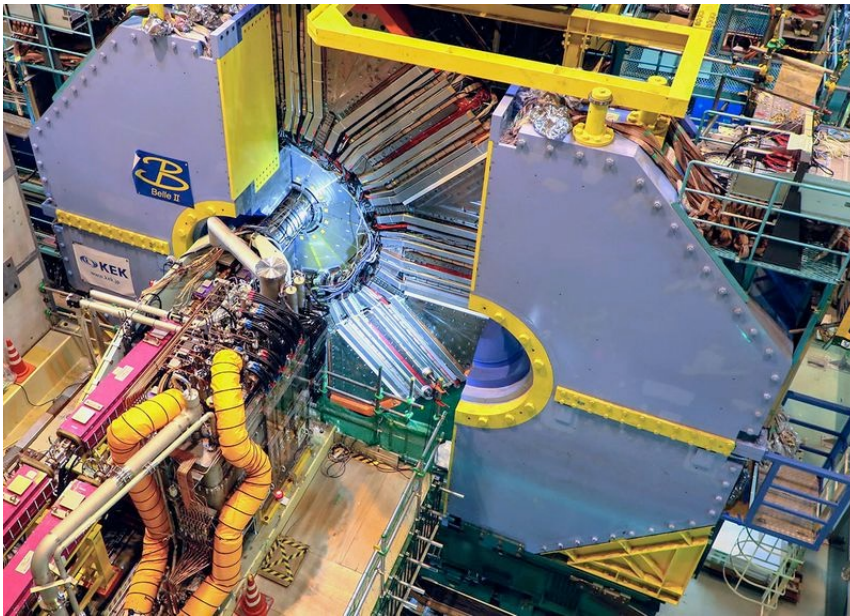


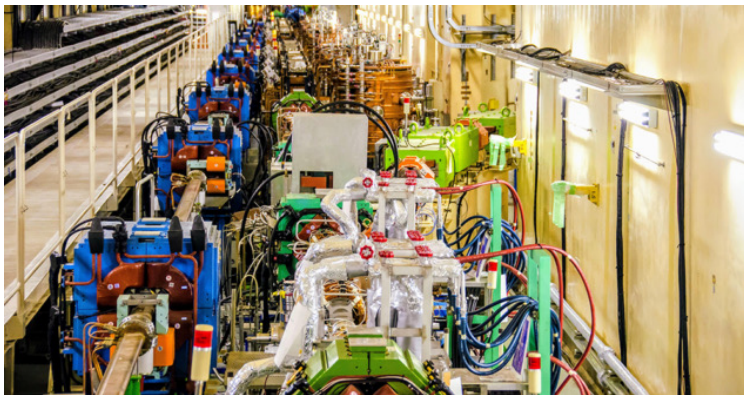
# Belle II: Opportunities for NP Discoveries in B physics



Tom Browder, University of Hawai'i at Manoa



The complex superconducting final focus is partially visible here (before closing the endcap).



Inside the SuperKEKB tunnel

Après Snowmass and the last Belle II Physics Run (Leo's talk) ( $L_{\text{peak}}=4.7 \times 10^{34}/\text{cm}^2/\text{sec}$ ,  $\text{Int}(L \text{ dt}) > 2.5 \text{ fb}^{-1}/\text{day}$ , which are new world records. A BaBar-sized data sample is now “on tape”.)

*A few early Physics Results from Belle II: B Physics*

Opportunities for *new physics discoveries* and the road ahead (the Belle II Physics Book, Snowmass Belle II Physics Whitepaper (WP) and other WPs)

Belle II/SuperKEKB Snowmass WPs:  
<https://confluence.desy.de/display/BI/Snowmass+2021>

# Snowmass 2022 (*International Physics Rodeo*)

Scenes from the actual Snowmass Rodeo in Colorado



N.B. Snowmass was *just held* in Seattle, Washington in summer of 2022. The last one was held in Minneapolis, Minnesota in 2013. It is unlikely that there will ever be another month-long planning meeting in Snowmass, CO.

Historical note: Young(ish) Scientist Pier Oddone (originally from Peru/Italy) introduced the concept and first proposal for an asymmetric energy B-factory to the broad HEP community at a Snowmass in 1988.



# Revisionist History and **Paradigm Shift**

The B factory experiments, Belle and BaBar, discovered large CP violation in the B system in 2001, compatible with the SM and provided a large range of CKM measurements. These provided the experimental foundation for the 2008 Nobel Prize to Kobayashi and Maskawa.

In the meantime, the LHC was constructed in 2008, ATLAS and CMS completely changed the nature of high energy physics. Of particular importance was the landmark discovery in 2012 of the Higgs boson.

This discovery was recognized by the 2013 Physics Nobel Prize to Englert and Higgs.

In addition, the high  $p_T$  experiments, established tight constraints on direct production of high mass particles (e.g.  $M(Z')$ ,  $M(W')$ )  $> 3$  TeV, vector-like fermions  $> 800$  GeV) and limits on SUSY. This *noble search* continues with the high luminosity LHC.

**Paradigm shift**: inspired by intriguing results from B factories, LHCb and the potential of Belle II, the possibility of finding new physics in flavor has emerged as a *complementary* route to the LHC.

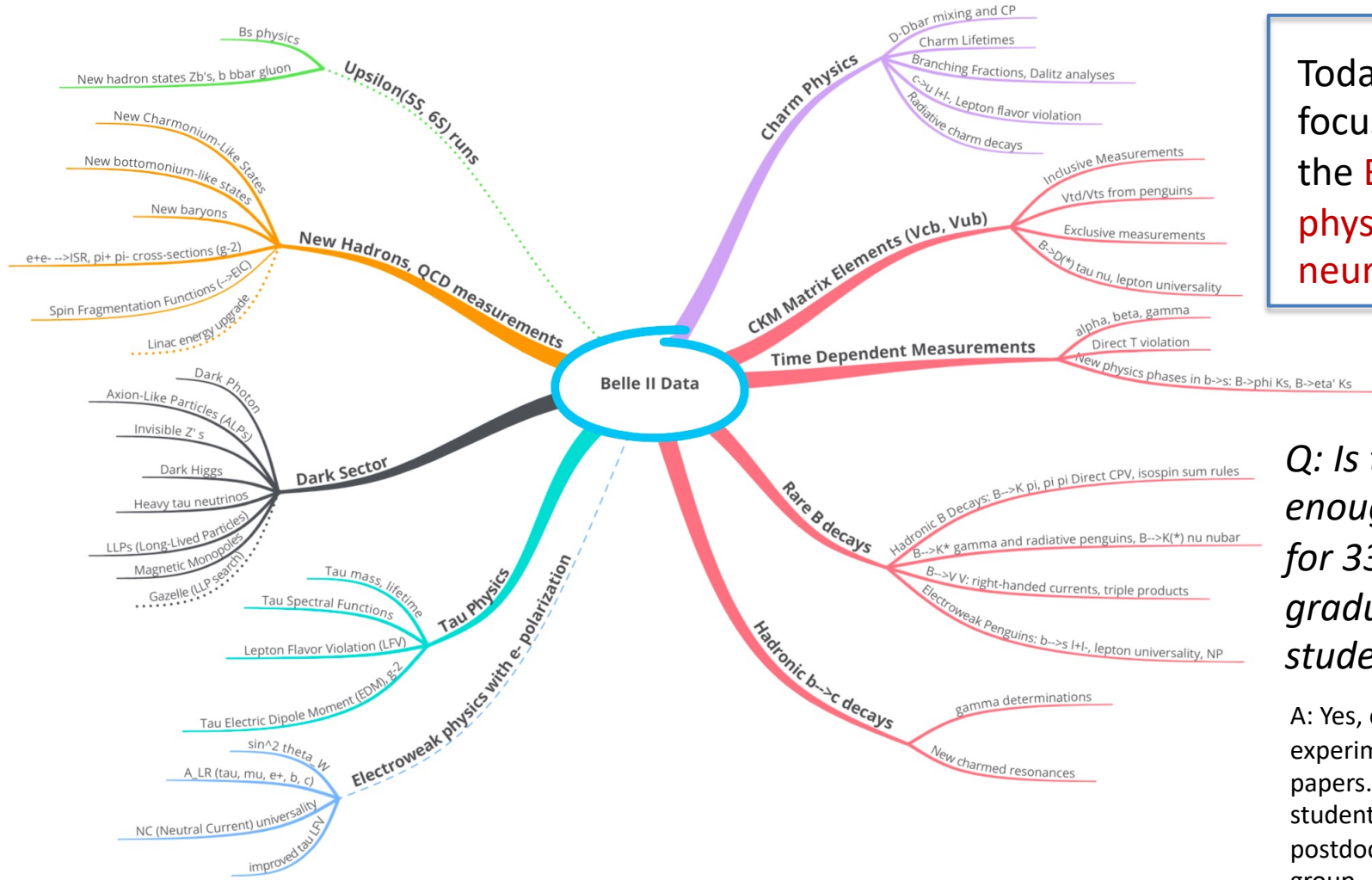


Younger theorists:  
**Dark Sector**  
may be another path.

# Belle II Physics “Mind Map” for Snowmass 2022



Wealth of new physics possibilities in different domains of HEP (weak, strong, electroweak interactions). Many opportunities for *initiatives* by **young scientists**.



Today, we focus on the **B physics neurons**.

*Q: Is there really enough physics for 330 graduate students ?*

*A: Yes, c.f. B factory experiments, >500 papers. Most by PhD student/advisor, postdoc or small group.*

*Dashed lines indicate extensions to SuperKEKB/Belle II that can enhance the physics reach of the facility. WP's <https://confluence.desy.de/display/BI/Snowmass+2021>*



# Steve Weinberg on crises in physics.

## I. INTRODUCTION

Physics thrives on crisis. We all recall the great progress made while finding a way out of various crises of the past: the failure to detect a motion of the Earth through the ether, the discovery of the continuous spectrum of beta decay, the  $\tau$ - $\theta$  problem, the ultraviolet divergences in electromagnetic and then weak interactions, and so on. Unfortunately, we have run short of crises lately. The “standard model” of electroweak and strong interactions currently faces neither internal inconsistencies nor conflicts with experiment. It has plenty of loose ends; we know no reason why the quarks and leptons should have the masses they have, but then we know no reason why they should not.

Perhaps it is for want of other crises to worry about that interest is increasingly centered on one veritable crisis: theoretical expectations for the cosmological constant exceed observational limits by some 120 orders of magnitude.<sup>1</sup> In these lectures I will first review the history of this problem and then survey the various attempts that have been made at a solution.



BTW can you identify the three Nobel Prize Winners ?

Do you know all the crises that Weinberg is referring to ?

BTW: now cosmology is stuck in its version of the Standard Model.

# FAQ: What is meant by “lepton universality” ?

This refers to the **weak interaction**

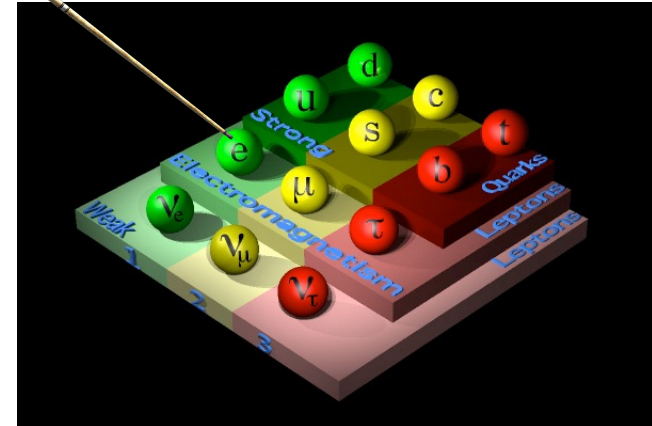
The weak couplings of leptons of different generations are the *same in the Standard Model*.

**Lepton universality** has been tested to O(1%) precision for lepton decays, pion and kaon decays....e.g.

$$g_{\mu} / g_{\tau} = 1.001 \pm 0.003$$

However in the  $b \rightarrow c$  *charged current weak interaction* and  $b \rightarrow s$  *neutral current weak interaction*, there are experimental hints of its breakdown ( $\sim 3\sigma$  level) at the 10-15% level.

QM Billiard Table  
Lepton row



For example,

$$R_D = \frac{\mathcal{B}(B \rightarrow D\tau\nu_{\tau})}{\mathcal{B}(B \rightarrow D\ell\nu_{\ell})}$$

$$R_D^* = \frac{\mathcal{B}(B \rightarrow D^*\tau\nu_{\tau})}{\mathcal{B}(B \rightarrow D^*\ell\nu_{\ell})}$$

$$l = e^{-}, \mu^{-}$$

Deviate from their SM expectations

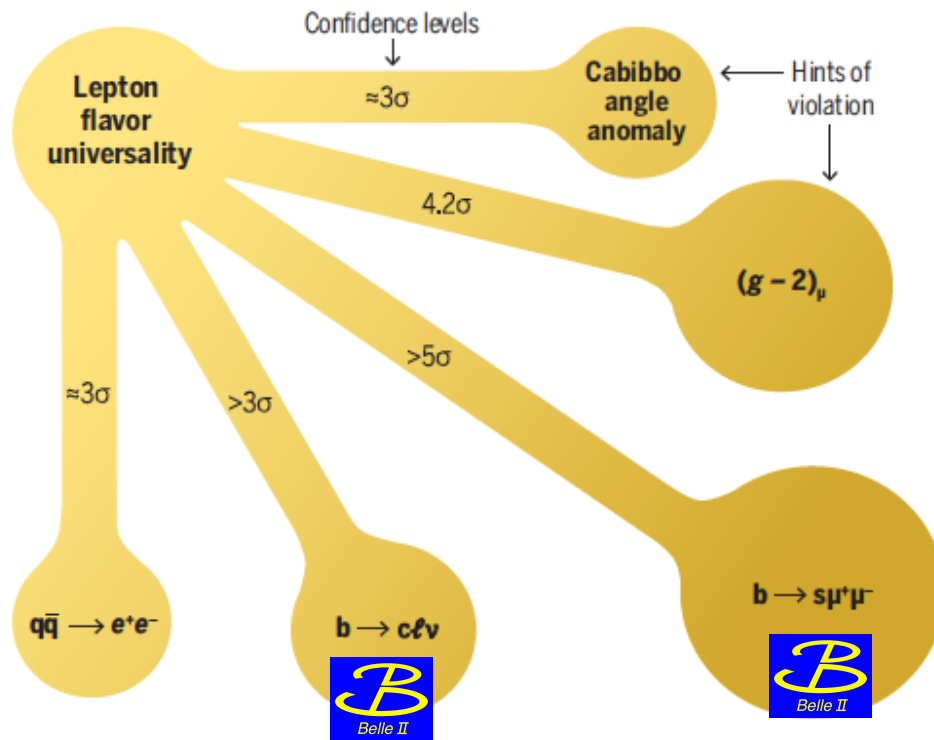


# An emerging crisis in High Energy Physics.



## Possible violations of lepton flavor universality are getting harder to ignore

Shown are five hints for the violation of lepton flavor universality from existing experimental data, with the size of each circle and length of each arm reflecting the level of confidence for the experimental data to break away from standard model predictions.



What are we doing to address this in Belle II ?

Is it real ?

Let's carry out a program of measurements at Belle II to find out.

From December 2021 SCIENCE magazine article by A. Crivellin and M. Hoferichter.






# Big Bang Theory (Flavor Changing Neutral Currents)

Sheldon, what about FCNCs ?

$t \rightarrow W^+ b$



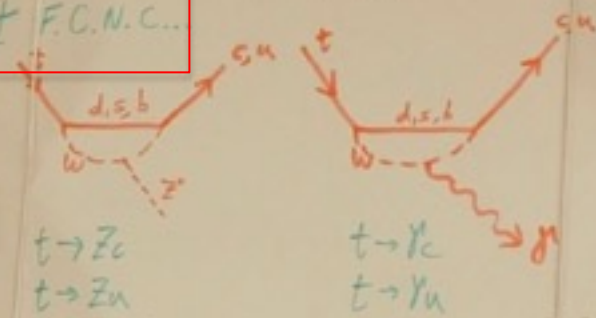
$$BR(t \rightarrow Wb) = \frac{\Gamma(t \rightarrow Wb)}{\Gamma(t \rightarrow Wq)}$$

$$= \frac{|V_{cb}|^2}{|V_{cd}|^2 + |V_{cs}|^2 + |V_{cb}|^2}$$

$$\approx \frac{(0.9745)^2}{(0.0094)^2 + (0.0410)^2 + (0.9745)^2}$$

$$= 99.82\%$$

but F.C.N.C...

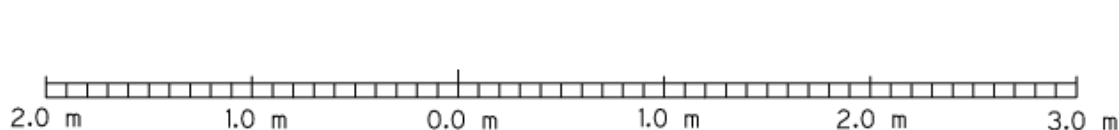
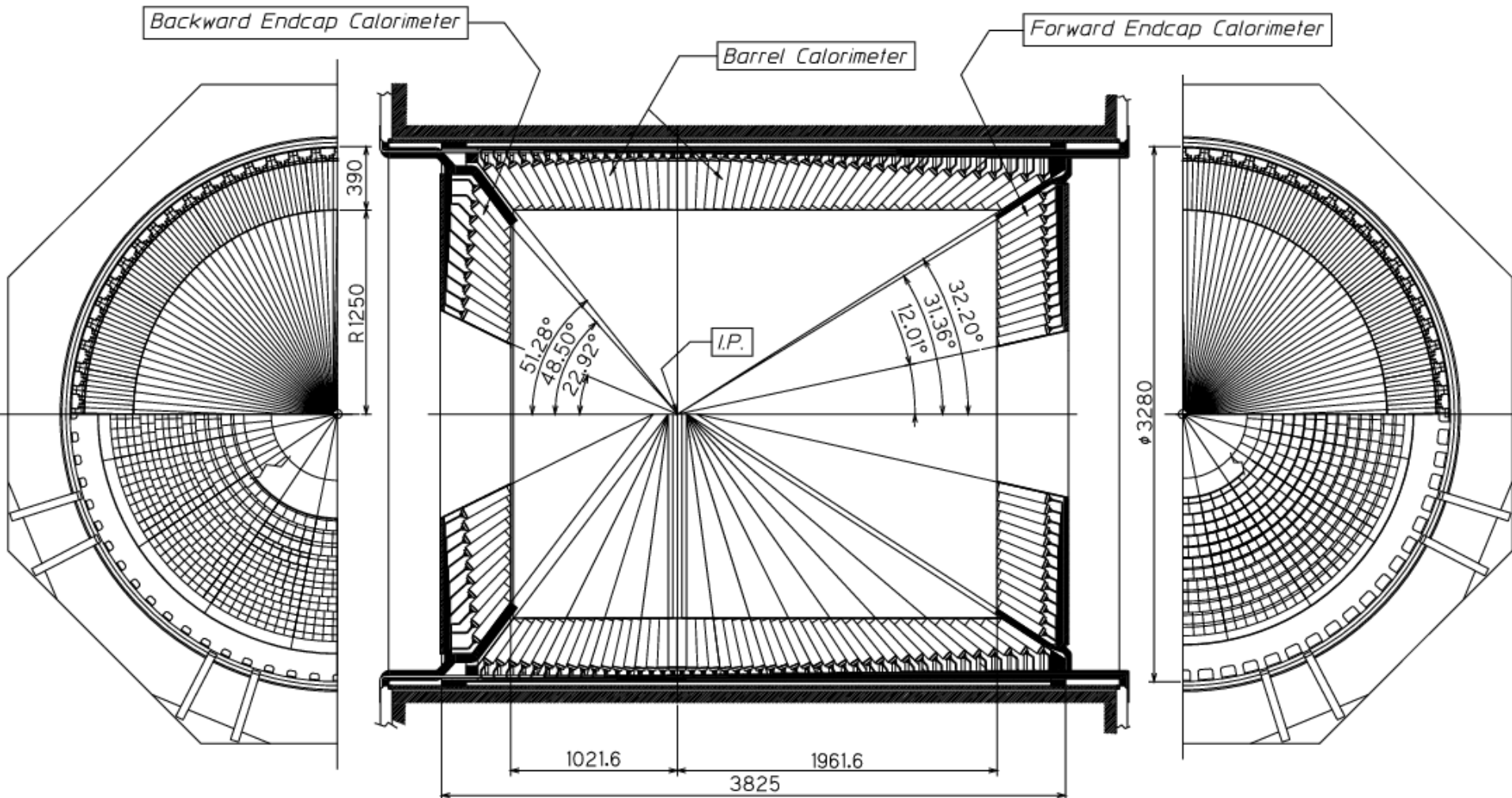


Remember FCNCs do not occur at 1<sup>st</sup> order in the SM. (only at 2<sup>nd</sup> order)

$$U_{CKM} = \begin{pmatrix} c_{12}c_{13} & & \dots \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & & \dots \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & & \dots \end{pmatrix}$$

Note the weak coupling constants are **complex**

Belle II's CsI(Tl) calorimeter (~Belle with improved waveform sampling and timing). 8736 crystals covering 90% of the solid angle.



unit (mm)

Belle's Neutral  
detection  
superpower

Fig. 69. Overall configuration of ECL.

See talk by  
Prof. Savino  
Longo.



# Re-discovery of Radiative Penguins at Belle II

1975: Vainshtein, Zakharov and Shifman



Examine the following  $b \rightarrow s \gamma$  decay modes in the Belle II Phase 3 dataset.

$$B^0 \rightarrow K^{*0} \gamma \rightarrow K^+ \pi^- \gamma$$

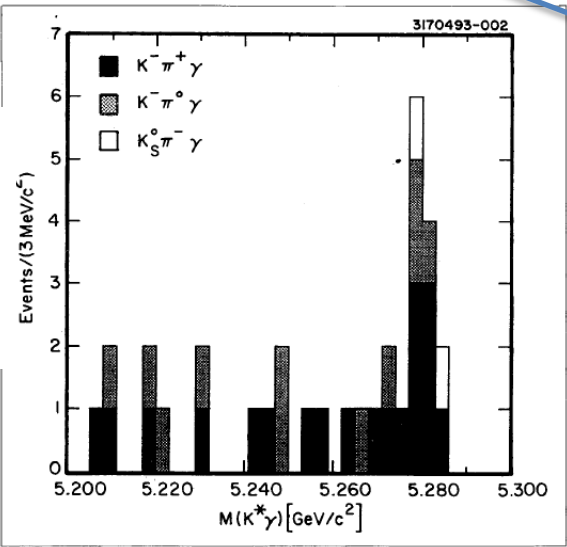
$$B^+ \rightarrow K^{*+} \gamma \rightarrow K^+ \pi^0 \gamma$$

$$B^+ \rightarrow K^{*+} \gamma \rightarrow K_S^0 \pi^+ \gamma$$

1993 CERN Courier:

CORNELL  
CLEO discovers  
B meson penguins

N.B. Using  $1.5 \times 10^6$   
B meson pairs



John Ellis, the CERN theorist who coined the name "Penguin".

Ed Thorndike,  
Rochester,  
CLEO



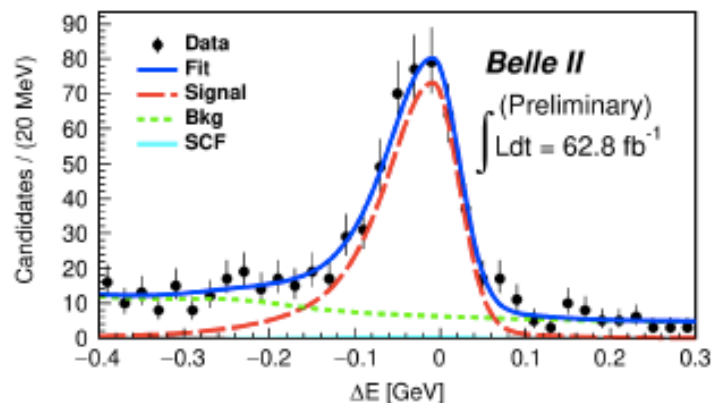
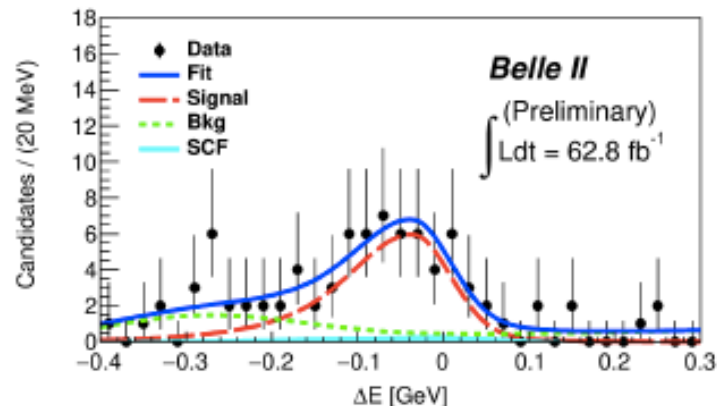
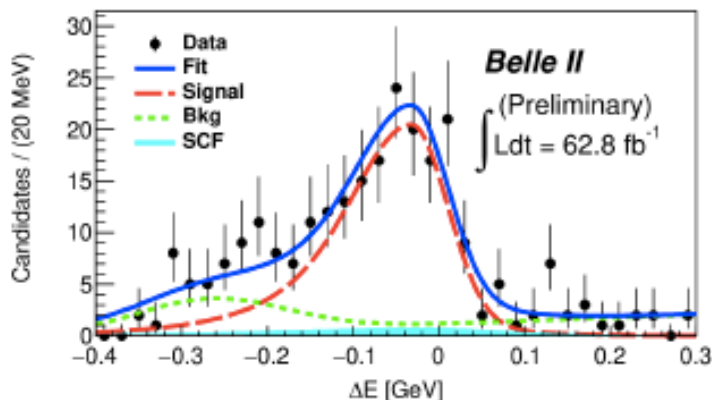
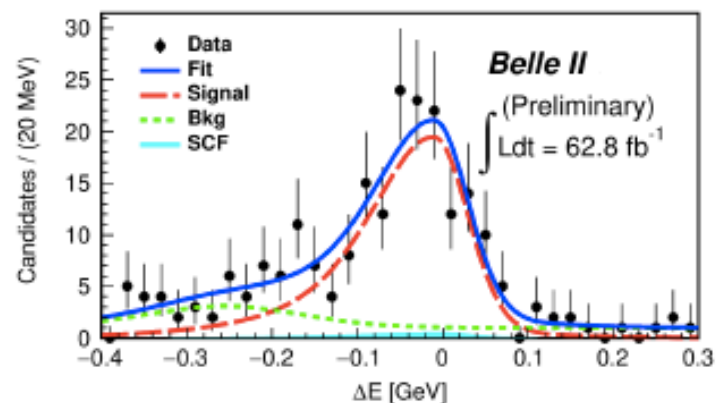
(a)  $B^0 \rightarrow K^{*0}[K^+\pi^-]\gamma$ (b)  $B^0 \rightarrow K^{*0}[K_S^0\pi^0]\gamma$ (c)  $B^+ \rightarrow K^{*+}[K^+\pi^0]\gamma$ (d)  $B^+ \rightarrow K^{*+}[K_S^0\pi^+]\gamma$ 

Figure 2.  $\Delta E$  distributions for each  $B \rightarrow K^*\gamma$  mode with the fit result superimposed. The black dots with error bars denote the data, the blue curve denotes the total fit, the dashed red curve is the signal component, the dotted green curve is the background component, and the filled cyan region is the misreconstructed signal component.



BELLE2-CONF-2021-028

Table I. Signal yield, efficiency and measured branching fraction for each mode. When two uncertainties are given, the first is statistical and the second is systematic. The world-average values reported by the PDG are given for comparison.

Mode	Signal yield	Signal efficiency (%)	B.F (Fit) $\times 10^{-5}$	B.F (PDG) $\times 10^{-5}$
$B^0 \rightarrow K^{*0}[K^+\pi^-]\gamma$	$454 \pm 28$	14.9	$4.5 \pm 0.3 \pm 0.2$	$4.18 \pm 0.25$
$B^0 \rightarrow K^{*0}[K_S^0\pi^0]\gamma$	$50 \pm 10$	1.7	$4.4 \pm 0.9 \pm 0.6$	$4.18 \pm 0.25$
$B^+ \rightarrow K^{*+}[K^+\pi^0]\gamma$	$169 \pm 18$	4.7	$5.0 \pm 0.5 \pm 0.4$	$3.92 \pm 0.22$
$B^+ \rightarrow K^{*+}[K_S^0\pi^+]\gamma$	$160 \pm 17$	4.1	$5.4 \pm 0.6 \pm 0.4$	$3.92 \pm 0.22$

Table II. Measured branching fractions for combined charged and neutral modes. The first uncertainty is statistical and the second is systematic. The world-average values reported by the PDG are given for comparison.

Mode	B.F (Fit) $\times 10^{-5}$	B.F (PDG) $\times 10^{-5}$
$B^0 \rightarrow K^{*0}\gamma$	$4.5 \pm 0.3 \pm 0.2$	$4.18 \pm 0.25$
$B^+ \rightarrow K^{*+}\gamma$	$5.2 \pm 0.4 \pm 0.3$	$3.92 \pm 0.22$

So far, branching fractions only.



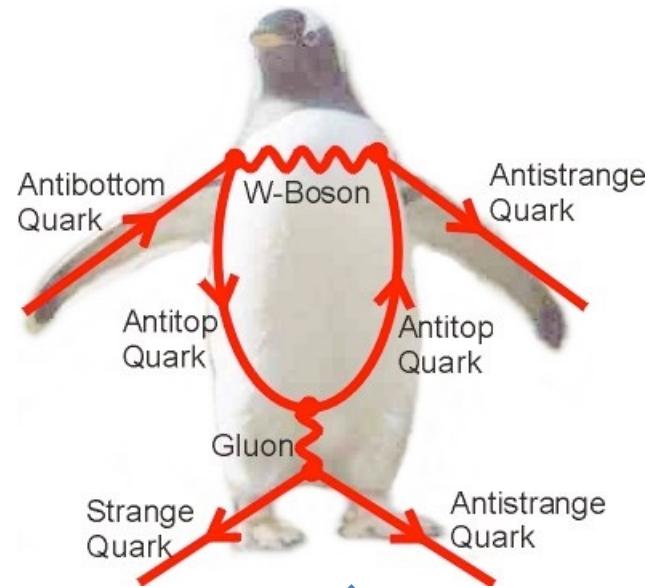


Move on to **gluonic penguins**:

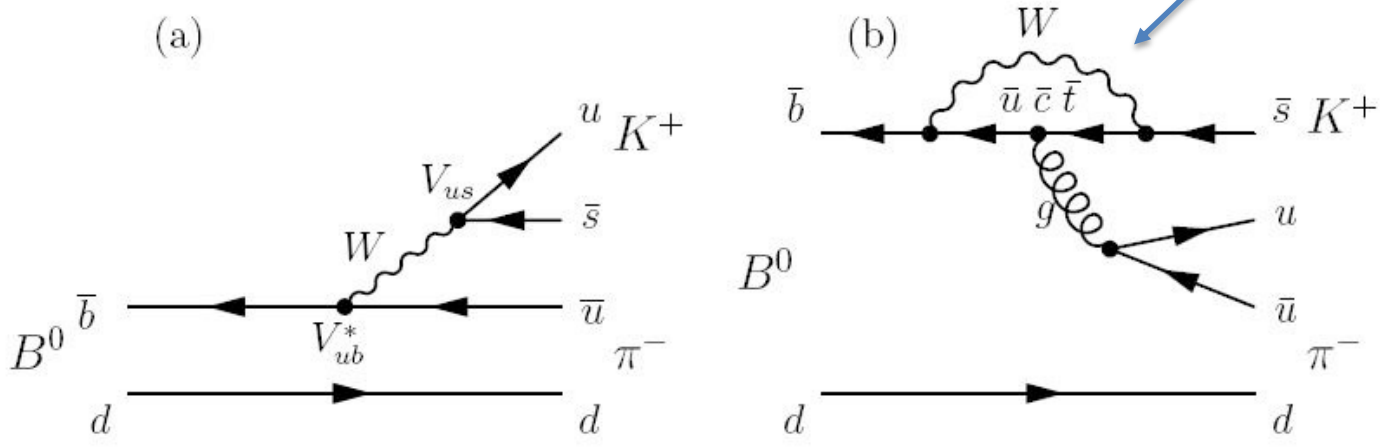
Rare Decay Mascot/Feynman Diagram

Let's consider

$$B^0 \rightarrow \pi^- K^+$$



Feynman diagrams for this process



Feynman tree (a) and penguin (b) diagrams for the  $B_d^0 \rightarrow K^+ \pi^-$  decay

N.B. Both amplitudes contribute, but the **(b→s) Penguin** is larger and has a different weak phase ( $V_{ts}$  vs  $V_{ub}$ ).

## LETTERS

In 2008, “the K pi puzzle” appeared in *Nature*. Charged and neutral  $A(\text{CP})$ 's for  $B \rightarrow K\pi$  penguins differ. Is this a sign of **new physics** ?

## Difference in direct charge-parity violation between charged and neutral $B$ meson decays

The Belle Collaboration\*

Also confirmed by BaBar

Mode	$A_{\text{CP}}$		
	BaBar	Belle	LHCb
$K^+\pi^-$	$-0.107 \pm 0.016^{+0.006}_{-0.004}$	$-0.069 \pm 0.014 \pm 0.007$	$-0.080 \pm 0.007 \pm 0.003$
$K^+\pi^0$	$0.030 \pm 0.039 \pm 0.010$	$0.043 \pm 0.024 \pm 0.002$	$0.025^{+0.015+0.006}$
$K^0\pi^+$	$-0.029 \pm 0.039 \pm 0.010$	$-0.011 \pm 0.021 \pm 0.006$	$-0.022 \pm 0.025 \pm 0.010$
$K^0\pi^0$	$-0.13 \pm 0.13 \pm 0.03$	$0.14 \pm 0.13 \pm 0.06$	

In summary, we have measured the CP asymmetries for  $B \rightarrow K^\pm \pi^\mp$ ,  $K^\pm \pi^0$  and  $\pi^\pm \pi^0$  using 535 million  $B\bar{B}$  pairs. Direct CP violation in  $B^\pm \rightarrow K^\pm \pi^\mp$  is observed, accompanied by a large deviation between  $\mathcal{A}_{K^\pm \pi^\mp}$  and  $\mathcal{A}_{K^\pm \pi^0}$ . Although this deviation could be due to our limited understanding of the strong interaction, the difference in direct CP asymmetries for charged versus neutral  $B$  decays may be an indication of new sources of CP violation beyond the standard model of particle physics.

# “Trapping” the Electroweak Penguin in $B \rightarrow K \pi$

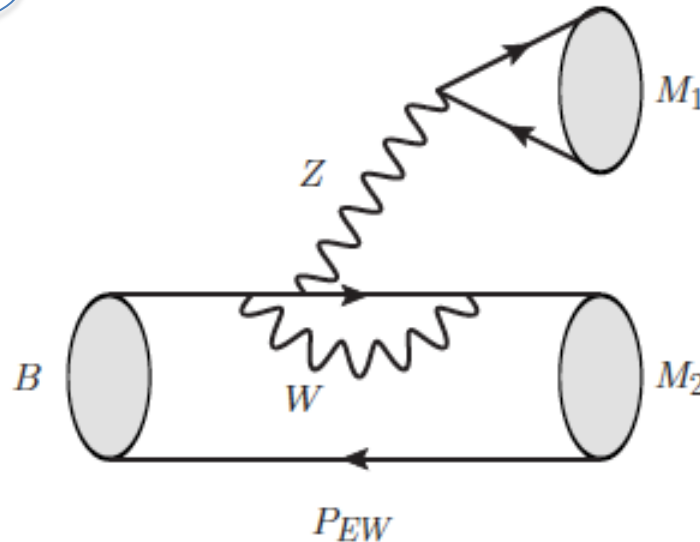
<https://arxiv.org/abs/hep-ph/0508047>

The isospin sum rule

$$I_{K\pi} = \mathcal{A}_{K^+\pi^-} + \mathcal{A}_{K^0\pi^+} \frac{\mathcal{B}(K^0\pi^+) \tau_{B^0}}{\mathcal{B}(K^+\pi^-) \tau_{B^+}} - 2\mathcal{A}_{K^+\pi^0} \frac{\mathcal{B}(K^+\pi^0) \tau_{B^0}}{\mathcal{B}(K^+\pi^-) \tau_{B^+}} - 2\mathcal{A}_{K^0\pi^0} \frac{\mathcal{B}(K^0\pi^0)}{\mathcal{B}(K^+\pi^-)}$$

(You may need to review isospin symmetry for the strong interaction)

**NP** can enter through this type of diagram, which would violate the sum rule



Michael Gronau

Have now observed the four  $B \rightarrow K \pi$  modes, needed for the *isospin sum rule test* of NP. This includes the difficult mode  $B \rightarrow K_S \pi^0$ . **Now have  $A_{CP}$  for all 4 modes and sensitivity estimates for the future.**







Examples of hadronic penguins ( $b \rightarrow s$  gluon) at Belle II.

$B \rightarrow K^- \pi^+$  and c.c.

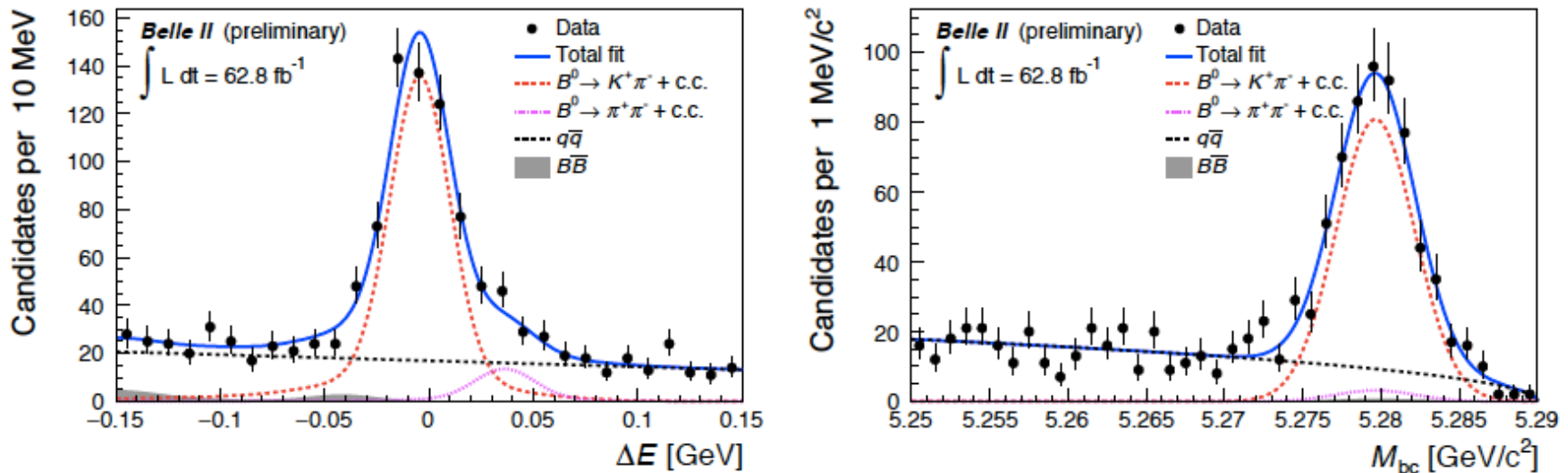


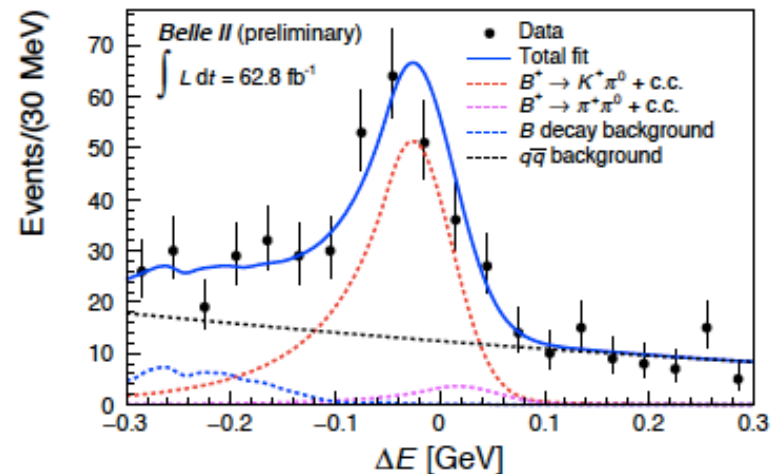
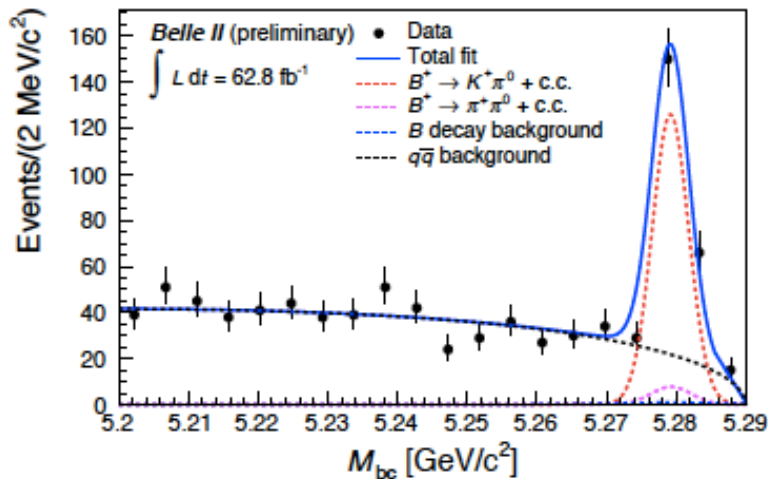
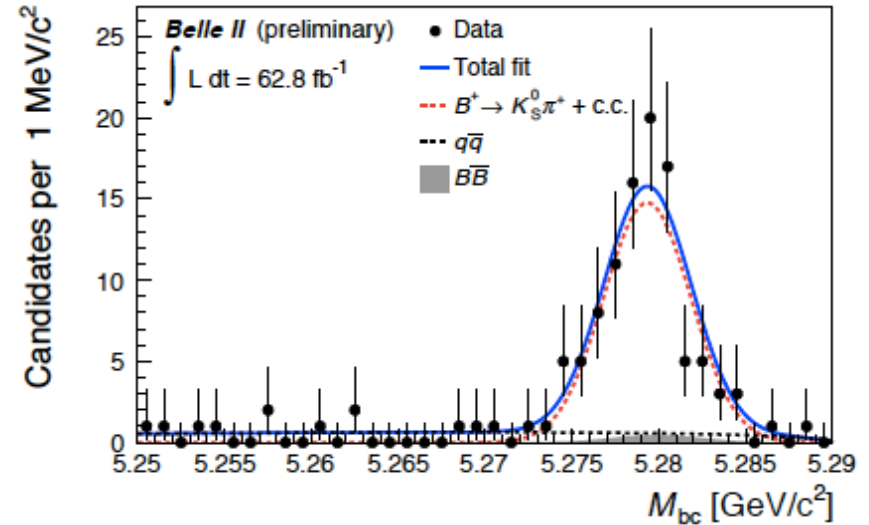
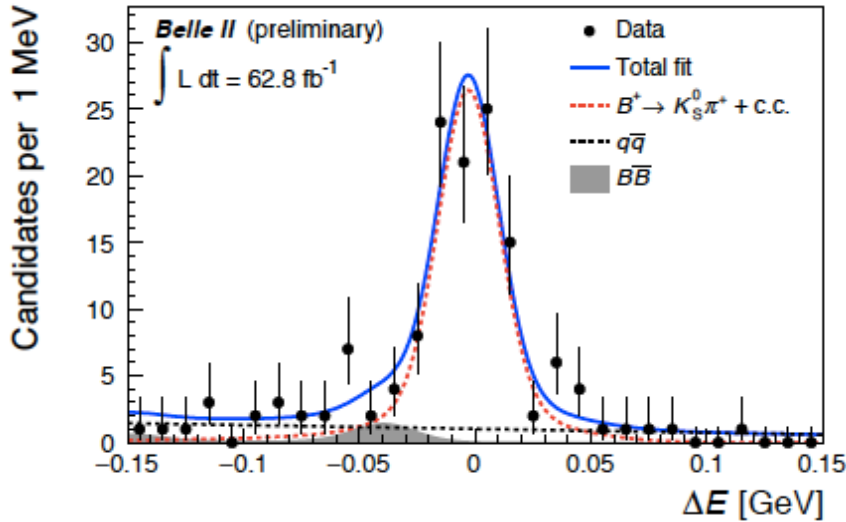
FIG. 2. Distributions of  $\Delta E$  (left) and  $M_{bc}$  (right) for  $B^0 \rightarrow K^+ \pi^-$  candidates reconstructed in 2019–2020 Belle II data, selected with an optimized continuum-suppression and kaon-enriching selection. The distributions are shown in signal-enriched regions of  $5.273 < M_{bc} < 5.286 \text{ GeV}/c^2$  and  $-0.04 < \Delta E < 0.03 \text{ GeV}$ , respectively. Fit projections are overlaid.

Details in <https://arxiv.org/abs/2106.03766>

These modes uses Belle II's other superpowers: tracking (Soeren Prell) and Particle ID (Alan Schwartz).



# More examples of **hadronic penguins** ( $b \rightarrow s$ gluon) at Belle II (modes with one Ks or one pi-zero)

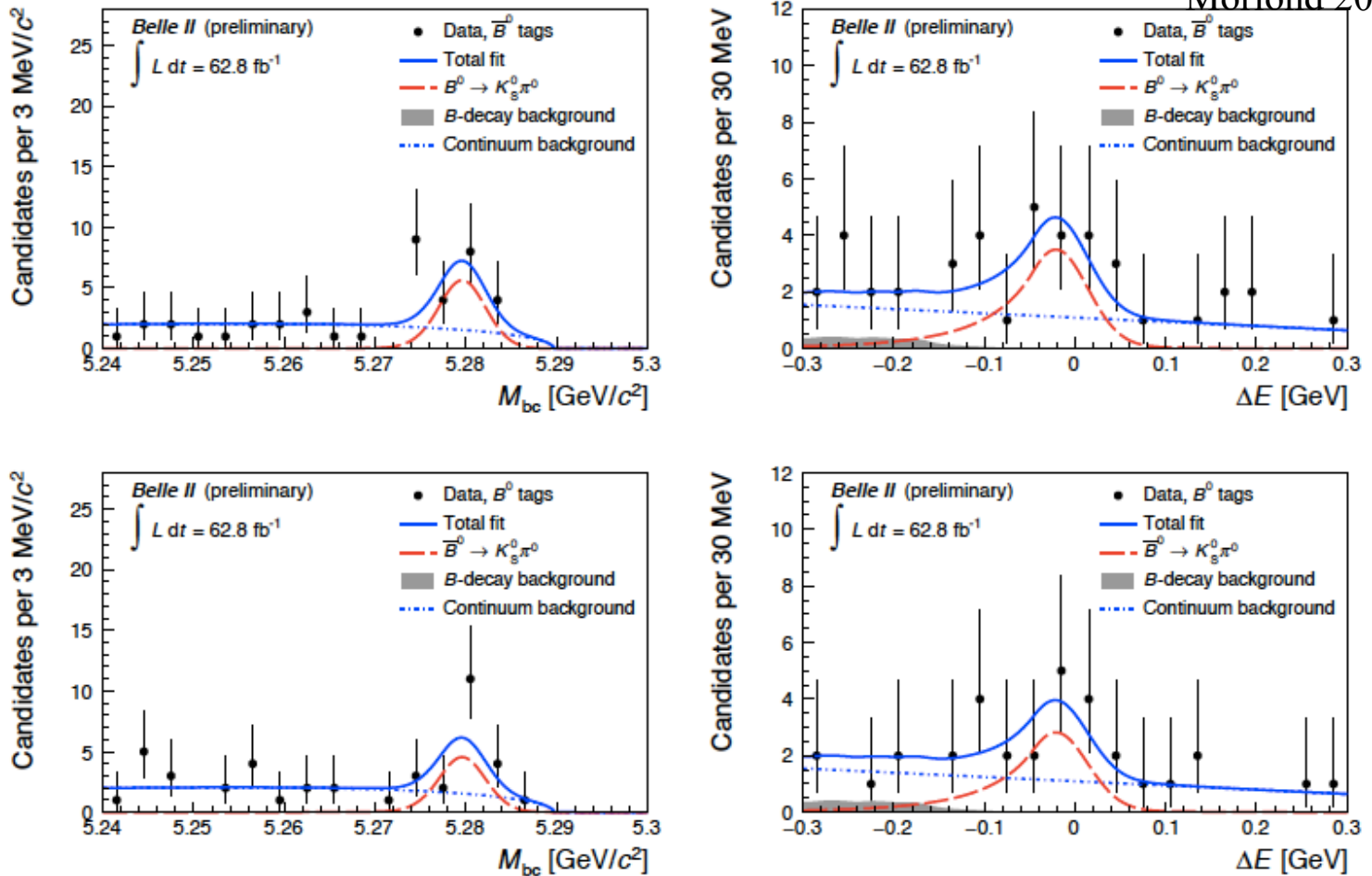


Details in <https://arxiv.org/abs/2106.03766> ; <https://arxiv.org/abs/2105.04111>



# Belle II's first result on $A_{CP}(B^0 \rightarrow K^0 \pi^0)$

Update with  
x3 data at  
Moriond 2022



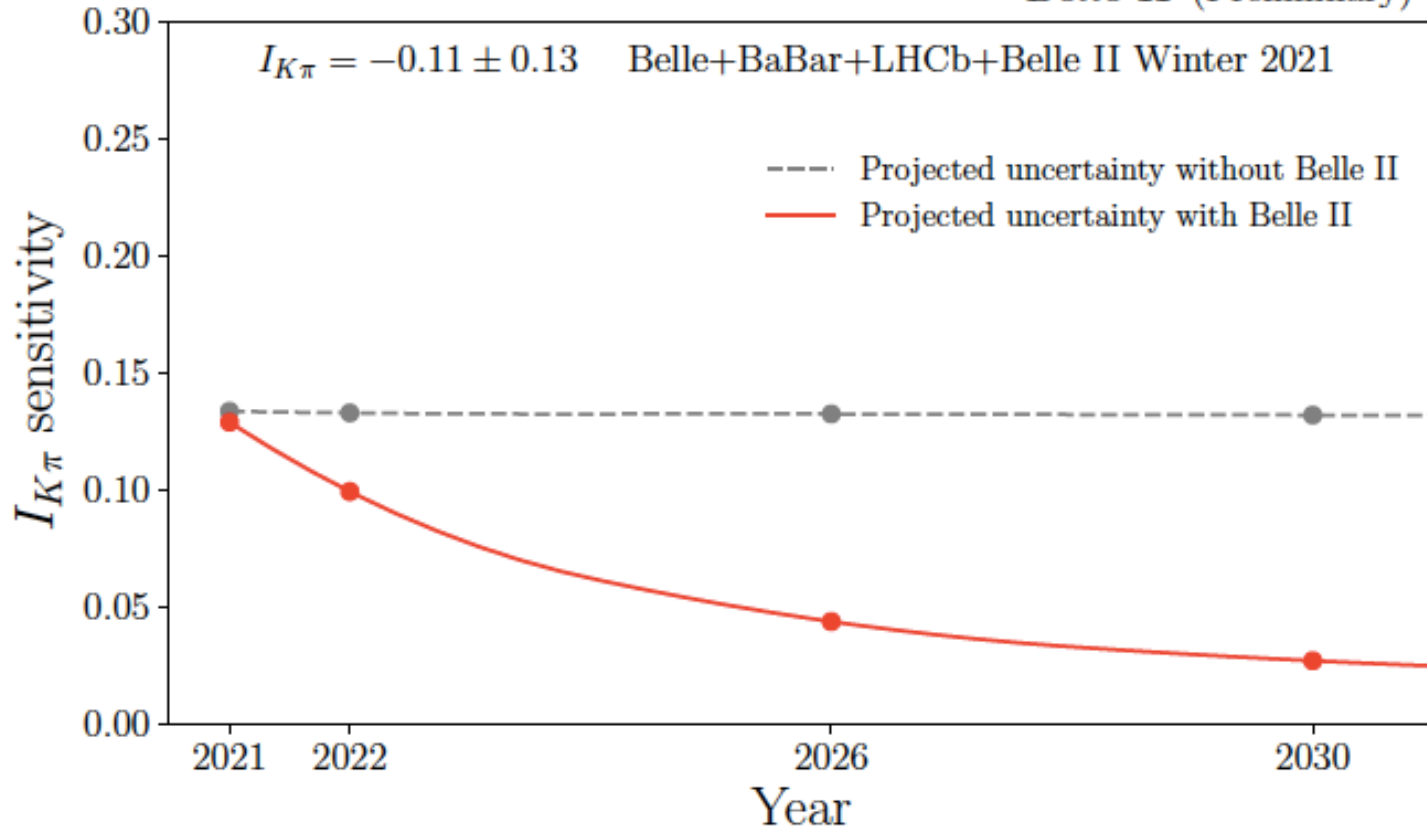
About 50  
events with  
62 fb<sup>-1</sup>

Difficult:  
Not self-  
tagging.

FIG. 3. Flavor-specific ( $M_{bc}$ ,  $\Delta E$ ) projections on 2019-2020 Belle II data. The top panel shows candidates where  $B_{tag}$  is tagged as a  $\bar{B}^0$  (signal-side:  $B^0$ ) and the bottom panel for candidates where  $B_{tag}$  is tagged as a  $B^0$  (signal-side:  $\bar{B}^0$ ). The distribution and fit are integrated over  $r$ -bin in the good tag region  $0.25 \leq r \leq 1$  and in the signal region (left panel:  $-0.16 < \Delta E < 0.08$  GeV, right panel:  $M_{bc} > 5.27$  GeV/c<sup>2</sup>).

Details in <https://arxiv.org/abs/2104.14871>





Without Belle II measurements of  $A_{CP}(B^0 \rightarrow K^0 \pi^0)$ , we are stuck.

*Need Belle II's neutral superpowers*

FIG. 4. The projected uncertainty on  $I_{K\pi}$  with and without Belle II inputs. The inputs for  $I_{K\pi}$  are averages of the estimated updates from ongoing LHCb and Belle II experiments with current world averages [10]. The red curve shows a projection when updates on the complete set of  $K\pi$  measurements are considered, and the grey curve is the case if only  $A_{K^+\pi^-}, A_{K^+\pi^0}, A_{K^0\pi^+}$  are updated by LHCb. The projection corresponds to the luminosity plans from LHCb and Belle II.

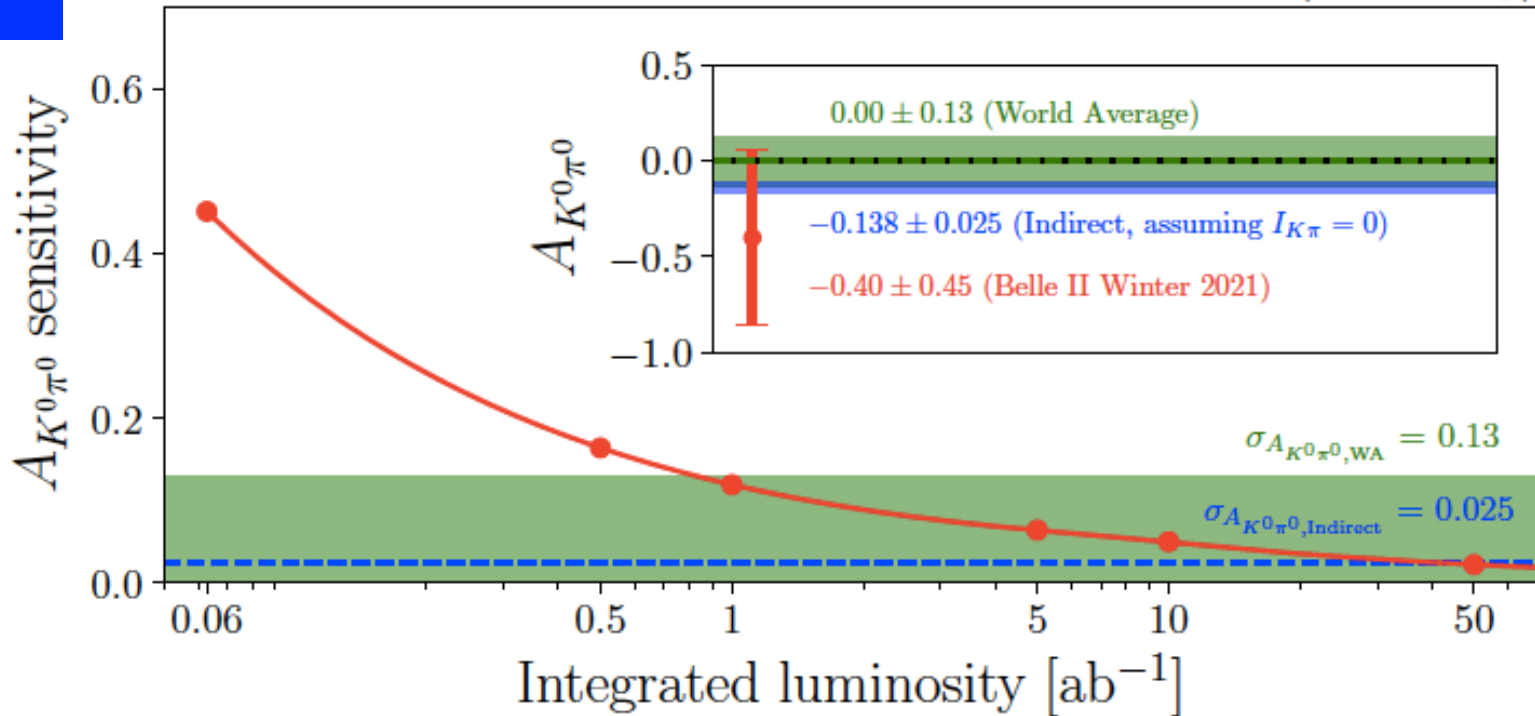
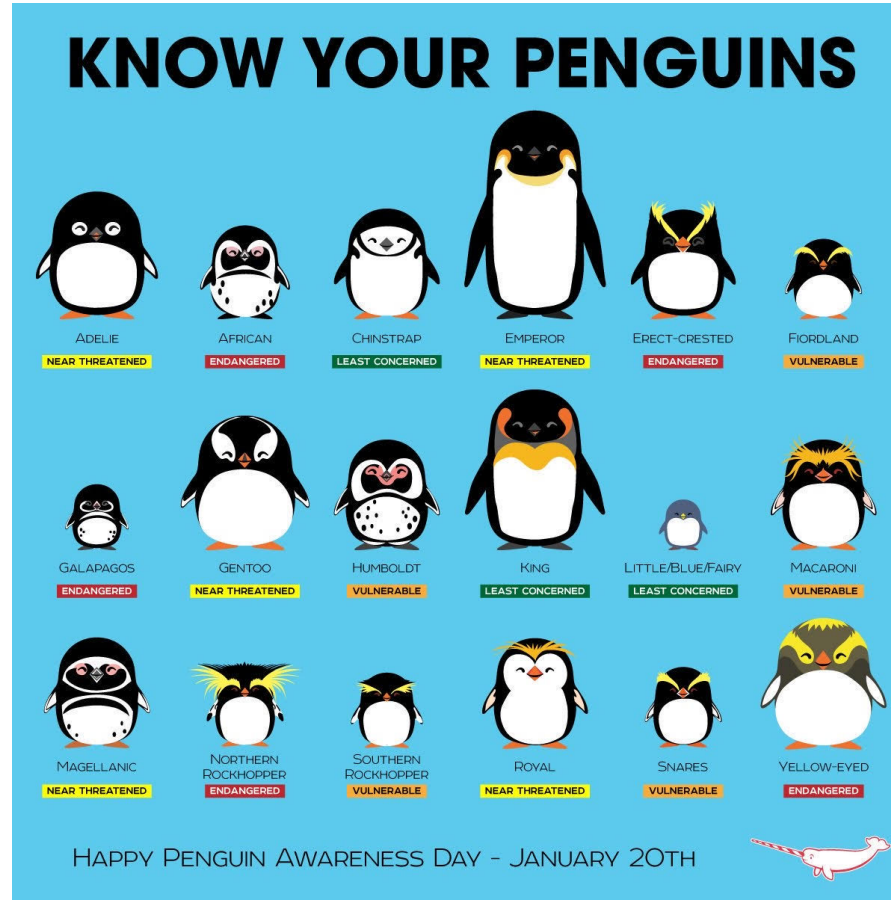
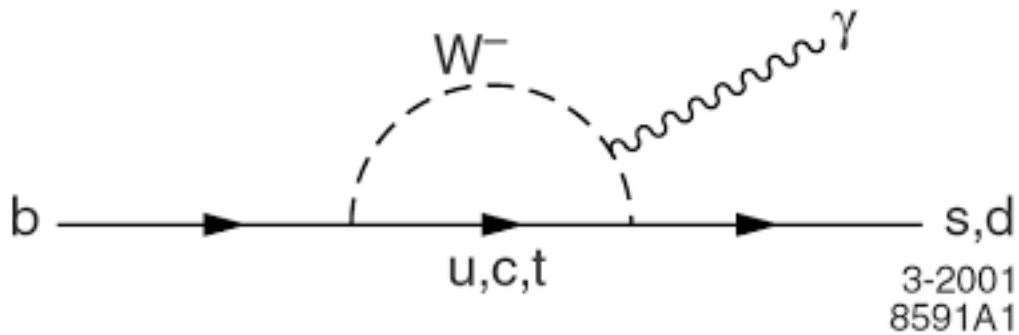


FIG. 5. The projected uncertainty on  $\mathcal{A}_{K^0 \pi^0}$  measurement. The inset panel shows the comparison of (red marker) the measurement reported here with (green band) the world average value, and (blue band) the indirect determination from Eq. 1 assuming  $I_{K\pi} = 0$  and world average values for the other inputs. The red curve in the main panel is Belle II's expected uncertainty on the  $\mathcal{A}_{K^0 \pi^0}$  measurement as a function of the integrated luminosity, while the green and blue dashed lines are the uncertainties of the world average value and of the indirect determination, respectively.

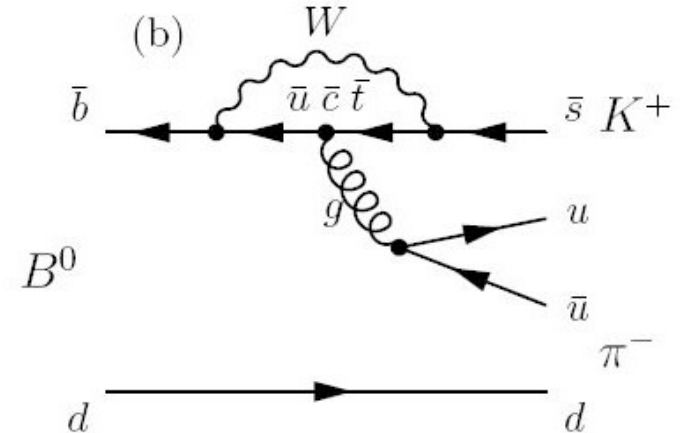
Recap:



“Radiative Penguin”



“Gluonic Penguin”

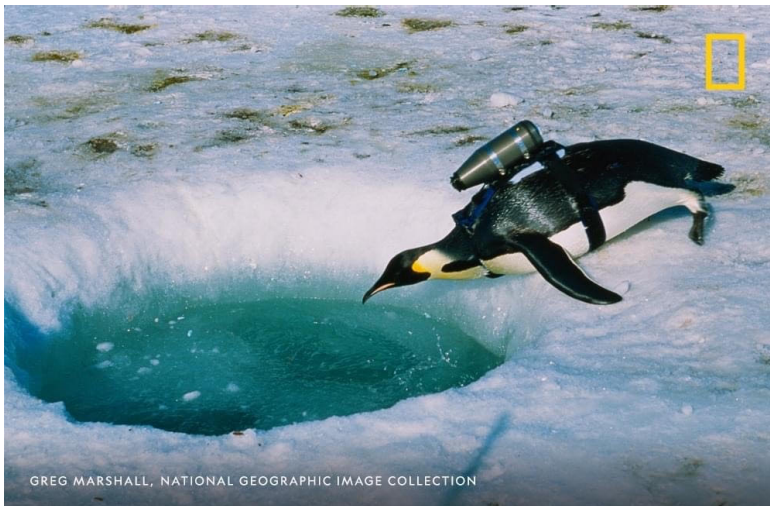




Now will describe some *speculations about how Belle II might discover new physics Beyond the SM (BSM)*

Research penguin

Photo Credit: National Geographic

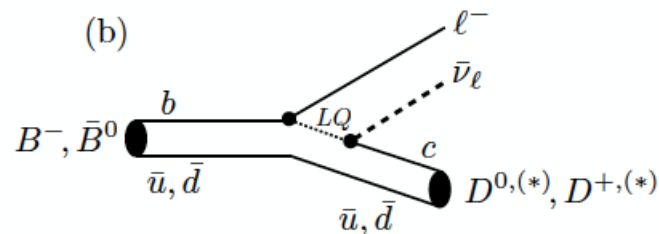
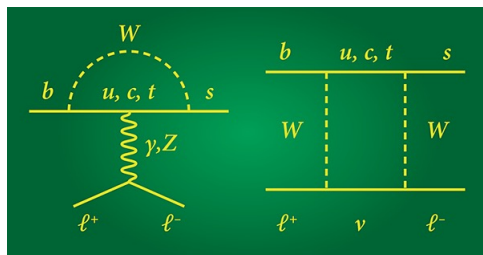


Sequoia National Forest



Exploring the unknown with  $b \rightarrow s$  “electroweak penguins”:  
(weak neutral current)

Discovering NP with  $b \rightarrow c \ell \nu$  “trees”:  
(weak charged current)





# A Snowmass Highlight (shown at Cincinnati and Seattle)



What happens at 1, 5, 50  $\text{ab}^{-1}$   
(or even 250  $\text{ab}^{-1}$  in the 2030's )?



**Belle II**  
Higher sensitivity to decays with photons and neutrinos (e.g.  $B \rightarrow K\nu\nu, \mu\nu$ ), inclusive decays, time dependent CPV in  $B_d, \tau$  physics.

**LHCb**  
Higher production rates for ultra rare B, D, & K decays, access to all b-hadron flavours (e.g.  $\Lambda_b$ ), high boost for fast  $B_s$  oscillations.

*Overlap in various key areas to verify discoveries.*

**Upgrades**  
*Most key channels will be stats. limited (not theory or syst.).*

Observable	2022 Belle(II), BaBar	2022 LHCb	Belle-II 5 $\text{ab}^{-1}$	Belle-II 50 $\text{ab}^{-1}$	LHCb 50 $\text{fb}^{-1}$	Belle-II 250 $\text{ab}^{-1}$	LHCb 300 $\text{fb}^{-1}$
$\sin 2\beta/\phi_1$	0.03	0.04	0.012	0.005	0.011	0.002	0.003
$\gamma/\phi_3$	11°	4°	4.7°	1.5°	1°	0.8°	0.35°
$\alpha/\phi_2$	4°	—	2°	0.6°	—	0.3°	—
$ V_{ub} / V_{cb} $	4.5%	6%	2%	1%	2%	< 1%	1%
$S_{CP}(B \rightarrow \eta' K_S^0)$	0.08	—	0.03	0.015	—	0.007	—
$A_{CP}(B \rightarrow \pi^0 K_S^0)$	0.15	—	0.07	0.04	—	0.018	—
$S_{CP}(B \rightarrow K^{*0} \gamma)$	0.32	—	0.11	0.035	—	0.015	—
$R(B \rightarrow K^{*} \ell^+ \ell^-)^\dagger$	0.26	0.12	0.09	0.03	0.022	0.01	0.009
$R(B \rightarrow D^{*} \tau \nu)$	0.018	0.026	0.009	0.0045	0.0072	<0.003	<0.003
$R(B \rightarrow D \tau \nu)$	0.034	—	0.016	0.008	—	<0.003	—
$\mathcal{B}(B \rightarrow \tau \nu)$	24%	—	9%	4%	—	2%	—
$\mathcal{B}(B \rightarrow K^{*} \nu \bar{\nu})$	—	—	25%	9%	—	4%	—
$\mathcal{B}(\tau \rightarrow e \gamma)$ UL	$42 \times 10^{-9}$	—	$22 \times 10^{-9}$	$6.9 \times 10^{-9}$	—	$3.1 \times 10^{-9}$	—
$\mathcal{B}(\tau \rightarrow \mu \mu \mu)$ UL	$21 \times 10^{-9}$	$46 \times 10^{-9}$	$3.6 \times 10^{-9}$	$0.36 \times 10^{-9}$	$1.1 \times 10^{-9}$	$0.07 \times 10^{-9}$	$5 \times 10^{-9}$

The dagger refers to a measurement in the range  $1 < q^2 < 6 \text{ GeV}^2/c^2$

Consideration of further luminosity upgrade and electron polarization capability of SuperKEKB are started for ultimate new physics searches with heavy flavor quarks and leptons including  $\tau$  lepton  $g - 2$  in the light of muon  $g - 2$  anomaly [28].

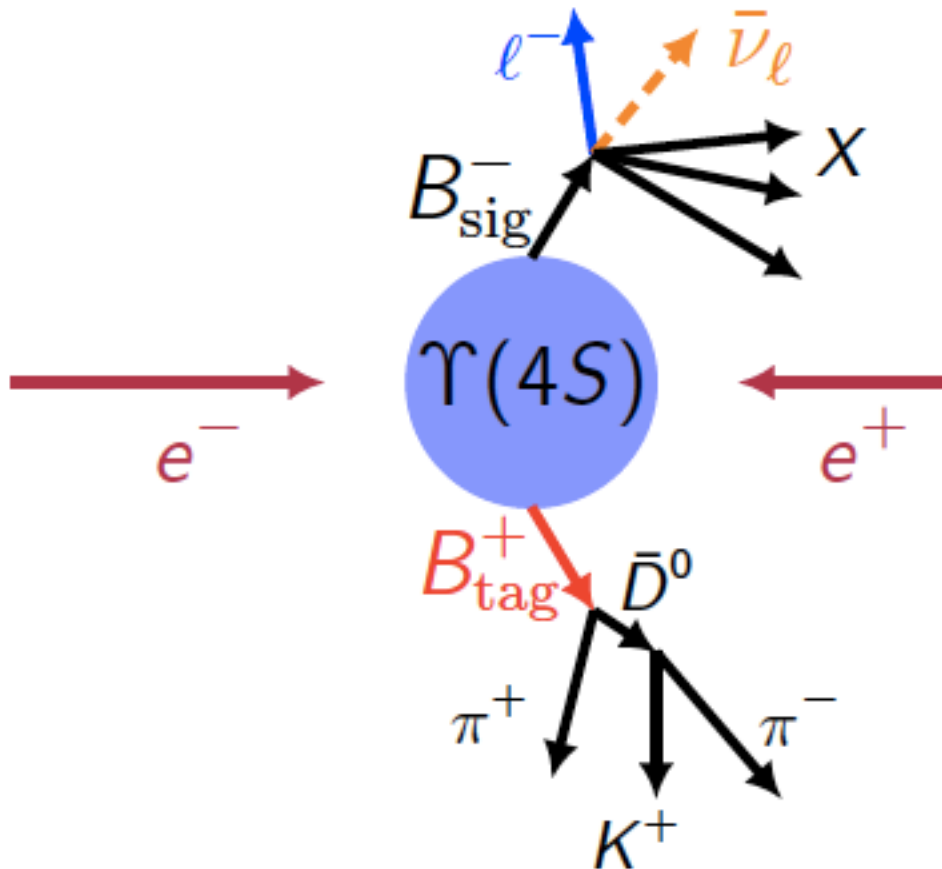
JAHEP report to Snowmass: Arxiv 2203:13979

# Some **critical** Belle II capabilities for flavor (B) physics

Full and equally strong capabilities for electrons and muons

**Photons**,  $K_S$ 's with excellent resolution and efficiency

Neutrinos via “**missing energy**” and missing momentum. **Hermeticity.**



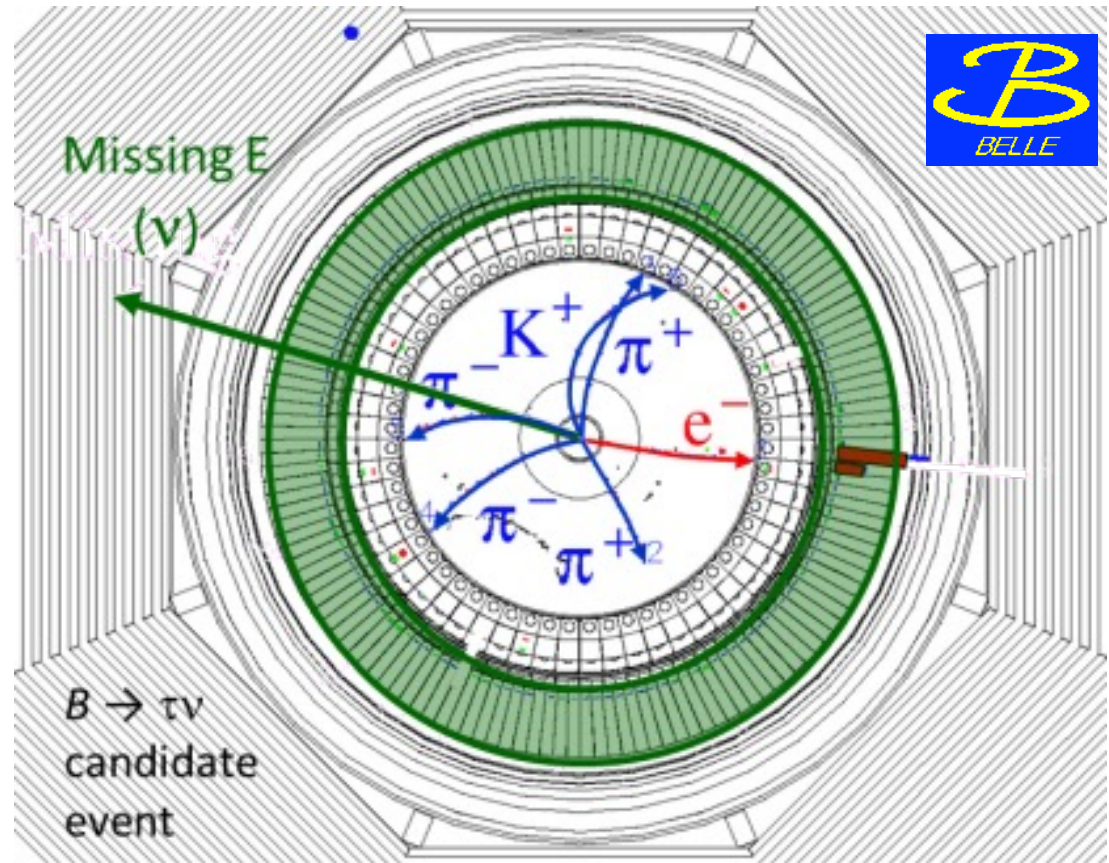
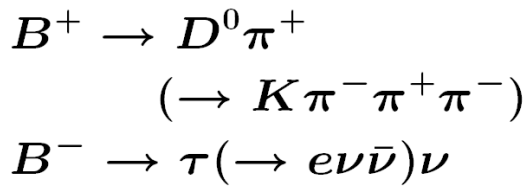
<https://arxiv.org/abs/2008.06096>

This is now called **FEI**  
“Full Event Interpretation”  
and uses large numbers of  
tag modes via a **BDT**  
(Boosted Decision Tree).

*Clean but efficiency  $\varepsilon \sim 0.5\%$*

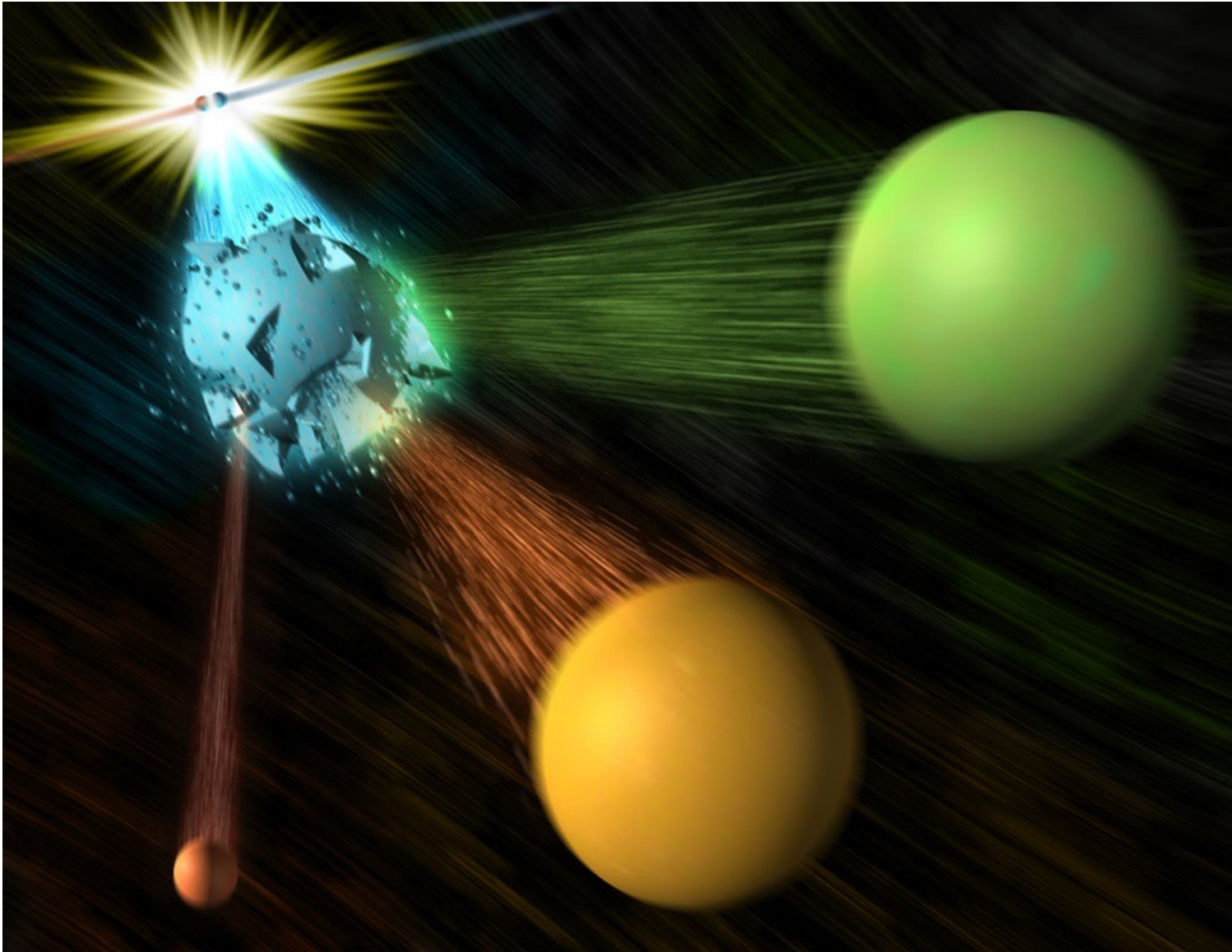
T. Keck et al., Comput. Softw. Big Sci. 3, 6  
(2019), arXiv:1807.08680 [hep-ex].

Example of a Missing Energy Decay ( $B \rightarrow \tau \nu$ ) in old Belle Data  
(recorded before 2010)



*The clean  $e^+e^-$  environment (and the CsI(Tl) crystal calorimeter) makes this possible.*

# Possible breakdown of lepton universality in $B \rightarrow D^{(*)} \tau \nu$



Let's try to understand this picture of the production process (EM) and a weak decay

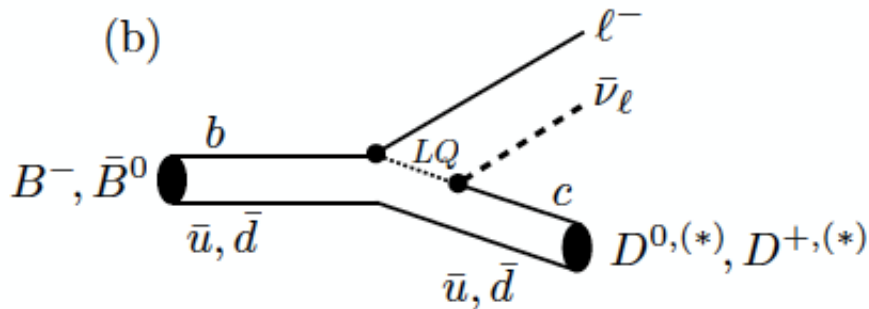
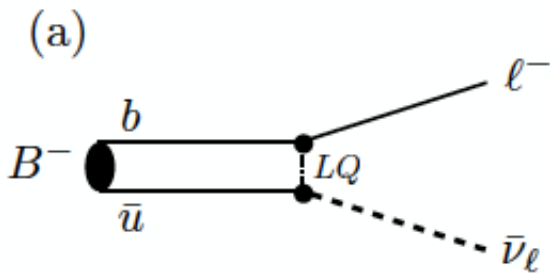
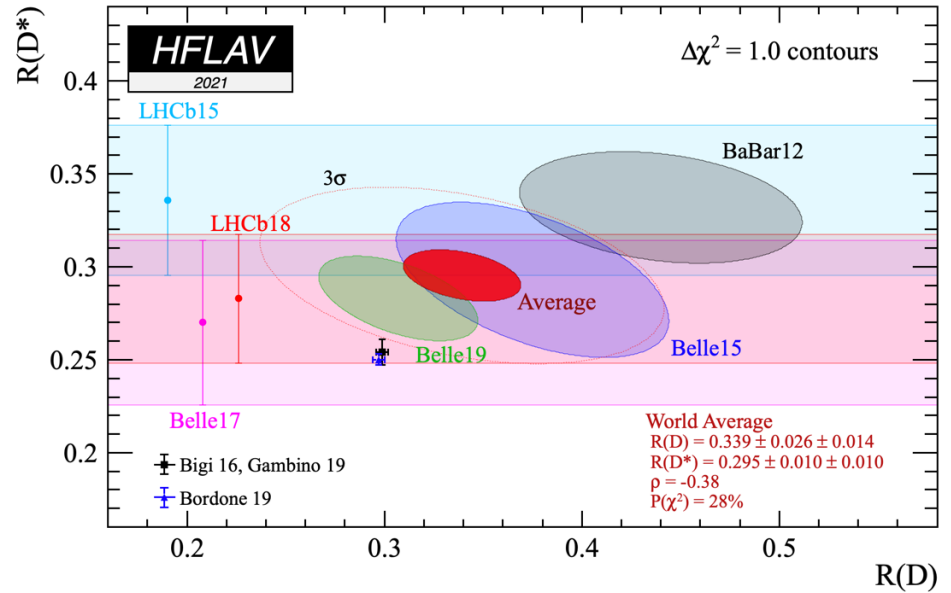


# $B \rightarrow D^{(*)} \tau \nu$ , possible breakdown of **lepton universality**

$$R_D^{(*)} = \frac{\mathcal{B}(B \rightarrow D^{(*)} \tau \nu_\tau)}{\mathcal{B}(B \rightarrow D^{(*)} \ell \nu_\ell)}$$

This is NP in the weak  $b \rightarrow c$  charged current

Some new physics possibilities  
(**leptoquarks (LQ)**, charged Higgs  
type 3 etc..):



With current data from Belle, LHCb and BaBar:  
 Evidence of **lepton universality breakdown** in semileptonic B decays with  $\tau$  leptons. Last Belle measurement (2019) with semileptonic tags brings down the WA discrepancy from  $4 \rightarrow 3.4\sigma$

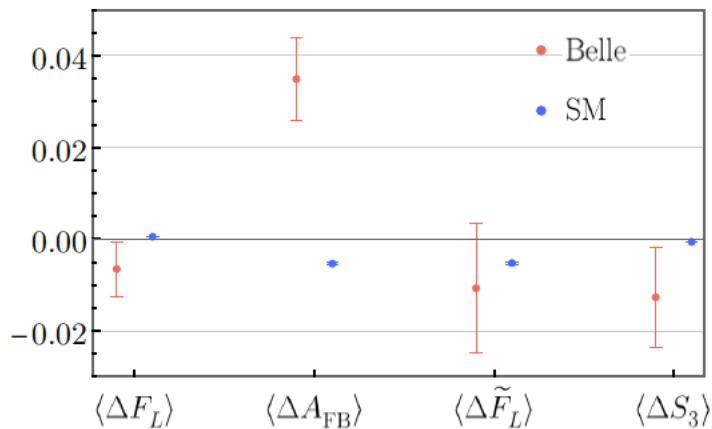
	$5 \text{ ab}^{-1}$	$50 \text{ ab}^{-1}$
$R_D$	$(\pm 6.0 \pm 3.9)\%$	$(\pm 2.0 \pm 2.5)\%$
$R_{D^*}$	$(\pm 3.0 \pm 2.5)\%$	$(\pm 1.0 \pm 2.0)\%$
$P_\tau(D^*)$	$\pm 0.18 \pm 0.08$	$\pm 0.06 \pm 0.04$



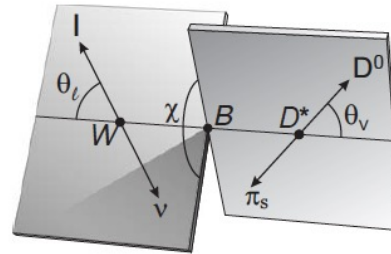
# Hot and fairly New: $\Delta A_{FB}$ in $b \rightarrow c \ell \nu$ (LFU violation)



$$\Delta A_{FB}(B \rightarrow D^{*+} \ell \nu) = A_{FB}(B \rightarrow D^{*+} \mu^{-} \bar{\nu}) - A_{FB}(B \rightarrow D^{*+} e^{-} \bar{\nu})$$

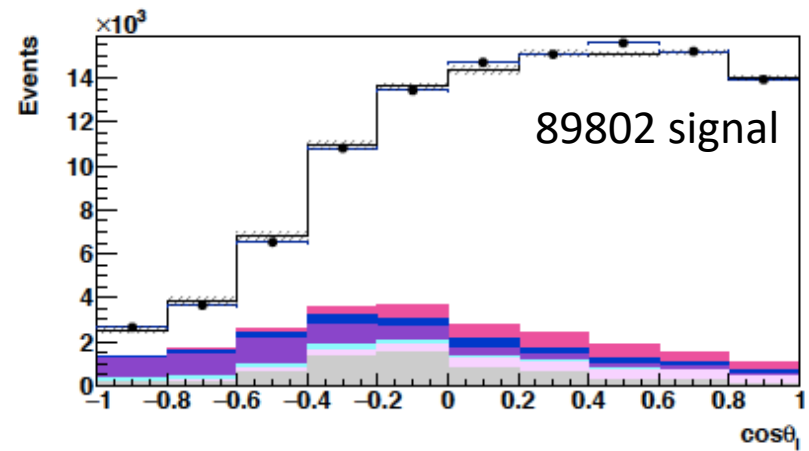
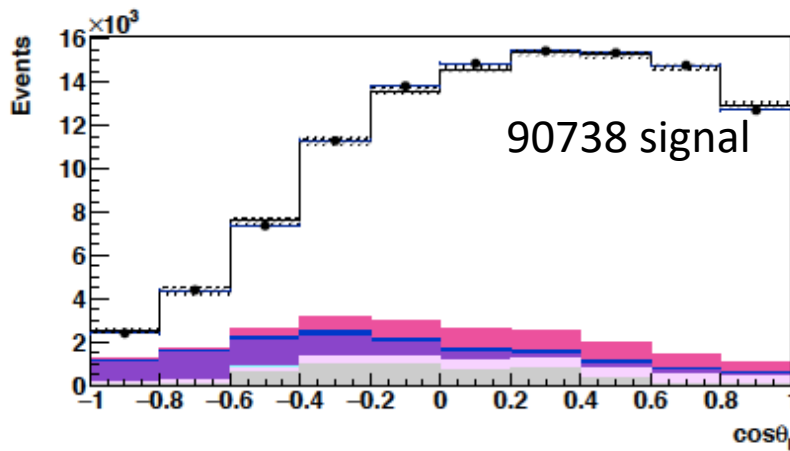


$\sim 4\sigma$  deviation in Eur. Phys. J.C. 81 (2021). Theoretical meta-analysis of  $0.71 \text{ ab}^{-1}$  of Belle data from Arxiv 1809.03290 (E. Waheed et al (Belle)).



“untagged”, no FEI used.

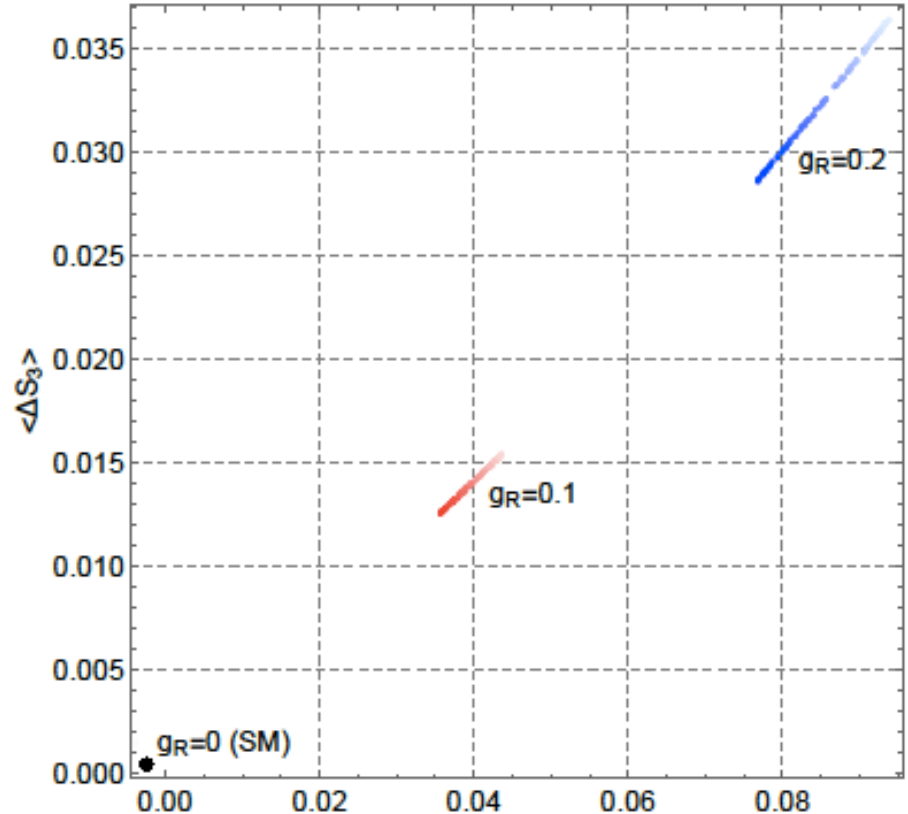
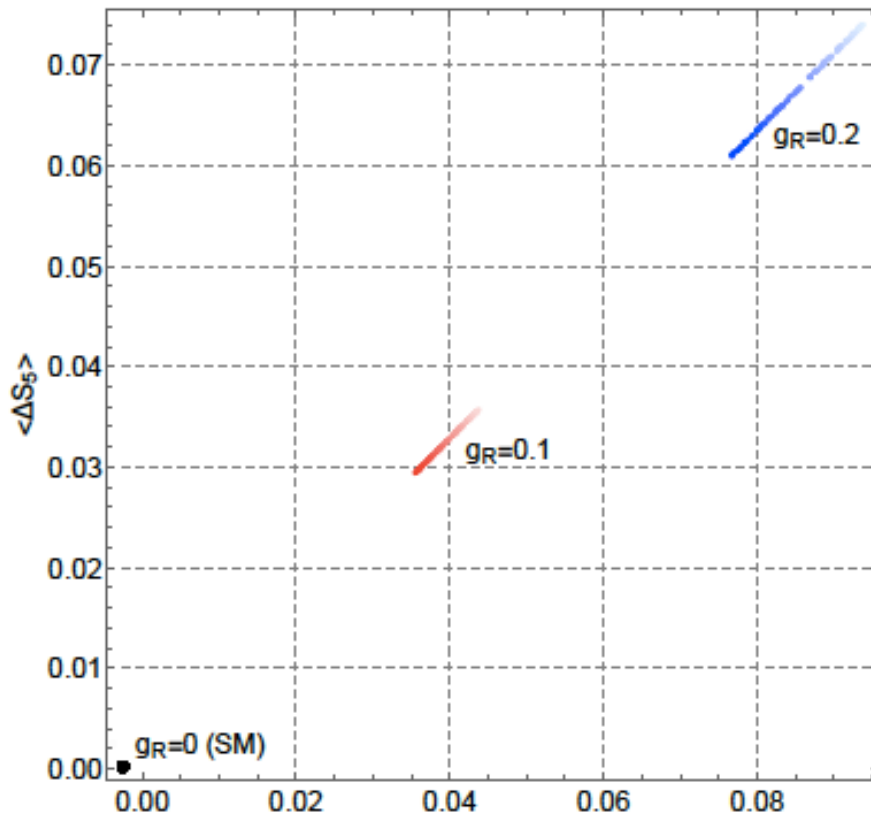
FIG. 3. Definition of the angles  $\theta_\ell$ ,  $\theta_\nu$  and  $\chi$  for the decay



Evidence for NP couplings in  $b \rightarrow c \mu \nu$  ?

$$\Delta A_{FB}(B \rightarrow D^{*+} \ell \nu) = A_{FB}(B \rightarrow D^{*+} \mu^{-} \nu) - A_{FB}(B \rightarrow D^{*+} e^{-} \nu)$$

NP implies **correlated** angular asymmetries ( $\Delta A_{FB}$  vs  $\Delta S_5$  or  $\Delta S_3$ ).



<https://arxiv.org/abs/2203.07189>

Snowmass WP

<https://arxiv.org/abs/2206.11283>

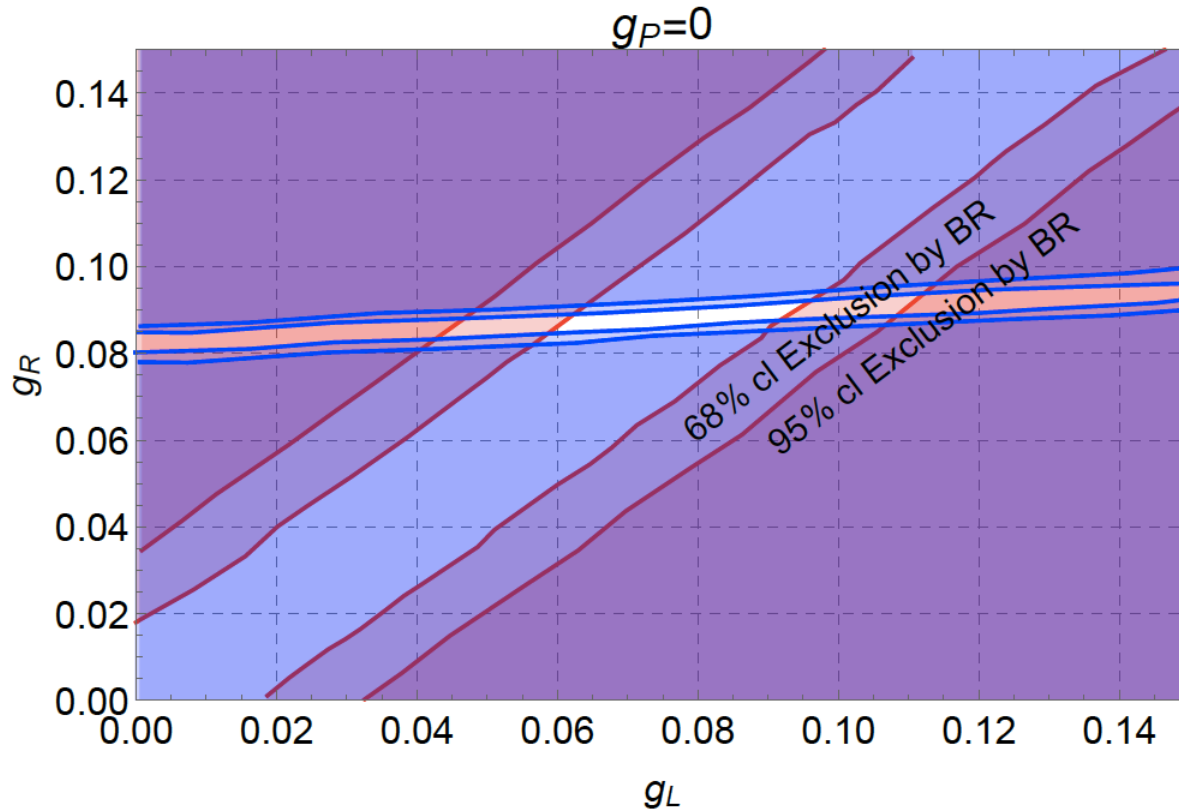
PRD version

# Snowmass Highlight:

$$\Delta A_{FB}(B \rightarrow D^{*+} \ell \nu) = A_{FB}(B \rightarrow D^{*+} \mu^{-} \nu) - A_{FB}(B \rightarrow D^{*+} e^{-} \nu)$$

N.B. Form Factor uncertainties cancel out in  $\Delta$  variables

+ constraints on NP coupling parameters @ 250  $\text{ab}^{-1}$



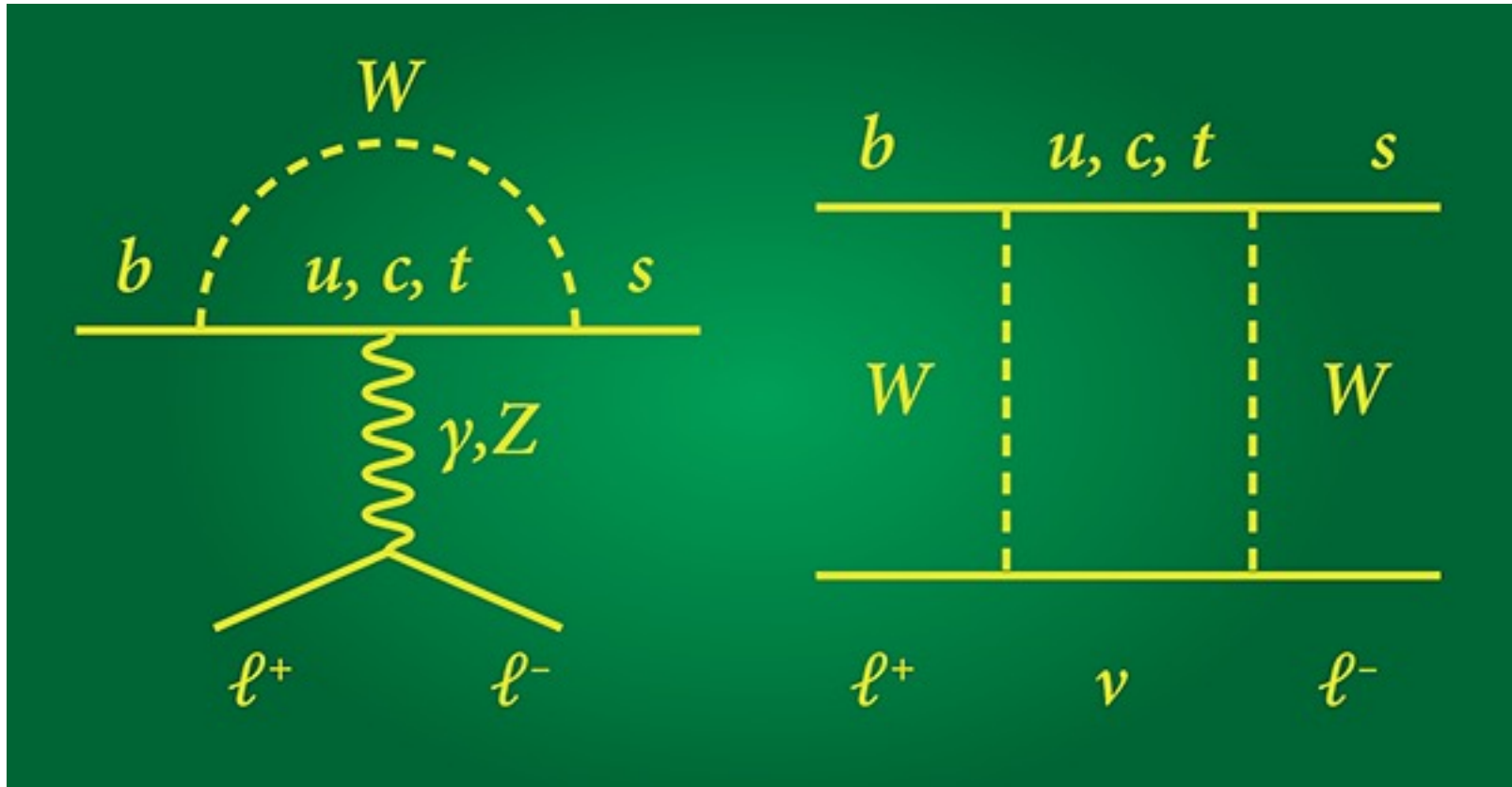
Angular asymmetries provide a tighter constraint on NP LFUV couplings (right-handed V+A, extra left-handed V-A and pseudo-scalar couplings).

<https://arxiv.org/abs/2203.07189>





# Lepton Universality Tests in $b \rightarrow s \ell^+ \ell^-$ transitions

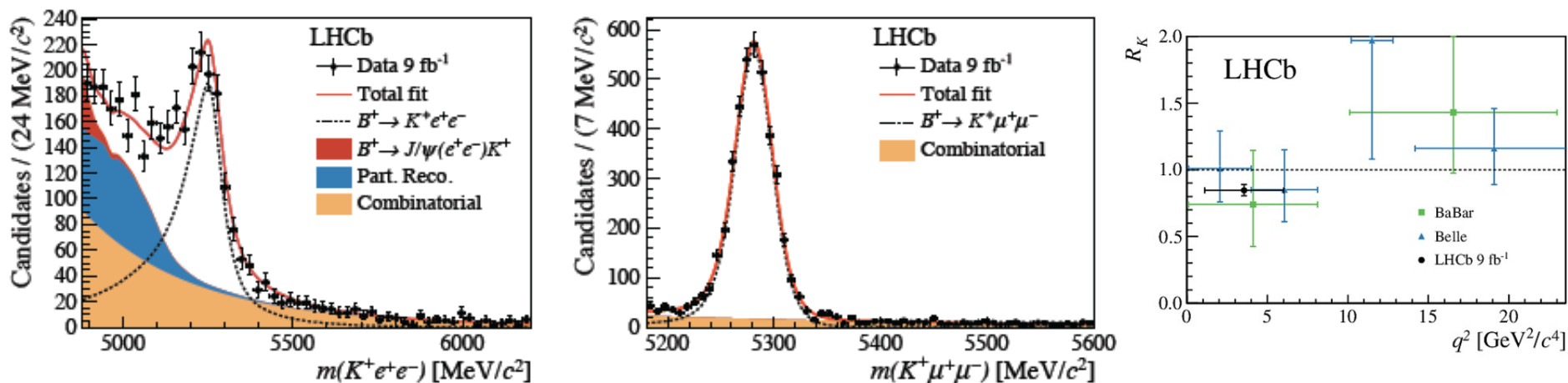


“Electroweak Penguin”

“Box”

# Possible breakdown of **Lepton Universality** in $b \rightarrow s l^+ l^-$ transitions by the LHCb experiment at CERN, reported in 2021.

<https://arxiv.org/abs/2103.11769>, published in Nature



$$R_K = \frac{BF(B^+ \rightarrow K^+ \mu^+ \mu^-)}{BF(B^+ \rightarrow J/\psi(\rightarrow \mu^+ \mu^-) K^+)} \bigg/ \frac{BF(B^+ \rightarrow K^+ e^+ e^-)}{BF(B^+ \rightarrow J/\psi(\rightarrow e^+ e^-) K^+)}$$

$$R_K(1.1 < q^2 < 6.0 \text{ GeV}^2/c^4) = 0.846^{+0.042+0.013}_{-0.039-0.012}$$

**<1 (lepton universality prediction)**

And thus this *might indicate* the breakdown of the Standard Model of Particle Physics ( $3.1 \sigma$ )

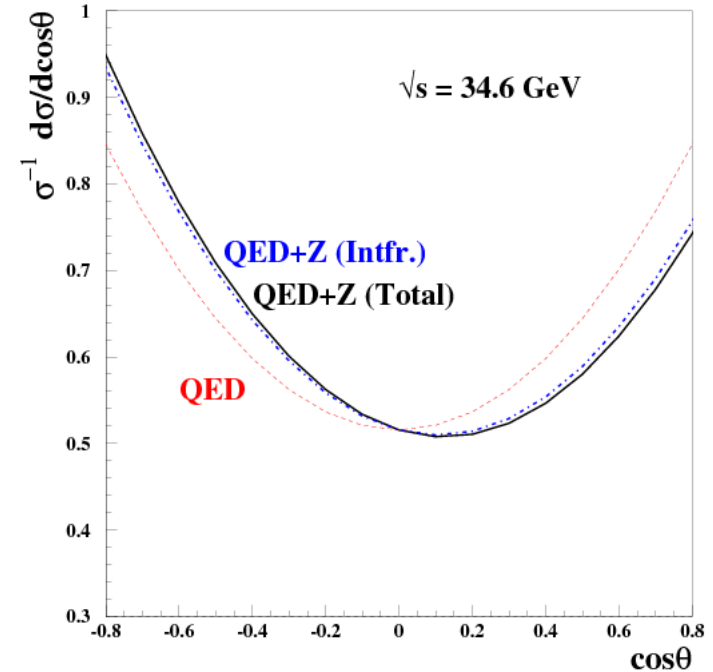
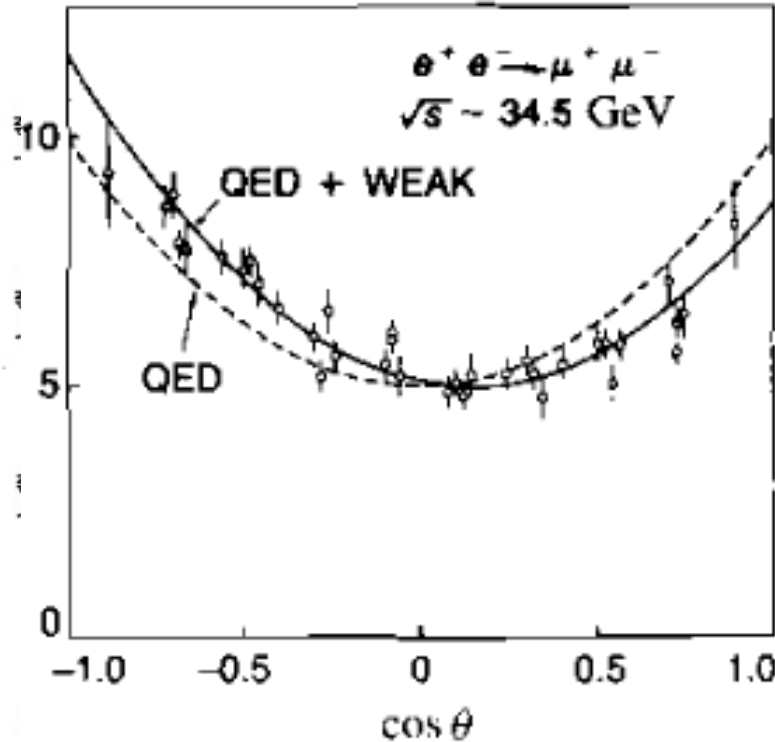
Note:  $q^2 = M^2(l^+ l^-)$



# High Energy Physics History: finding NP in $A_{FB}$ (using interference,



SLAC

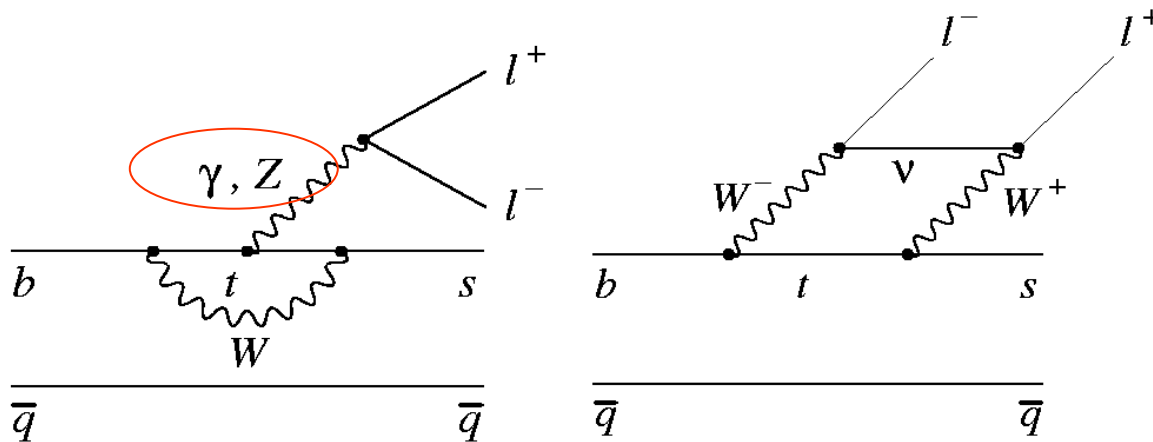
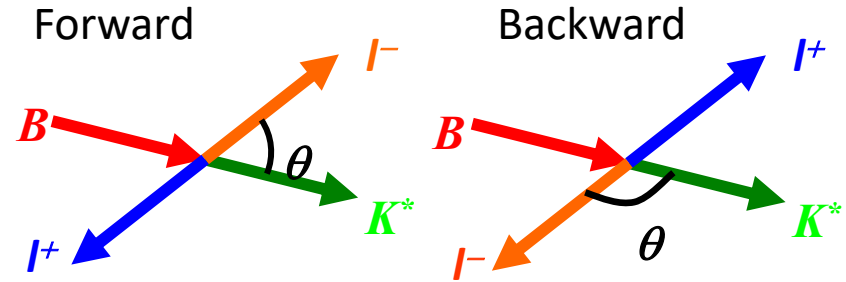


Conclusion: There is a Z boson at higher energy *even though* colliders of the time did not have enough  $\sqrt{s}$  to produce it



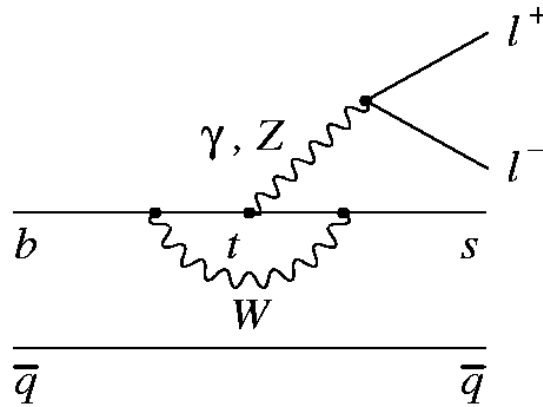
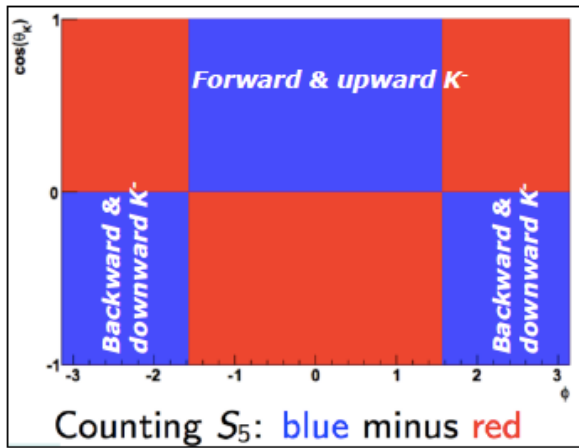
$$A_{\text{FB}}(B \rightarrow K^* l^+ l^-)(q^2)$$

The SM forward-backward asymmetry in  $b \rightarrow s l^+ l^-$  can arise from the interference between  $\gamma$  and  $Z^0$  contributions.



Note that all the heavy particles of the SM (W, Z, top) enter in this decay.

# More on $A_{FB}(B \rightarrow K^* l^+ l^-)(q^2)$ and $S_5(q^2)$



Can in effect vary  $v_s$  for NP

$A_{FB}$  depends on  $q^2 = M^2(l^+ l^-)$

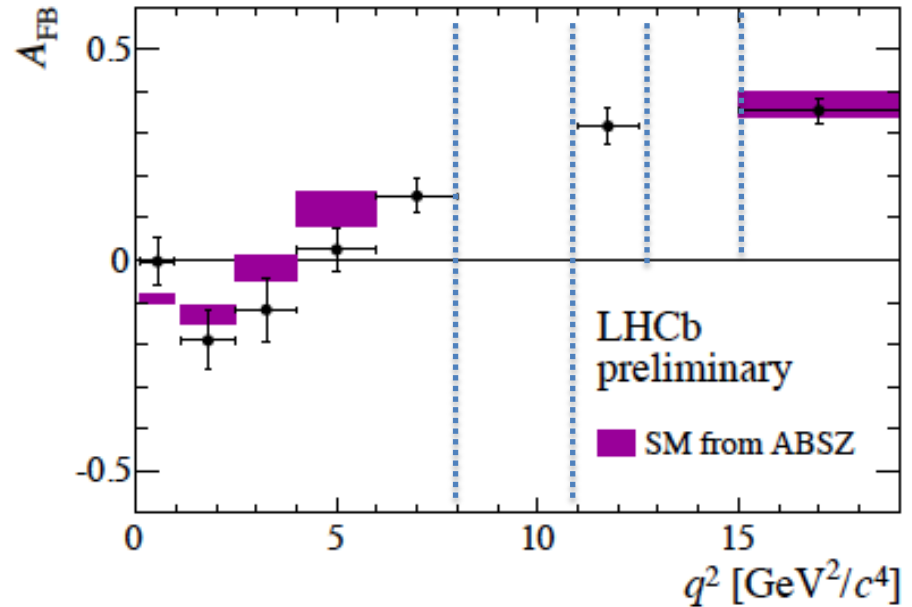
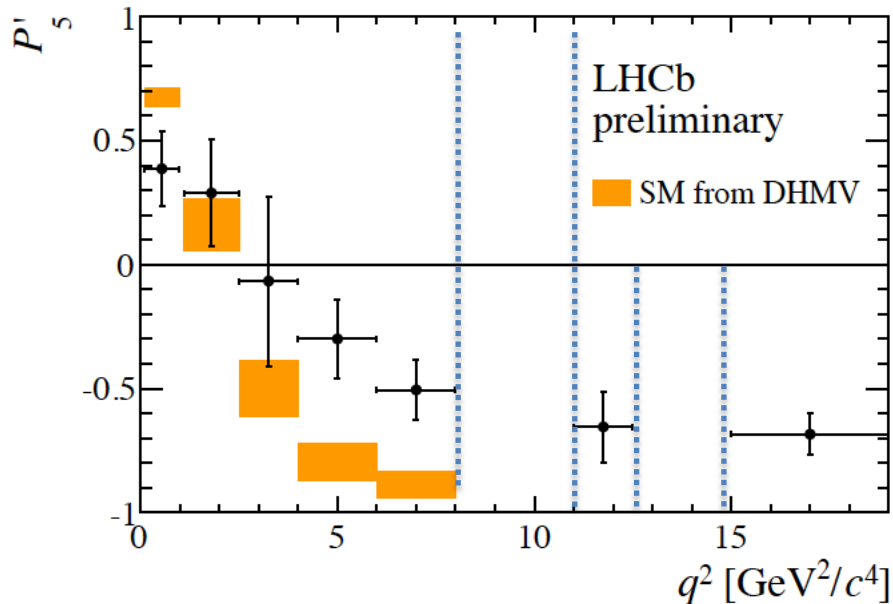
$$A_{FB}(B \rightarrow K^* l^+ l^-) = -C_{10} \xi(q^2) \left[ \text{Re}(C_9) F_1 + \frac{1}{q^2} C_7 F_2 \right]$$

G. Burdman, Phys.Rev. D57 (1998) 4254

The “zero-crossing” of  $A_{FB}$  depends only on a ratio of form factors and is a relatively *clean* observable.

# LHCb $3fb^{-1}$ results on $B \rightarrow K^* \mu^+ \mu^- (q^2)$

A different angular asymmetry, involving  $\chi$



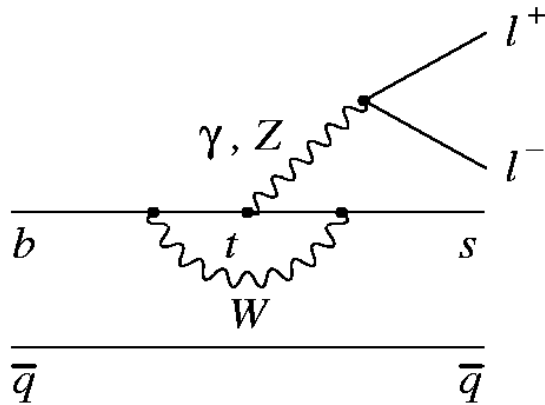
“The  $P_5'$  measurements are only compatible with the SM prediction at a level of  $3.7\sigma$ ....A mild tension can also be seen in the  $A_{FB}$  distribution, where the measurements are systematically  $\leq 1\sigma$  below the SM prediction in the region  $1.1 < q^2 < 6.0$  GeV<sup>2</sup>”

*These angular asymmetries persist in 2022*

*Comment on  $B \rightarrow K^* \mu^+ \mu^- (q^2)$  ( $b \rightarrow s l^+ l^-$ )*

*Is HEP History repeating itself?* [*but make sure this is not a SM resonance/ non-factorizable/long distance effect.*]

Why would NP appear first in this mode (and not others) ?



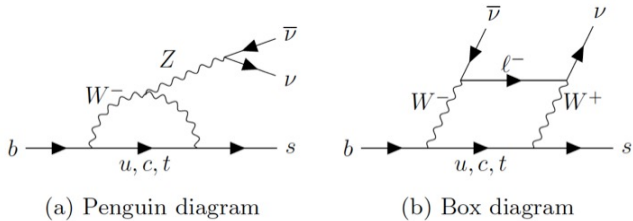
Possible answer: All the heavy particles of the SM ( $t$ ,  $W$ ,  $Z$ ) and maybe NP (except the Higgs) appear here. Sensitive to NP **via QM interference** (linear effects). ++*Lepton Flavor Universality Violation*



# Feynman Diagrams and Model Building



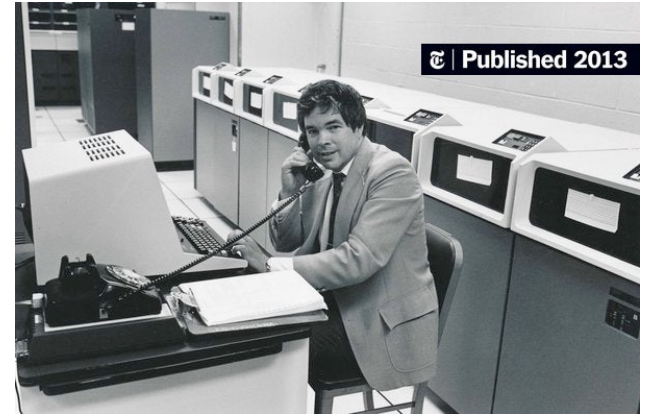
Feynman family and diagrams



# Paradigm shift



# Effective Field Theory → Wilson Coefficients



Ken Wilson ("Wilson coefficients")



$C_7, C_9, C_{10}$

## New Physics Couplings in $b \rightarrow s$

The effective Hamiltonian for  $b \rightarrow s$  transitions can be written as

$$\mathcal{H}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \frac{e^2}{16\pi^2} \sum_i (C_i O_i + C'_i O'_i) + \text{h.c.}$$

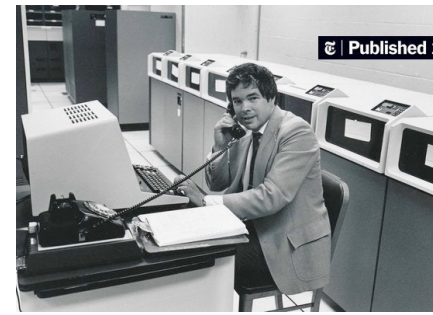
and we consider NP effects in the following set of dimension-6 operators,

$$\begin{aligned} O_9 &= (\bar{s}\gamma_\mu P_L b)(\bar{\ell}\gamma^\mu \ell), & O'_9 &= (\bar{s}\gamma_\mu P_R b)(\bar{\ell}\gamma^\mu \ell), \\ O_{10} &= (\bar{s}\gamma_\mu P_L b)(\bar{\ell}\gamma^\mu \gamma_5 \ell), & O'_{10} &= (\bar{s}\gamma_\mu P_R b)(\bar{\ell}\gamma^\mu \gamma_5 \ell). \end{aligned}$$

The primes are NP right-handed couplings.



# New Physics Couplings in $b \rightarrow s$



Ken Wilson

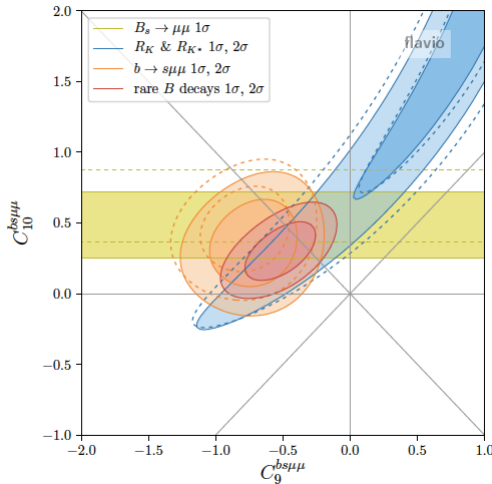
The effective Hamiltonian for  $b \rightarrow s$  transitions can be written as

$$\mathcal{H}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}}V_{tb}V_{ts}^* \frac{e^2}{16\pi^2} \sum_i (C_i O_i + C'_i O'_i) + \text{h.c.}$$

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The primes are right-handed couplings.

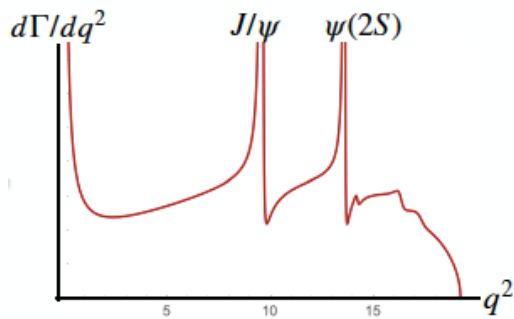


		$b \rightarrow s\mu\mu$		LFU, $B_s \rightarrow \mu\mu$		all rare $B$ decays	
Wilson coefficient		best fit	pull	best fit	pull	best fit	pull
NP errors	$C_9^{bs\mu\mu}$	$-0.75^{+0.22}_{-0.23}$	$3.4\sigma$	$-0.74^{+0.20}_{-0.21}$	$4.1\sigma$	$-0.73^{+0.15}_{-0.15}$	$5.2\sigma$
	$C_{10}^{bs\mu\mu}$	$+0.42^{+0.23}_{-0.24}$	$1.7\sigma$	$+0.60^{+0.14}_{-0.14}$	$4.7\sigma$	$+0.54^{+0.12}_{-0.12}$	$4.7\sigma$
	$C_9^{bs\mu\mu} = -C_{10}^{bs\mu\mu}$	$-0.53^{+0.13}_{-0.13}$	$3.7\sigma$	$-0.35^{+0.08}_{-0.08}$	$4.6\sigma$	$-0.39^{+0.07}_{-0.07}$	$5.6\sigma$
SM errors	$C_9^{bs\mu\mu}$	$-0.88^{+0.22}_{-0.21}$	$3.7\sigma$	$-0.74^{+0.20}_{-0.21}$	$4.1\sigma$	$-0.78^{+0.15}_{-0.15}$	$5.3\sigma$
	$C_{10}^{bs\mu\mu}$	$+0.44^{+0.21}_{-0.21}$	$2.1\sigma$	$+0.60^{+0.14}_{-0.14}$	$4.7\sigma$	$+0.54^{+0.12}_{-0.12}$	$4.8\sigma$
	$C_9^{bs\mu\mu} = -C_{10}^{bs\mu\mu}$	$-0.58^{+0.17}_{-0.18}$	$3.6\sigma$	$-0.35^{+0.08}_{-0.08}$	$4.6\sigma$	$-0.39^{+0.07}_{-0.07}$	$5.5\sigma$

$C_9 : >5\sigma$   
from the SM

Altmanshofer, Stangel fit to all data (mostly LHCb)  
<https://arxiv.org/pdf/2103.13370.pdf>

*Be very careful about  $5\sigma$  New Physics (NP) claims, leftmost column assumes minimal QCD, resonance effects in angular asymmetries and  $q^2$  distribution.*

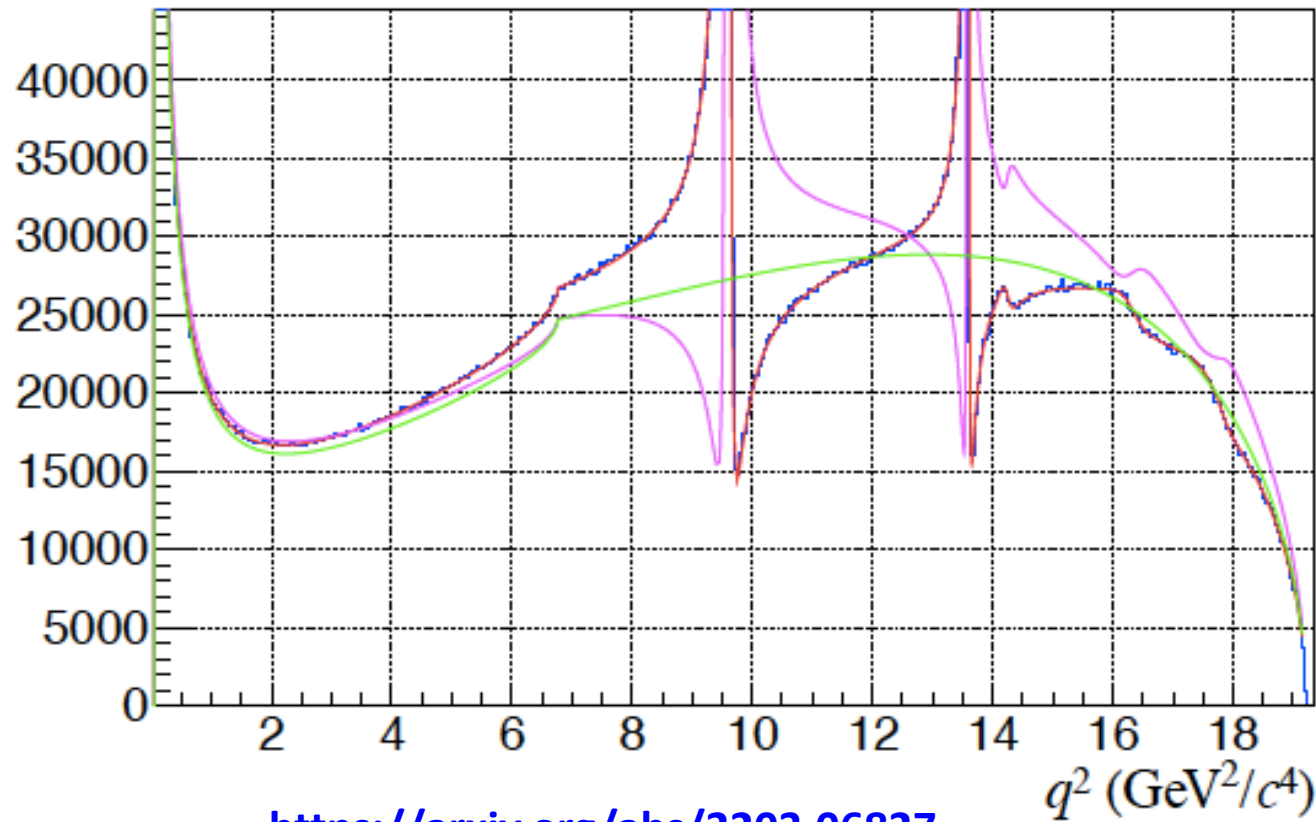


The green curve is the short-distance  $b \rightarrow s l^+ l^-$  contribution. The non-factorizable phase is an uncertainty.

There are also uncertainties in  $B \rightarrow K^*$  form factors.



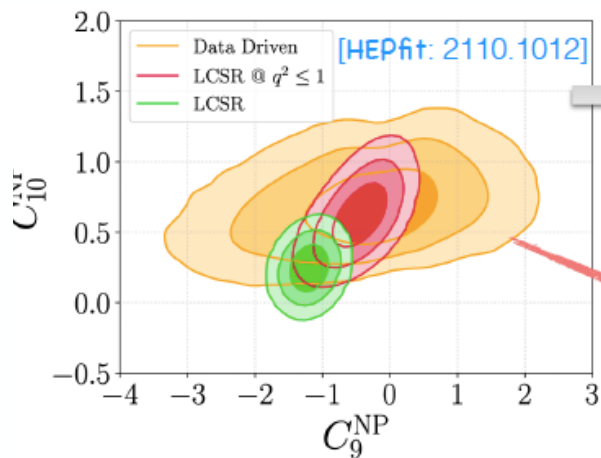
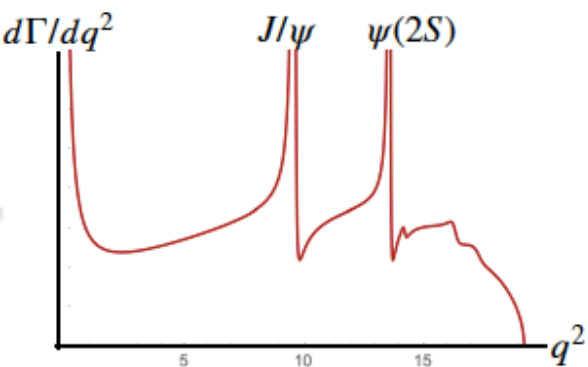
Alexei Sibidanov



<https://arxiv.org/abs/2203.06827>

FIG. 21. The  $q^2$  distribution of  $\bar{B} \rightarrow \bar{K}^* \mu^+ \mu^-$  decay in the presence of  $c\bar{c}$  resonances. The histogram is the result from the **EvtGen** generator, the green curve shows the result of the likelihood integration without resonances, and the red curve is the result of the likelihood integration when resonances are included. The contribution of these resonances (and non-factorizable effects) will be a limiting uncertainty in the extraction of NP Wilson coefficients from  $B \rightarrow K^* \mu^+ \mu^-$ .

# Angular analysis



Global fits to current  $b \rightarrow s\ell\ell$  data with three different treatments to parametrize charm-loop effects

Hadronic uncertainties due to long distance physics can overshadow new physics

Further cancellation required for angular observables

Resonant  $b \rightarrow c \bar{c} s$  contributions

☑ use different lepton flavors:  $\frac{S_i^{b \rightarrow s\mu\mu}(q^2)}{\Gamma_f^{b \rightarrow s\mu\mu}(q^2)} - \frac{S_i^{b \rightarrow see}(q^2)}{\Gamma_f^{b \rightarrow see}(q^2)}$

Q-observables  
[Capdevila et. al, '16]  
[Belle: '16]

☑ Directly extract  $\Delta C_9 = \delta C_9^{b \rightarrow s\mu\mu} - \delta C_9^{b \rightarrow see}$

The solution to the problem is the Delta ( $\Delta$ ) Observables.





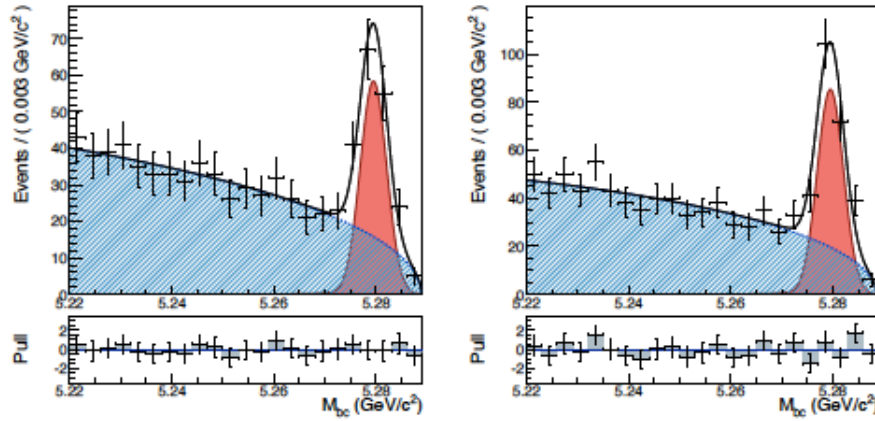


FIG. 1. Distribution of the beam-energy constrained mass for selected  $B \rightarrow K^* e^+ e^-$  (left) and  $B \rightarrow K^* \mu^+ \mu^-$  (right). Combinatorial background (shaded blue), signal (red filled) and total (solid) fit functions are superimposed on the data points

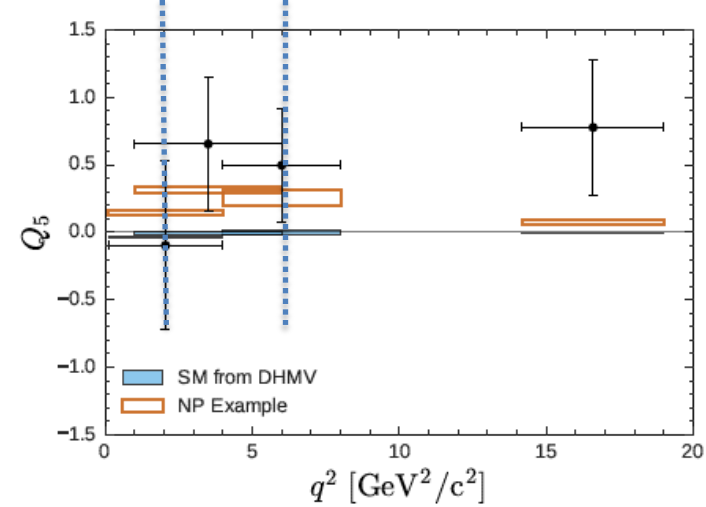
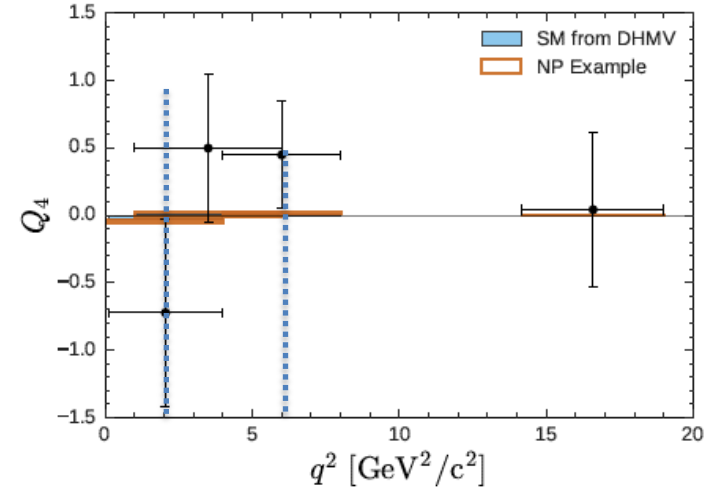
$$\Delta P'_4 = P'_4(B \rightarrow K^* \mu^+ \mu^-) - P'_4(B \rightarrow K^* e^+ e^-)$$

a.k.a.  $Q_4$

$$\Delta P'_5 = P'_5(B \rightarrow K^* \mu^+ \mu^-) - P'_5(B \rightarrow K^* e^+ e^-)$$

a.k.a.  $Q_5$

Belle has tried out some of the  $\Delta$  Observables with  $0.7 \text{ ab}^{-1}$

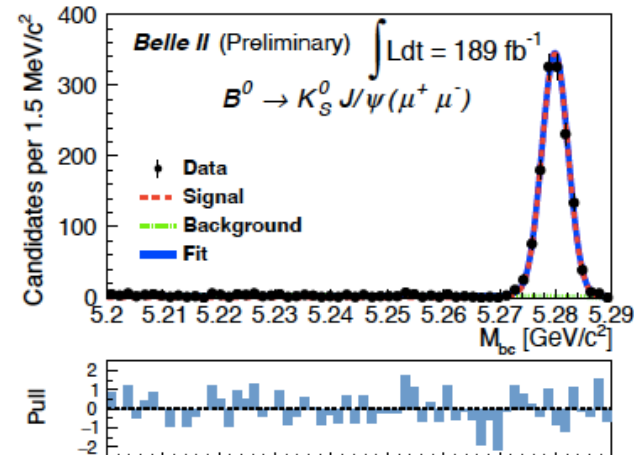
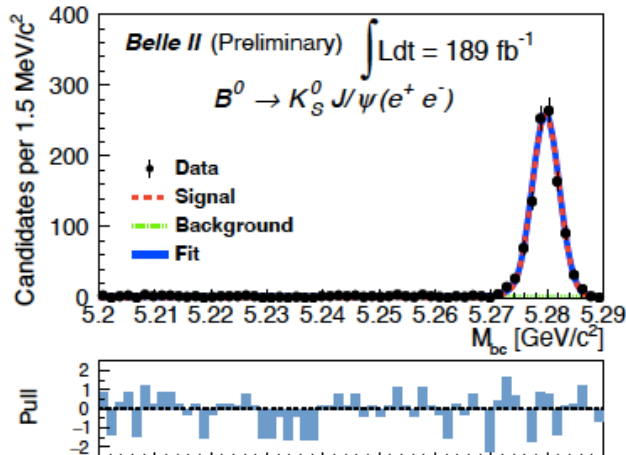
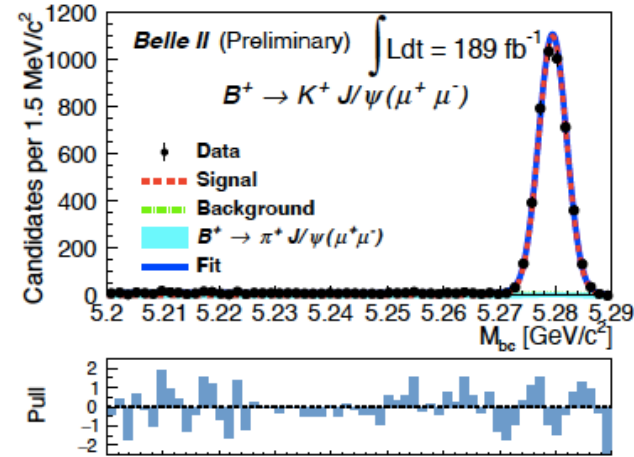
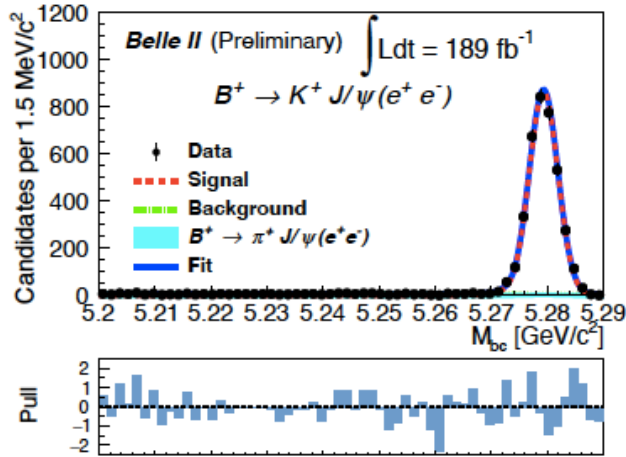


S. Wehle, C. Niebuhr, S. Yashchenko, et al.  
(Belle Collaboration), [PRL118, 111801 \(2017\)](#)



# Belle II is gearing up for e vs $\mu$ lepton universality tests (e.g. $B \rightarrow K J/\psi$ , $\psi \rightarrow l^+ l^-$ from recent data, $190 \text{ fb}^{-1}$ )

Includes  
brems  
recovery  
for  
electrons



<https://arxiv.org/abs/2207.11275>

$$R_{K^+}(J/\psi) = 1.009 \pm 0.022 \pm 0.008$$
$$R_{K^0}(J/\psi) = 1.042 \pm 0.042 \pm 0.008$$

## Reminder and Motivation:

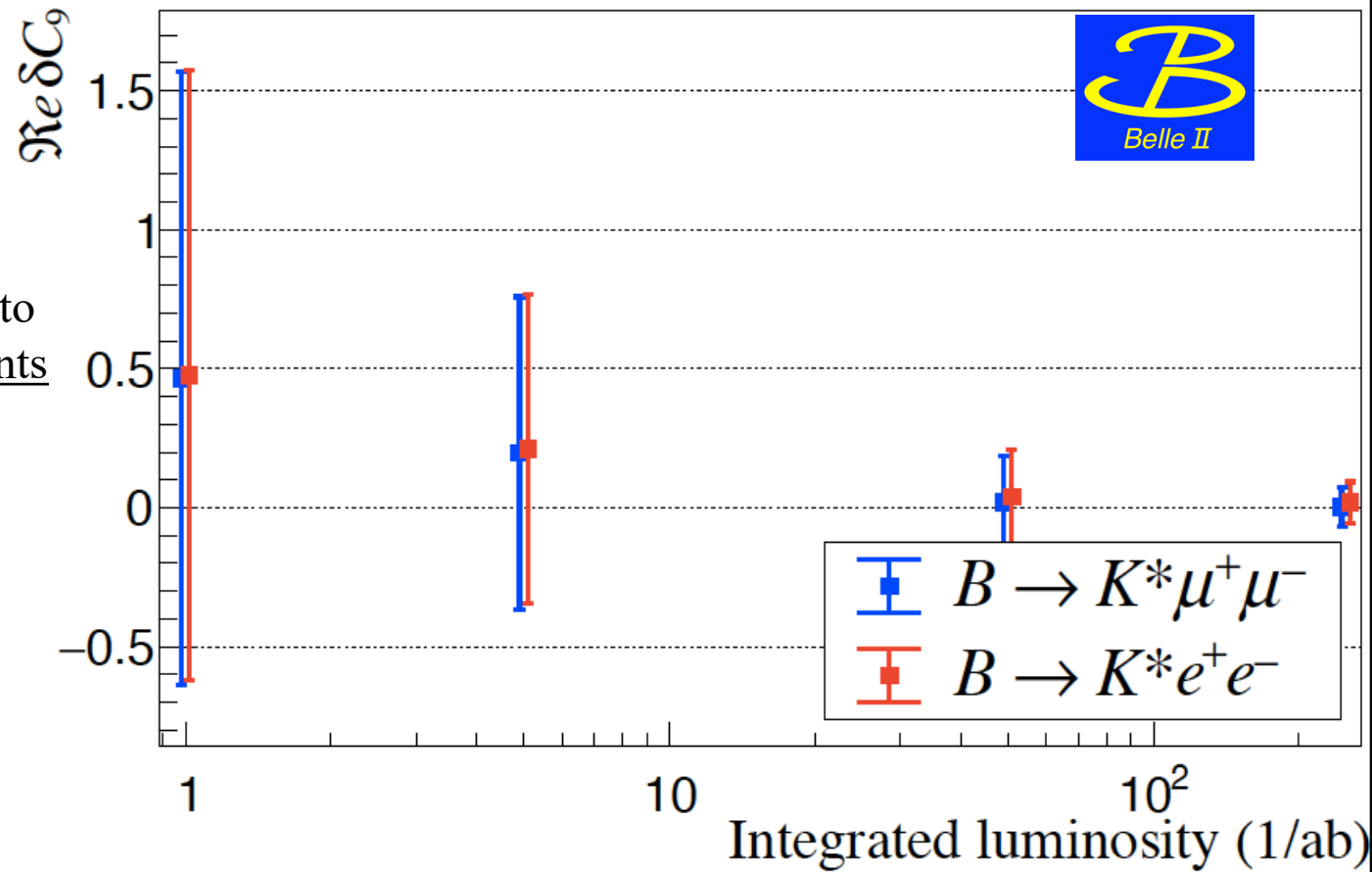
$C_9$  : Global fit to world  $b \rightarrow s$  data gives a  $>5\sigma$  deviation from the SM

## What about the future ?

Estimates use pseudo-experiments with **4-D unbinned maximum likelihood fits to 4 variables** in  $B \rightarrow K^* l^+ l^-$  to extract Wilson coefficients  $C_i$  directly from data.

Use  $q^2 > 1 \text{ GeV}^2$  and  $|q^2 - M^2| < 0.25 \text{ GeV}^2$  and assume 25% Belle efficiency

A. Sibidanov et al.



# What's Ahead for Belle II ?

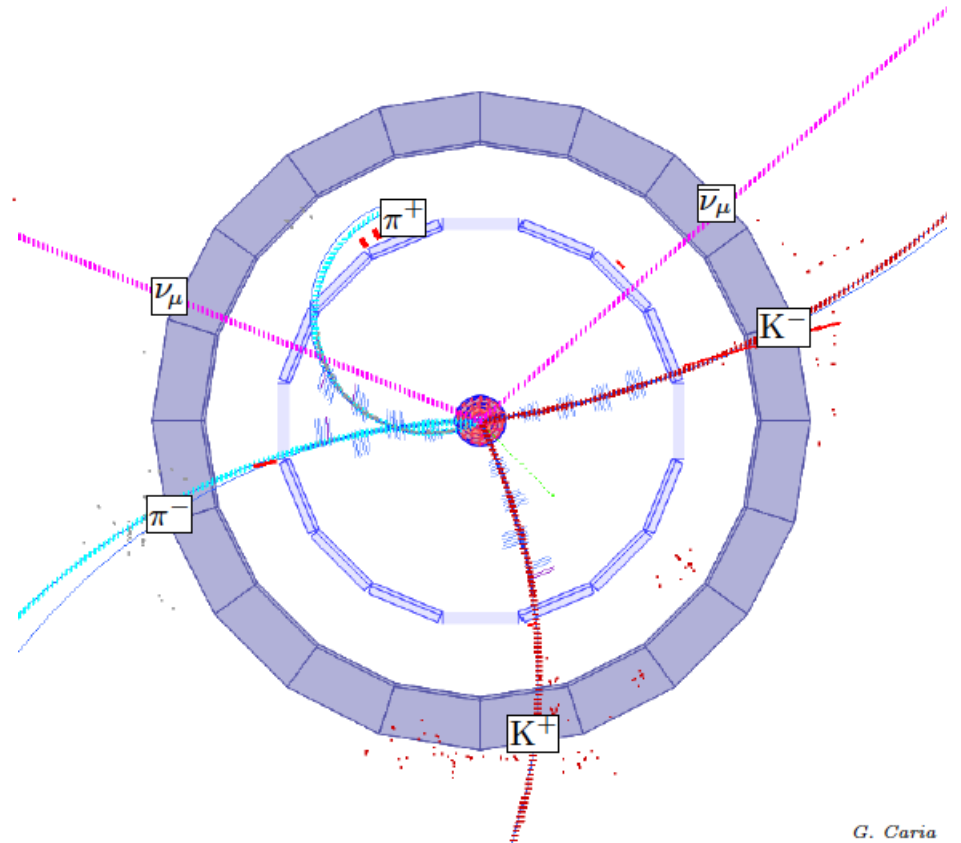
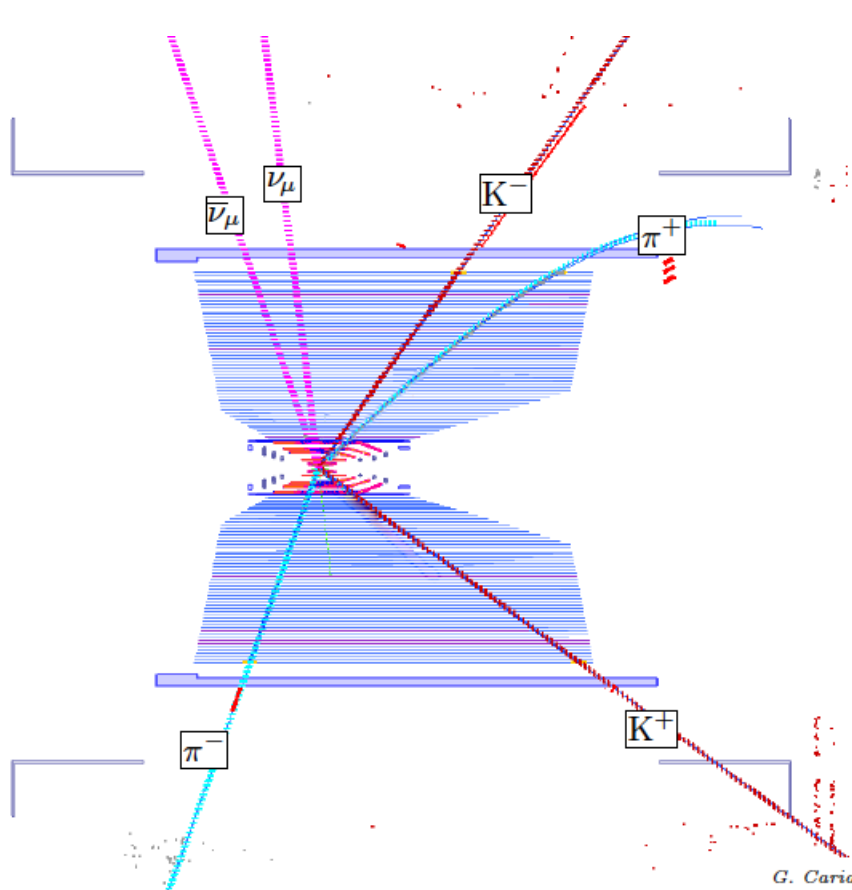
“Missing Energy Decay” in a Belle II GEANT4 MC simulation

Signal:  $B \rightarrow K \nu \nu$

tag mode:  $B \rightarrow D\pi$ ;  $D \rightarrow K\pi$

View in r-z

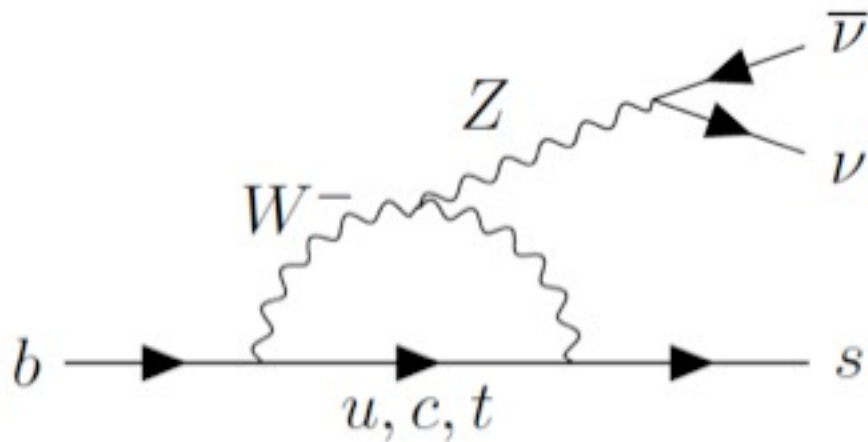
Zoomed view of the vertex region in r--phi



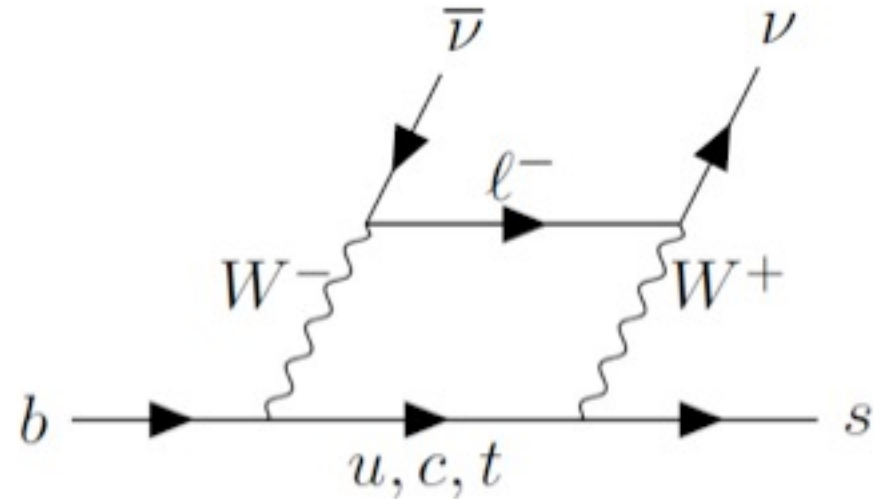




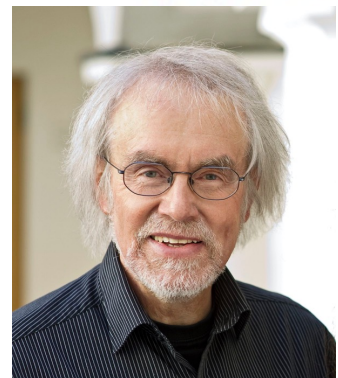
# $B \rightarrow K \nu \bar{\nu}$ : NP without hadronic uncertainties



(a) Penguin diagram



(b) Box diagram



Andrezej  
Buras

Note that in contrast to  $B \rightarrow K^{(*)} l^+ l^-$  angular asymmetries, there are **NO** long distance (charm annihilation) contributions from  $B \rightarrow J/\psi K^{(*)}$  and  $B \rightarrow \psi(2S) K^{(*)}$

For example, <https://arxiv.org/abs/1409.4557>

The  $B \rightarrow K^{(*)} \nu \bar{\nu}$  modes are accessible to **Belle II** (and Belle), but might be hard at a hadron experiment.



# Calibration Mode for $B \rightarrow K$ nu nubar

$$B^+ \rightarrow J/\psi K^+, J/\psi \rightarrow \mu^+ \mu^-$$

$$B \rightarrow K \nu \bar{\nu}$$

Hadronic FEI or semileptonic FEI, require full reconstruction of individual decay modes, effective efficiency is 1% at best.

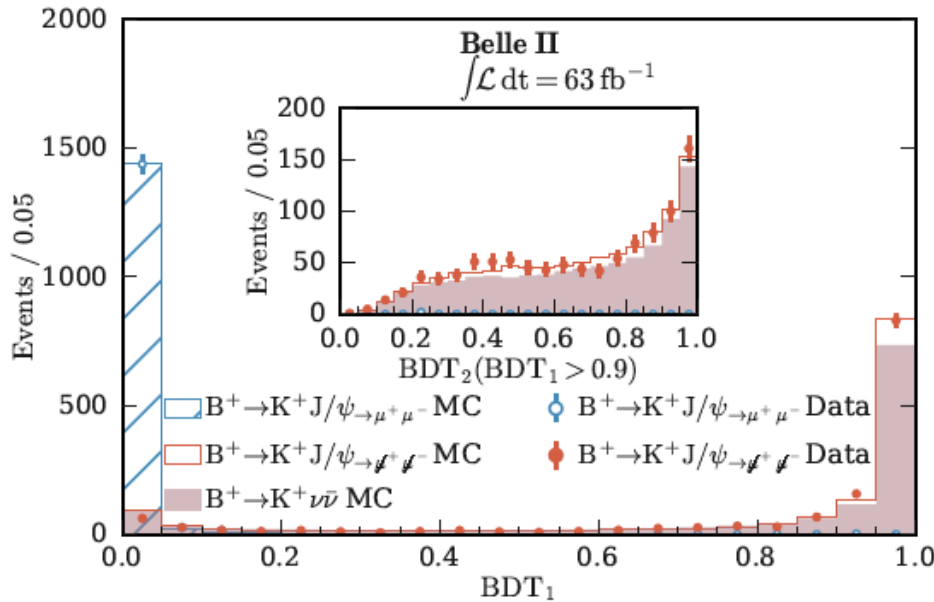


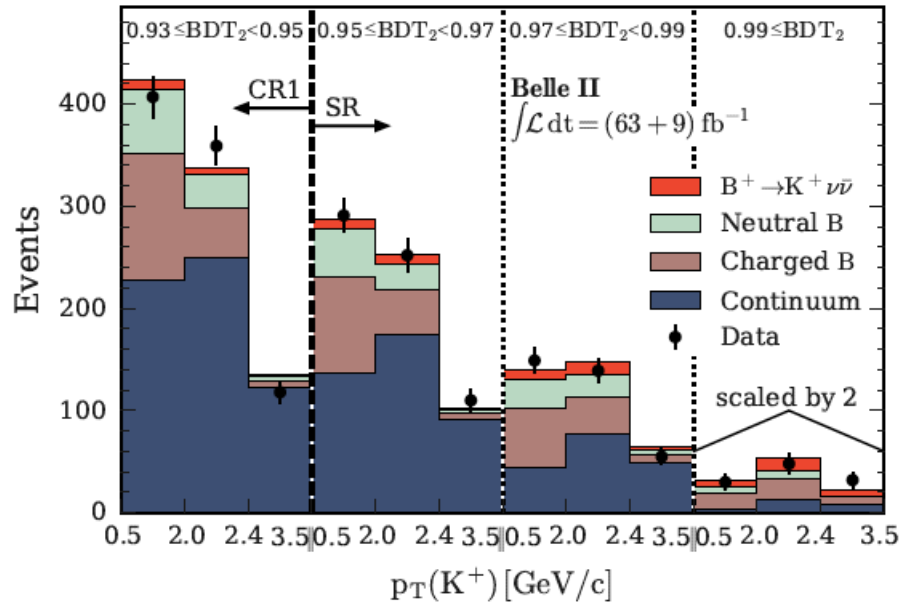
FIG. 2: Distribution of the classifier output  $BDT_1$  (main figure) and  $BDT_2$  for  $BDT_1 > 0.9$  (inset). The distributions are shown before ( $J/\psi \rightarrow \mu^+ \mu^-$ ) and after ( $J/\psi \rightarrow \mu^+ \mu^-$ ) the muon removal and update of the kaon-candidate momentum of selected  $B^+ \rightarrow K^+ J/\psi$  events in simulation and data. As a reference, the classifier outputs directly obtained from simulated  $B^+ \rightarrow K^+ \nu \bar{\nu}$  signal events are overlaid. The simulation histograms are scaled to the total number of  $B^+ \rightarrow K^+ J/\psi$  events selected in data.

**New Idea** (Sasha Glazov et al):  
 Try *inclusive ROE (Rest Of the Event)* tagging and improve efficiency by a factor of 5-10. Backgrounds are higher but manageable by fitting.

Use two BDTs:  $BDT_1$  for continuum bkg suppression and then  $BDT_2$  to distinguish  $B \bar{B}$  bkg from signal.



# B → K ν nubar candidates: p<sub>T</sub>(K) distribution in BDT2 bins



*inclusive ROE (Rest Of the Event) tagging*

$$B \rightarrow K \nu \bar{\nu}$$

There is an excess from a 2D histogram fit, which corresponds to

$$\mu = [4.2^{+2.9+1.8}_{-2.8-1.6}] \times SM$$

FIG. 3: Yields in on-resonance data and as predicted by the simultaneous fit to the on- and off-resonance data, corresponding to an integrated luminosity of  $63 \text{ fb}^{-1}$  and  $9 \text{ fb}^{-1}$ , respectively. The predicted yields are shown individually for charged and neutral  $B$ -meson decays and the sum of the five continuum contributions. The leftmost three bins belong to CR1 with  $\text{BDT}_2 \in [0.93, 0.95]$  and the other nine bins correspond to the SR, three for each range of  $\text{BDT}_2 \in [0.95, 0.97, 0.99, 1.0]$ . Each set of three bins is defined by  $p_T(K^+) \in [0.5, 2.0, 2.4, 3.5] \text{ GeV}/c$ . All yields in the rightmost three bins are scaled by a factor of two.



# B → K ν ν̄: NP without hadronic uncertainties

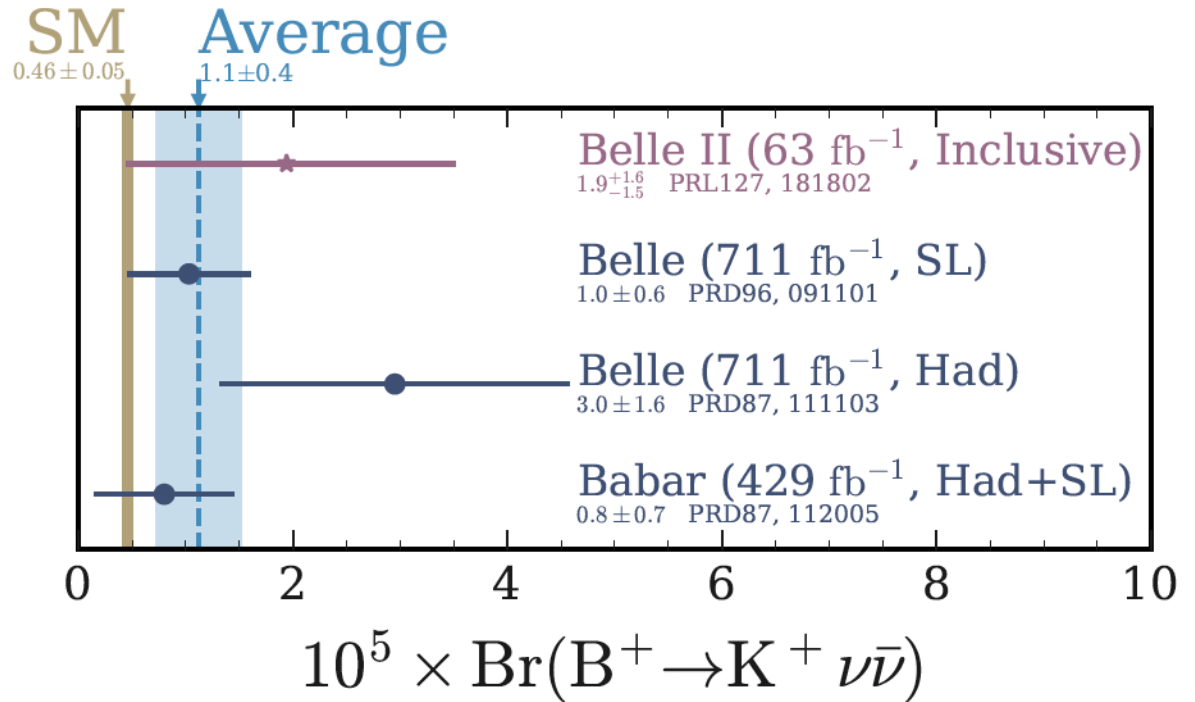
An emerging anomaly ???

$$B \rightarrow K \nu \bar{\nu}$$

New Technique from Belle II with inclusive ROE (Rest of the Event) tagging.

Phys. Rev. Lett. 127, 181802, (20

Now apply to old Belle and new Belle II data. Stay tuned.



But it is also possible that NP shows up in  $b \rightarrow s l^+ l^-$  but not in  $b \rightarrow s \nu \bar{\nu}$  or vice-versa

Dark matter could also play a major role.

>>> This is one way that Belle II could discover New Physics soon <<<

More details in this theory preprint (TEB, N. Deshpande, R. Mandal, R. Sinha):

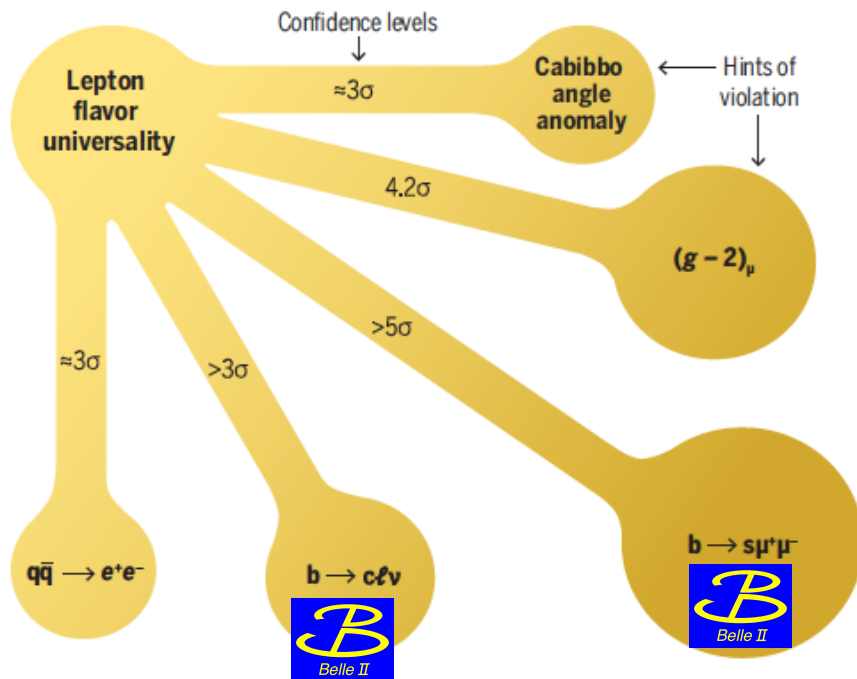
<https://arxiv.org/abs/2107.01080>, published as Phys. Rev. D. 104, 053007 (2021)



Conclusion: Here are some more examples of how Belle II might find New Physics in the coming years.

### Possible violations of lepton flavor universality are getting harder to ignore

Shown are five hints for the violation of lepton flavor universality from existing experimental data, with the size of each circle and length of each arm reflecting the level of confidence for the experimental data to break away from standard model predictions.



From December 2021 SCIENCE magazine article by A. Crivellin and M. Hoferichter.

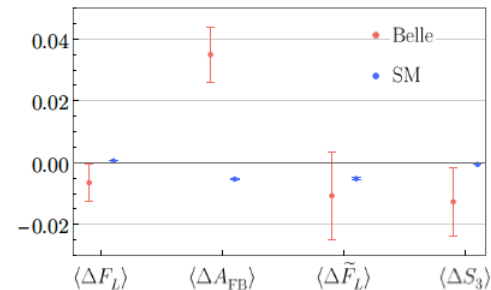
Belle II's Strong and Unique Capabilities for **New Physics and resolution of HEP Anomalies:**



$$\Delta A_{FB}(B \rightarrow D^{*+} \ell \nu) = A_{FB}(B \rightarrow D^{*+} \mu^{-} \nu) - A_{FB}(B \rightarrow D^{*+} e^{-} \nu)$$

$\Delta A_{FB}$

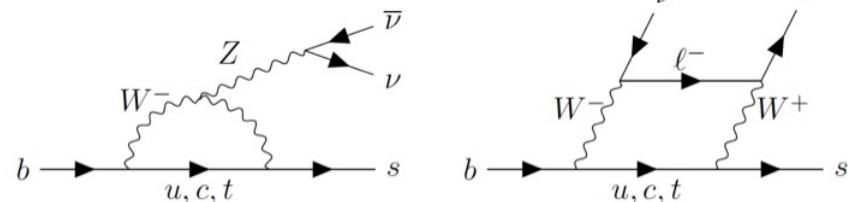
A 4σ deviation from the SM is found in Belle data (0.71 ab<sup>-1</sup>)



Eur. Phys. J.C. 81 (2021).

<https://arxiv.org/abs/2104.02094>

Belle II's unique inclusive and **missing energy** capabilities. The current WA for B → K ν nubar (2.4 +1.1) SM.



(a) Penguin diagram

(b) Box diagram

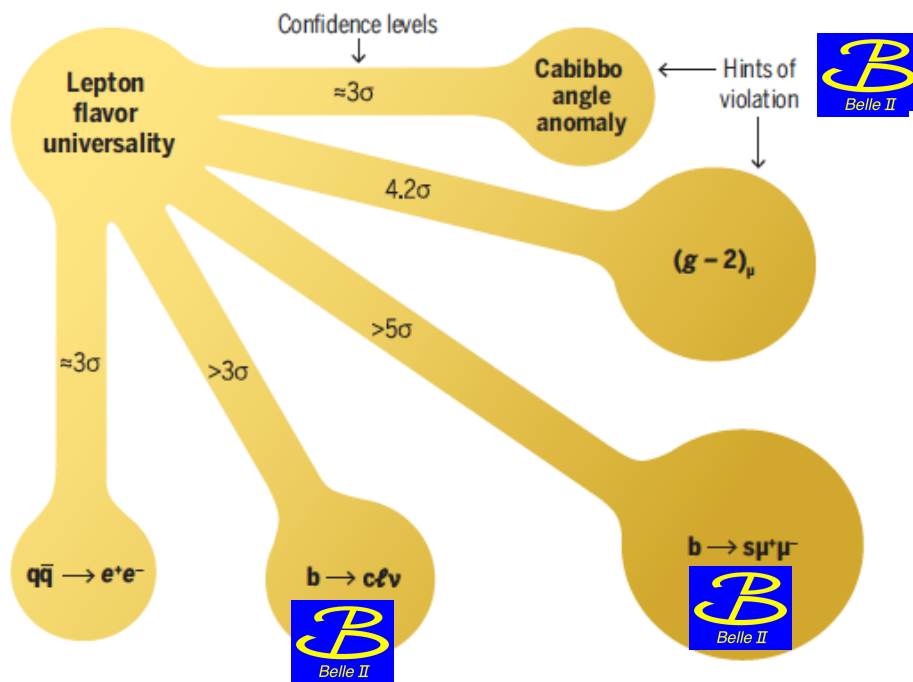
Belle II; Phys. Rev. Lett. 127, 181802 (2021)

But these modes require lots of data...."There is no royal road to new physics" (to paraphrase Euclid).

# But wait there's more.....

## Possible violations of lepton flavor universality are getting harder to ignore

Shown are five hints for the violation of lepton flavor universality from existing experimental data, with the size of each circle and length of each arm reflecting the level of confidence for the experimental data to break away from standard model predictions.



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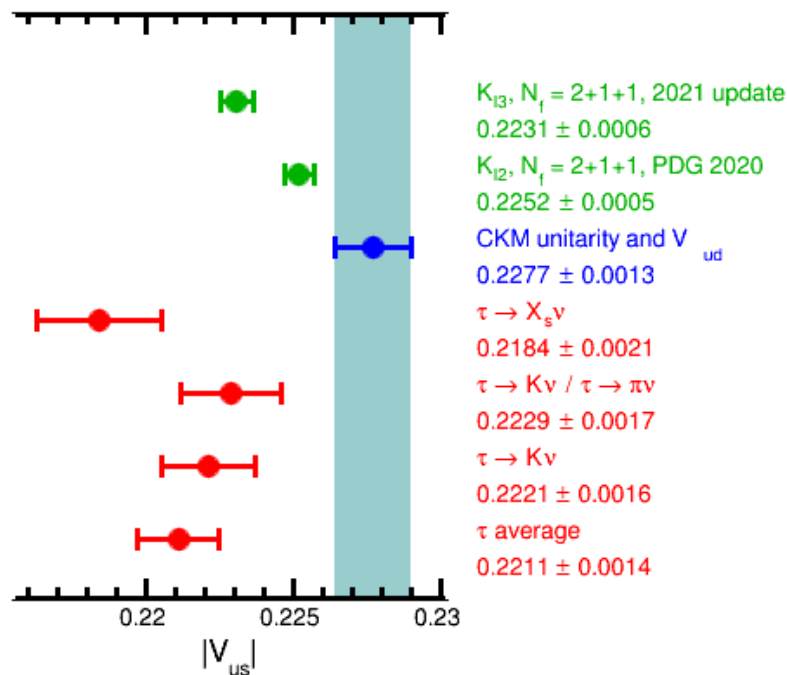
A major supporting role of Belle II in the resolution of two more of the other HEP anomalies.

## Belle II can contribute to the resolution of the Cabibbo Angle Anomaly (CAA)



There is a  $\sim 3\sigma$  discrepancy between  $|V_{us}|$  measured from tau and kaon semileptonic decays.

*Belle II will measure  $|V_{us}|$  in inclusive tau decays to high precision*



The CAA could be another hint of lepton flavor universality violation

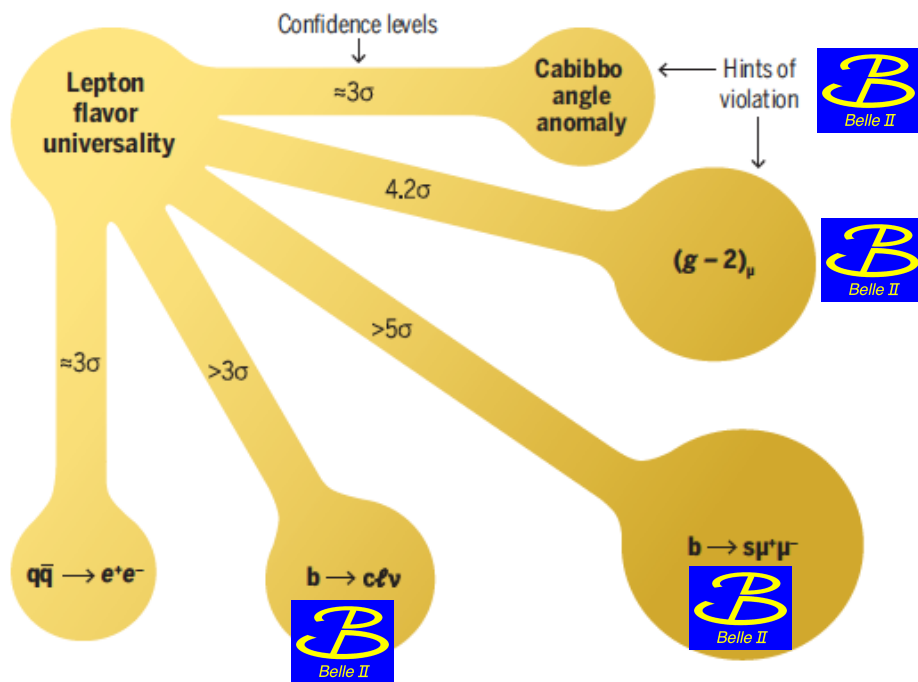


+But wait there's more.....

Belle II can contribute to  $g-2$

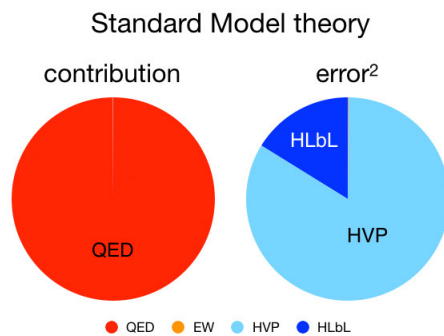
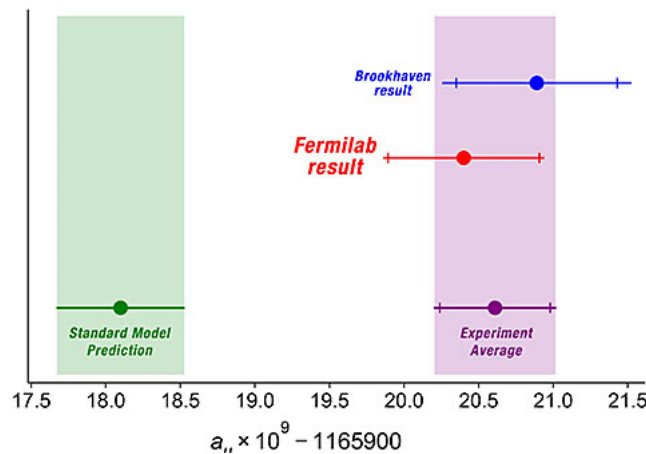
### Possible violations of lepton flavor universality are getting harder to ignore

Shown are five hints for the violation of lepton flavor universality from existing experimental data, with the size of each circle and length of each arm reflecting the level of confidence for the experimental data to break away from standard model predictions.



From December 2021 SCIENCE magazine article by A. Crivellin and M. Hoferichter.

A major supporting role of Belle II in the resolution of two more of the other major HEP anomalies



KLOE and BaBar data disagree

*Belle II can measure the cross-section for  $e^+e^- \rightarrow \pi\pi$  vs  $\sqrt{s}$  and reduce the hadronic vacuum polarization error in  $g-2$  (dominant theory uncertainty). This could help to determine whether there is really New Physics in  $g-2$  (muon).*





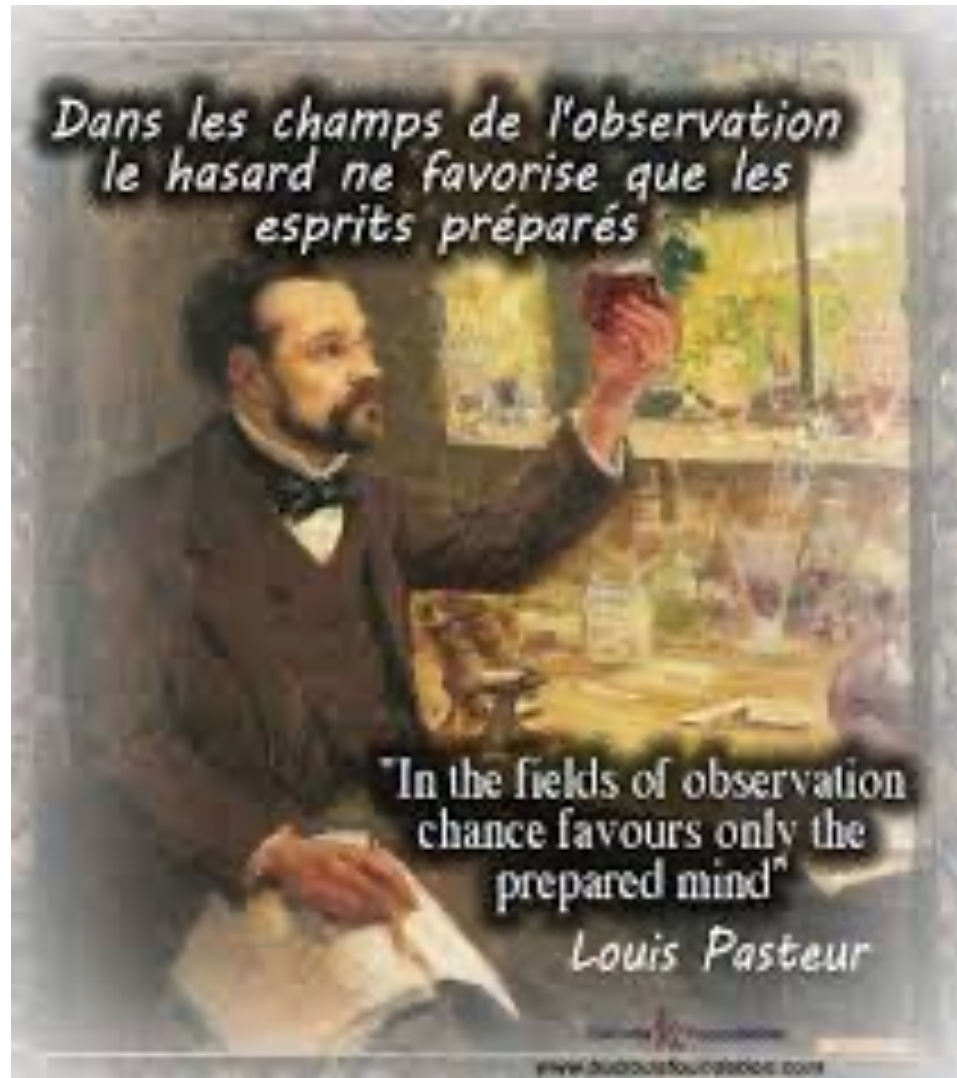
# New Physics Opportunities

Belle II Executive Summary for Snowmass

<https://arxiv.org/abs/2203.10203>

- Apologies for only covering a small range of possibilities.
- Leverage Belle II's unique **photon, electron,  $\pi^0$ , missing energy capabilities** for rare B decays. From Snowmass: *Use  $\Delta$  Observables to find LFU violation in angular asymmetries (ideally suited for Belle II at high luminosity).*
- A number of  $b \rightarrow s$  and  $b \rightarrow c$  processes have hints of NP. **(New: pay attention to  $B \rightarrow K \nu \bar{\nu}$  as new data comes in, Belle II has demonstrated improved sensitivity).** Along with  $B \rightarrow D^{(*)} \tau \nu$  and  $D^{(*)} l \nu$  ( $b \rightarrow c$ ), the anomalies will be studied in detail at Belle II in the near future.
- A lot to learn at this workshop to allow **you** to make the next round of discoveries. Belle II has strong and unique capabilities for **New Physics discoveries and resolution of the major** high energy physics anomalies (and not just B physics).

# Backup slides





# More Belle II Superpowers

4. [arXiv:2206.08280](#) [pdf, other] [hep-ex](#)

Measurement of the branching fraction of the  $B^0 \rightarrow K_S^0 \pi^0 \gamma$  decay using  $190 \text{ fb}^{-1}$  of Belle II data

**Authors:** Belle II Collaboration, F. Abudinén, I. Adachi, K. Adamczyk, L. Aggarwal, P. Ahlburg, H. Ahmed, J. K. Ahn, H. Aihara, N. Akopov, A. Aloisio, F. Ameli, L. Andricsek, N. Anh Ky, D. M. Asner, H. Atmacan, V. Aulchenko, T. Aushev, V. Aushev, T. Aziz, V. Babu, S. Bacher, H. Bae, S. Baehr, S. Bahinipati, et al. (570 additional authors not shown)

**Abstract:** We report the measurement of the branching fraction of the  $B^0 \rightarrow K_S^0 \pi^0 \gamma$  decay in  $e^+ e^- \rightarrow \Upsilon(4S) \rightarrow B\bar{B}$  data recorded by the Belle II experiment at the SuperKEKB asymmetric-energy collider and corresponding to  $190 \text{ fb}^{-1}$  of integrated luminosity. The signal yield is measured to be  $121 \pm 29$  (stat.), leading to the branching fraction... [▽ More](#)

**Submitted** 16 June, 2022; **originally announced** June 2022.

**Comments:** 10 pages, 3 figures

**Report number:** BELLE2-CONF-2022-008

6. [arXiv:2206.07453](#) [pdf, other] [hep-ex](#)

First decay-time-dependent analysis of  $B^0 \rightarrow K_S^0 \pi^0$  at Belle II

**Authors:** Belle II Collaboration, F. Abudinén, I. Adachi, R. Adak, K. Adamczyk, L. Aggarwal, P. Ahlburg, H. Ahmed, J. K. Ahn, H. Aihara, N. Akopov, A. Aloisio, F. Ameli, L. Andricsek, N. Anh Ky, D. M. Asner, H. Atmacan, V. Aulchenko, T. Aushev, V. Aushev, T. Aziz, V. Babu, S. Bacher, H. Bae, S. Baehr, et al. (569 additional authors not shown)

**Abstract:** We report measurements of the branching fraction ( $\mathcal{B}$ ) and direct  $CP$ -violating asymmetry ( $A_{CP}$ ) of the charmless decay  $B^0 \rightarrow K^0 \pi^0$  at Belle II. A sample of  $e^+ e^-$  collisions, corresponding to  $189.8 \text{ fb}^{-1}$  of integrated luminosity, recorded at the  $\Upsilon(4S)$  resonance is used for the first decay-time-dependent analysis of these decays within the experiment. We reconst...

[▽ More](#)

**Submitted** 15 June, 2022; **originally announced** June 2022.



# $B \rightarrow K \nu \bar{\nu}$ : NP *without* hadronic uncertainties !

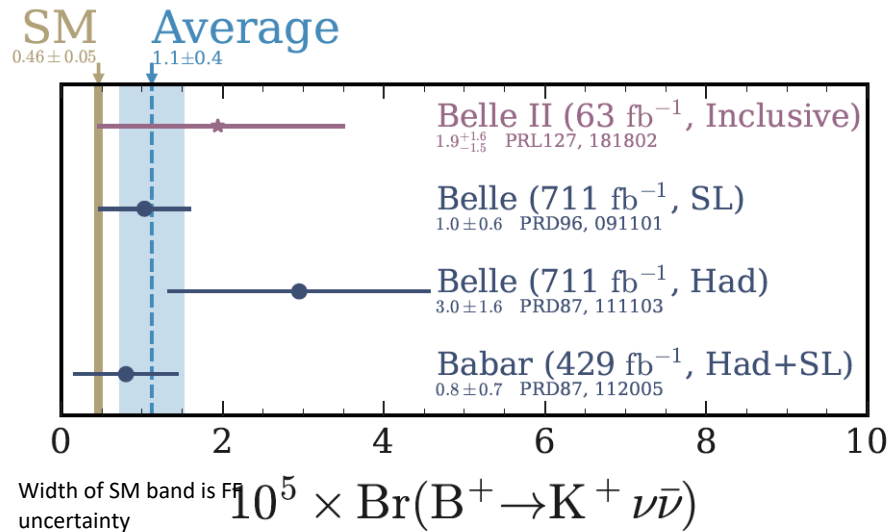
4% experimental error on  $B \rightarrow K^* \nu \bar{\nu}$  with Belle II@250  $ab^{-1}$

$$B \rightarrow K \nu \bar{\nu}$$

New Technique from Belle II with inclusive ROE (Rest of the Event) tagging improves sensitivity.

Phys. Rev. Lett. 127, 181802, (2021)

An emerging anomaly ???



Andrezej Buras

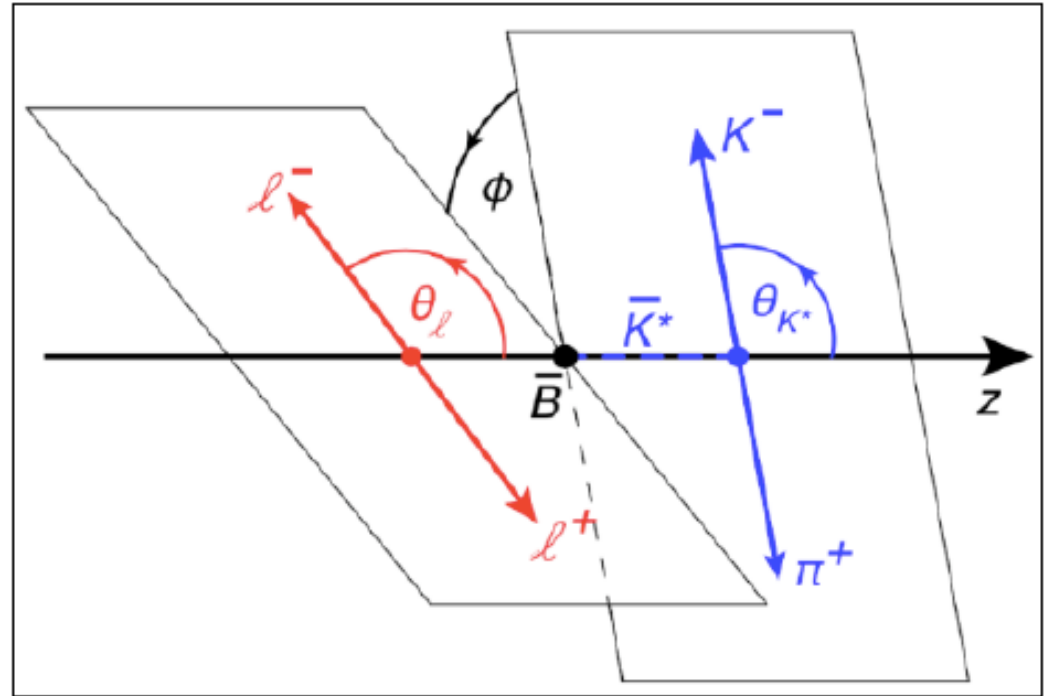
“Note there are no charm loops here”- Wolfgang A.

*But it is also possible that NP shows up only in  $b \rightarrow s l^+ l^-$  but not in  $b \rightarrow s \nu \bar{\nu}$  or vice-versa. The two classes of EWPs are related but distinct.*

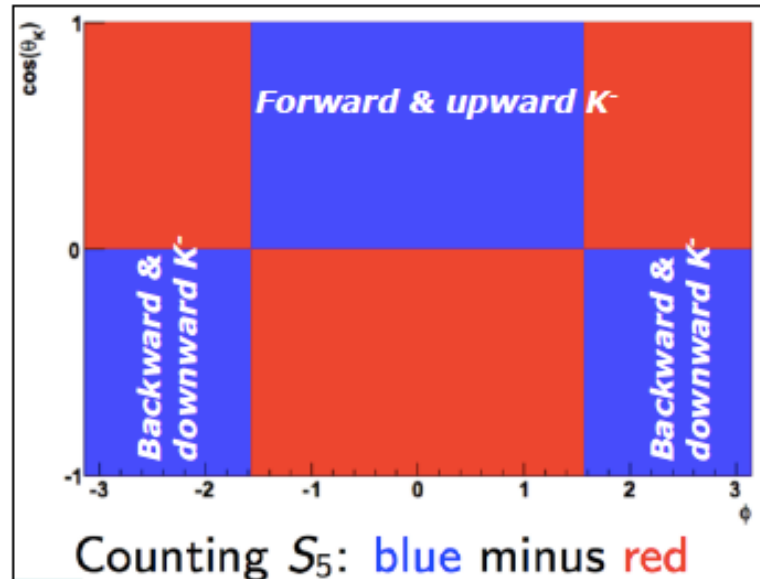
This is one way that Belle II could discover New Physics soon. For example: <https://arxiv.org/abs/2107.01080>, Phys. Rev. D. 104, 053007 (2021)

Dark matter could also play a major role.

# Angular asymmetries

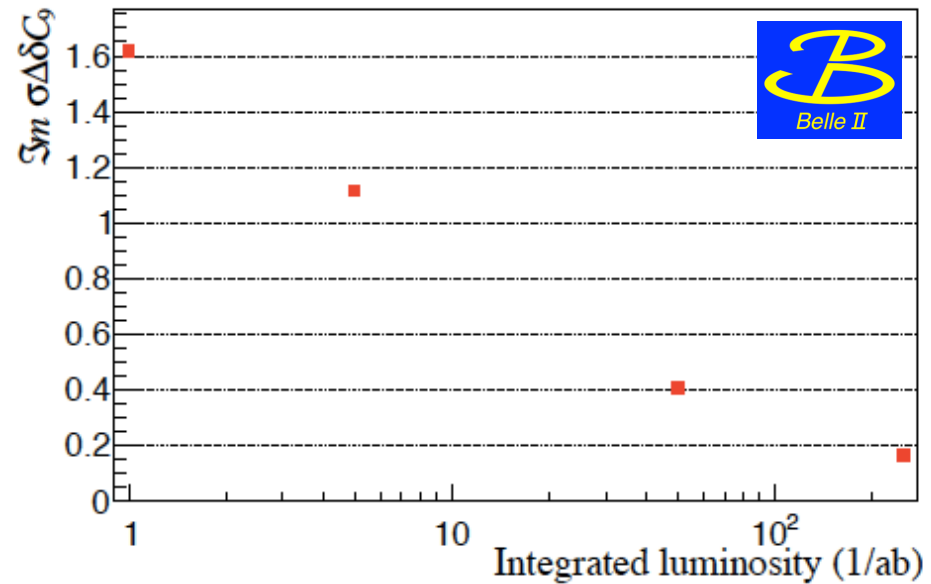
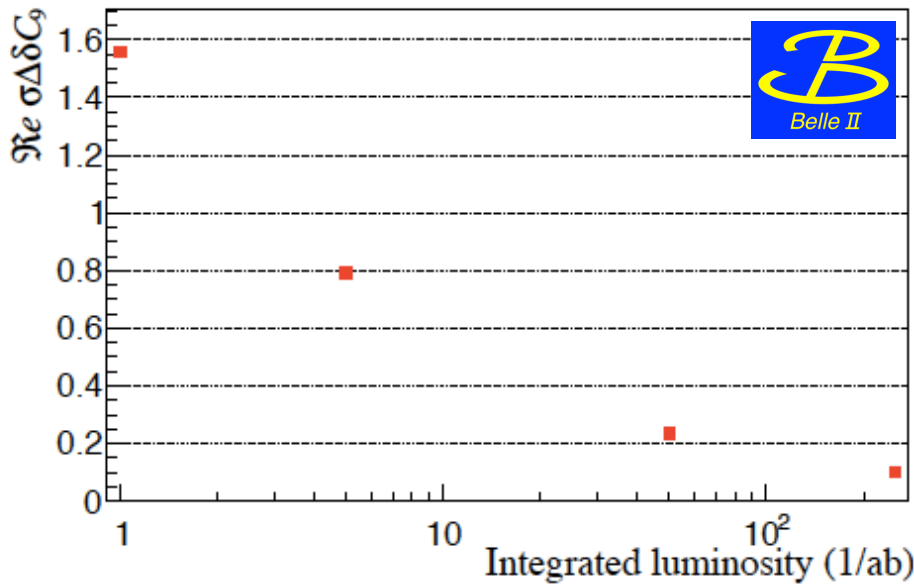


$P_5'$  :



Snowmass Bullet Point:

Use the  $\Delta$  Observables in  $B \rightarrow K^* 1^+ 1^-$  to discover New Physics at Belle II without QCD and hadronic uncertainties.

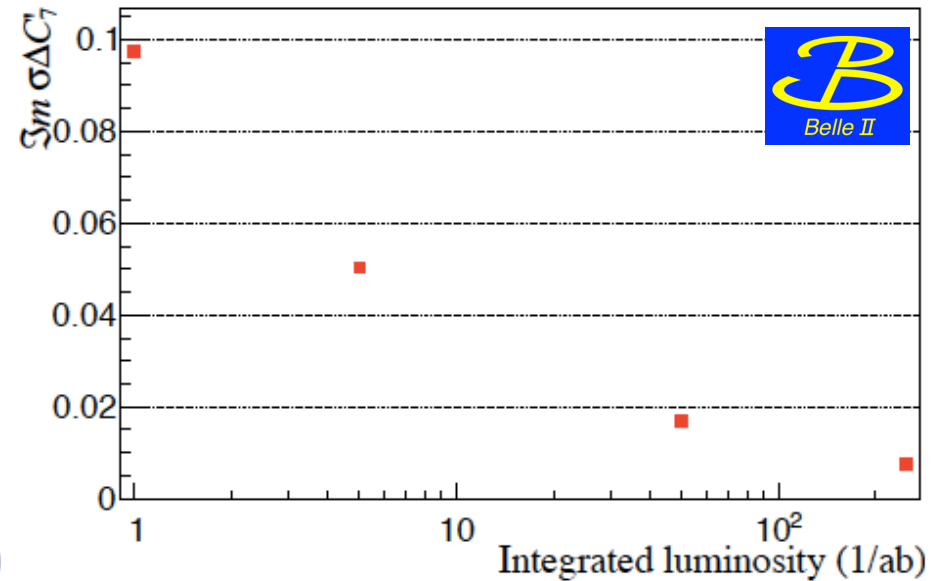
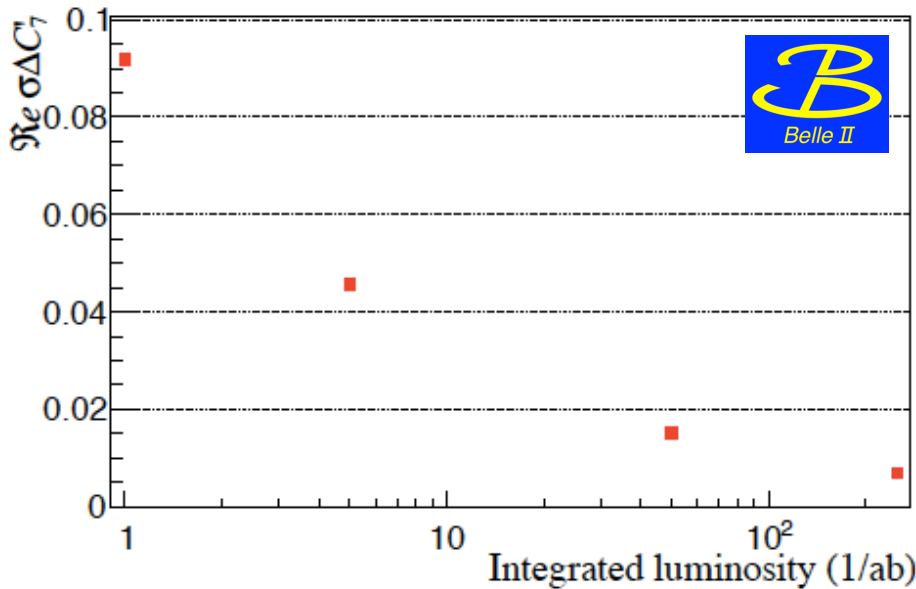


[A. Sibidanov et al., https://arxiv.org/abs/2203.07189](https://arxiv.org/abs/2203.07189)



# Belle II Sensitivity to NP Right-Handed Currents, ( $C_7'$ )

A. Sibidanov et al., <https://arxiv.org/abs/2203.07189>



Snowmass Bullet Point:

Use the  $\Delta$  Observables in  $B \rightarrow K^* l^+ l^-$  to discover New Physics at Belle II without QCD and hadronic uncertainties.



# FAQ: How do Belle II at KEK and LHCb at CERN capabilities compare ?

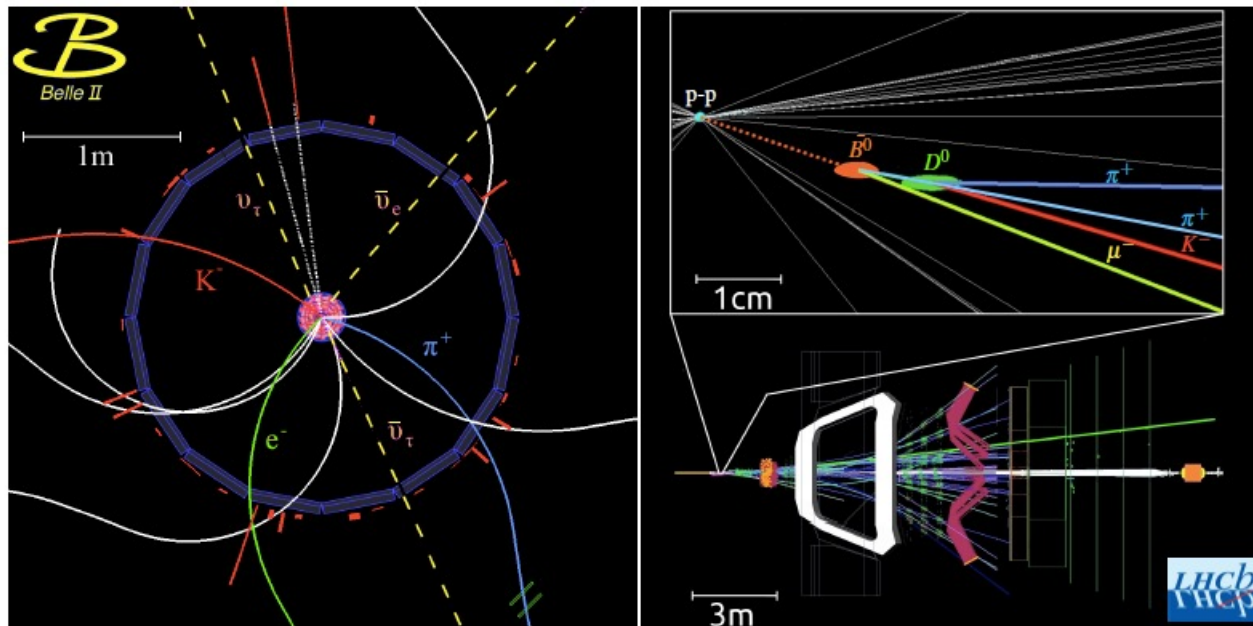


Figure credit:  
G. Ciezarak et al,  
Nature  
546, 227 (2017)

++Belle II can  
do the dark  
sector

1. LHCb has a large  $b\bar{b}$  cross-section (hundreds of microbarns versus nanobarns) and good sensitivity, signal to background, for modes with dimuons, and all charged final states using vertexing. Triggering and flavor tagging effs. are much lower than in  $e^+e^-$ .
  2. Belle II has a simple and clean event environment with  $B^0$  - $\bar{B}^0$  pairs produced in *a coherent QM state with no additional particles*.
  3. Belle II can measure *inclusive processes*
  4. Belle II can measure *electrons* just as well as muons. (important for lepton universality checks).
  5. Belle II can measure final states with  $\gamma$ 's, *Kshorts and missing neutrinos well*.
- Rule of thumb for statistics in this case:  
1  $\text{fb}^{-1}$  at LHCb is 1  $\text{ab}^{-1}$  at Belle II.  
( $\rightarrow$  Need good **SuperKEKB performance** and long runs in the coming years)

The  $B^0$ -anti  $B^0$  meson pairs at the Upsilon(4S) are produced in a coherent, entangled quantum mechanical state.

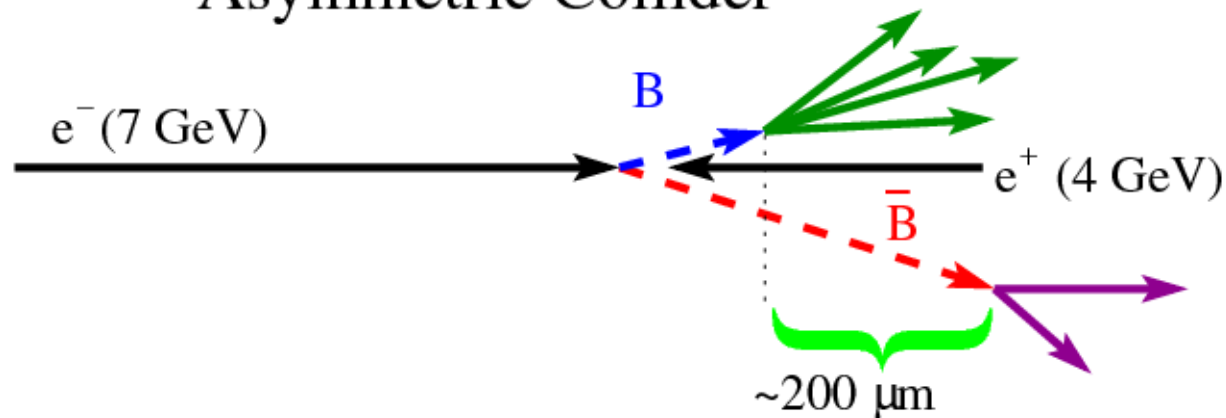
(Exercise: why is there a minus sign ?)

$$|\Psi\rangle = |B^0(t_1, f_1)\bar{B}^0(t_2, f_2)\rangle - |B^0(t_2, f_2)\bar{B}^0(t_1, f_1)\rangle$$

Need to measure decay times to observe CP violation (particle-antiparticle asymmetry).

One B decays  $\rightarrow$  collapses the flavor wavefunction of the other anti-B. (N.B. One B must decay before the other can mix) [*exercise: explain*]

### Asymmetric Collider

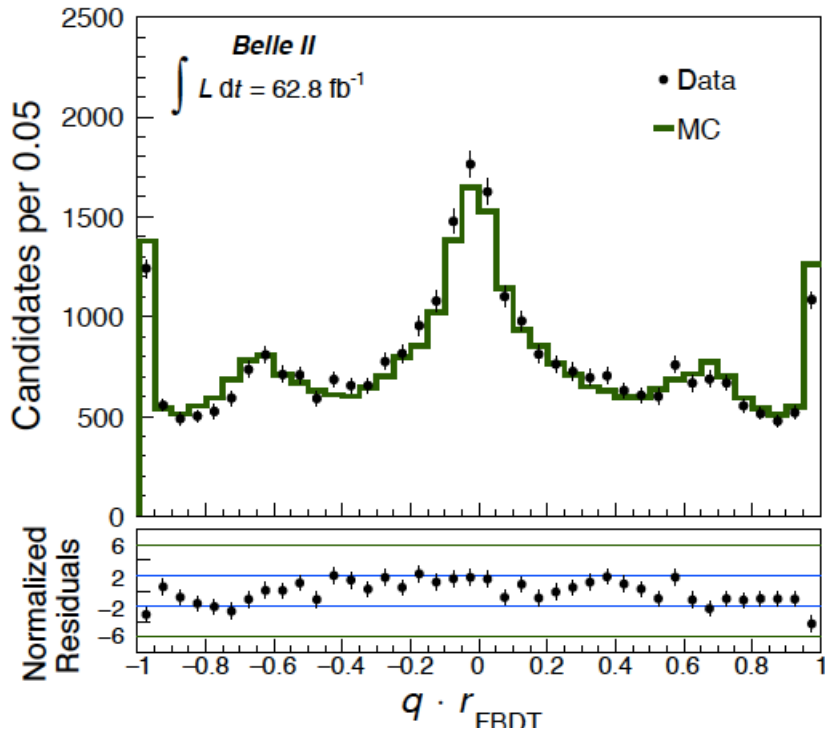


Not to scale

The beam energies are asymmetric (7 on 4 GeV)

The decay distance is increased by around a factor  $\sim 7$

# 2021 update: Flavor Tagging (b quark or anti-b quark ?)



Categories	Targets for $\bar{B}^0$	Underlying decay modes
Electron	$e^-$	$\bar{B}^0 \rightarrow D^{*+} \bar{\nu}_\ell \ell^-$
Intermediate Electron	$e^+$	$\hookrightarrow D^0 \pi^+$
Muon	$\mu^-$	$\hookrightarrow X K^-$
Intermediate Muon	$\mu^+$	
Kinetic Lepton	$l^-$	$\bar{B}^0 \rightarrow D^+ \pi^- (K^-)$
Intermediate Kinetic Lepton	$l^+$	$\hookrightarrow K^0 \nu_\ell \ell^+$
Kaon	$K^-$	
Kaon-Pion	$K^-, \pi^+$	
Slow Pion	$\pi^+$	
Maximum P*	$l^-, \pi^-$	$\bar{B}^0 \rightarrow \Lambda_c^+ X^-$
Fast-Slow-Correlated (FSC)	$l^-, \pi^+$	$\hookrightarrow \Lambda \pi^+$
Fast Hadron	$\pi^-, K^-$	$\hookrightarrow p \pi^-$
Lambda	$\Lambda$	

*Time-independent method with  $62.8 \text{ fb}^{-1}$*

r- Interval	$\epsilon_{\text{eff},i} \pm \delta\epsilon_{\text{eff},i}$		
	FBDT	DNN	Belle
0.000 – 0.100	$0.1 \pm 0.1$	$0.0 \pm 0.1$	0.0
0.100 – 0.250	$0.5 \pm 0.2$	$0.3 \pm 0.1$	$0.4 \pm 0.1$
0.250 – 0.500	$3.3 \pm 0.5$	$2.3 \pm 0.4$	$2.3 \pm 0.1$
0.500 – 0.625	$3.3 \pm 0.5$	$4.2 \pm 0.5$	$3.5 \pm 0.1$
0.625 – 0.750	$6.1 \pm 0.6$	$3.8 \pm 0.5$	$4.6 \pm 0.2$
0.750 – 0.875	$5.4 \pm 0.5$	$5.9 \pm 0.6$	$5.5 \pm 0.1$
0.875 – 1.000	$11.3 \pm 0.6$	$12.3 \pm 0.7$	$13.8 \pm 0.3$
Total	$30.0 \pm 1.3$	$28.8 \pm 1.3$	$30.1 \pm 0.4$



So far, the fast BDT does better than deep learning neural net.

We obtain  $\epsilon_{\text{eff}} = \epsilon(1-2w)^2 = 30.0 \pm 1.2 \pm 0.4 \%$ , which is similar to the Belle result of  $30.1 \pm 0.4 \%$

<https://arxiv.org/abs/2008.02707>,  
 submitted to EPJC



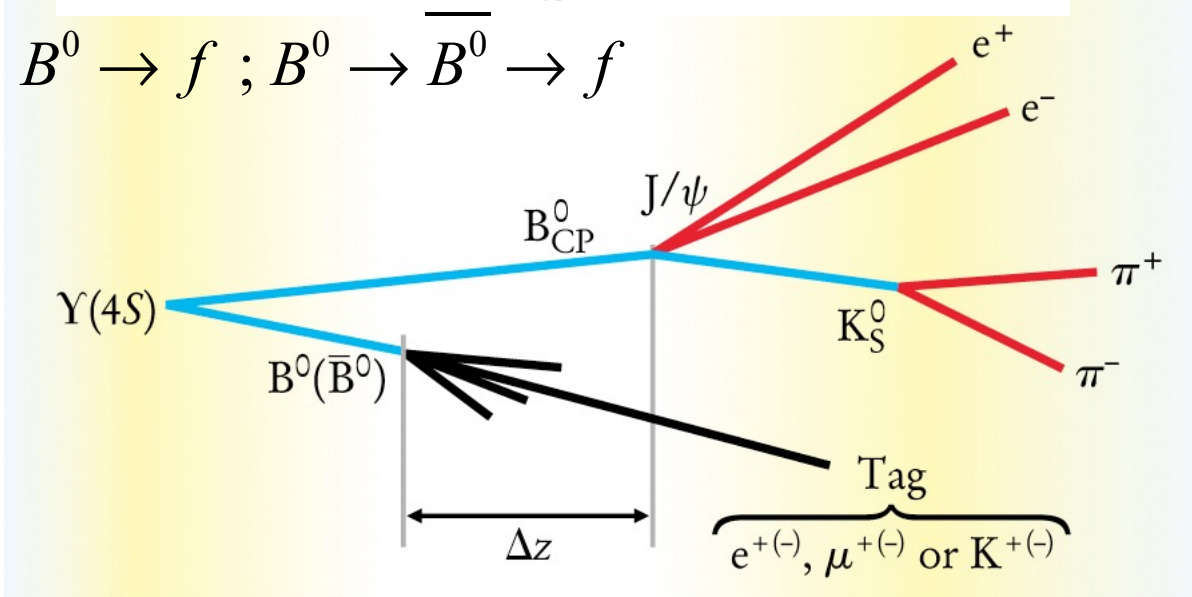
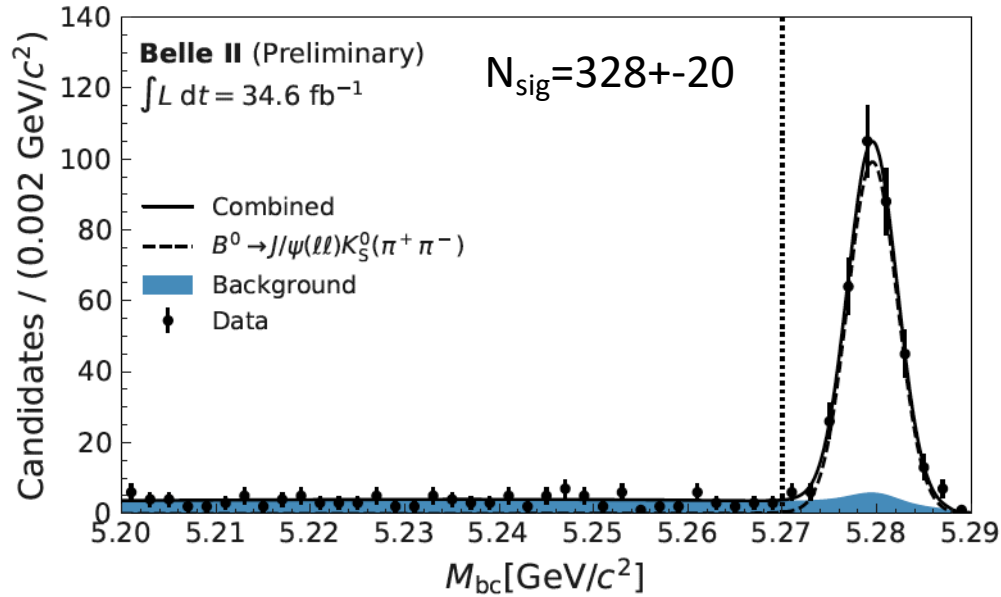
# Observation of $B \rightarrow J/\psi K_S$ and the road to CPV

A "Golden" CP Eigenstate

Test with 17% of the Phase 3 data sample.

Now apply a *simplified analysis*:

- 1) Only one CP eigenstate
- 2) No beam spot constraint
- 3) Flavor tagging does not separate r-bins



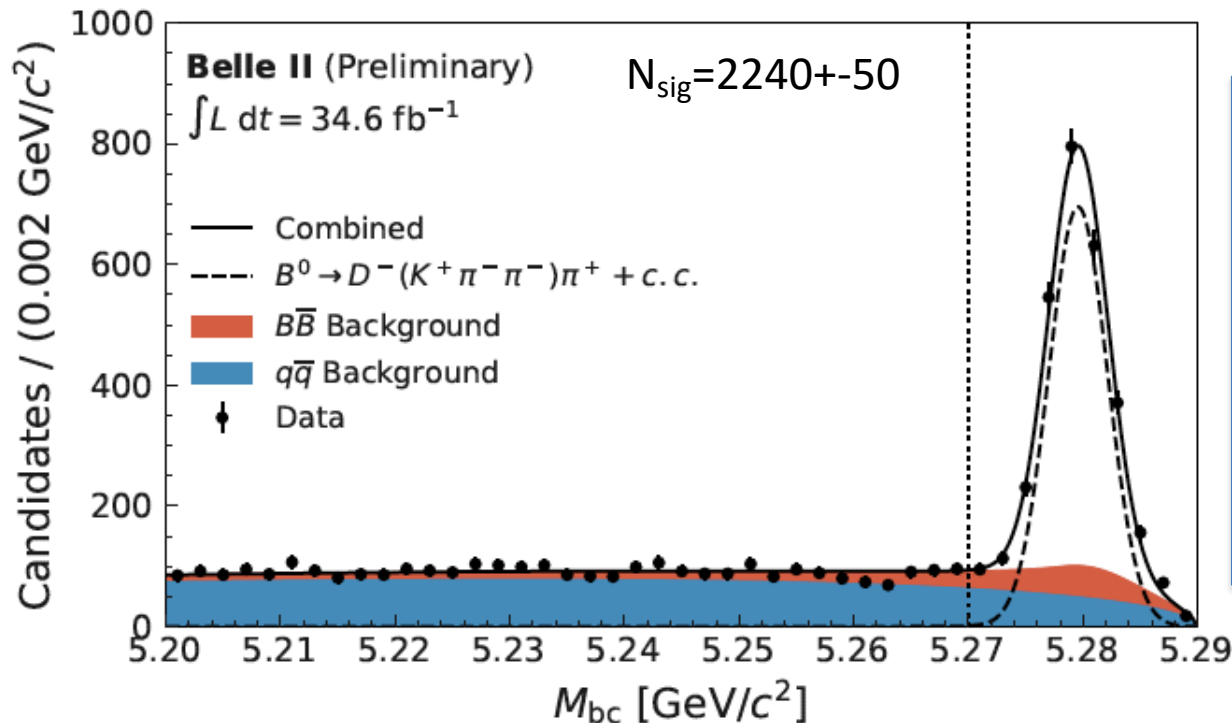
$$\Delta t \approx \frac{\Delta z}{\beta \gamma}$$

Figure credit: Physics Today



This is a flavor-specific B decay mode with a charged track topology similar to the  $B \rightarrow J/\psi K_S$  signal.

$B^0 \rightarrow D^- \pi^+$  is not self-conjugate and is **not** a CP eigenstate (but can be used to check time-dependence of B-Bbar mixing).



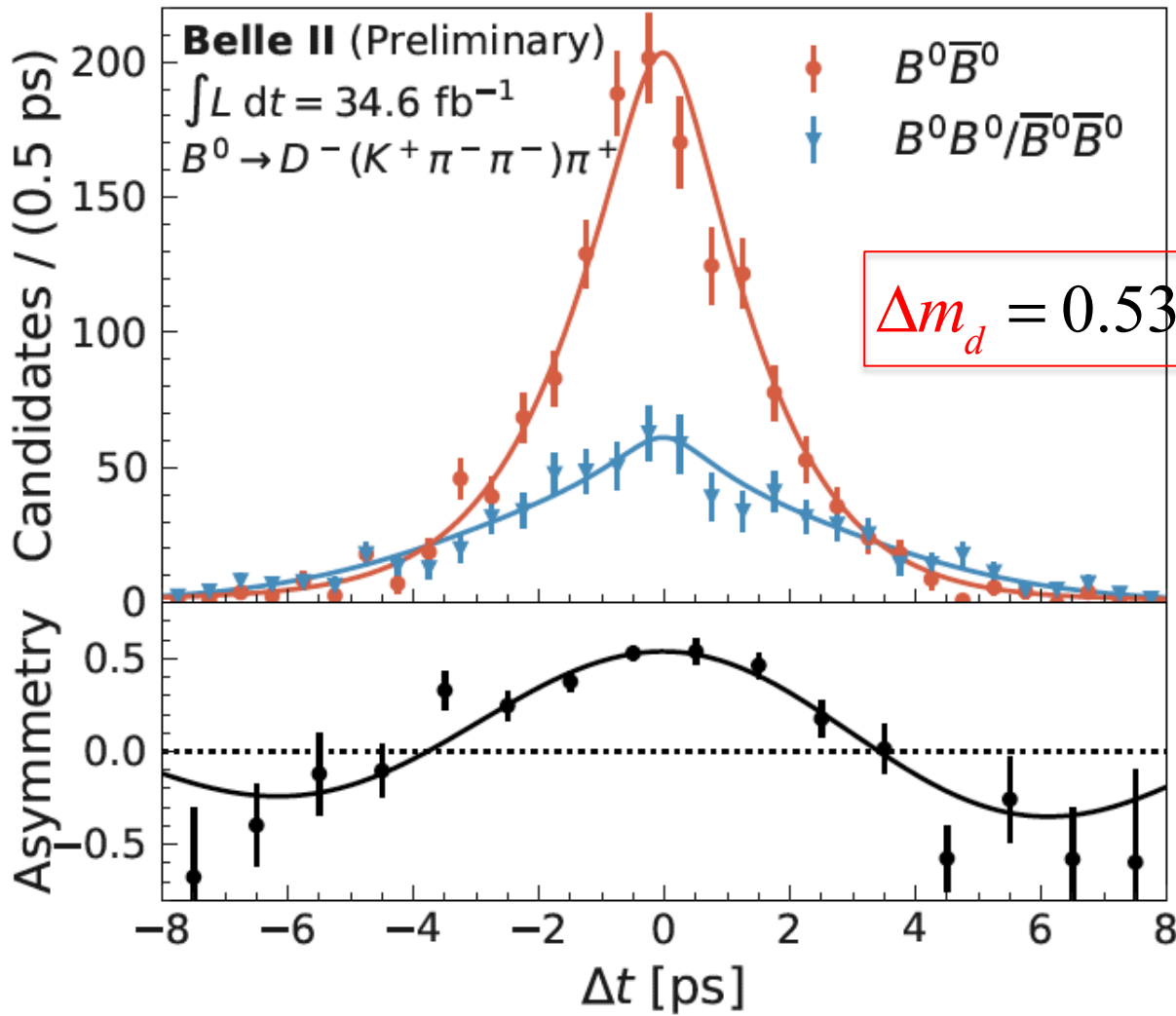
Start with a  $B^0$  (wait a while,  $\sim a \text{ few } \times 10^{-12} \text{ sec}$ ).

There is a large probability that the  $B^0$  will turn into its anti-particle, an anti- $B^0$  (discovered by ARGUS at DESY in 1987)

*The variable on the x-axis is beam-constrained mass (CM energy/2 or beam energy is used instead of reconstructed energy)*



# Time Dependent Mixing asymmetry (not CPV)

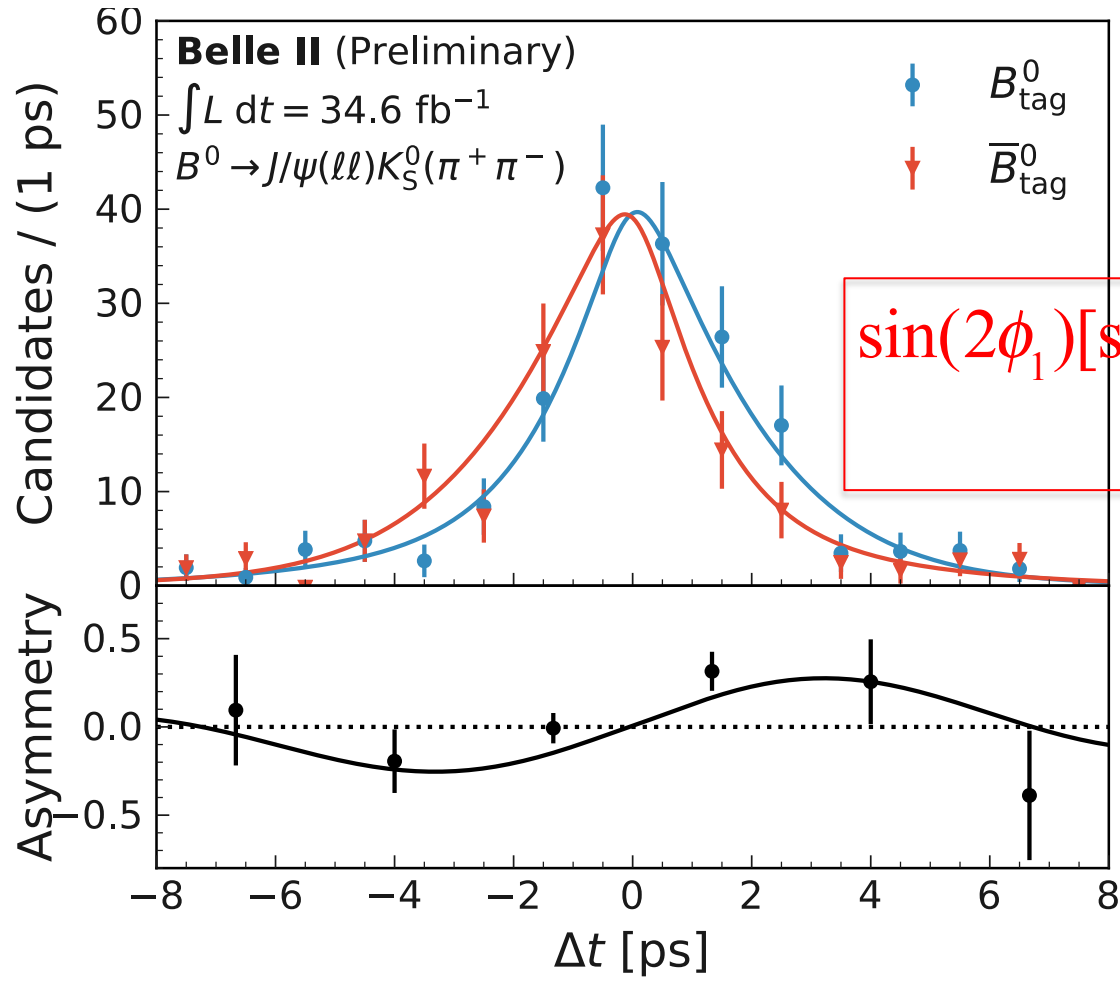


$$\text{Asym}(\text{mixing}) = \frac{OF - SF}{OF + SF}$$

$$N_{SF/OF} \sim \frac{\exp(-|\Delta t|/\tau)}{4\tau} [1 \pm (1 - 2w) \cos(\Delta m_d \Delta t)] \otimes R(\Delta t)$$

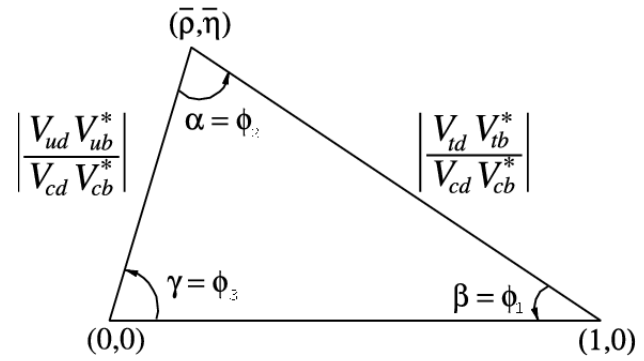


# Hint of **time-dependent CPV** from Belle II ( $2.7\sigma$ significance)



$$\sin(2\phi_1) [\sin(2\beta)] = 0.55 \pm 0.21 \pm 0.04$$

(WA=0.685±0.019)

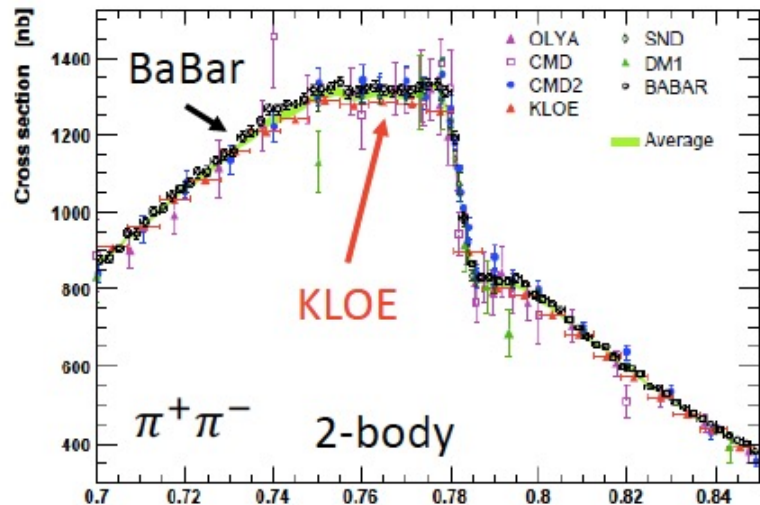


Based on the QM interference of  
 $B^0 \rightarrow f_{CP} ; B^0 \rightarrow \overline{B^0} \rightarrow f_{CP}$

Expect updates in summer 2022 (*x6 more data*)

$$N_{+/-} \sim \frac{\exp(-|\Delta t|/\tau)}{4\tau} \left\{ 1 \pm (1 - 2w) \sin(2\phi_1) \sin(\Delta m_d \Delta t) \right\} \otimes R(\Delta t)$$

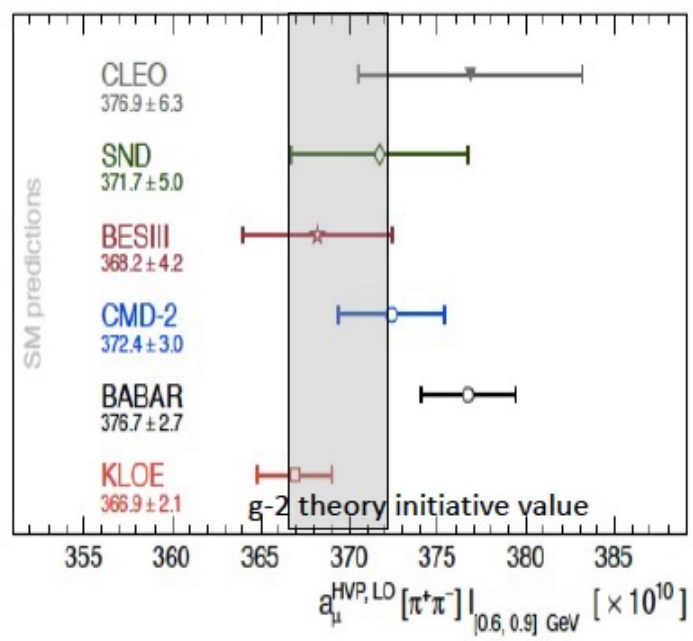
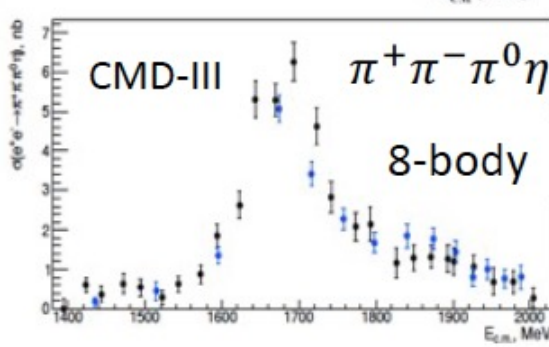
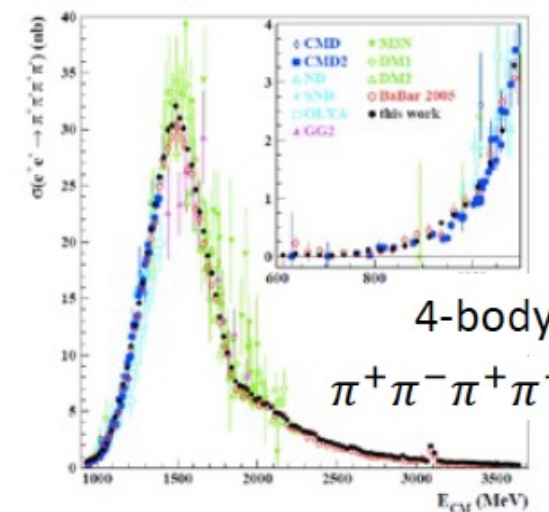
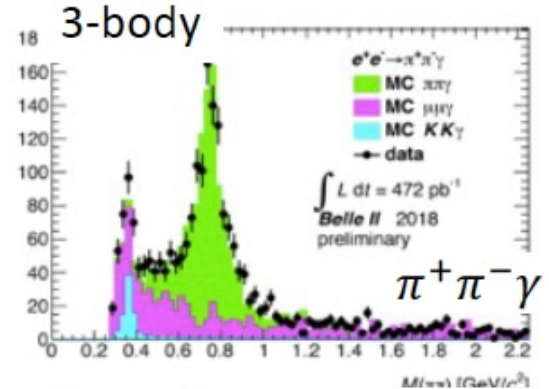
# Gold Standard: $e^+e^- \rightarrow \text{hadrons}$



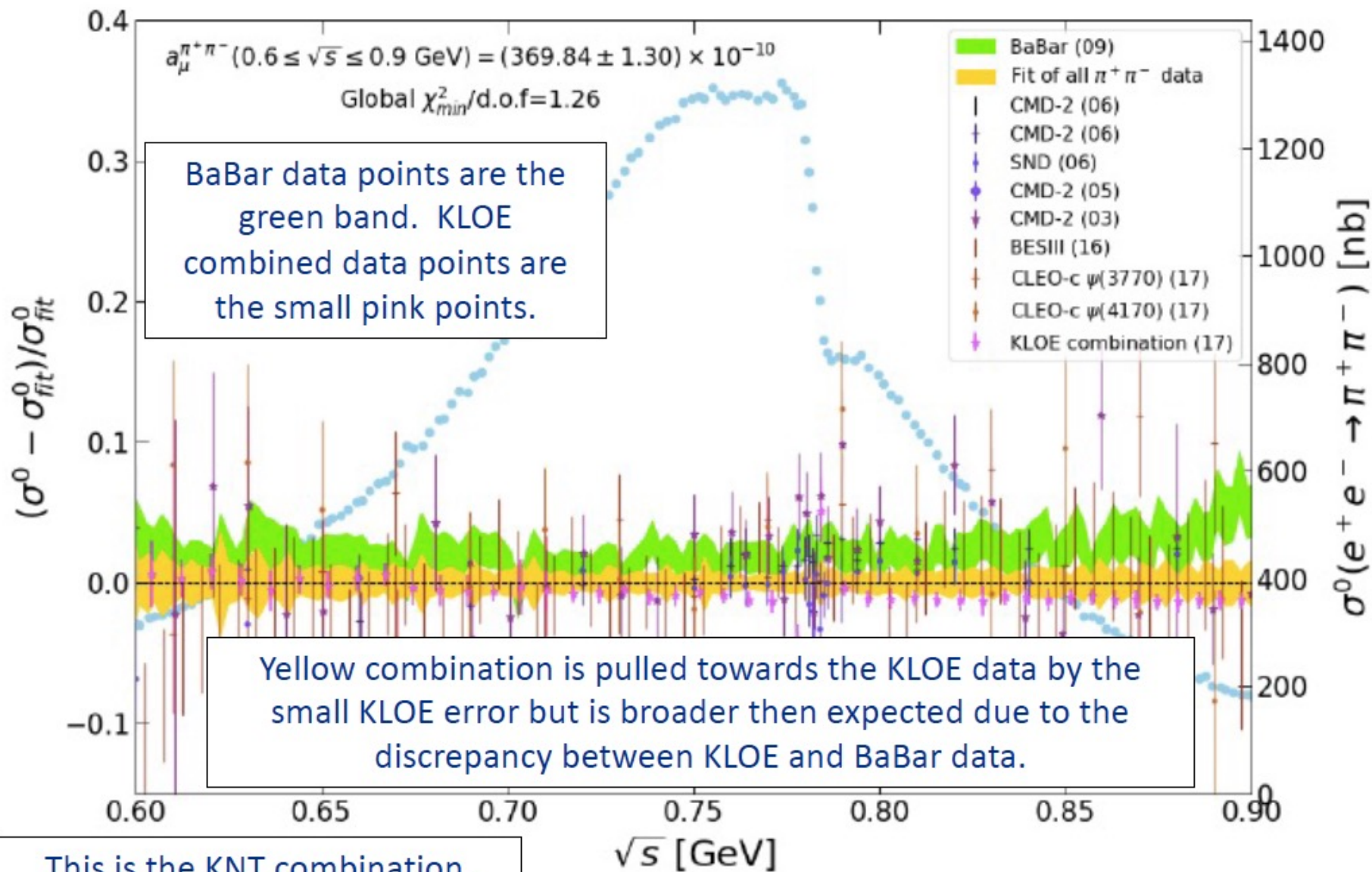
Channel	$a_e^{\text{had, LOVP}} \times 10^{14}$
Chiral perturbat	
$\pi^0\gamma$	$0.04 \pm 0.00$
$\pi^+\pi^-$	$0.31 \pm 0.01$
$\pi^+\pi^-\pi^0$	$0.00 \pm 0.00$
$\eta\gamma$	$0.00 \pm 0.00$

Excl	
$\pi^0\gamma$	$1.19 \pm 0.03$
$\pi^+\pi^-$	$138.59 \pm 0.54$
$\pi^+\pi^-\pi^0$	$12.29 \pm 0.25$
$\pi^+\pi^-\pi^+\pi^-$	$3.67 \pm 0.05$
$\pi^+\pi^-\pi^0\pi^0$	$4.80 \pm 0.19$
$(2\pi^+2\pi^-\pi^0)_{\text{no } \eta\omega}$	$0.24 \pm 0.02$
$(\pi^+\pi^-3\pi^0)_{\text{no } \eta}$	$0.15 \pm 0.03$
$(3\pi^+3\pi^-)_{\text{no } \omega}$	$0.06 \pm 0.00$
$(2\pi^+2\pi^-\pi^0)_{\text{no } \eta}$	$0.33 \pm 0.04$
$(\pi^+\pi^-4\pi^0)_{\text{no } \eta}$	$0.05 \pm 0.05$
$(3\pi^+3\pi^-\pi^0)_{\text{no } \eta\omega}$	$0.00 \pm 0.00$
$K^+K^-$	$5.86 \pm 0.06$
$K_S^0K_L^0$	$3.33 \pm 0.05$
$KK\pi$	$0.66 \pm 0.03$
$KK2\pi$	$0.47 \pm 0.02$
$KK3\pi$	$0.01 \pm 0.00$
$\eta\gamma$	$0.18 \pm 0.01$
$\eta\pi^+\pi^-$	$0.33 \pm 0.01$
$(\eta\pi^+\pi^-\pi^0)_{\text{no } \omega}$	$0.17 \pm 0.02$
$\eta2\pi^+2\pi^-$	$0.02 \pm 0.00$
$\eta\pi^+\pi^-\pi^0\pi^0$	$0.03 \pm 0.00$
$\eta\omega$	$0.07 \pm 0.01$
$\omega(\rightarrow \pi^0\gamma)\pi^0$	$0.22 \pm 0.00$
$\omega(\rightarrow \text{npp})2\pi$	$0.03 \pm 0.00$
$\omega(\rightarrow \text{npp})3\pi$	$0.04 \pm 0.01$
$\omega2\pi^+2\pi^-$	$0.00 \pm 0.00$
$\eta\phi$	$0.10 \pm 0.00$
$\omega\eta\pi^0$	$0.06 \pm 0.01$
$\omega(\rightarrow \text{npp})KK$	$0.00 \pm 0.00$
$\eta(\rightarrow \text{npp})KK_{\text{no } \phi \rightarrow KK}$	$0.00 \pm 0.00$
$\phi \rightarrow \text{unaccounted}$	$0.01 \pm 0.01$
$p\bar{p}$	$0.01 \pm 0.00$
$n\bar{n}$	$0.01 \pm 0.00$

Other	
Inclusive channel	$10.38 \pm 0.16$
$J/\psi$	$1.49 \pm 0.05$
$\psi'$	$0.37 \pm 0.01$
$\Upsilon(1S)$	$0.01 \pm 0.00$
$\Upsilon(2S)$	$0.00 \pm 0.00$
$\Upsilon(3S)$	$0.00 \pm 0.00$
$\Upsilon(4S)$	$0.00 \pm 0.00$
pQCD ( $\sqrt{s} > 11.199 \text{ GeV}$ )	$0.48 \pm 0.00$
Total ( $< \infty \text{ GeV}$ )	$186.08 \pm 0.66$



$$e^+e^- \rightarrow \pi^+\pi^-$$



This is the KNT combination.  
 The white paper combination  
 has about a 60% boarder error.



# Leading order QCD: HVP

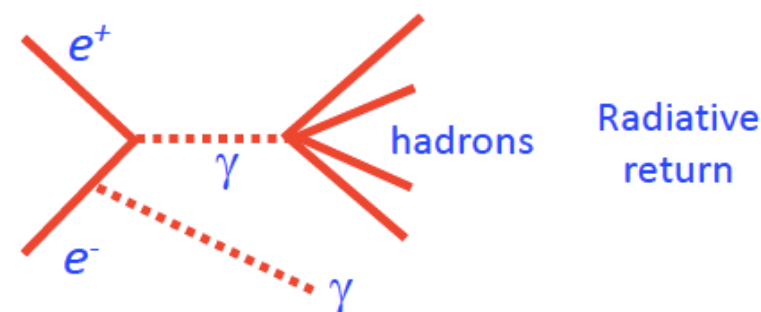
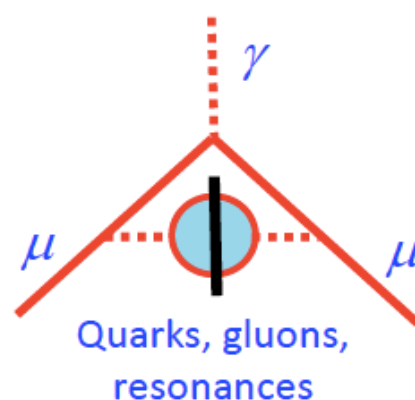
The leading order QCD contribution comes from Hadronic Vacuum Polarization

This can be taken directly from data by the measurement of the differential cross section  $e^+e^- \rightarrow \text{hadrons}$ .

The assumptions are analyticity and the optical theorem.

This is considered the gold standard and the Muon g-2 Theory initiative only uses this data in their prediction.

The  $1/s$  scaling puts significant weight to the two-pion low energy region around the  $\rho$  and  $\omega$  but data from all regions and all final states needs to be included.

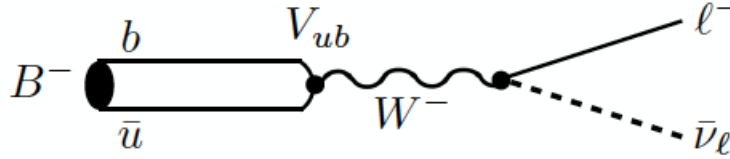


$$\frac{g-2}{2} HVP = \frac{\alpha^2}{3\pi^2} \int_{s_0}^{\infty} \frac{ds}{s} R(s) K_l(s)$$

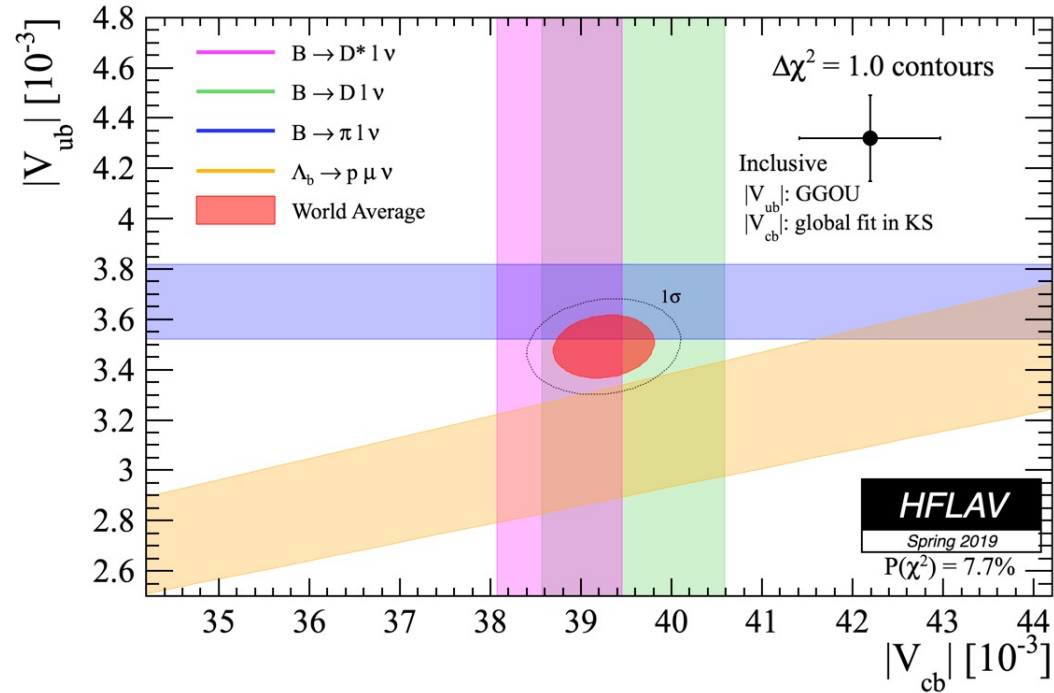
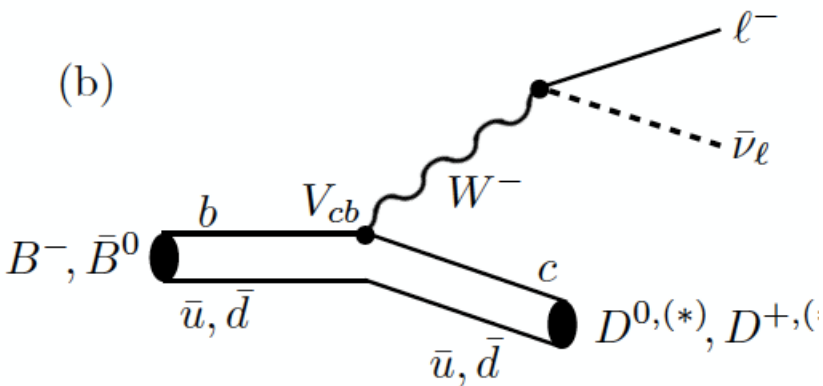


# Motivation for semileptonic decays: $V_{cb}$ , $V_{ub}$

(a)



(b)



a) Purely leptonic decays e.g.

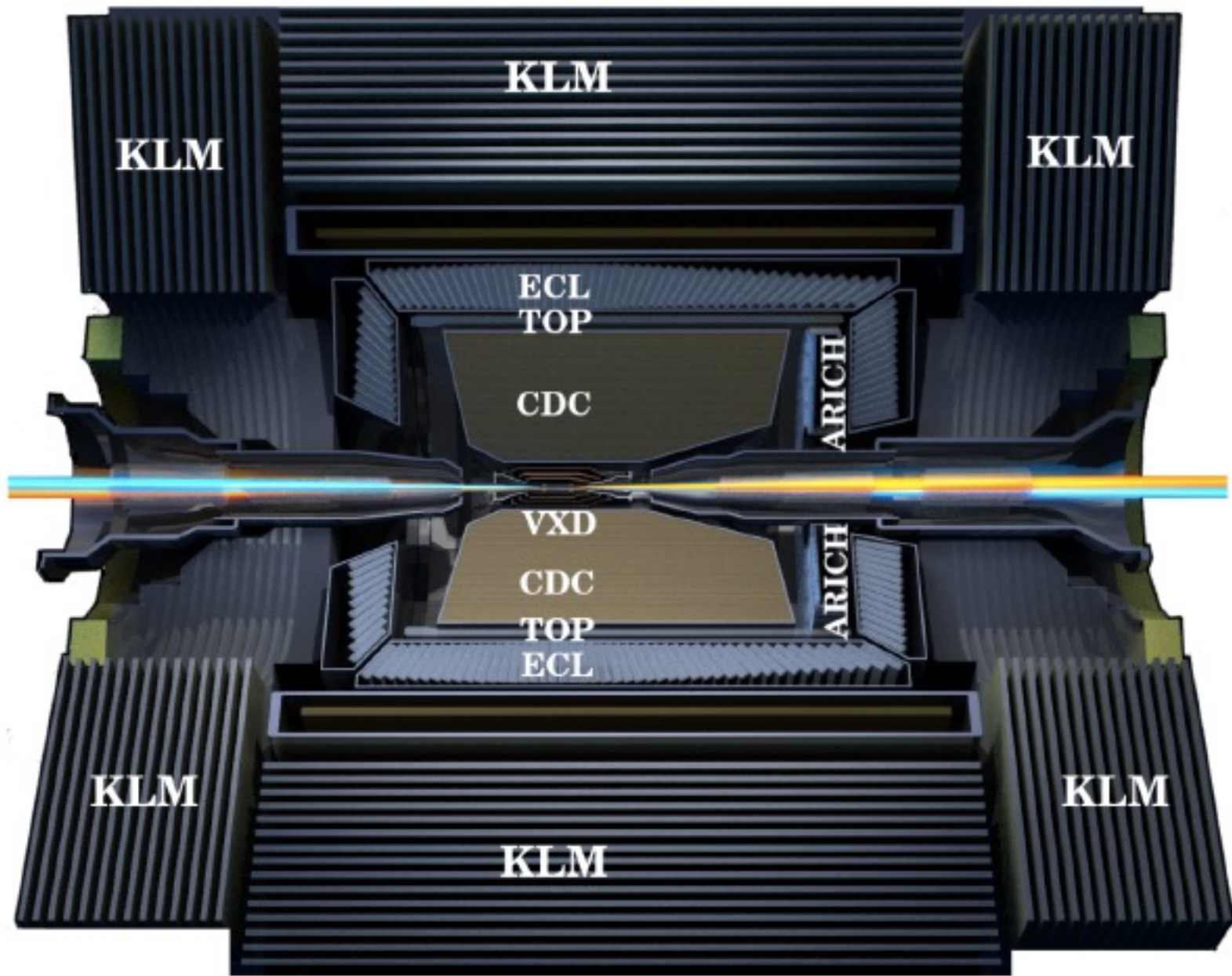


b) Semileptonic decays e.g.



Figure credit:

Tensions persist between **exclusive** and **inclusive** (e+e-) measurements of fundamental CKM elements  $|V_{cb}|$ ,  $|V_{ub}|$



# B → K\* 1+1-(q<sup>2</sup>) bootcamp at B2TIP

Angular dependence



(-) means the term is only in  $\Gamma - \bar{\Gamma}$

$$\frac{1}{d(\Gamma + \bar{\Gamma}) / dq^2} \frac{d^3(\Gamma + \bar{\Gamma})}{d\Omega} =$$

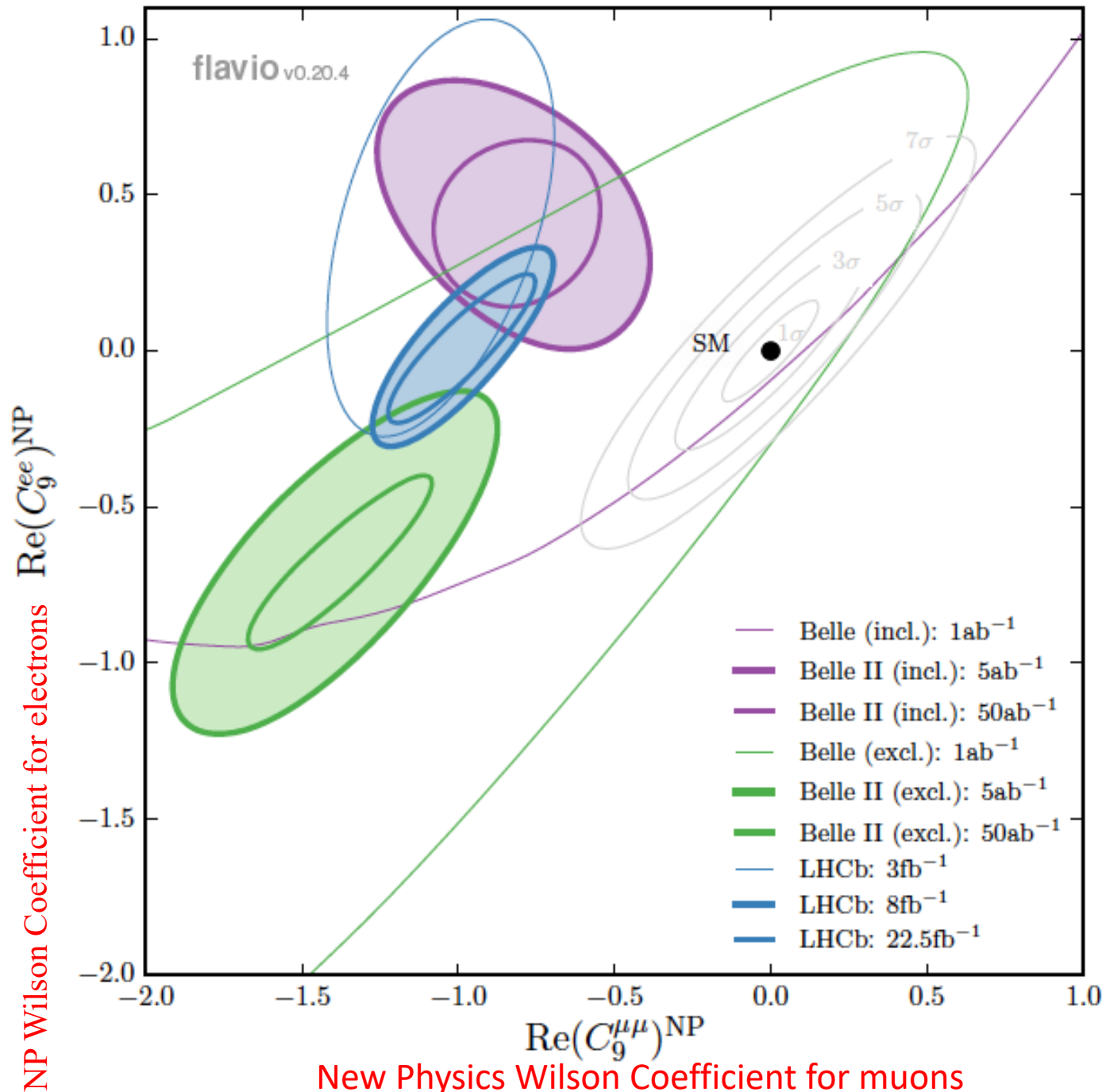
$F_L$  is the longitudinal polarization fraction.

$$\frac{9}{32\pi} \left[ \begin{aligned} & \frac{3}{4}(1 - F_L) \sin^2 \vartheta_K + F_L \cos^2 \vartheta_K \\ & + \frac{1}{4}(1 - F_L) \sin^2 \vartheta_K \cos 2\vartheta_L \\ & - F_L \cos^2 \vartheta_K \cos 2\vartheta_L + S_3 \sin^2 \vartheta_K \sin^2 \vartheta_L \cos 2\phi \\ & + S_4 \sin 2\vartheta_K \sin 2\vartheta_L \cos \phi + \boxed{\phantom{S_5 \sin 2\vartheta_K \sin 2\vartheta_L \sin \phi}} \\ & + \boxed{\phantom{S_6 \sin 2\vartheta_K \sin 2\vartheta_L \sin \phi}} + S_7 \sin 2\vartheta_K \sin \vartheta_L \sin \phi \\ & + \boxed{\phantom{S_8 \sin 2\vartheta_K \sin 2\vartheta_L \sin \phi}} \end{aligned} \right]$$

*Introduce  $P_{4,5}' = S_{4,5} / \text{sqrt}[F_L(1 - F_L)]$  to reduce dependence on form factors*

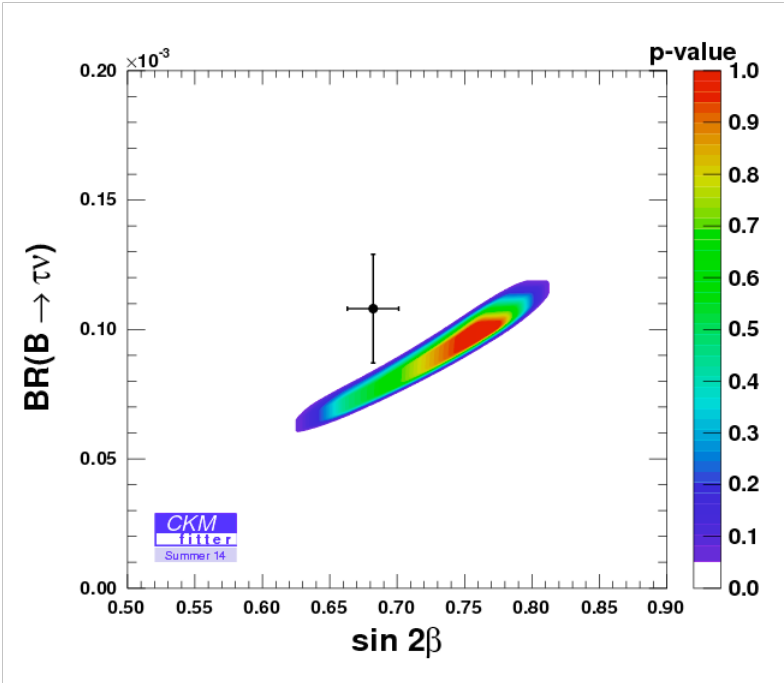
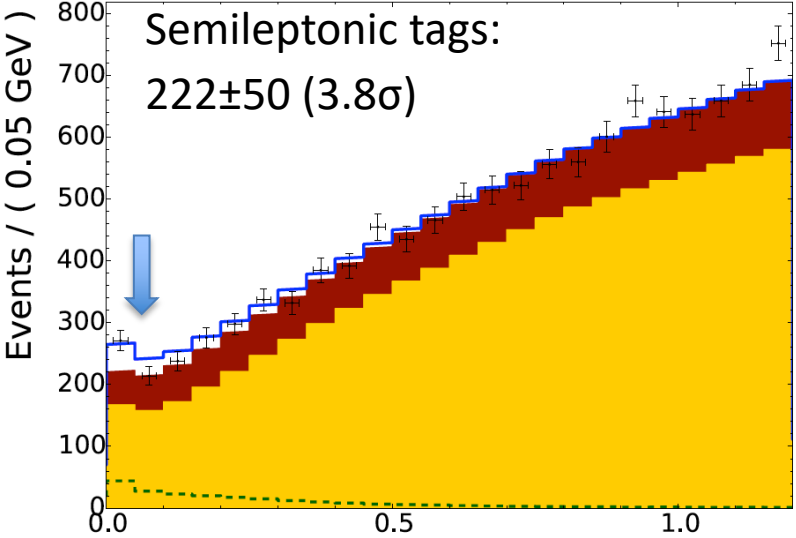
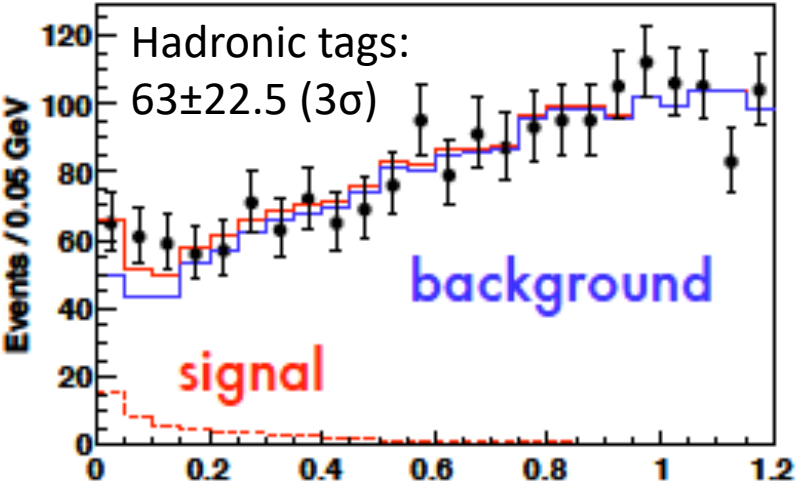
# NP in $b \rightarrow s |l^+l^-$

Prepared by D. Straub et al. for the Belle II Physics Book (edited by P. Urquijo and E. Kou)



Belle II can do both inclusive and exclusive. Equally strong capabilities for electrons and muons.

Example: **old Belle  $B \rightarrow \tau \nu$  results** with full *reprocessed* data sample: either hadronic or semileptonic tags (PRD 92, 051102 (2015))



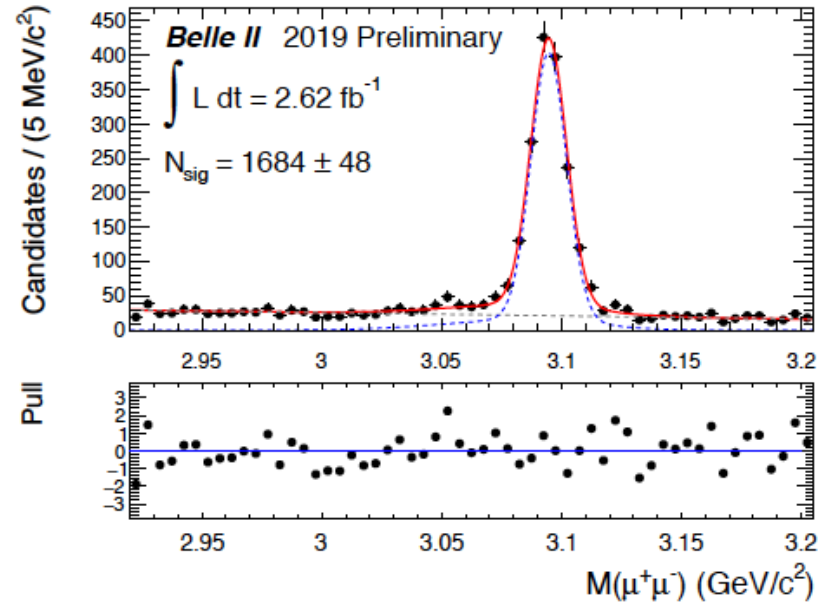
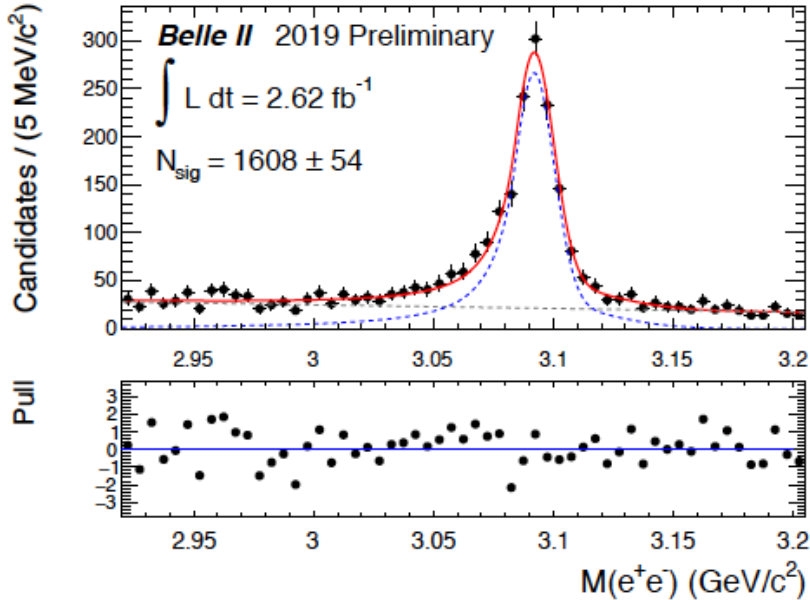
With the full B factory statistics only “evidence”. No single observation from either Belle or BaBar.

➔ The horizontal axis is the “Extra Calorimeter Energy” or  $E_{ECL}$





# Signals for $B \rightarrow J/\psi X$ in Phase 3 data



Clear signals for  $B \rightarrow J/\psi X$  in  $\sim 1/2$  of Phase 3 data. Note the small radiative tail on the di-electrons (does include bremsstrahlung recovery).

*$\rightarrow$  Belle II has equally strong capabilities for electrons and muons.*