

# Looking for New Physics in B decays at Belle II

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# Introduction

A variety of BSM scenarios have interesting implications for B-physics.

Here I will focus on **BSM models for the B anomalies**,

## $b \rightarrow s$ anomalies

$\mu$  vs  $e$  universality in  $b \rightarrow sl\ell$   
+ ang. obs. and rates in  $b \rightarrow s\mu\mu$   
 $\sim 4\sigma$

## $b \rightarrow c$ anomalies

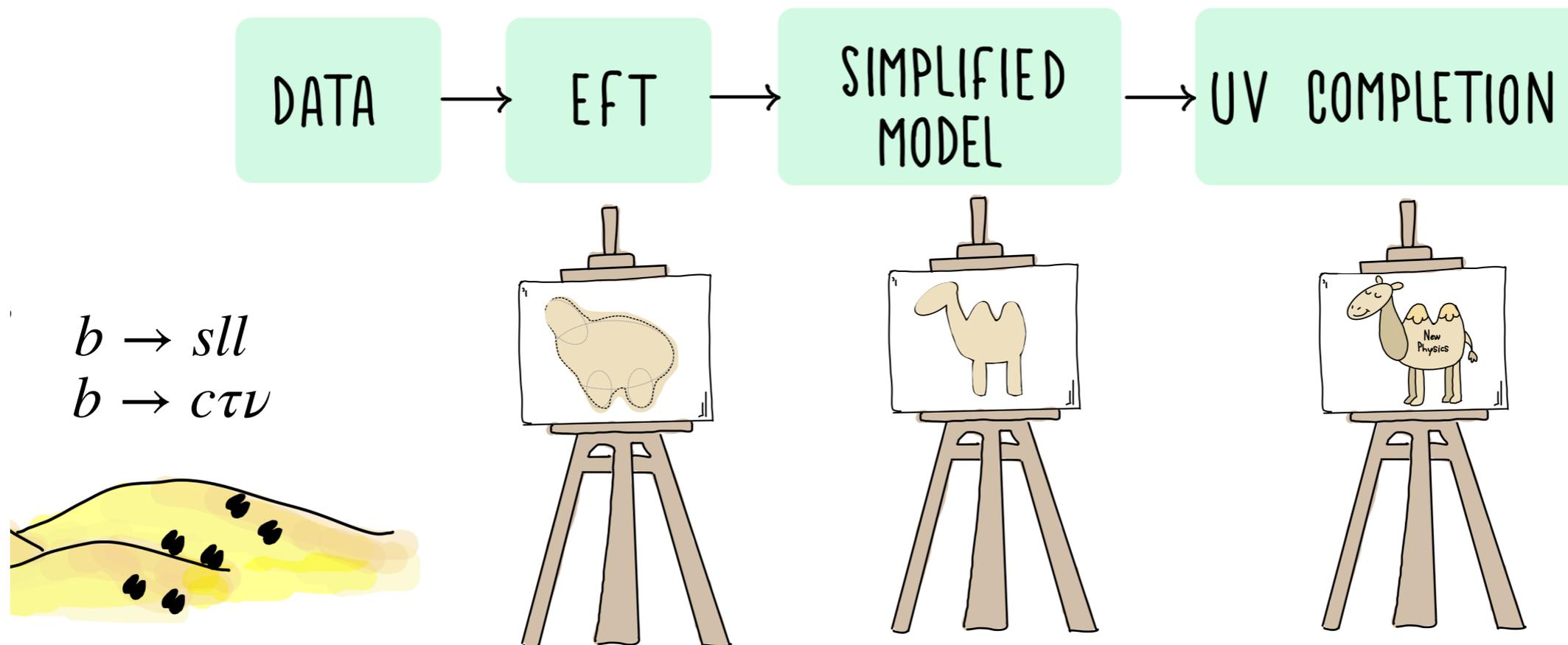
$\tau$  vs  $\mu, e$  universality in  $b \rightarrow cl\nu$   
 $\sim 3\sigma$

at present the only direct experimental motivation to expect NP in (other) B decays.

# Reading the footprints of the B anomalies

? If the anomalies in  $b \rightarrow sll$  and  $b \rightarrow c\tau\nu$  are true NP signals, which NP could be responsible for them, and where else should we see it?

Usual strategy:



- at each step, we can investigate connections with other observables
- strength of connections becomes more model-dependent going from left to right)

# EFT lessons

Minimal assumptions: B anomalies are real, and are due to some heavy NP.

Under these assumptions, the right framework to parametrize NP contributions to observables is an **Effective Theory**.

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The first place to look for similar effects are **observables in the same partonic transition** (ratios, BRs, angular distributions...):

## $b \rightarrow s$ anomalies

- $B \rightarrow X_s \mu \mu$
- $B_s \rightarrow \mu \mu$
- $R_{pK}^{\mu/e}$

## $b \rightarrow c$ anomalies

- Ang. dist. of  $B \rightarrow D^{(*)} \tau \nu$
- $F_L(D^*), P_\tau(D^*)$
- $R_{D^+}, R_{D^*}, R_{D^{*+}}, R_{D_s^*}, R_{\Lambda_c} \dots$

On top of assessing the “size” (scale) of the NP effect, these observables help us pin down its Lorentz structure.

# Observables in $b \rightarrow c\tau\nu$

$$\mathcal{L}_{\text{eff}} = -2\sqrt{2}G_F V_{cb} \left[ (1+g_{V_L})(\bar{c}_L\gamma^\mu b_L)(\bar{\tau}_L\gamma_\mu\nu_L) + g_{V_R}(\bar{c}_R\gamma^\mu b_R)(\bar{\tau}_L\gamma_\mu\nu_L) + g_{S_R}(\bar{c}_L b_R)(\bar{\tau}_R\nu_L) \right. \\ \left. + g_{S_L}(\bar{c}_R b_L)(\bar{\tau}_R\nu_L) + g_T(\bar{c}_R\sigma^{\mu\nu} b_L)(\bar{\tau}_R\sigma_{\mu\nu}\nu_L) \right]$$

Switching on only one WC at a time, only  $g_{V_L}$  is able to account for all  $b \rightarrow c$  data.  
But different solutions, e.g.  $V_L + S_R, S_L + T\dots$  are still possible.

Looking at different obs. in  $b \rightarrow c\tau\nu$  can help us disentangle these scenarios.

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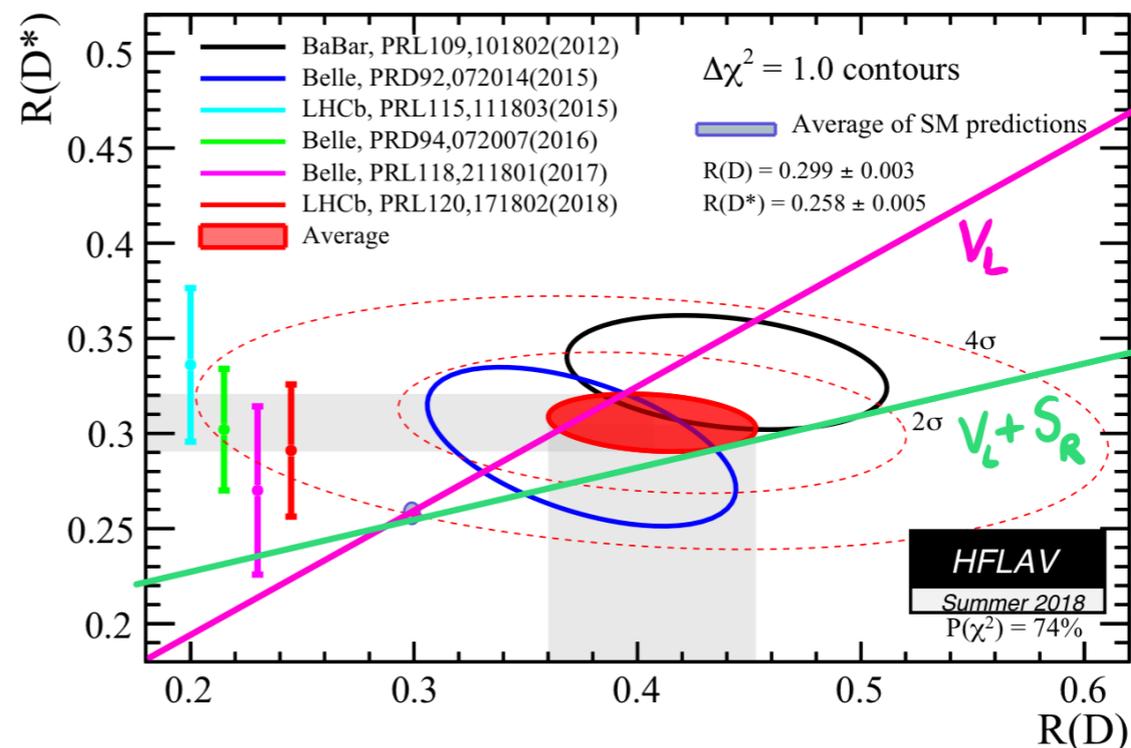
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## LFU ratios

$$V_L \text{ only} \quad \Delta R_D = \Delta R_D^*$$

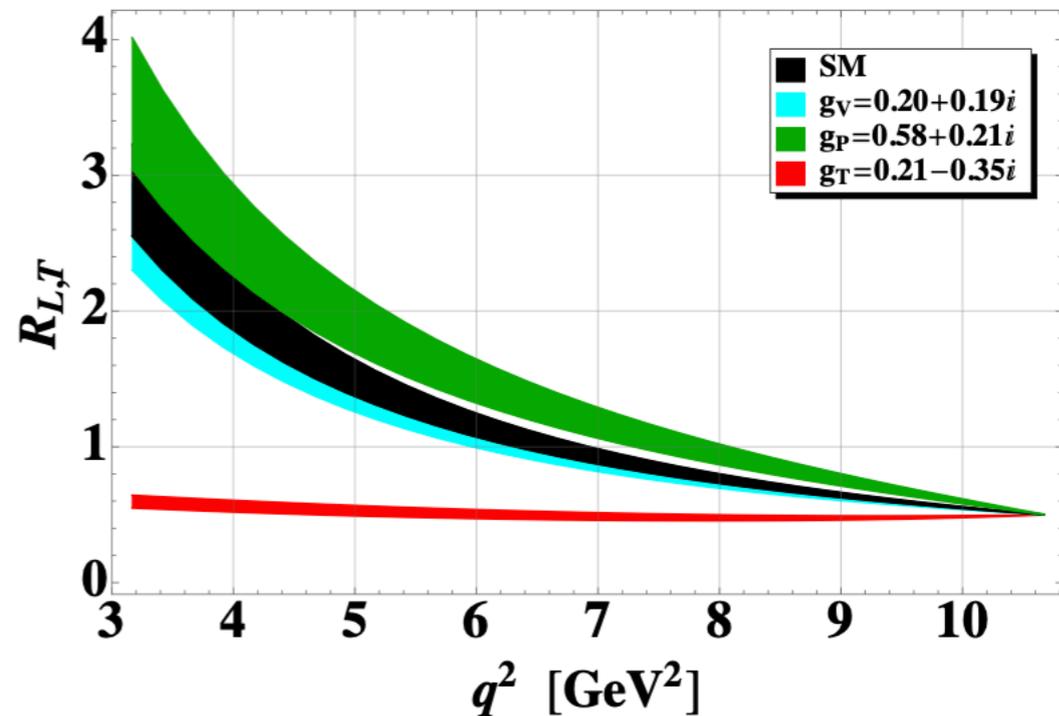
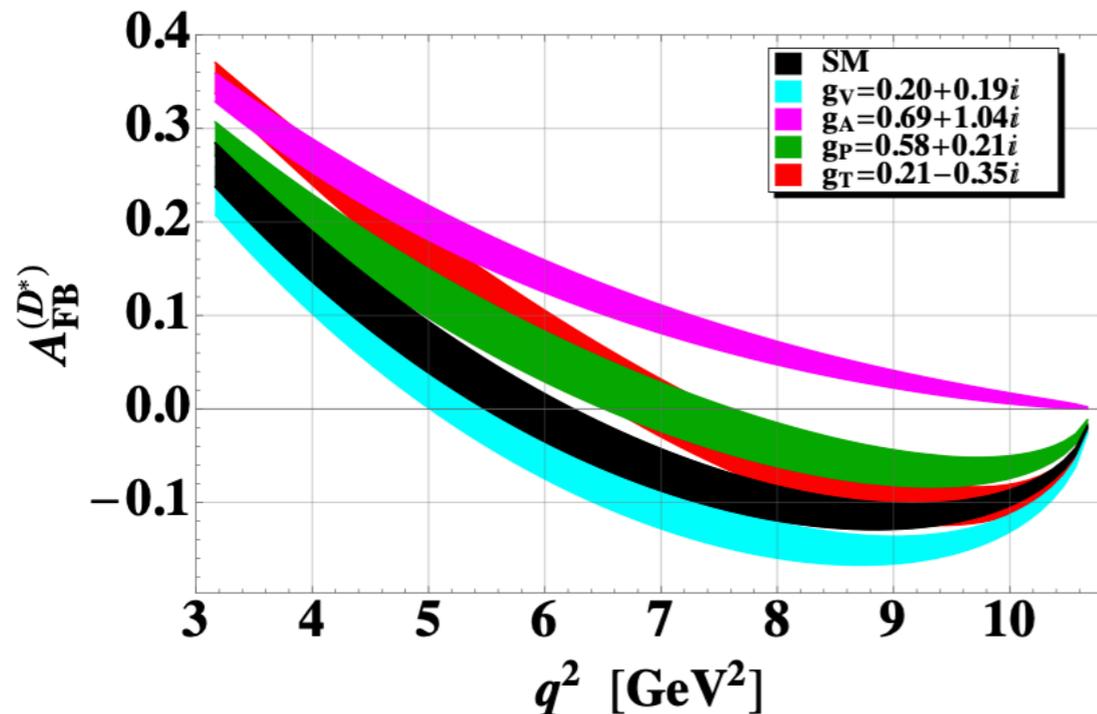
$$V_L + S_R \quad \Delta R_D > \Delta R_D^*$$

$$\Delta R_i \equiv \frac{R_i}{R_i^{\text{SM}}} - 1$$



# Observables in $b \rightarrow c\tau\nu$

Ang. observables, polarizations...



See also [Murgui et. al \(2019\)](#) for a similar analysis

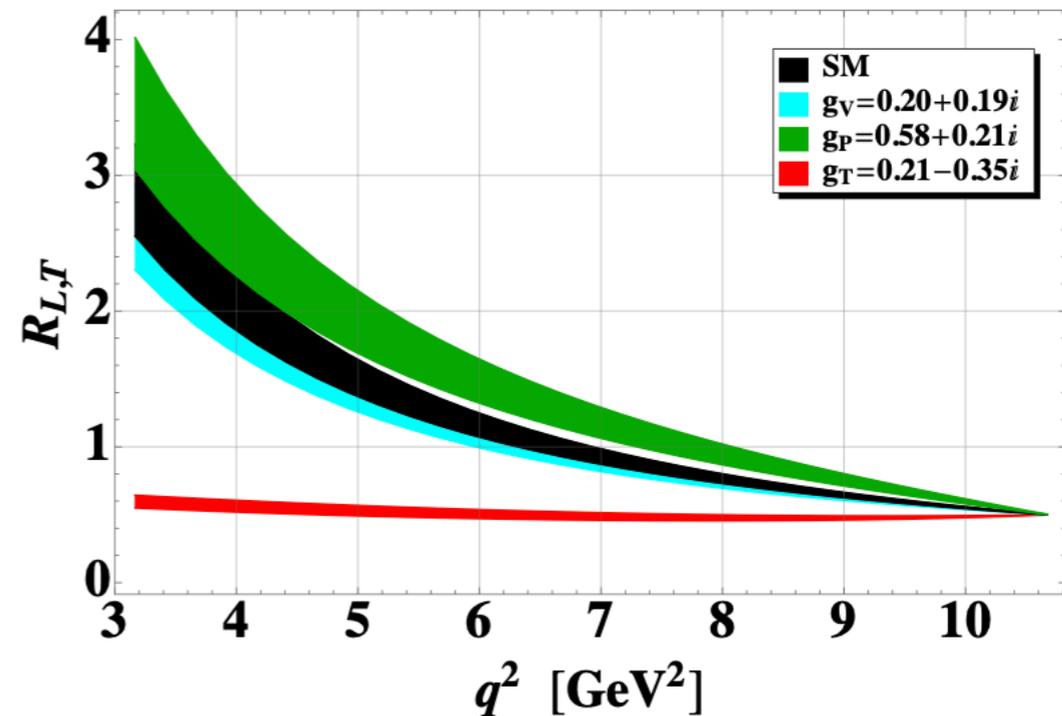
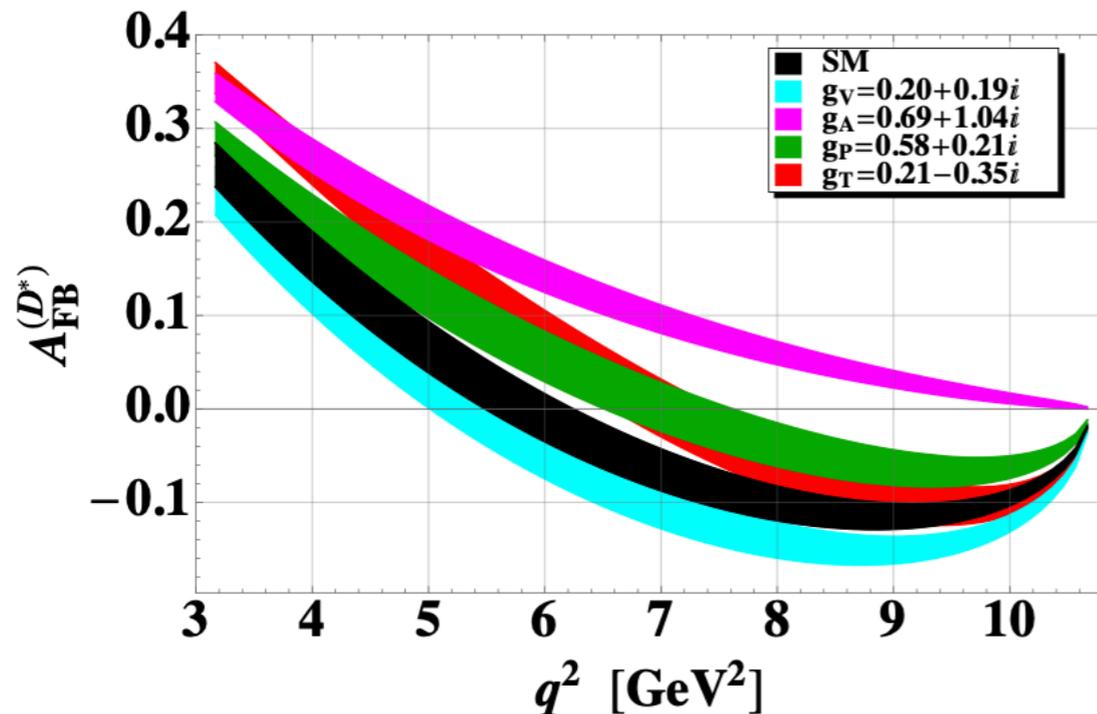
More precise measurements of ratios, angular correlations, polarizations, and asymmetries are crucial to disentangle these scenarios.

Many of these are more easily accessible at Belle II w.r.t. LHCb.

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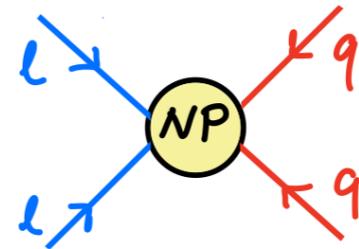
# Combined explanation of the B anomalies

The two anomalies fit well in the SMEFT:

$$\begin{array}{ccc}
 b \rightarrow sll & & b \rightarrow cl\nu \\
 (\bar{s}_L \gamma^\mu b_L)(\bar{l}_L \gamma_\mu l_L) & \xleftrightarrow{SU(2)_L} & (\bar{c}_L \gamma^\mu b_L)(\bar{\nu}_L \gamma_\mu \nu_L)
 \end{array}$$

⇒ Minimal solution: **left-handed, TeV scale NP in semi-leptonic operators:**

$$\mathcal{L}_{\text{EFT}}^{\text{NP}} = -\frac{1}{v^2} \left( C_{lq}^{(3)} (\bar{l}_L \gamma^\mu \tau^a l_L)(\bar{q}_L \gamma^\mu \tau^a q_L) + C_{lq}^{(1)} (\bar{l}_L \gamma^\mu l_L)(\bar{q}_L \gamma^\mu q_L) \right) + \dots$$



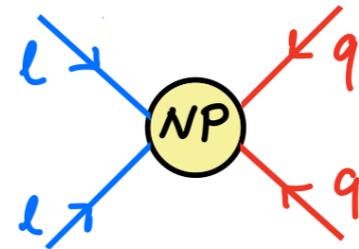
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The only viable tree-level mediators are **leptoquarks**:

no  $4\ell$  and  $4q$  processes at tree level, and no resonant production at LHC.

Three possibilities (for a combined explanation):

$S_1 + S_3$  [[Crivellin et al. \(2017\)](#); [Buttazzo et al. \(2017\)](#); [Marzocca \(2018\)](#)...]

$S_3 + R_2$  [[Bečirević et al. \(2018, 2022\)](#)]

$U_1$  [[di Luzio et al. \(2017\)](#); [Calibbi et al. \(2017\)](#); [Bordone, CC, et al. \(2017\)](#); [Barbieri, Tesi \(2017\)](#); [Heck, Teresi \(2018\)](#)...]

# General phenomenological consequences

1. Large  $b \rightarrow s\tau\tau$
2. Large  $\tau/\mu$  violation in B and  $\tau$  decays
3. Enhancement of  $B \rightarrow K\nu\bar{\nu}$

Belle II plays an important role in assessing all these effects.

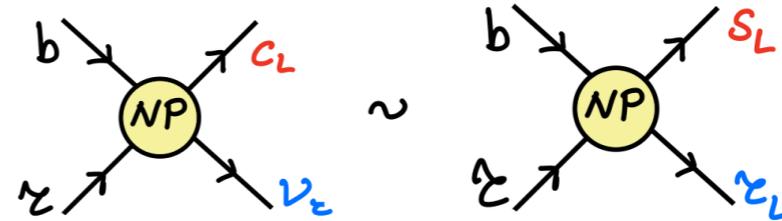
# Large $b \rightarrow s\tau\tau$

[(\*)Exception: R2 + S3]

Driven by the CC anomaly (when explained via the triplet(\*))

$$l_L^3 = \begin{bmatrix} \nu_e \\ \tau_L \end{bmatrix} \quad q_L^2 = \begin{bmatrix} c_L \\ s_L \end{bmatrix}$$

$SU(2)_L$  invariance  $\Rightarrow$



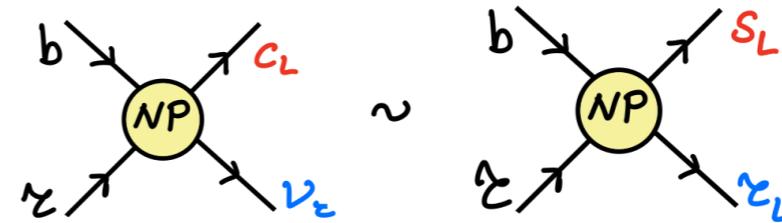
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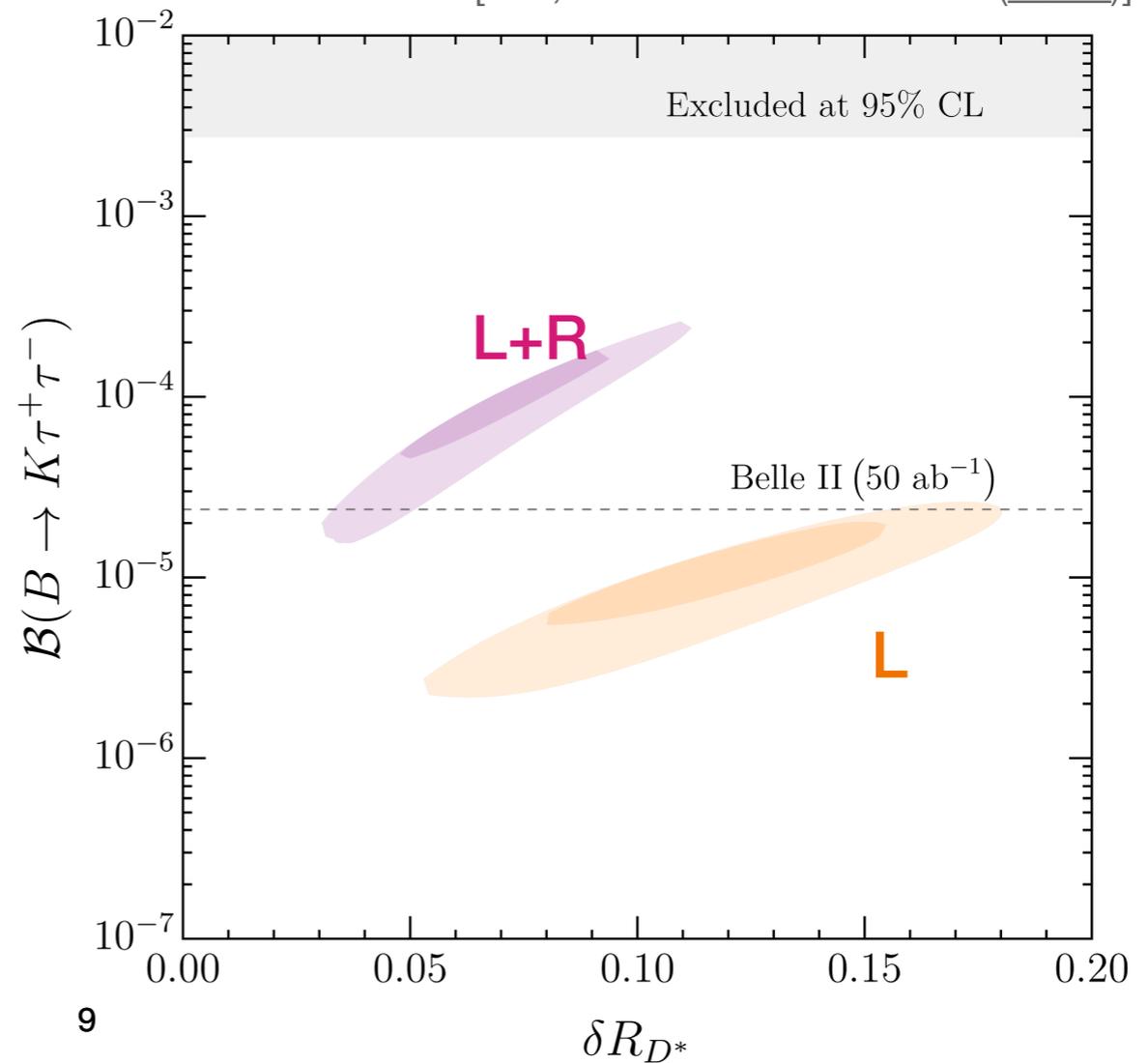
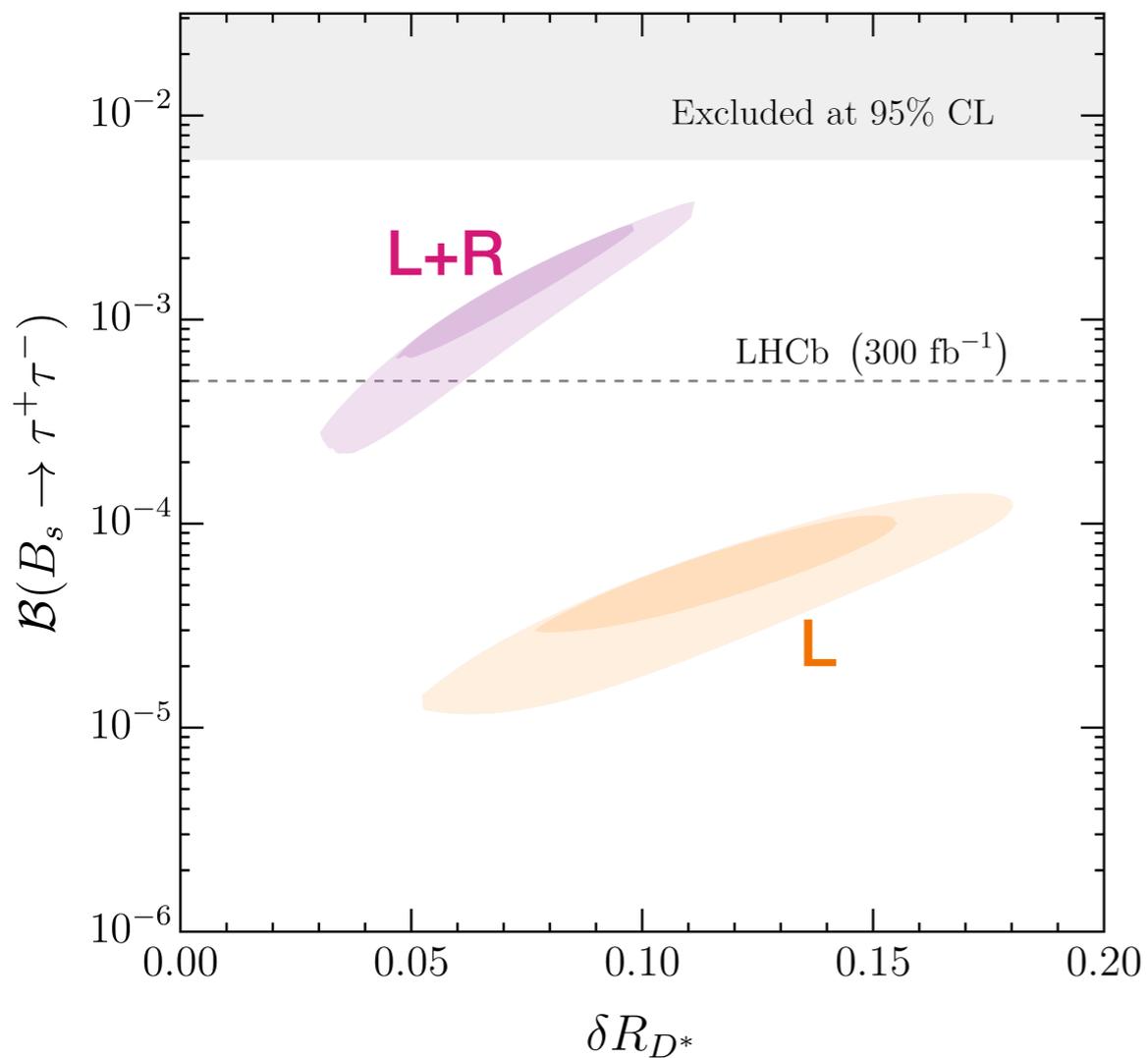
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Projections for the  $U_1$ :

[CC, Fuentes Martin et al. (2021)]



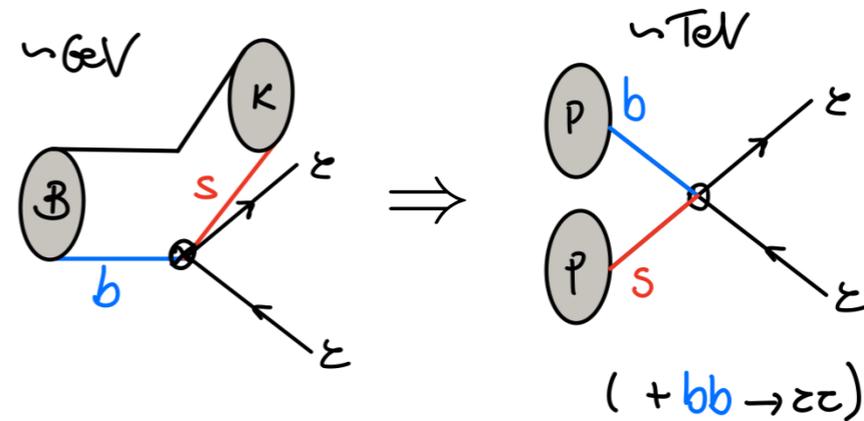
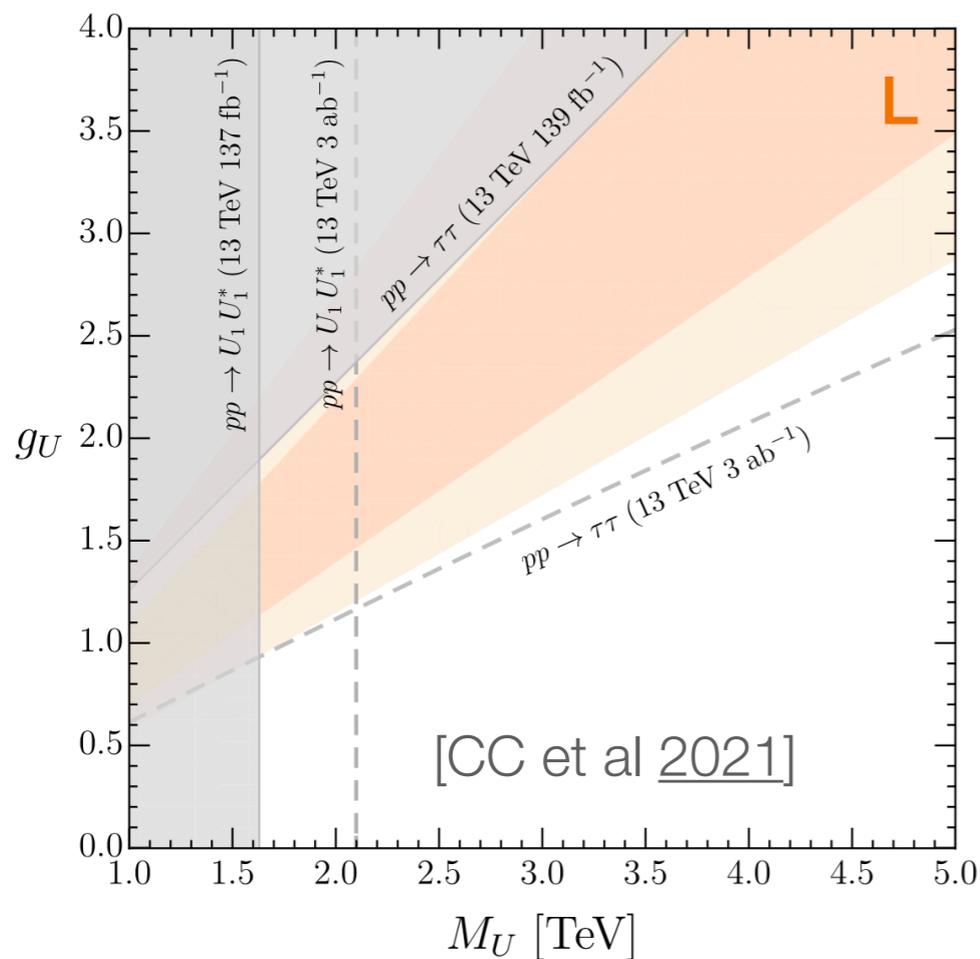
# High-pT bounds from $pp \rightarrow \tau\tau$

[Faroughy, Greljio, Kamenik (2016);  
Fuentes-Martin et al. (2020)]

The same interaction can be probed in **di-tau tails** at the LHC.

The obtained bounds are generally stronger than the low-energy ones.

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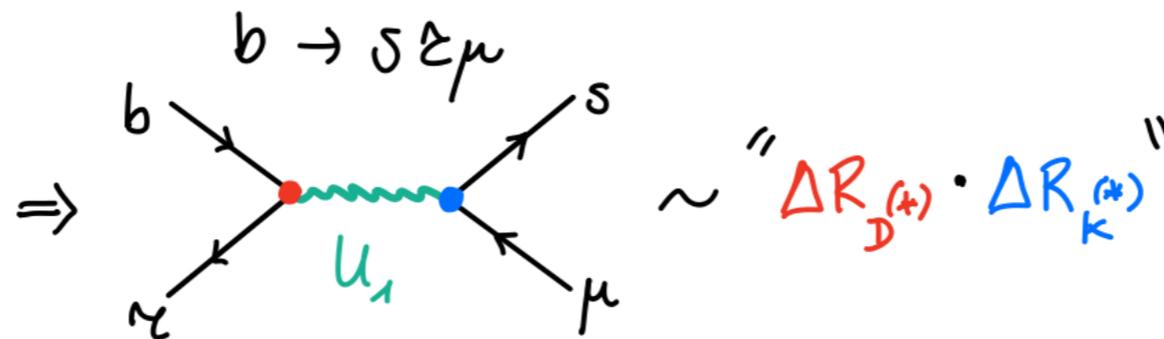
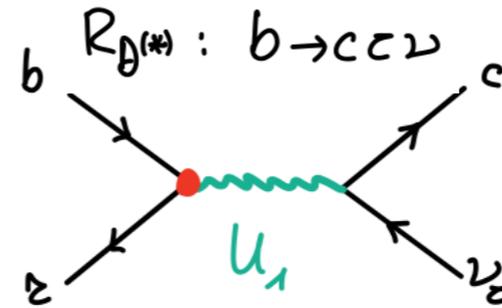
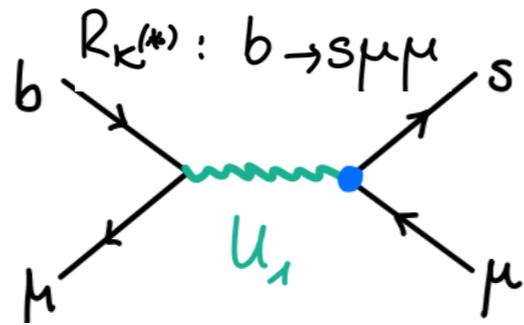
- $U_1$  solution is completely falsifiable at HL-LHC (or we will find a  $U_1$ !)
- same for  $R_2 + S_3$ ,
- still space left for  $S_1 + S_3$

Models for  $R_D^{(*)}$  only yield similar enhancements in  $B \rightarrow K\tau\tau$ ,  $B_s \rightarrow \tau\tau$  and  $pp \rightarrow \tau\tau$ .

**Any change in  $R_D^{(*)}$  will alter these conclusions significantly.**

# Lepton Flavour Violation in $b \rightarrow s\tau\mu$ and $\tau$ decays

Driven by the presence of both CC and NC anomaly:



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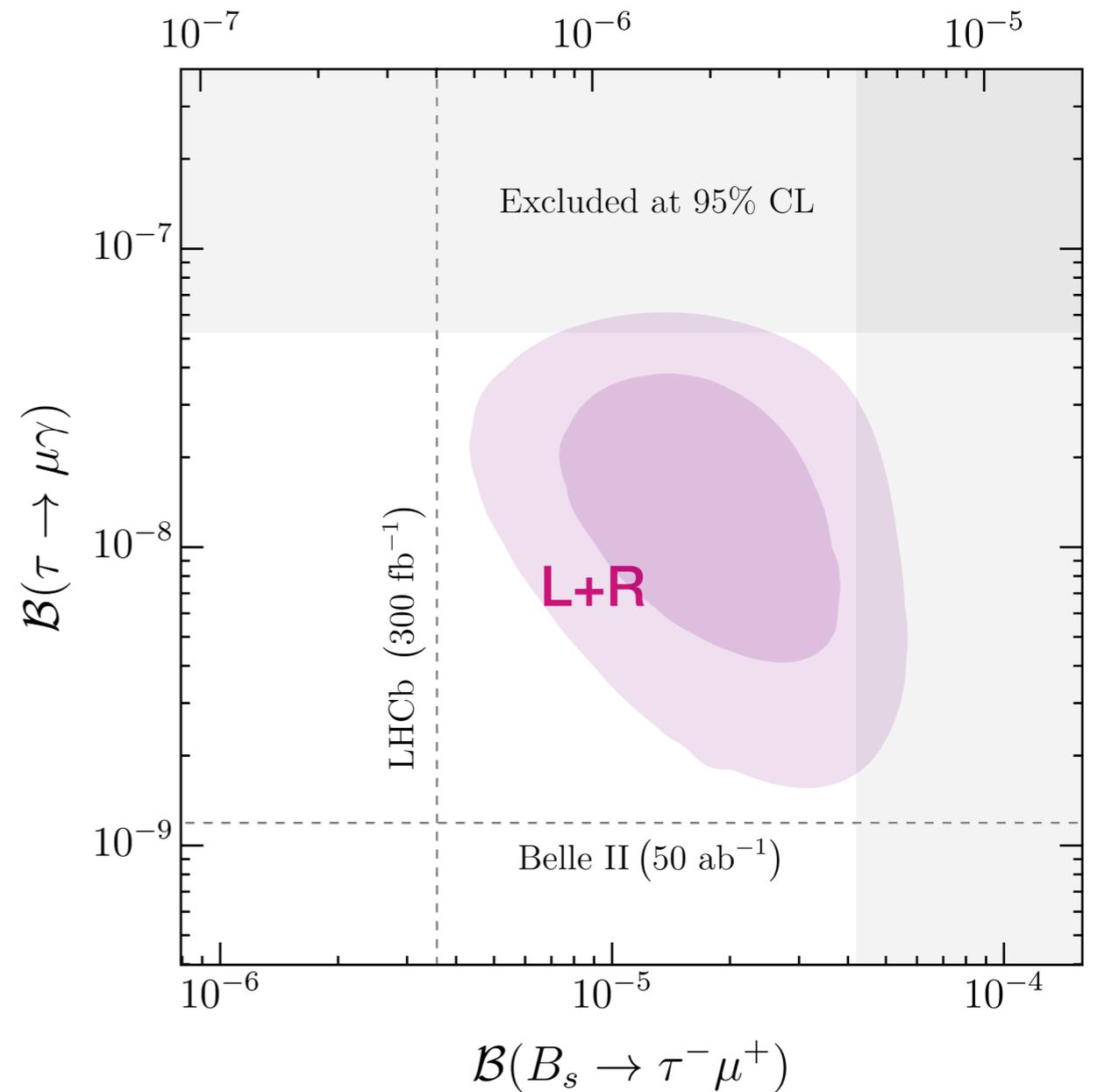
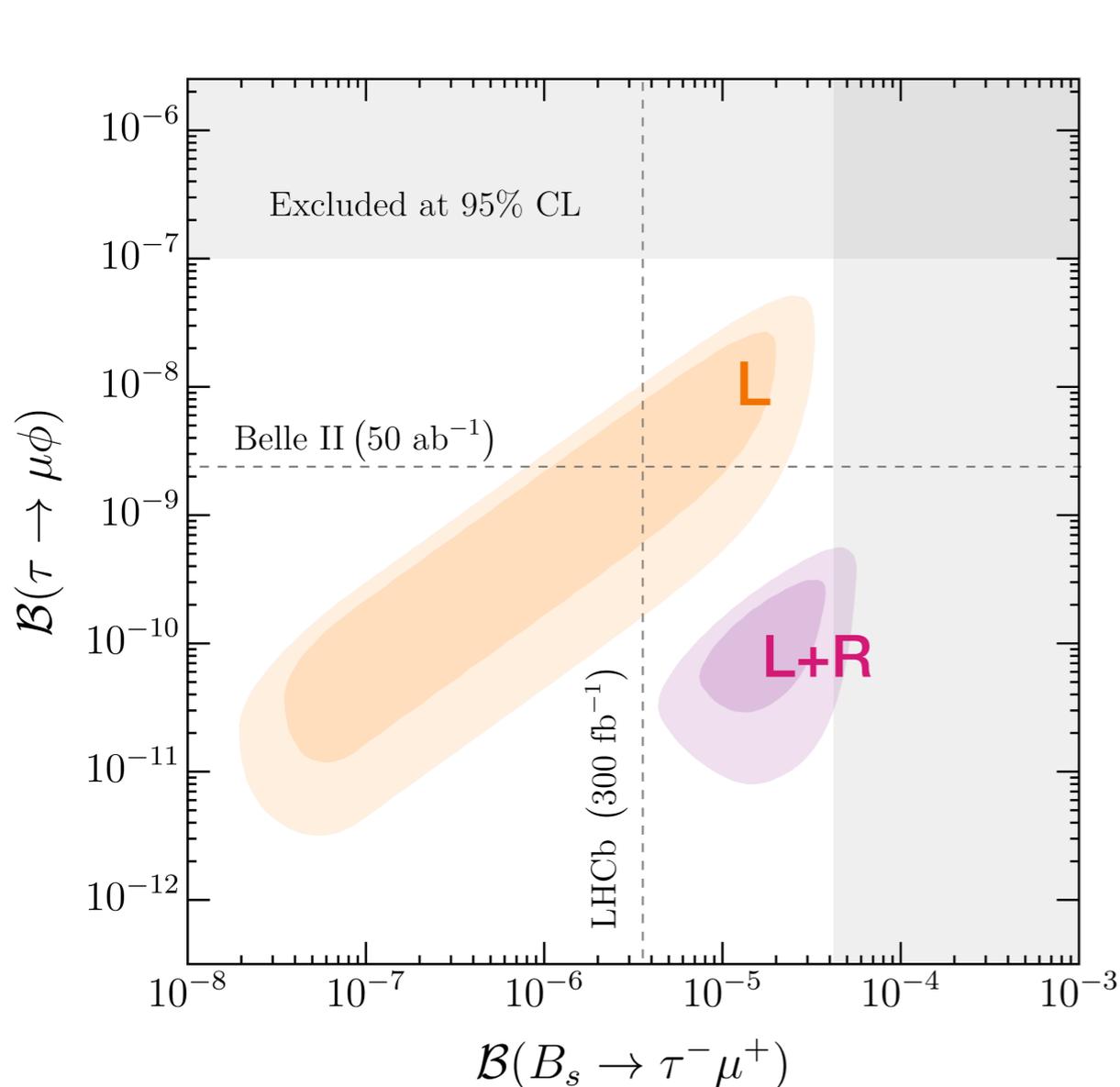
$$\mathcal{B}(B_s \rightarrow \tau\mu) \approx \mathcal{B}(B \rightarrow K\tau\mu) \approx 10^{-7} - 10^{-6}$$

$$\mathcal{B}(\tau \rightarrow \mu\phi) \approx 10^{-10} - 10^{-8}$$

$$\mathcal{B}(B_s \rightarrow \tau\mu) \approx 1 \times 10^{-5}$$

$$\mathcal{B}(B \rightarrow K\tau\mu) \approx 1 \times 10^{-6}$$

$$\mathcal{B}(\tau \rightarrow \mu\gamma) \approx 1 \times 10^{-8}$$



[CC, Fuentes Martin, Faroughy Isidori, Neubert, 2021]

# Enhancement of $B \rightarrow K\nu\nu$

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$$\mathcal{A}(b \rightarrow s\nu\bar{\nu}) \propto \left( \underbrace{C_{lq}^{(3)}}_{\text{does } R_{D^{(*)}} \text{ (} \rightarrow \text{ large)}} - C_{lq}^{(1)} \right)_{2333} \Rightarrow \text{need } C_{lq}^{(3)} \approx C_{lq}^{(1)} \text{ at 1\% level not to overshoot present bounds!}$$

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At tree level, this is achieved automatically for the  $U_1$  and can be imposed for  $S_1 + S_3$ , .....but is anyway spoiled by radiative effects, resulting in a **20-50% enhancement** over the SM. [Fuentes-Martin et al [2020](#) ,[2021](#) , Gherardi et al [2008](#)]

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→ **Belle II will probe the parameter space preferred by all these models entirely.**

Usual caveat: the size of the effect is driven by the CC anomaly.

Other interesting tests:  $b \rightarrow d$  and  $b \rightarrow u$

Belle II has the potential to test LFU in other quark transitions, like  $b \rightarrow u$  and  $b \rightarrow d$ .

A priori no obvious connection between these and  $b \rightarrow c$  and  $b \rightarrow s$ .

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$$\frac{b \rightarrow c\ell\nu}{b \rightarrow u\ell\nu} = \frac{b \rightarrow c\ell\nu}{b \rightarrow u\ell\nu} \Bigg|_{\text{SM}}$$

$$\frac{b \rightarrow s\ell\ell}{b \rightarrow d\ell\ell} = \frac{b \rightarrow s\ell\ell}{b \rightarrow d\ell\ell} \Bigg|_{\text{SM}}$$

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**Universality tests in  $B \rightarrow \pi$ :**

$$\frac{\mathcal{B}(B \rightarrow \pi\tau\nu)}{\mathcal{B}(B \rightarrow \pi\ell\nu)} \approx \frac{\mathcal{B}(B \rightarrow \pi\tau\nu)}{\mathcal{B}(B \rightarrow \pi\ell\nu)} \Bigg|_{\text{SM}} \left( 0.75 \frac{R_D}{R_D^{\text{SM}}} + 0.25 \frac{R_{D^*}}{R_{D^*}^{\text{SM}}} \right)$$

$$\frac{\mathcal{B}(B \rightarrow \pi\mu\mu)}{\mathcal{B}(B \rightarrow \pi ee)} \approx R_{K^{(*)}}$$

# Conclusions

BSM models for B anomalies predict a variety of signatures relevant for Belle II:

- Modifications of obs. In  $b \rightarrow sl\ell$  and  $b \rightarrow c\tau\nu$
- Enhancement of  $B \rightarrow K\tau\tau$
- $B \rightarrow K\tau\mu, \tau \rightarrow \mu\phi, \tau \rightarrow \mu\gamma$
- Enhancement of  $B \rightarrow K\nu\bar{\nu}$
- Possibly effects in  $b \rightarrow u$  and  $b \rightarrow d$  transitions