## **Particle Identification**



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US Belle II Summer School Iowa State University 1 August 2022

Cerenkov radiation

- Belle II iTOP
- LHCb RICH detectors
- dE/dx and muons



### The Cerenkov effect:

- When a charged particle traverses a medium with a velocity exceeding that of light in that medium, photons are radiated. The condition is v > c/n, or  $\beta > 1/n$ . the greater the index of refraction *n*, the lower the threshold velocity  $\beta$ .
- the photons are emitted at an angle:

60

50

40

30

20

10

 $p(\pi) = 0.08$ 

0.5

CHERENKOV ANGLE @ (DEG)

$$\cos heta_c = rac{1}{neta}$$

0.9

0.29

0.8

RELATIVE VELOCITY  $\beta = v/c$ 



velocity

 $\Delta \theta$ 

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0.7

0.14

0.5

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 $\beta = 0.995$ 

1.4 GeV



The light can be collected in several ways:

### Measuring light yield:



### Measuring Cerenkov angle (RICH):



By measuring either light yield or  $\theta_c$ , one measures  $\beta$ , the velocity of the charged particle. Given the momentum of the particle, i.e., as measured in a spectrometer (tracking detector in a magnetic field), one determines the mass (the particle identity):  $m = p/(\gamma\beta)$ 

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Charged particles must be above the 1.0 Cerenkov threshold (1/n) to be detected: 99.9 0.8 GAS-CERENKOV EFFICIENCY FREON 114 (n=1,0014) Freon 114: n = 1.0014,  $\beta_{thresh} = 0.9986$ ,  $\gamma_{thresh} = 18.92$ , 0.6 0.4 Ifrom Fernow.  $p_{thresh}(\pi) = 2.637 \; GeV$  $\Rightarrow$ Introduction to Experimental 0.2 Particle Physics] 0.0 2.0 2.5 3.0 3.5 PION MOMENTUM (GeV/c)

Most photons are emitted in the UV, which is hard to detect; but a generous number of photons are emitted in the visible, which matches well to mirrors, windows, and PMTs. The greater  $\beta$ , the greater  $\theta_c$ , and the more light is radiated:

$$\begin{aligned} \frac{dN}{dx \, d\lambda} &= 4\pi \alpha \frac{1}{\lambda^2} \left( 1 - \frac{1}{\beta^2 n^2} \right) \\ \Rightarrow \quad \frac{dN}{dx} &= 2\pi \alpha \int_{n\beta} \left( 1 - \frac{1}{\beta^2 n^2} \right) \frac{d\lambda}{\lambda^2} \\ &\approx 2\pi \alpha \sin^2 \theta_c \left( \frac{1}{\lambda_1} - \frac{1}{\lambda_2} \right) \\ &\approx 390 \sin^2 \theta_c \qquad (\lambda = 350 - 500 \text{ nm}) \end{aligned}$$

Note: detecting photons down to  $\lambda = 330$  nm increases photon yield by 20% (390  $\rightarrow$  470)

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### Cerenkov Radiation: Belle II quartz (n=1.466)

Photon yield per cm of radiator

Cerenkov angle:



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### **Cerenkov Radiation: LHCb**





## **Cerenkov Radiation: LHCb**





### **Cerenkov Radiators**

N<sub>PE</sub>/cm versus Refractive Index for Various Radiators

 $\begin{array}{c} \gamma_{\text{threshold}} \, \text{versus Refractive Index for} \\ & \text{Various Radiators} \end{array}$ 



[from Ratcliffe, Imaging Rings in Ring Imaging Cherenkov Counters, presented at RICH 2002, Pylos, Greece]

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## Belle II iTOP ("imaging time-of-propagation") Detector



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## **Belle II iTOP Detector: how it works**



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## **Belle II iTOP Detector: how it works**



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## **Belle II iTOP Detector: mirror focusing**

A spherical mirror focuses parallel rays to common point:



This can "remove" the thickness of the radiator bar:



![](_page_13_Picture_0.jpeg)

# **Belle II iTOP Detector: time-of-propagation**

Time of propagation (TOP) depends on path length and group velocity:  $v_g = \frac{c}{n_g}$ 

This "group index"  $n_g$  differs from the "phase index" n that governs  $\cos \theta_c = 1/(n\beta)$ .

Also, the group index has much larger chromatic dispersion than the phase index:

dispersion

$$n(\lambda) = n(\lambda_0) - \lambda rac{dn}{d\lambda}$$

- the Cerenkov angle θ<sub>c</sub> has very small (but non-zero) variation with photon wavelength: UV photons are emitted at slightly larger angles than IR photons
- the TOP has notable variation with photon wavelenth: UV photons propagate slower than IR photons; this degrades time resolution

![](_page_13_Figure_9.jpeg)

![](_page_13_Figure_10.jpeg)

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![](_page_14_Picture_0.jpeg)

## **Belle II iTOP Detector: time-of-propagation**

	<b>Refractive Ind</b> Data in 22°C in	<b>ex and Dispersion</b> 760mm Hg dry niti							
Corning 7980	Wavelength [air] λ [nm]	Refractive Index <sup>*2</sup> n		lfro	m Ratcliff	e Imaging	Rinas in		
(quartz)	1128 64	1 448870		Rin	g Imaging	Cherenko	V		
	1064.00	1 449633		Co	unters, pre	esented at	RICH		
~1.4% effect for n ~2.6% effect for TOP	1060.00	1 449681		200	02, Pylos,	Greece]			
	1013 98 n	1 450245							
	852 11 n	1 452469							
	$\frac{0.052.11 \text{ m}_{s}}{706.52 \text{ n}}$	1.152109	Refractive	e Indices ar	nd Dispersi	on versus W	/avelength	for SiO <sub>2</sub>	
	656.27 n	1.456370	2				<u> </u>	<b></b>	10
	643.85 n	1.456707							
	632 80 n	1.457021							
	589 29 n-	1.458406	× 1.9 -						
	587 56 n	1 458467	ğn						
	546 07 n	1 460082				n (group)			
	486.13 n	1.100002	. <u>≥</u> 1.8 -			— Dispersion (n)		-	1
	479 99 n	1.463509	g I			Disporsion [n (	(group)]		
	435.83 n	1.466701					group)]		4
	404 66 n	1 469628	<b>č</b> 1.7 -			Dispersion [n (	group)]/ Dispei	rsion (n)	ž
	365.01 n	1 474555							ģ
	334.15	1 479785							-
	312.57	1.484514	1.6 -					+	0.1
	308.00	1.485663							
	248.30	1.508433							
	248.00	1.508601	1.5 -						
	214.44	1.533789							
	206.20	1.542741							
	194.17	1.559012	1.4	I	I	I		<b>_</b>	0.01
	193.40	1.560208	0 19	0 29	0.39	0 49	0.59	0.69	
	193.00	1.560841	0.10	0.20	0.00	0.10	0.00	0.00	
	184.89	1.575131		Photo	on Wavele	ength $\lambda$ (mic	crons)		

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![](_page_15_Picture_0.jpeg)

## **Belle II iTOP Detector: light yield**

![](_page_15_Figure_2.jpeg)

MC tuning includes measured quantum efficiency (~25%) and collection efficiency (~55%) of PMTs

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![](_page_16_Picture_0.jpeg)

## **Belle II iTOP Detector: performance**

- *iTOP particle likelihood has a contribution from*  $\gamma$  *yield and a contribution from position + time (* $\theta_c$ *)*
- Evaluate performance with  $D^{*+} \rightarrow D^{0}(\rightarrow K^{-}\pi^{+})\pi^{+}$  or  $D^{*-} \rightarrow D^{0}(\rightarrow K^{+}\pi^{-})\pi^{-}$  decays

Effic(K) = (# K tracks identified as K)/(# K tracks) MisID( $\pi$ ) = (#  $\pi$  tracks identified as K)/(#  $\pi$  tracks)

![](_page_16_Figure_5.jpeg)

"ROC" curve

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![](_page_17_Picture_0.jpeg)

## **Belle II iTOP Detector: performance**

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![](_page_17_Figure_5.jpeg)

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![](_page_18_Picture_0.jpeg)

Bethe-Bloch formula (1933):

$$\left\langle -rac{dE}{dx}
ight
angle = K\left(rac{Z}{A}
ight)\left(rac{z^2}{eta^2}
ight)\left[\ln\left(rac{2meta^2\gamma^2}{I}
ight) -eta^2 -rac{\delta(eta\gamma)}{2}
ight]$$

 $K = 0.307 \text{ MeV mol}^{-1} \text{ cm}^2$ 

#### Complicated, but note:

- depends on material only via (Z/A), which varies little over a range of materials
- depends on incident particle via z<sup>2</sup>, i.e, higher z gives much more ionization
- also depends on βγ of incident particle: broad minimum is at βγ ~ 3, sharp drop at lower values, slow rise at larger values. This corresponds to ~2 MeV/(g/cm<sup>2</sup>)
- does not depend on m of incident particle, only  $\beta\gamma$  and, knowing p, that is how one uses dE/dx for particle ID.
- the 8/2 term is a small correction (the "density effect")

[from D. Groom, PDG Review "Passage of Particles Through Matter"]

![](_page_18_Figure_12.jpeg)

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![](_page_19_Picture_0.jpeg)

Bethe-Bloch formula (1933):

### Finer points:

denominator "dx" is a distance, but units are grams/cm<sup>2</sup> (not cm)

![](_page_19_Picture_5.jpeg)

- $1/\beta^2$  drop is due to less time for incident particle to interact with electrons
- "relativistic rise" is due to greater range of electric field of incident particle, and greater maximum energy transfer
- when plotting vs. particle momentum, the  $\beta\gamma$  curve gets shifted; shift depends on particle mass

![](_page_19_Figure_9.jpeg)

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![](_page_20_Picture_0.jpeg)

Bethe-Bloch formula (1933):

$$\left\langle -rac{dE}{dx}
ight
angle = K\left(rac{Z}{A}
ight)\left(rac{z^2}{eta^2}
ight)\left[\ln\left(rac{2meta^2\gamma^2}{I}
ight) -eta^2 -rac{\delta(eta\gamma)}{2}
ight]$$

- This formula gives only the **mean energy loss** it is dominated by huge-energy-loss (tinyimpact-parameter) collisions, which liberate electrons (called "delta rays").
- ⇒ PDG recommends using "most probable energy loss" formula for calculations

![](_page_20_Figure_6.jpeg)

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![](_page_21_Picture_0.jpeg)

dE/dx curve for CDC (all charged w/ CC)

![](_page_21_Figure_3.jpeg)

 $\Rightarrow$  good  $\pi/K$  separation up to p ~1.1 GeV/c

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![](_page_22_Picture_0.jpeg)

## **Particle ID homework problem #1:**

## *In the forward endcap region of Belle II is the ARICH:*

4.0 cm of aerogel radiator, followed by an array of finely segmented photodetectors. The separation is 20.0 cm (see figure) and n(aerogel) = 1.050.

What is the difference in radii of the "Cerenkov rings" at the photodetectors for:

- a)  $\pi^+$  and  $K^+$  with p=3.0 GeV/c?
- b) for  $\pi^+$  and  $K^+$  with p = 4.0 GeV/c?
- c) what are the mean γ yields for these four cases?

![](_page_22_Picture_8.jpeg)

![](_page_22_Figure_9.jpeg)

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![](_page_23_Picture_0.jpeg)

## **Particle ID homework problem #2:**

![](_page_23_Figure_2.jpeg)

For a  $\mu^-$  to reach the KLM, it must pass through (at least) the following material:

- a) iTOP quartz: 2.0 cm @  $\rho$  = 2.201 g/cm<sup>3</sup>
- **b) ECL CsI crystal:** 30 cm @  $\rho$  = 4.51 g/cm<sup>3</sup>
- **c)** magnet coil Alum.: 10 cm @  $\rho$  = 2.710 g/cm<sup>3</sup>
- **d) KLM iron (1 layer):** 4.7 cm @ ρ = 7.874 g/cm<sup>3</sup>

What is the minimum energy required for a "minimum-ionizing"  $\mu^ \langle -dE/dx \rangle \approx 2 \text{ MeV} / (q/cm^2)$ 

to reach the first scintillator layer of the KLM?