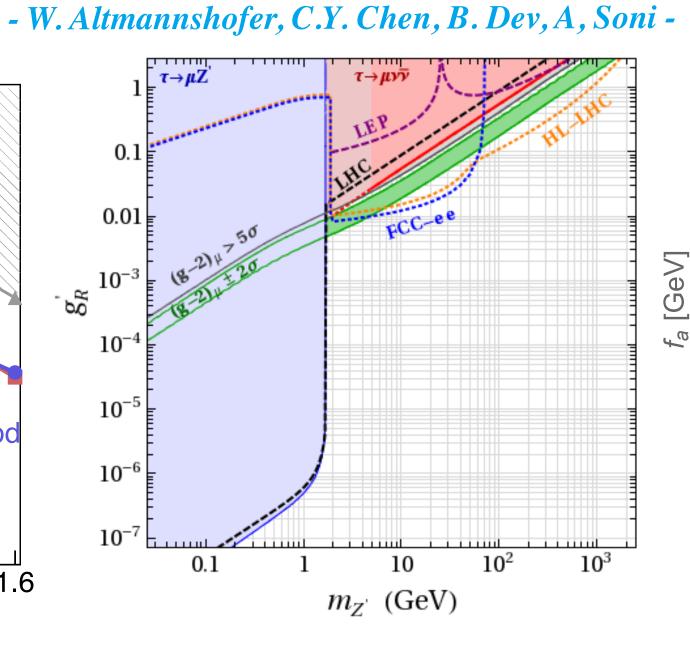
Status of the analysis:

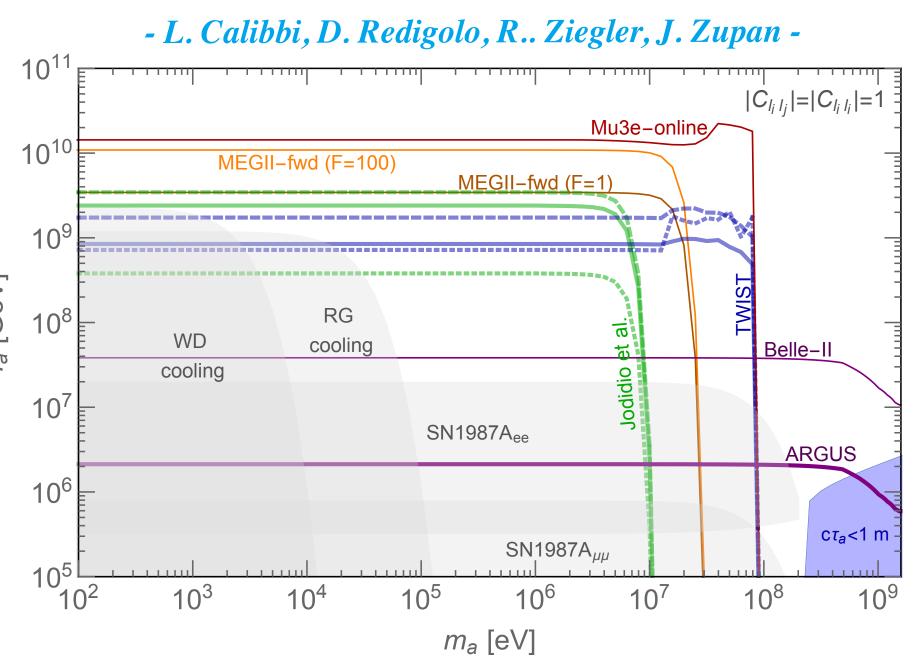
- background suppression already quite effective
- ongoing work to further suppress BG using BDT
- → UL estimation using the frequentist profile-likehood method using asymptotic approach
- alternative test using the Bayesian approach

Various NP scenarios:

- → LFV Z′
- strong bound from ARGUS measurement
- \rightarrow light ALP a
 - exploring regions in parameter space not reachable by other experiments



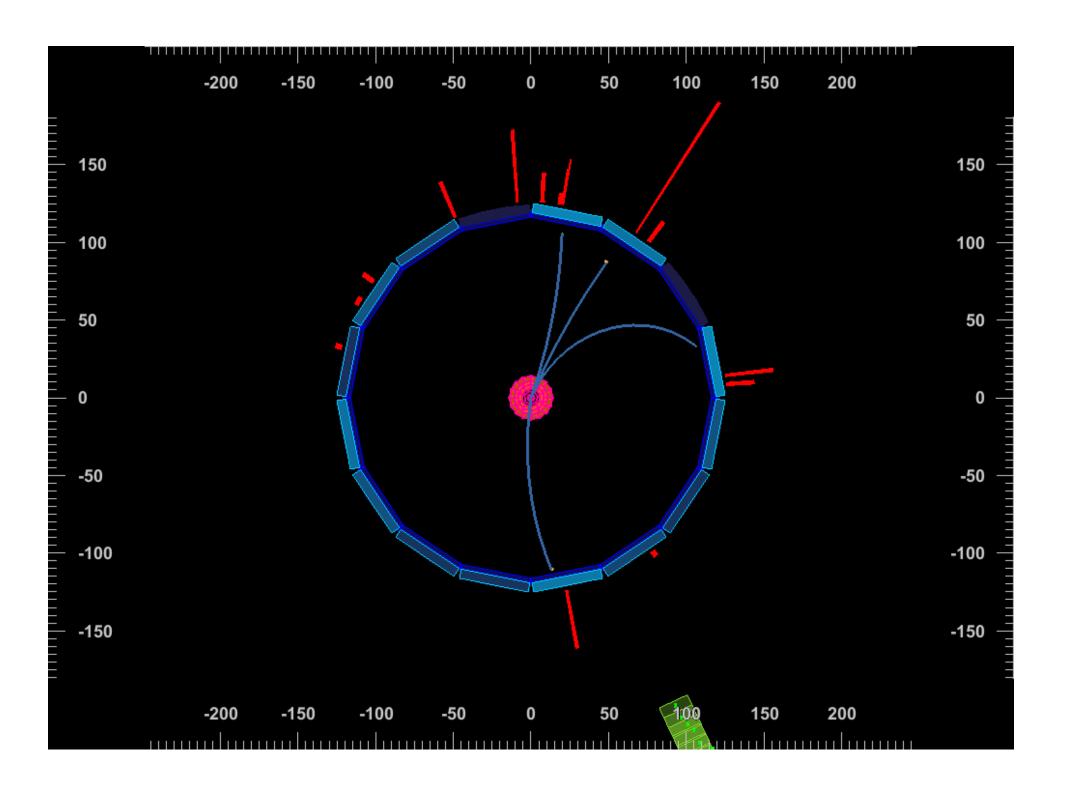




DESY.

Summary

e^+e^- annihilation data is ideal for precision measurements and NP searches!



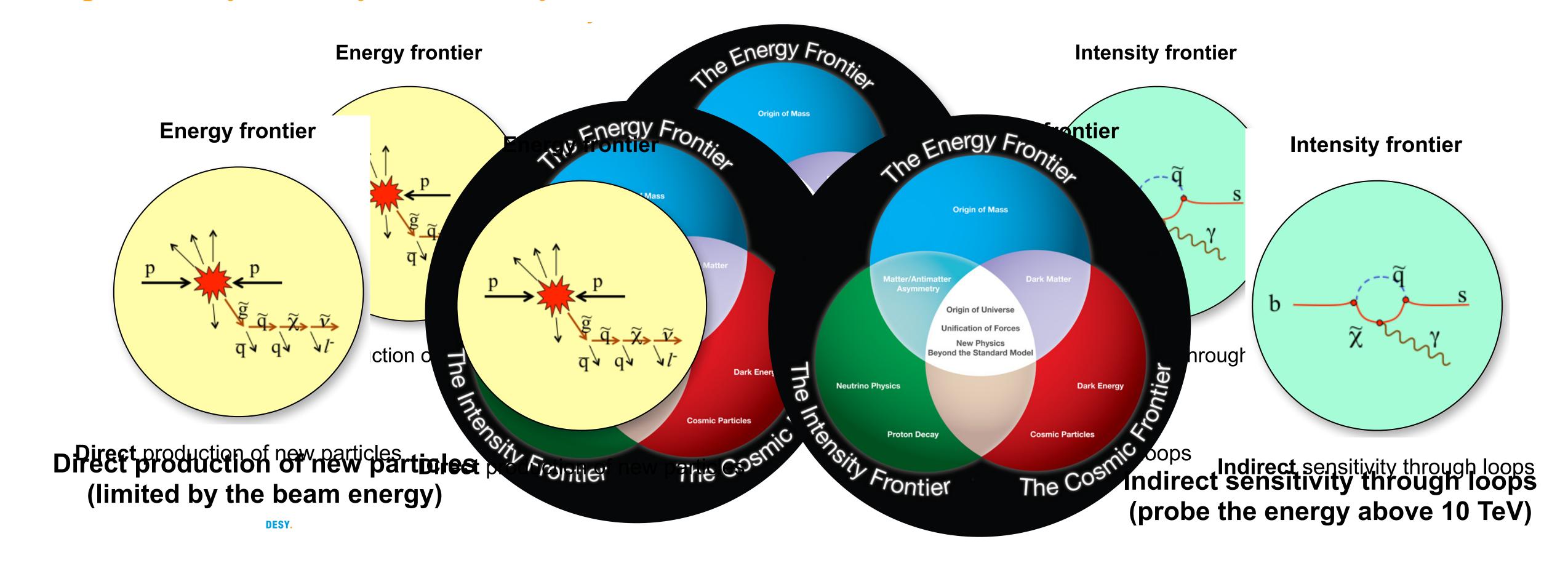
If got interested, join us! Contact me!

→ Belle II experiment started

- → Achieved world record luminosity $L = 2.4 \times 10^{35} \text{cm}^{-2} \text{s}^{-1}$
- → Accelerator tuning is ongoing; more data will be recorded soon
- → **T mass** and **lifetime** measurements with the early data are very promising and show the potential of Belle II precision measurements
- → LFU and V_{us} (exclusive) analysis started
- $\rightarrow \tau \rightarrow \mu\mu\mu$ indicates the potential of LFV searches
- **-**
- Belle II will provide the world largest number (5x1010) of $e^+e^- \to \tau^+\tau^-$ events
 - → t precision measurements and NP searches will reach higher sensitivity w.r.t. the previous experiments

Outlook

Complementary Pathways to New Physics



If NP is seen by one frontier, the confirmation by the other would be important!



Jeniffer Summer School 2021

DESY.