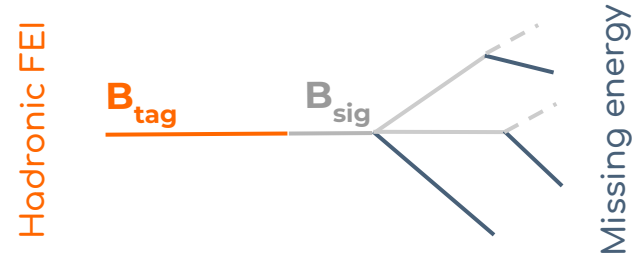


Digging into hadronic B decays

BAW India 2022 @ IISER Mohali

Vidya Sagar Vobbiliseti, Karim Trabelsi

17 December 2022



Motivation for studying B-tagging

$B \rightarrow K \tau l/\tau$ searches rely on the purity of B-tagging

$B^+ \rightarrow K^+ \tau l$ has 1 - 2 neutrinos in the final state

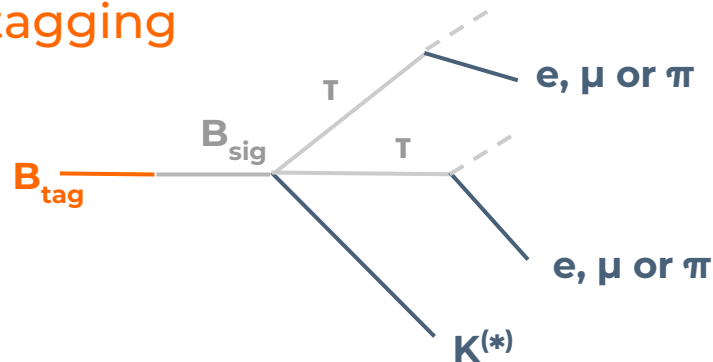
$B^+ \rightarrow K^+ \tau \tau$ has 2 - 4 neutrinos in the final state

⇒ Huge background

⇒ Requires high purity in the tag-side

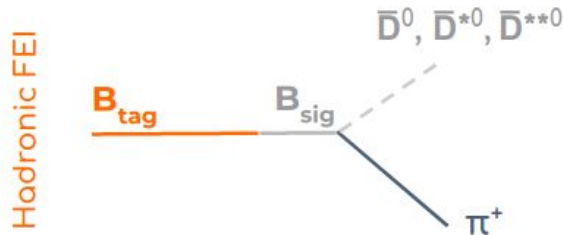
For hadronic B_{tag} : ϵ_{tag} (<1%) is a limiting factor.

Many interesting B-physics studies involve missing energy: $D^{(*)} \tau \nu$, $K^{(*)} \tau l$, $K^{(*)} \tau \tau$, $K^{(*)} \nu \nu$, $\pi l \nu$, τl , $\tau \nu$, $\mu \nu$... which require B-tagging.



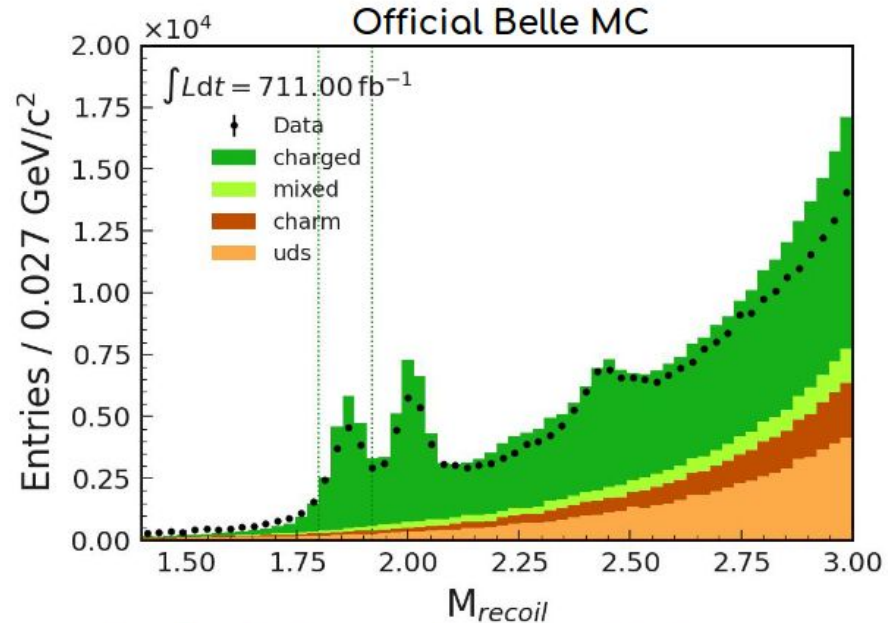
Irrespective of tagging strategy, optimal MC modeling is essential for good performance of ML techniques (NN/BDT).

Partial reconstruction for more statistics!



We can look for D^0, D^{*0} and even D^{**0} in the recoil mass of a fully reconstructed B and a π^\pm

Within a narrow region around the peak, we know that one B decays to $D^0\pi^+$ and we can study the other B (decaying hadronically)

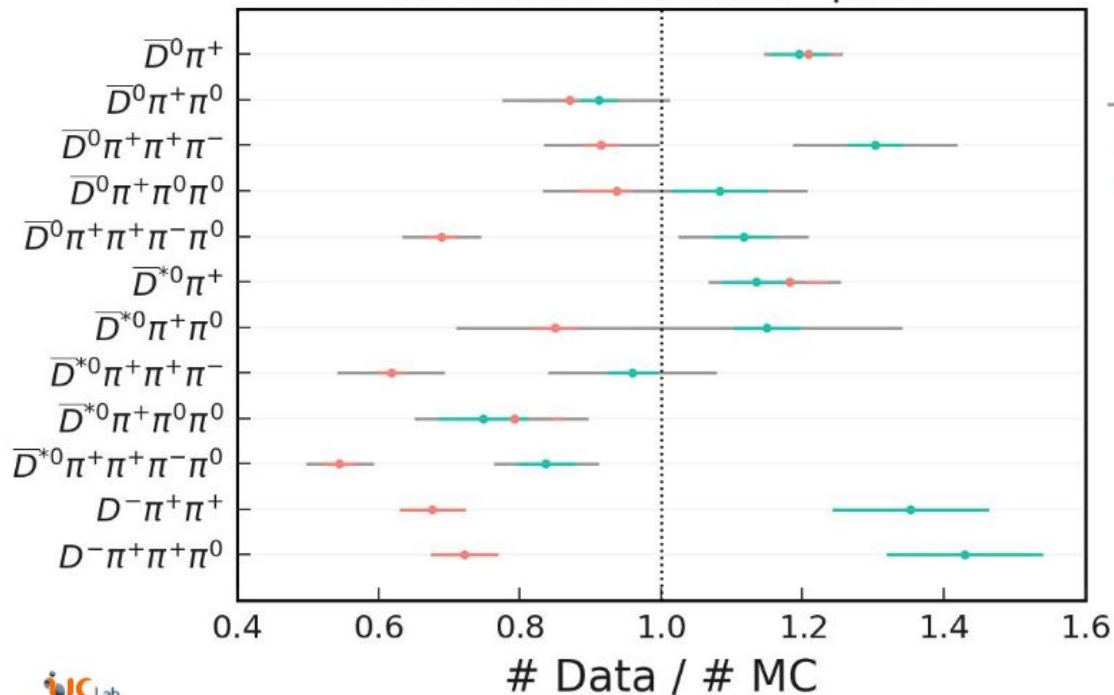


-16k events in a 3σ window around each peak in data.
Roughly $\frac{1}{3}$ statistics of $X_c \text{ lv}$ sample, but much smaller systematic.

[BELLE2-NOTE-PH-2021-029, Belle note bn1615]

Updated calibration factors per mode

3σ window around D^0 peak



$\int Ldt = 711 \text{ fb}^{-1}$

- PDG uncertainty
- Official MC
- Proposed MC

Overall calibration factor:

$(82.6 \pm 0.9)\%$

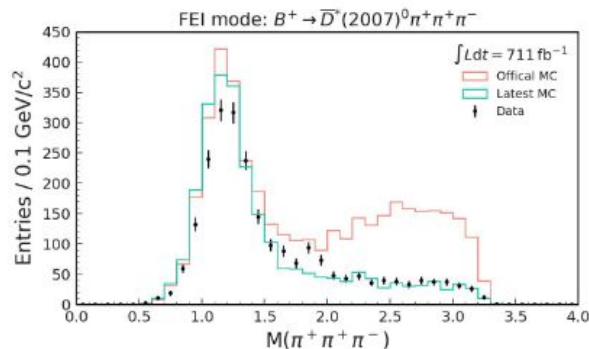
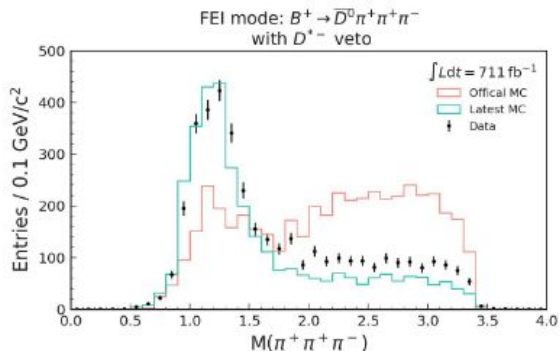


$(104.2 \pm 1.2)\%$

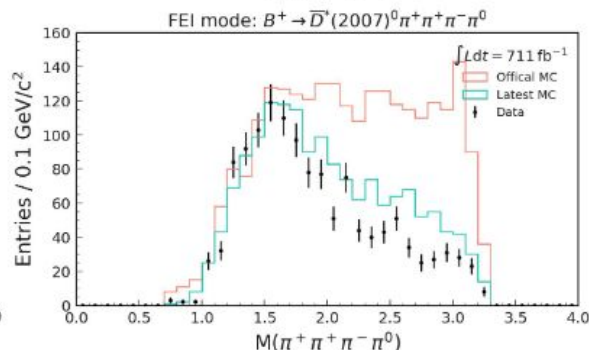
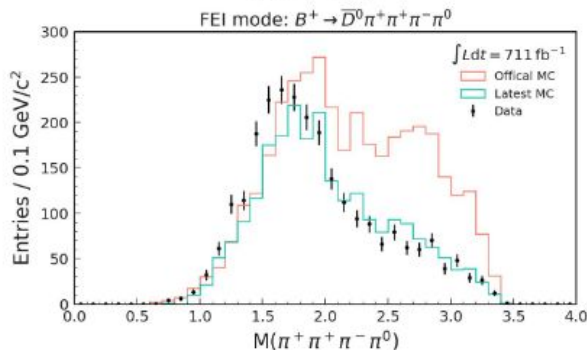
Decay description is improved!

The improvement is not limited to calibration factors, but more importantly in the invariant masses (of intermediate particles), which are used as training variables in FEI

$3\pi^\pm$ case:



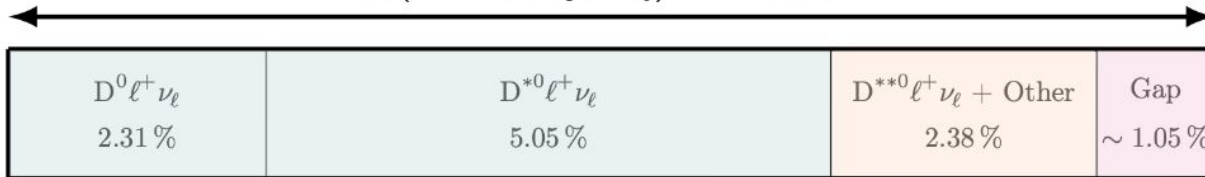
$3\pi^\pm \pi^0$ case:



Semi-Leptonic gap

[Raynette van Tonder]

$$\mathcal{B}(B^+ \rightarrow X_c^0 \ell^+ \nu_\ell) \approx 10.79\%$$



Decay	$\mathcal{B}(B^+)$	$\mathcal{B}(B^0)$
$B \rightarrow D \ell^+ \nu_\ell$	$(2.4 \pm 0.1) \times 10^{-2}$	$(2.2 \pm 0.1) \times 10^{-2}$
$B \rightarrow D^* \ell^+ \nu_\ell$	$(5.5 \pm 0.1) \times 10^{-2}$	$(5.1 \pm 0.1) \times 10^{-2}$
$B \rightarrow D_1 \ell^+ \nu_\ell$	$(6.6 \pm 0.1) \times 10^{-3}$	$(6.2 \pm 0.1) \times 10^{-3}$
$B \rightarrow D_2^* \ell^+ \nu_\ell$	$(2.9 \pm 0.3) \times 10^{-3}$	$(2.7 \pm 0.3) \times 10^{-3}$
$B \rightarrow D_0^* \ell^+ \nu_\ell$	$(4.2 \pm 0.8) \times 10^{-3}$	$(3.9 \pm 0.7) \times 10^{-3}$
$B \rightarrow D_1' \ell^+ \nu_\ell$	$(4.2 \pm 0.9) \times 10^{-3}$	$(3.9 \pm 0.8) \times 10^{-3}$
$B \rightarrow D \pi \pi \ell^+ \nu_\ell$	$(0.6 \pm 0.9) \times 10^{-3}$	$(0.6 \pm 0.9) \times 10^{-3}$
$B \rightarrow D^* \pi \pi \ell^+ \nu_\ell$	$(2.2 \pm 1.0) \times 10^{-3}$	$(2.0 \pm 1.0) \times 10^{-3}$
$B \rightarrow X_c \ell \nu_\ell$	$(10.8 \pm 0.4) \times 10^{-2}$	$(10.1 \pm 0.4) \times 10^{-2}$



Fairly well known.
Some iso-spin tension.



Broad states based on
3 measurements.
(BaBar, Belle, DELPHI)



Some hints from
the BaBar result.

This gap leads to up to 3σ difference in V_{cb} measured from inclusive vs exclusive.

[1507.08303]

Semi-Leptonic gap: Filled with η ?

[Raynette van Tonder]

Model 1:

Equidistribution of all final state particles in phase space

Decay	$\mathcal{B}(B^+)$	$\mathcal{B}(B^0)$
$B \rightarrow D \ell^+ \nu_\ell$	$(2.4 \pm 0.1) \times 10^{-2}$	$(2.2 \pm 0.1) \times 10^{-2}$
$B \rightarrow D^* \ell^+ \nu_\ell$	$(5.5 \pm 0.1) \times 10^{-2}$	$(5.1 \pm 0.1) \times 10^{-2}$
$B \rightarrow D_1 \ell^+ \nu_\ell$	$(6.6 \pm 0.1) \times 10^{-3}$	$(6.2 \pm 0.1) \times 10^{-3}$
$B \rightarrow D_2^* \ell^+ \nu_\ell$	$(2.9 \pm 0.3) \times 10^{-3}$	$(2.7 \pm 0.3) \times 10^{-3}$
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$B \rightarrow D \pi \pi \ell^+ \nu_\ell$	$(0.6 \pm 0.9) \times 10^{-3}$	$(0.6 \pm 0.9) \times 10^{-3}$
$B \rightarrow D^* \pi \pi \ell^+ \nu_\ell$	$(2.2 \pm 1.0) \times 10^{-3}$	$(2.0 \pm 1.0) \times 10^{-3}$
$B \rightarrow D \eta \ell^+ \nu_\ell$	$(4.0 \pm 4.0) \times 10^{-3}$	$(4.0 \pm 4.0) \times 10^{-3}$
$B \rightarrow D^* \eta \ell^+ \nu_\ell$	$(4.0 \pm 4.0) \times 10^{-3}$	$(4.0 \pm 4.0) \times 10^{-3}$
$B \rightarrow X_c \ell \nu_\ell$	$(10.8 \pm 0.4) \times 10^{-2}$	$(10.1 \pm 0.4) \times 10^{-2}$

Model 2:

Decay via intermediate broad D^{**} state

Decay	$\mathcal{B}(B^+)$	$\mathcal{B}(B^0)$
$B \rightarrow D_0^* \ell^+ \nu_\ell$ ($\hookrightarrow D \pi \pi$)	$(0.03 \pm 0.03) \times 10^{-2}$	$(0.03 \pm 0.03) \times 10^{-2}$
$B \rightarrow D_1^* \ell^+ \nu_\ell$ ($\hookrightarrow D \pi \pi$)	$(0.03 \pm 0.03) \times 10^{-2}$	$(0.03 \pm 0.03) \times 10^{-2}$
$B \rightarrow D_0^* \pi \pi \ell^+ \nu_\ell$ ($\hookrightarrow D^* \pi \pi$)	$(0.108 \pm 0.051) \times 10^{-2}$	$(0.101 \pm 0.048) \times 10^{-2}$
$B \rightarrow D_1^* \pi \pi \ell^+ \nu_\ell$ ($\hookrightarrow D^* \pi \pi$)	$(0.108 \pm 0.051) \times 10^{-2}$	$(0.101 \pm 0.048) \times 10^{-2}$
$B \rightarrow D_0^* \ell^+ \nu_\ell$ ($\hookrightarrow D \eta$)	$(0.396 \pm 0.396) \times 10^{-2}$	$(0.399 \pm 0.399) \times 10^{-2}$
$B \rightarrow D_1^* \ell^+ \nu_\ell$ ($\hookrightarrow D^* \eta$)	$(0.396 \pm 0.396) \times 10^{-2}$	$(0.399 \pm 0.399) \times 10^{-2}$

The current workaround to explain the SL gap is to fill it with $D^{(*)} \eta \ell \nu$, either as a non-resonant state or through $(D^{(*)} \eta)$ resonance.

But never seen.

Source of η : D^{**} ?

D^{**}	decay channel	branching ratio
$D_0(2300)^0$	$D^0\pi^0$	0.3333
	$D^+\pi^-$	0.6667
$D_1(2420)^0$	$D^{*0}\pi^0$	0.1997
	$D^{*+}\pi^-$	0.3994
	$D^0\pi^+\pi^-$	0.1719
	$D^0\pi^0\pi^0$	0.1145
	$D^+\pi^-\pi^0$	0.1145
$D_1(2430)^0$	$D^{*+}\pi^-$	0.6667
	$D^{*0}\pi^0$	0.3333
$D_2^*(2460)^0$	$D^{*0}\pi^0$	0.1334
	$D^{*+}\pi^-$	0.2669
	$D^0\pi^0$	0.1999
	$D^+\pi^-$	0.3998

TABLE XIX: Decay channels of D^{**}

Model 2:

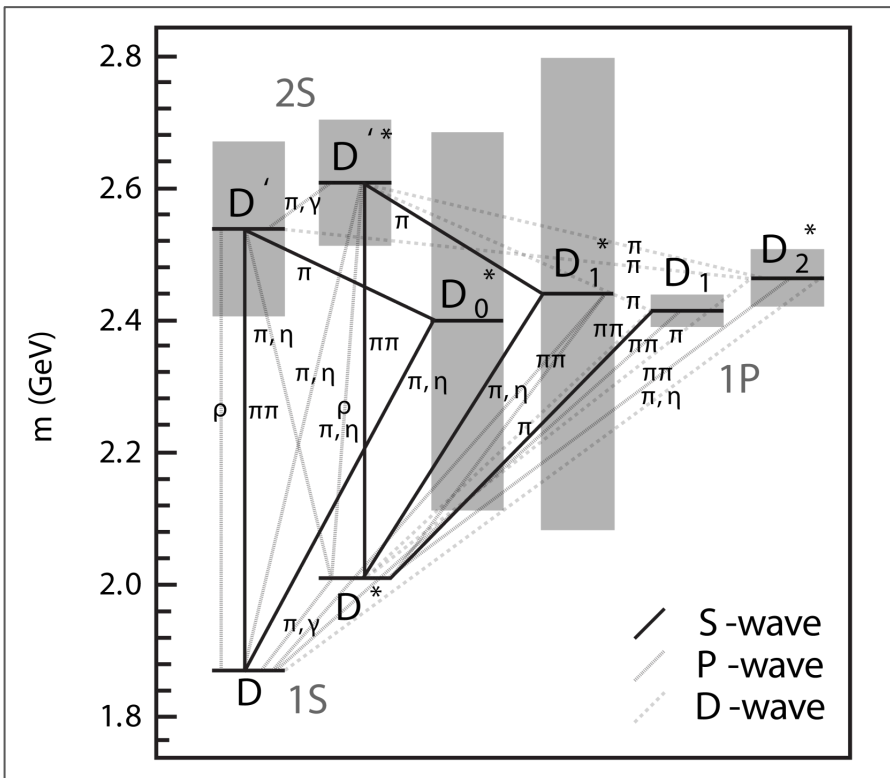
Decay via intermediate broad D^{**} state

Decay	$\mathcal{B}(B^+)$	$\mathcal{B}(B^0)$
$B \rightarrow D_0^* \ell^+ \nu_\ell$ ($\leftrightarrow D\pi\pi$)	$(0.03 \pm 0.03) \times 10^{-2}$	$(0.03 \pm 0.03) \times 10^{-2}$
$B \rightarrow D_1^* \ell^+ \nu_\ell$ ($\leftrightarrow D\pi\pi$)	$(0.03 \pm 0.03) \times 10^{-2}$	$(0.03 \pm 0.03) \times 10^{-2}$
$B \rightarrow D_0^* \pi\pi \ell^+ \nu_\ell$ ($\leftrightarrow D^* \pi\pi$)	$(0.108 \pm 0.051) \times 10^{-2}$	$(0.101 \pm 0.048) \times 10^{-2}$
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$B \rightarrow D_0^* \ell^+ \nu_\ell$ ($\leftrightarrow D\eta$)	$(0.396 \pm 0.396) \times 10^{-2}$	$(0.399 \pm 0.399) \times 10^{-2}$
$B \rightarrow D_1^* \ell^+ \nu_\ell$ ($\leftrightarrow D^* \eta$)	$(0.396 \pm 0.396) \times 10^{-2}$	$(0.399 \pm 0.399) \times 10^{-2}$

The decays of D^{**} are not well measured, and the Belle II model does not consider η .

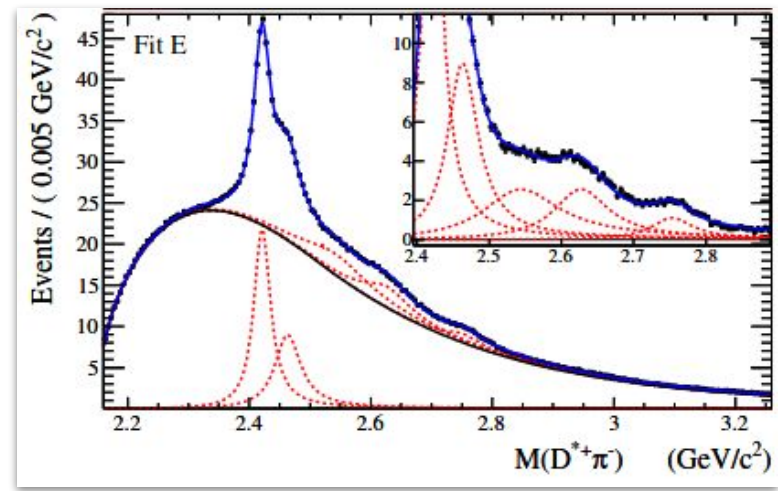
D^{**} decays and $B \rightarrow D^{**} X$ decays needs further studies.

Source of η : D(2S)?



In 2010, BaBar observed even higher D resonances, consistent with L=2.

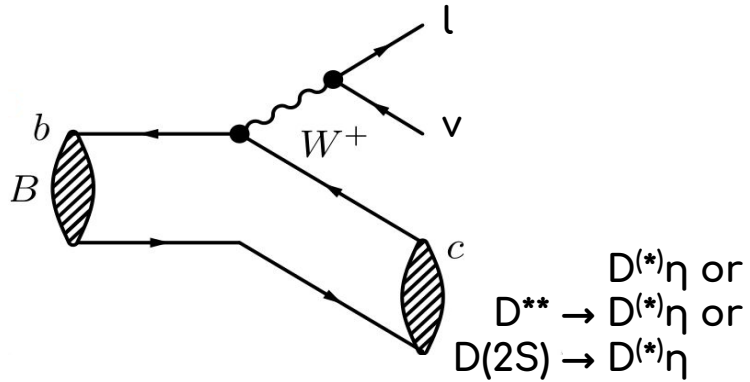
[1009.2076]



These D(2S) resonances have higher mass, and are potential candidates for sources of η filling the SL gap.

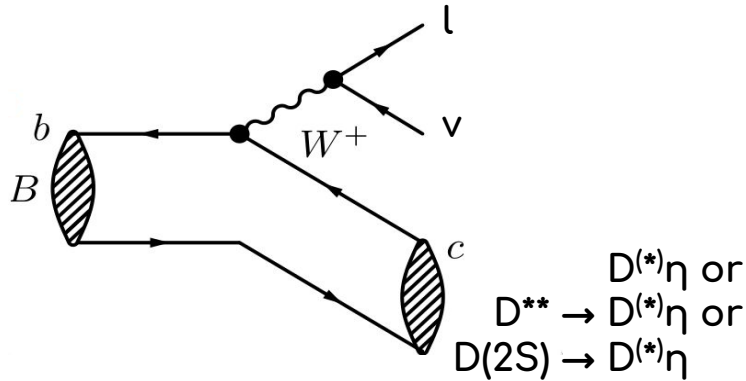
[arXiv:1202.1834]

SL $D^{(*)}\eta lv$

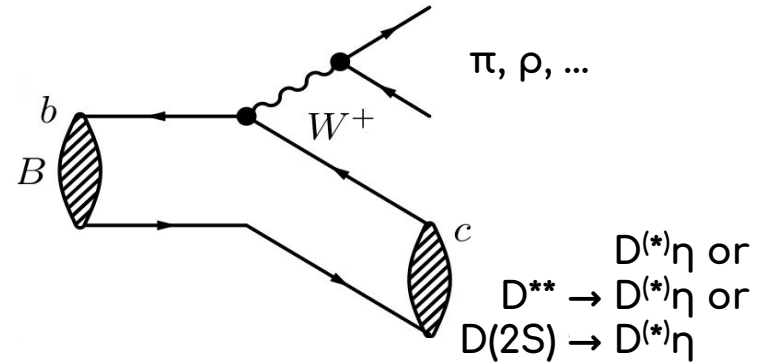


Signals of these SL decays are difficult to search for.

SL $D^{(*)}\eta lv \Rightarrow$ Hadronic $D^{(*)}\eta\pi, D^{(*)}\eta\rho$



Signals of these SL decays are difficult to search for.

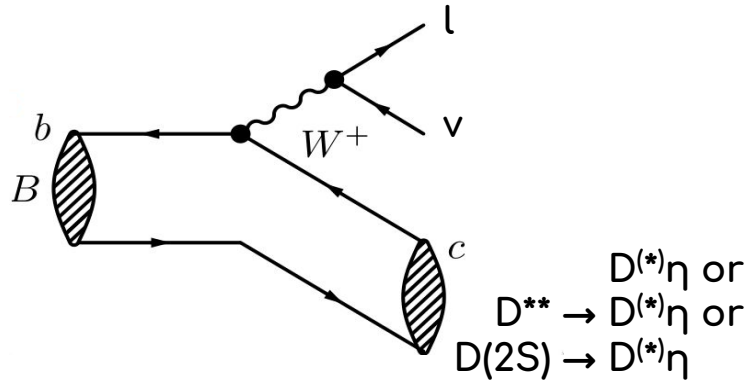


But the hadronic counterparts (changing lv with π/ρ) are easier to search.

The presence of $D^{(*)}\eta\pi$ can validate the assumption of η filling the SL-gap and can also describe the source of η .

Vismaya will talk more about the status.

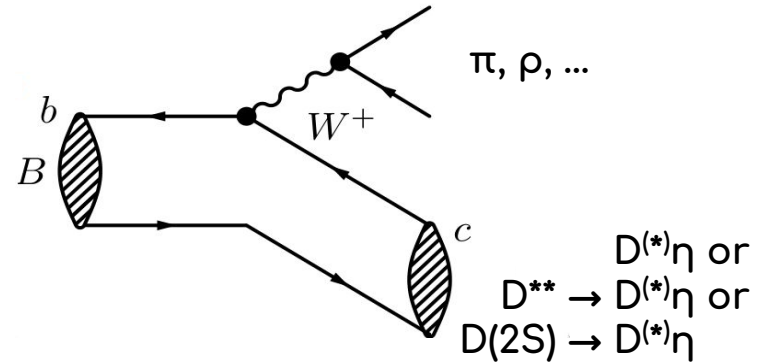
SL $D^{(*)}\eta lv \Rightarrow$ Hadronic $D^{(*)}\eta\pi, D^{(*)}\eta\rho$



Signals of these SL decays are difficult to search for.

$B \rightarrow D^*\pi$ is 1/10 of $B \rightarrow D^*lv$.

\Rightarrow A limit of $BF(B \rightarrow D^*\eta\pi) < 4 \times 10^{-4}$ is enough to invalidate η as a candidate for SL gap.

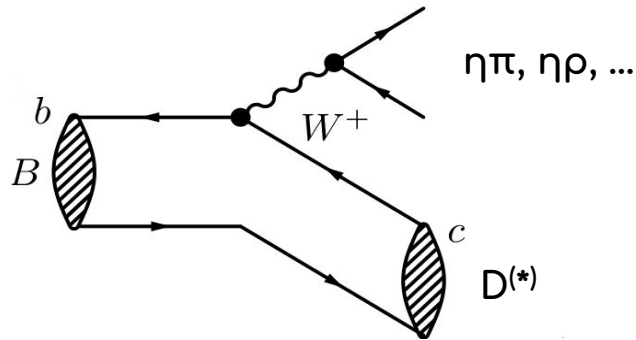
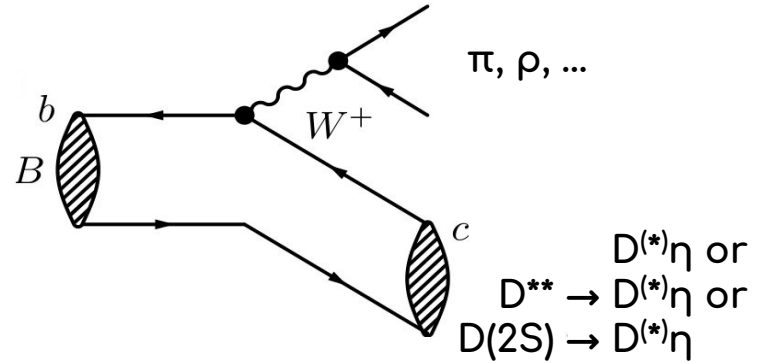


But the hadronic counterparts (changing lv with π/ρ) are easier to search.

The presence of $D^{(*)}\eta\pi$ can validate the assumption of η filling the SL-gap and can also describe the source of η .

Vismaya will talk more about the status.

Hadronic $D^{(*)}\eta\pi$ vs $D^{(*)}\eta\rho$

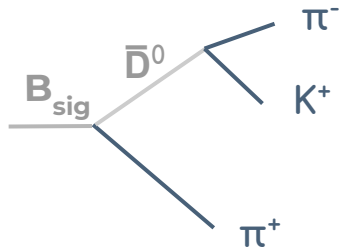


In the alternative way of producing η through W , the $\eta\pi$ contribution is suppressed.
 G-parity violation \Rightarrow Second class current.
 (also seen in τ decays)

But $\eta\rho$ is still possible.

So, studying both $D^{(*)}\eta\pi$ vs $D^{(*)}\eta\rho$ simultaneously can also shed light on the source of η .

Exclusive reconstruction

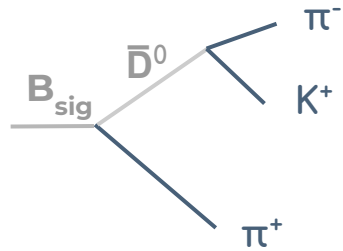


Reconstruct all the final state particles from the B
 \Rightarrow Calculate the 4-momentum of B.
And apply selection using ΔE (and M_{bc})

Efficiency =

$$BR_{\bar{D}^0 \rightarrow K^+ \pi^-} \times \epsilon_K \times \epsilon_{\pi^-} \times \epsilon_{\pi^+}$$

Exclusive vs Partial reconstruction

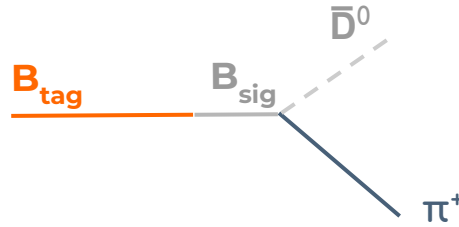


Reconstruct all the final state particles from the B
 \Rightarrow Calculate the 4-momentum of B.
 And apply selection using ΔE (and M_{bc})

Efficiency =

$$\text{BR}_{\bar{D}^0 \rightarrow K^+ \pi^-} \times \epsilon_K \times \epsilon_{\pi^-} \times \epsilon_{\pi^+}$$

Hadronic FEI



Popular when there are neutrinos which cannot be reconstructed, like in $B \rightarrow K \tau \bar{\nu}_\tau$

Instead of reconstructing the D exclusively, one could reconstruct the other B in the event fully. And look for the D in the recoil mass.

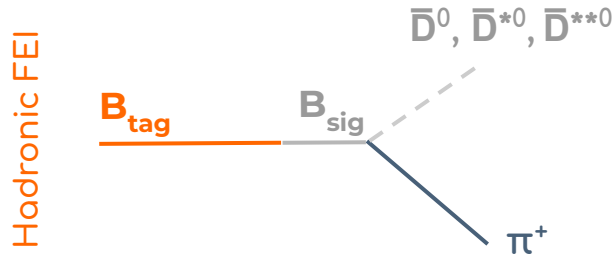
In CM frame of $Y(4S)$:

$$\begin{aligned} \vec{p}_{B_{sig}} &= -\vec{p}_{B_{tag}} \\ \vec{p}_X &= \vec{p}_{B_{sig}} - \vec{p}_{\pi^+} \\ E_X &= E_{beam} - E_{\pi^+} \\ M_X &= \sqrt{E_X^2 - \vec{p}_X^2} \end{aligned}$$

Here, efficiency =

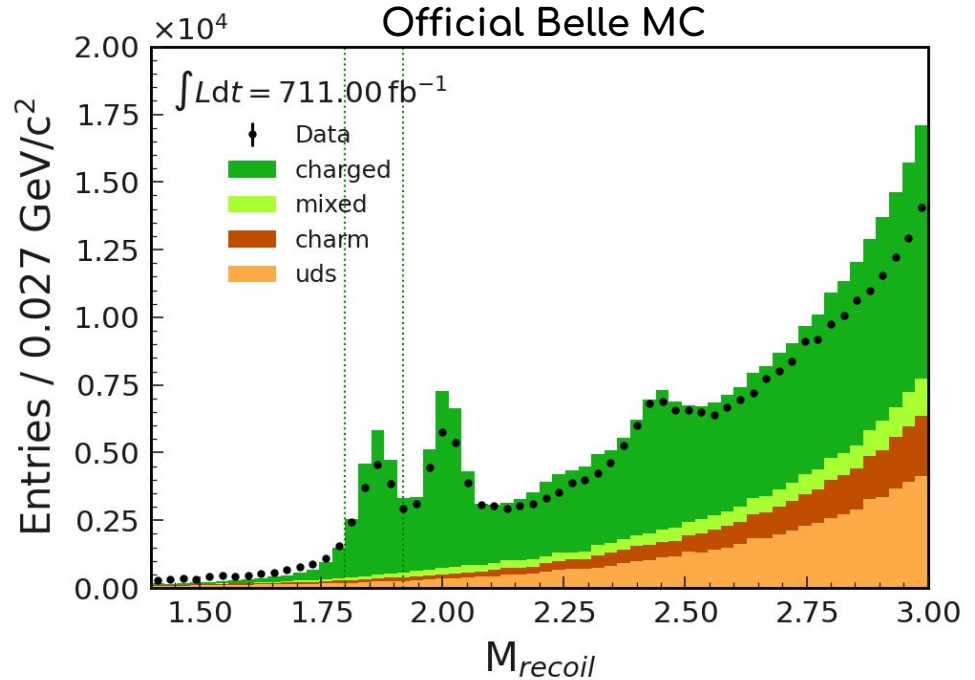
$$\epsilon_{B\text{-tag}} \times \epsilon_{\pi}$$

Recoil with π

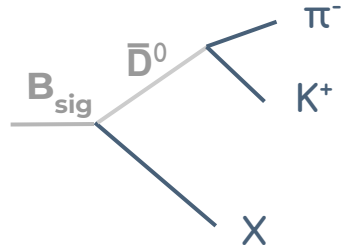


We can look for D^0, D^{*0} and even D^{**0} in the recoil mass of a fully reconstructed B and a π^\pm

Within a narrow region around the peak, we know that one B decays to $D^0\pi^+$ and we can study the other B (decaying hadronically)



Exclusive vs Partial reconstruction



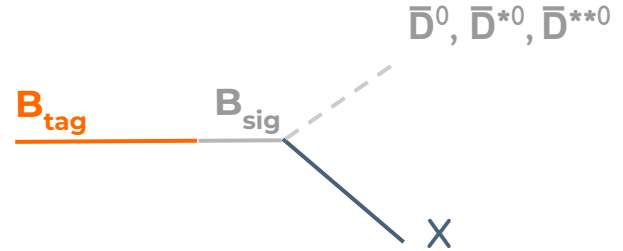
Efficiency =

for D^0 : $(BR_{\bar{D}^0 \rightarrow K \pi} \times \epsilon_K \times \epsilon_\pi) \times \epsilon_X$

for D^{*0} : $(BR_{\bar{D}^{*0} \rightarrow \bar{D}^0 \pi^0} \times \epsilon_{\pi^0} \times BR_{\bar{D}^0 \rightarrow K \pi} \times \epsilon_K \times \epsilon_\pi) \times \epsilon_X$

Here, D^* has lower efficiency than D .

Hadronic FEI

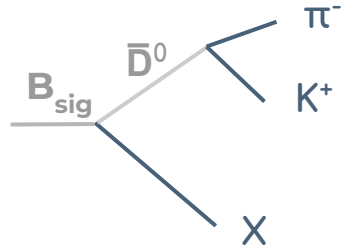


Efficiency =

$\epsilon_{B\text{-tag}} \times \epsilon_X$

Here D^* and D have same efficiency!

Exclusive vs Partial reconstruction



To extend on this idea, we are not limited to π .

X can be anything like $\pi\pi^0$ (ρ), $\pi\pi\pi$ (ω), $\eta\pi$, $\eta\rho$, $\omega\pi$, KK_S , KK^*?!

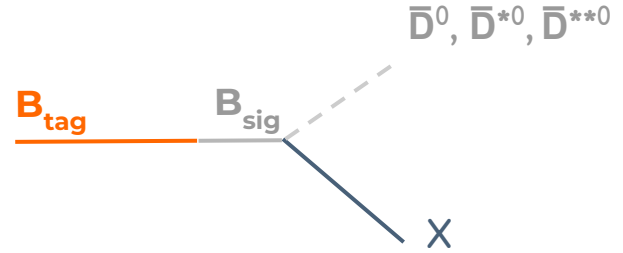
Efficiency =

for D^0 : $(BR_{D^0 \rightarrow K \pi} \times \epsilon_K \times \epsilon_\pi) \times \epsilon_X$

for D^{*0} : $(BR_{D^{*0} \rightarrow D^0 \pi^0} \times \epsilon_{\pi^0} \times BR_{D^0 \rightarrow K \pi} \times \epsilon_K \times \epsilon_\pi) \times \epsilon_X$

Here, D^* has lower efficiency than D .

Hadronic FEI

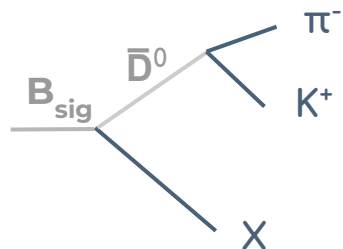


Efficiency =

$\epsilon_{B\text{-tag}} \times \epsilon_X$

Here D^* and D have same efficiency!

Exclusive vs Partial reconstruction

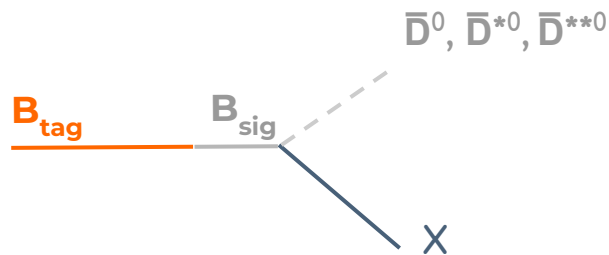


To extend on this idea, we are not limited to π .

X can be anything like $\pi\pi^0$ (ρ), $\pi\pi\pi$ (ω), $\eta\pi$, $\eta\rho$, $\omega\pi$, KK_S , KK^*?!

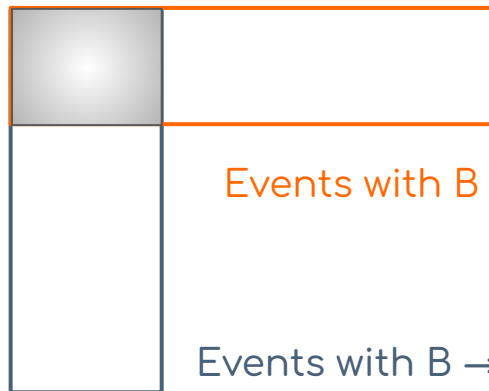
Here, D^* has lower efficiency than D .

Hadronic FEI



Here D^* and D have same efficiency!

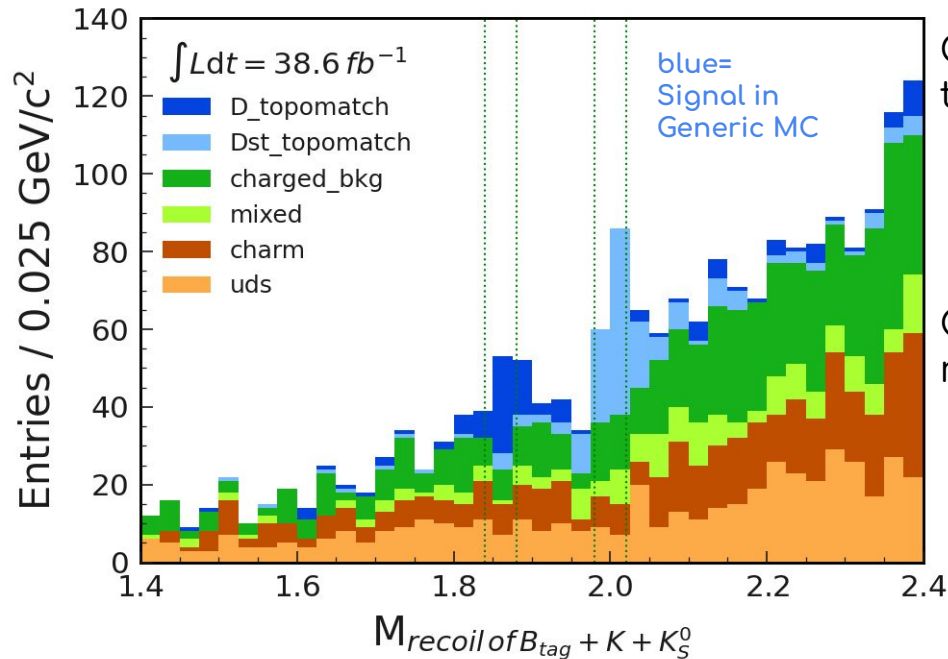
Both procedure look at different events:



Events with $B \rightarrow D^{(*)} X$ where the other $B \rightarrow \text{Had } B\text{-tag}$

Events with $B \rightarrow DX$ where $D \rightarrow K\pi$

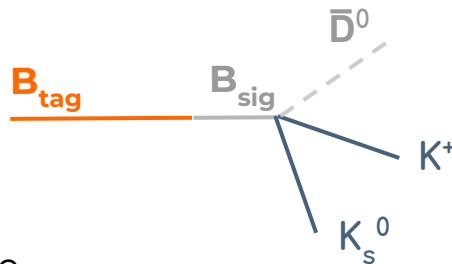
Example: DKK partial reconstruction



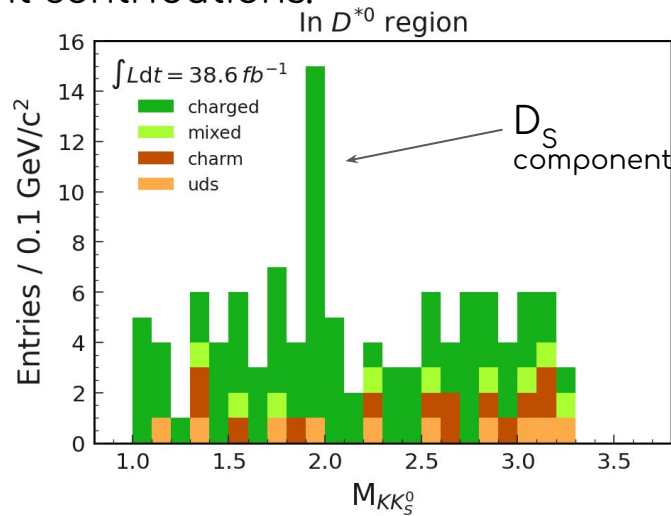
One can fit to get the BR

Hadronic FEI

[Valerio Bertacchi is working on exclusive approach]



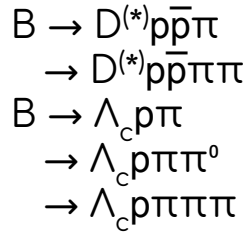
Can also look at the resonant contributions:



Efficiency (%)	hep-ex/0207041	
	Exclusive*	Partial
D	~0.3	~0.3
D*	~0.1	~0.3

Same efficiency!

Baryonic decays with recoil?

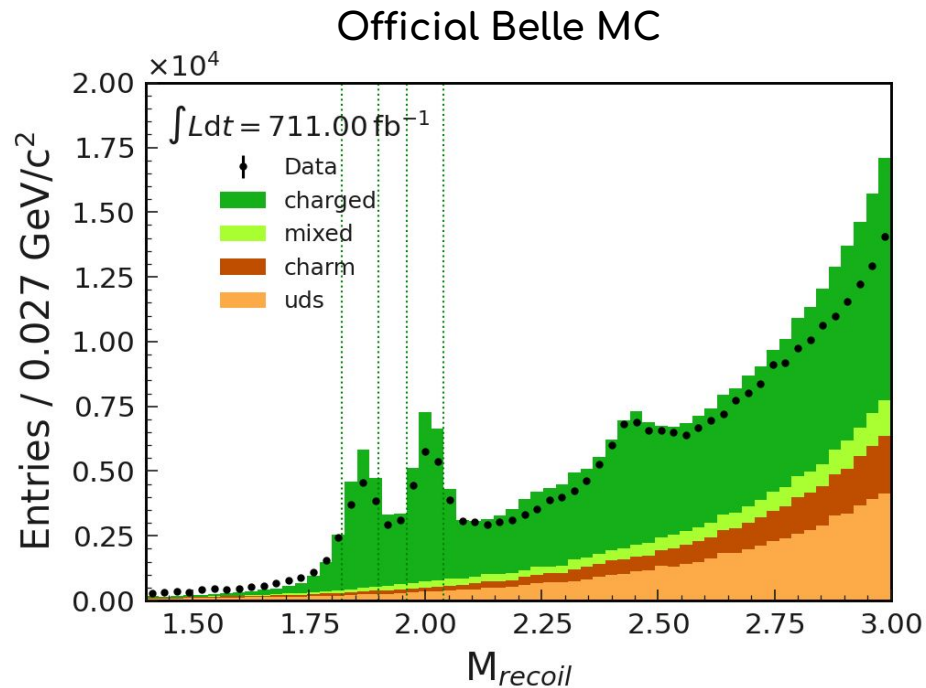
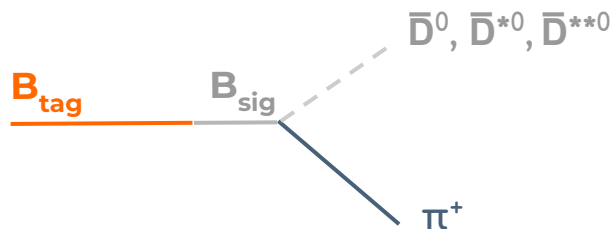


are the baryonic decays of B with the largest branching fractions
(some based on 20 year old CLEO measurements).

Clean enough to study using recoil method i.e., without reconstructing $D^{(*)}$ and Λ_c .

D** is more difficult

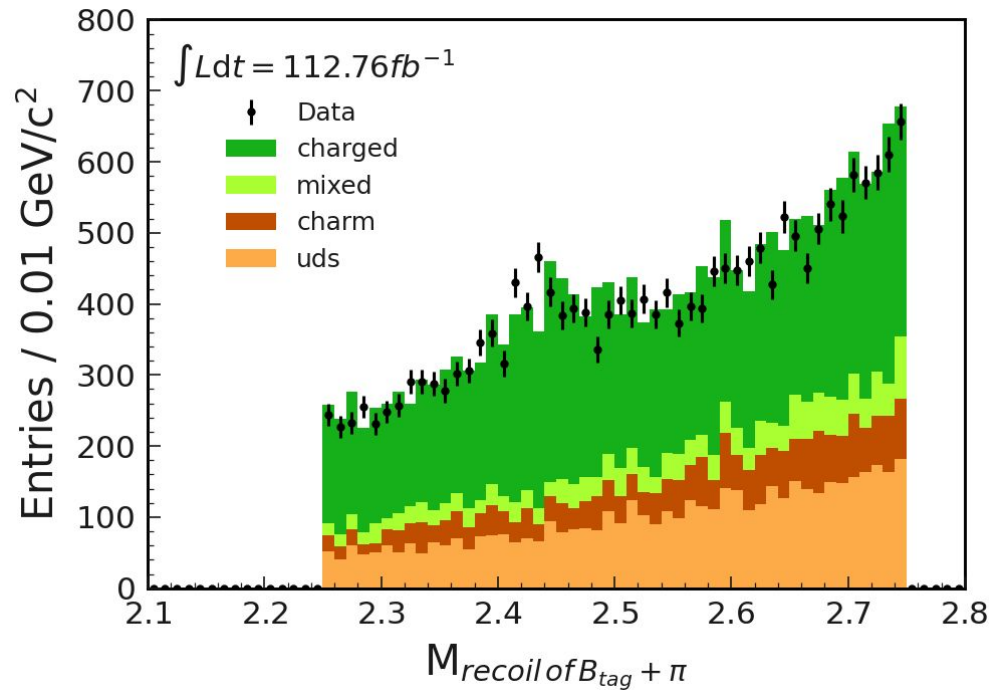
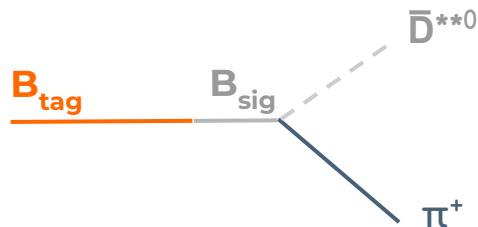
Hadronic FEI



The D and D* peaks are narrow and at the low-background region, but D** is more difficult to study here.

D** in recoil

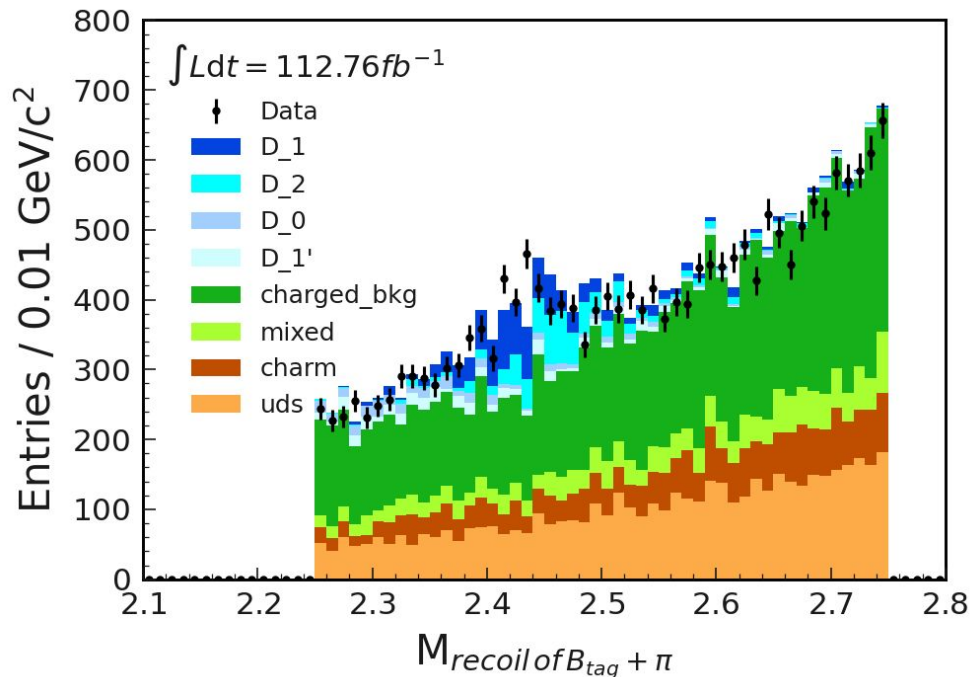
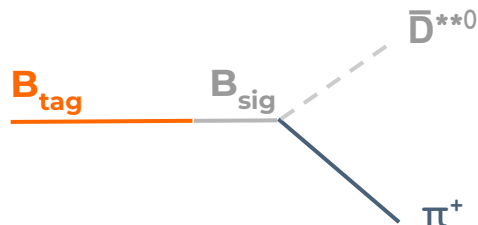
Hadronic FEI



We can first zoom into the D^{**} region.

D** in recoil

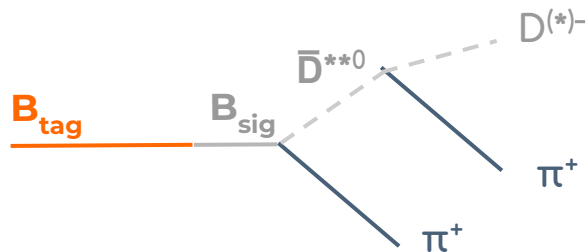
Hadronic FEI



We can first zoom into the D** region.
And focus on the “narrow” D**s: D₁ and D₂

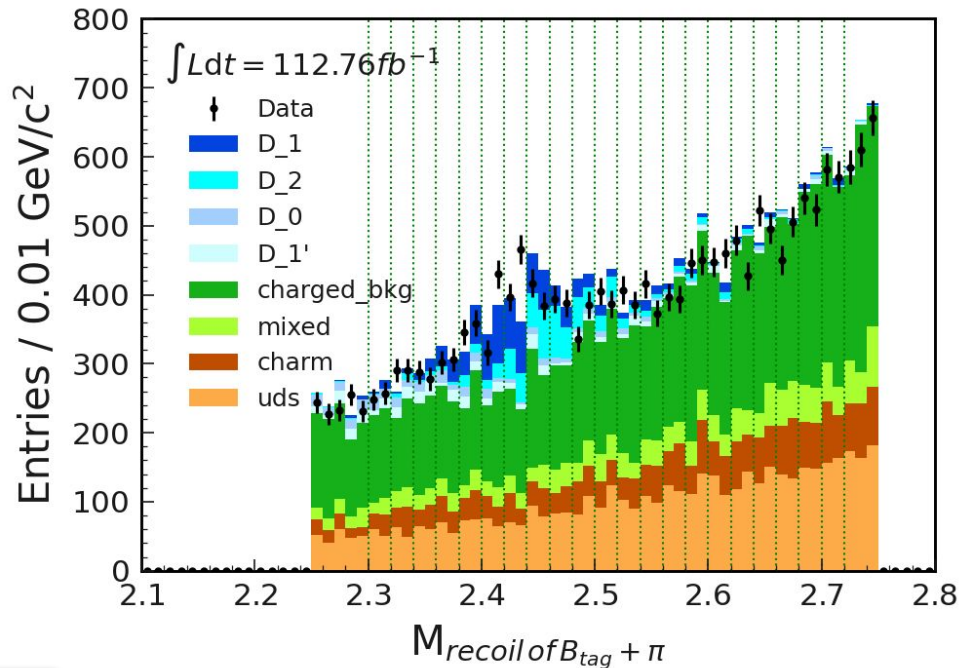
Double-recoil with D^{**} sample

Hadronic FEI



In these events, we can do a “double-recoil” by adding another π^+

D_1 can only decay to $D^{*-} \pi^+$, but
 D_2 can decay to both $D^- \pi^+$ and $D^{*-} \pi^+$



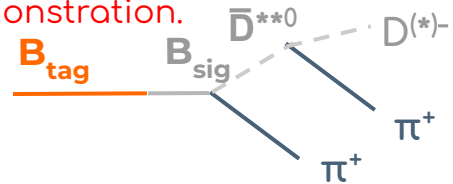
$D_1(2420)^0$	$D^{*0}\pi^0$	0.1997
	$D^{*+}\pi^-$	0.3994
	$D^0\pi^+\pi^-$	0.1719
	$D^0\pi^0\pi^0$	0.1145
	$D^+\pi^-\pi^0$	0.1145

$D_2^*(2460)^0$	$D^{*0}\pi^0$	0.1334
	$D^{*+}\pi^-$	0.2669
	$D^0\pi^0$	0.1999
	$D^+\pi^-$	0.3998

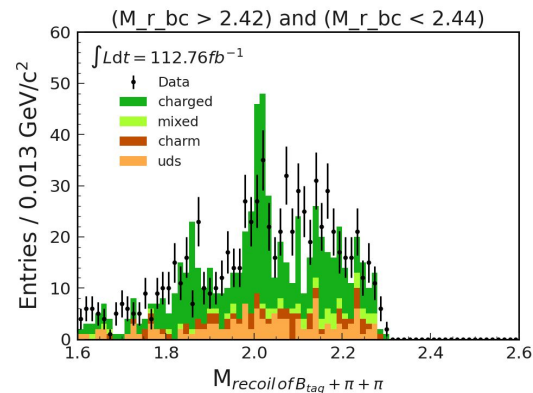
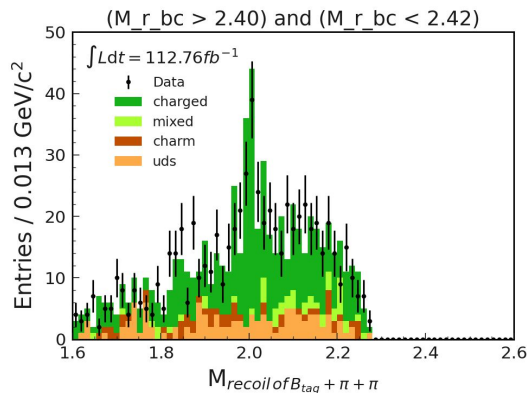
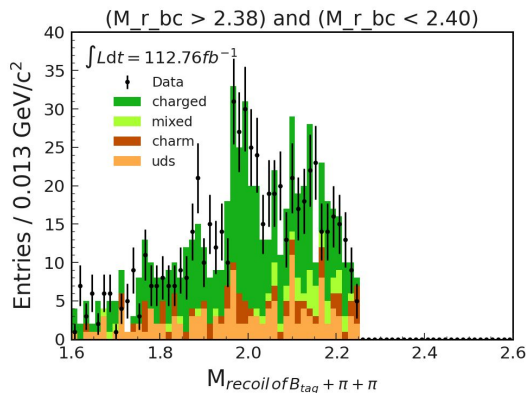
Only 1/10th of data; not optimized, just a demonstration.

Only 1/10th of data; not optimized, just a demonstration.

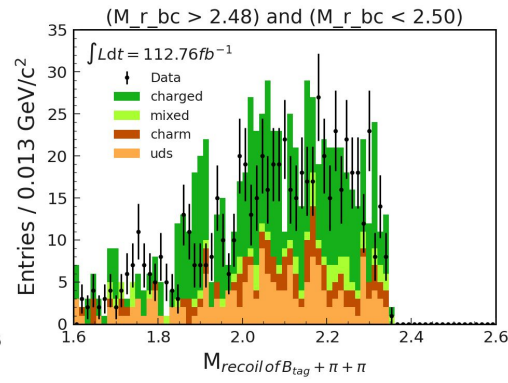
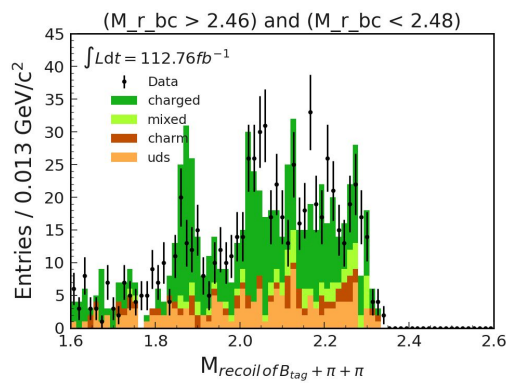
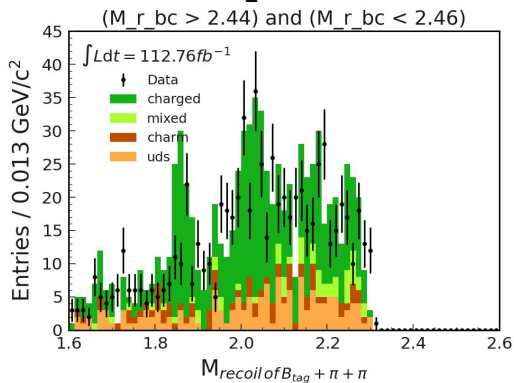
Double-recoil with D^{*-} sample



As expected, in the region of D_1 , we see mostly D^{*-} :

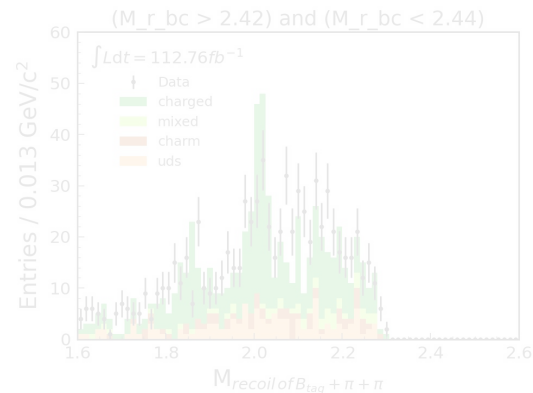
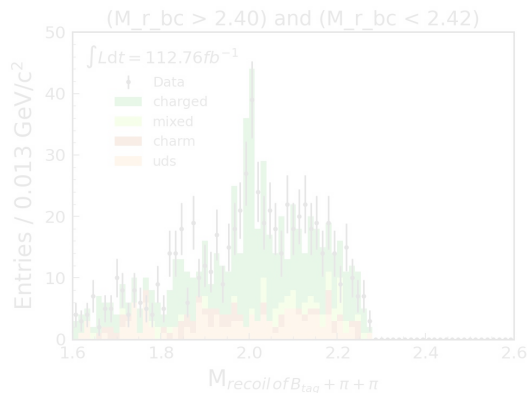
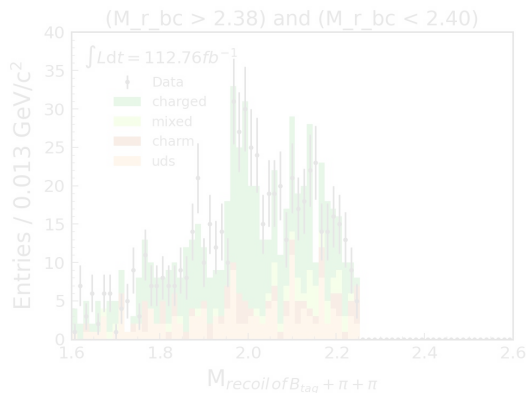
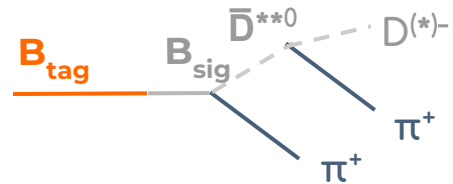


And in the region of D_2 , we see both D^- and D^{*-} :

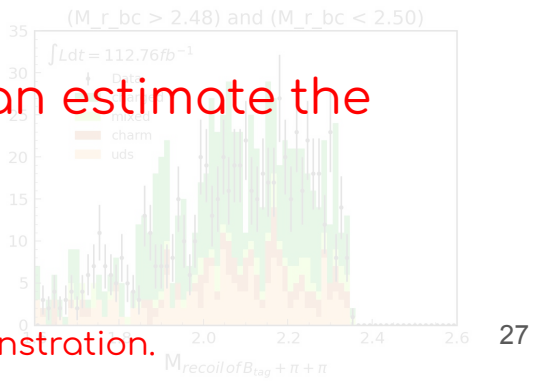
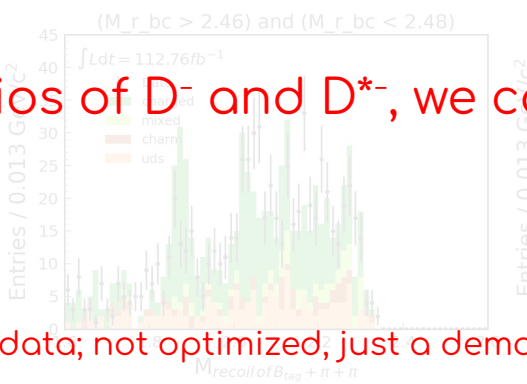
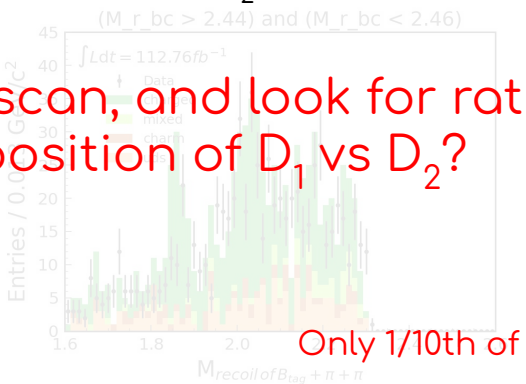


Double-recoil with D^{*-} sample

As expected, in the region of D_1 , we see mostly D^{*-} :



And in the region of D_2 , we see both D^- and D^{*-} :



If we scan, and look for ratios of D^- and D^{*-} , we can estimate the composition of D_1 vs D_2 ?

Only 1/10th of data; not optimized, just a demonstration.

Summary

We don't need to reconstruct the $D^{(*)}$ or Λ_c exclusively.

- There are many problems other than anomalies.
- Studying $B \rightarrow D^{(*)}\eta\pi$ and $B \rightarrow D^{(*)}\eta\rho$ along with possible intermediate resonances like D^{**} or $D(2S)$ will be a crucial input for understanding SL-gap and V_{cb} .
- Studying the decays of D^{**} and $D(2S)$ is also essential (charm physics)
- Demonstrated the performance of reconstruction $B \rightarrow D^{(*)}\pi$ with recoil-mass method.
- Many more exciting possibilities with recoil:
 - $B \rightarrow D^{(*)}\pi\pi^0$ (ρ), $B \rightarrow D^{(*)}\pi\pi\pi$ (a_1),
 - $B \rightarrow D^{(*)}\eta\pi$, $B \rightarrow D^{(*)}\eta\rho$, $B \rightarrow D^{(*)}\omega\pi$,
 - $B \rightarrow D^{(*)}KK_S$, $B \rightarrow D^{(*)}KK^*$
 - $B \rightarrow D^{(*)}p\bar{p}\pi$, $B \rightarrow D^{(*)}p\bar{p}\pi\pi$
 - $B \rightarrow \Lambda_c\rho\pi$, $B \rightarrow \Lambda_c\rho\pi\pi^0$, $B \rightarrow \Lambda_c\rho\pi\pi\pi$

Backup

Calibration factors per mode

with PDG uncertainties

3σ window around D^0 peak

Overall calibration factor:

$(82.6 \pm 0.9)\%$

$$\int L dt = 711 \text{ fb}^{-1}$$

— PDG uncertainty

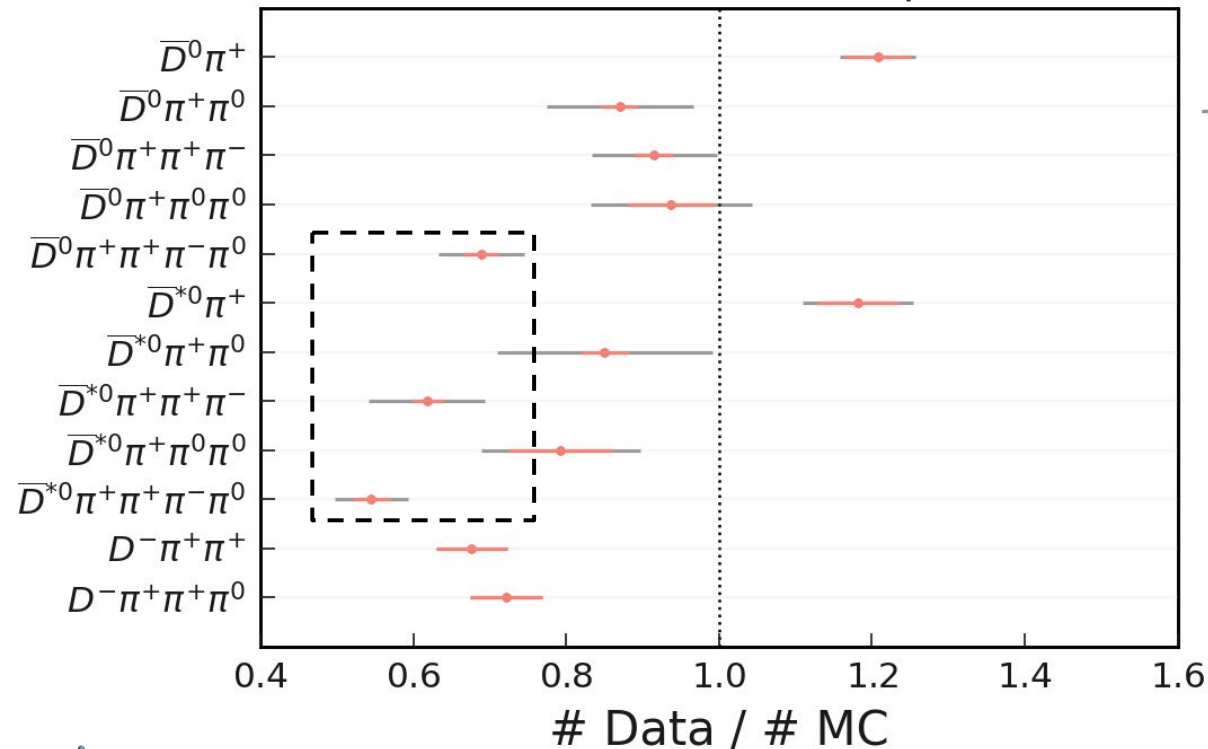
— Official MC

Modes with high multiplicity have large calibration factors! Even ~50%!

Even after considering PDG uncert, MC is clearly overestimating.

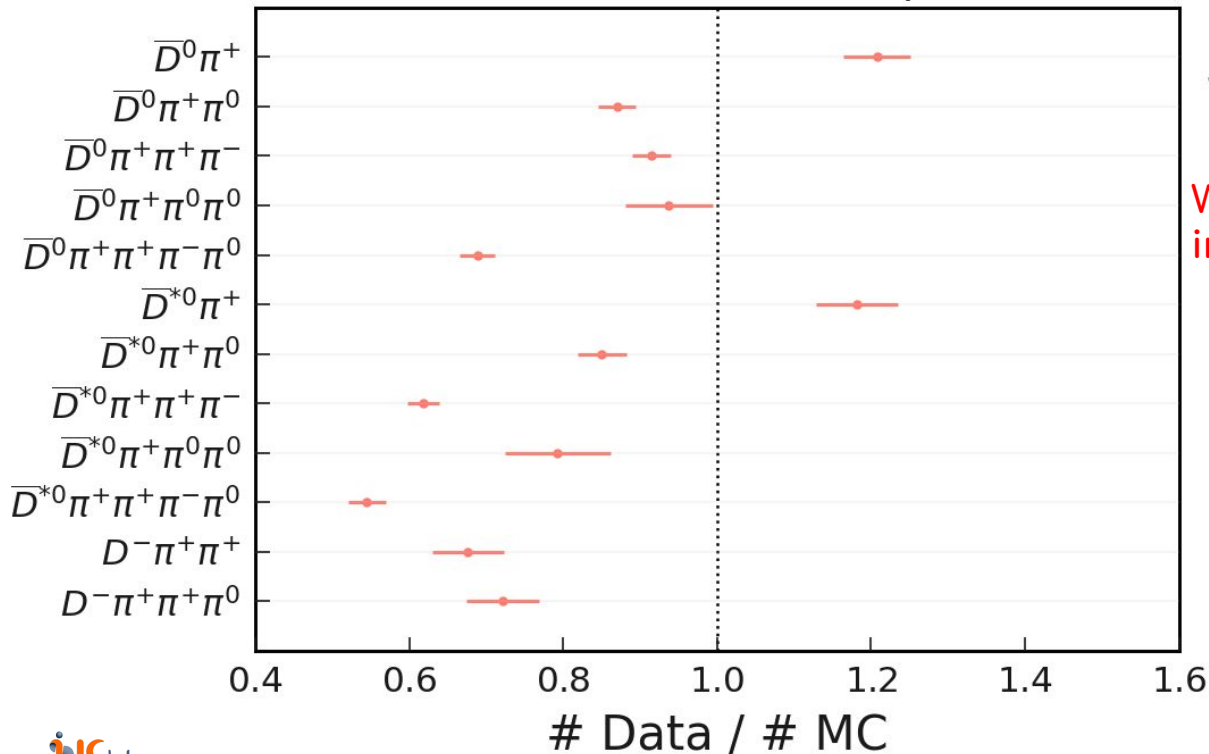
But the issue is not just in scaling, but also in the intermediate resonances to get to these final states.

⇒ We need a model for Hadronic B decays ! (a well educated and coherent update of DECAY table)



Systematics on calibration factors?

By counting the number of events per FEI mode:
 3σ window around D^0 peak



Overall calibration factor:
 $(82.6 \pm 0.9)\%$

$\int L dt = 711 \text{ fb}^{-1}$

Official MC

What systematics should be included?:

- π reco? (tracking, PID)
- BCS (highest ρ_π in CMS)
- Extraction method (Counting; includes background events)
→ Move to fitting?

Case study: $B^+ \rightarrow \bar{D}^0 \pi^+ \pi^+ \pi^-$

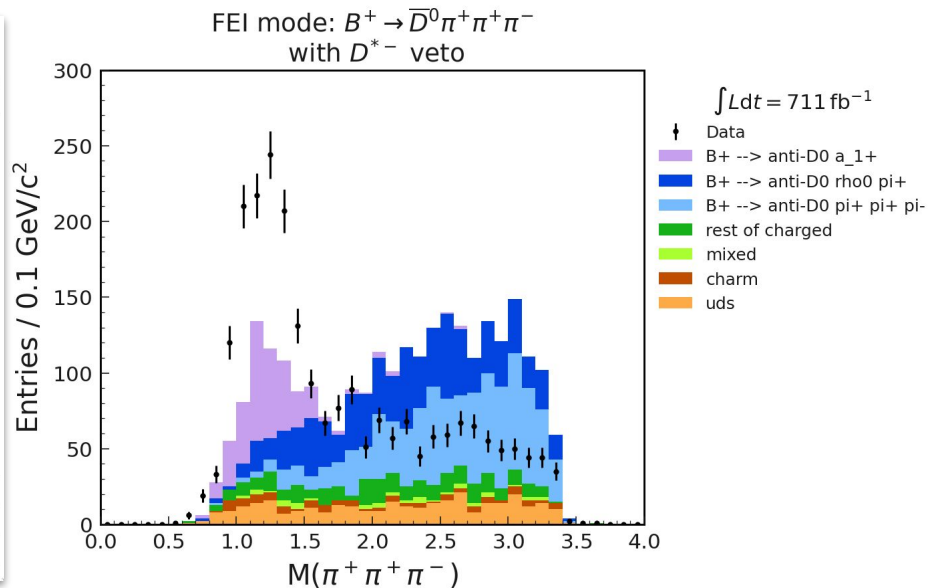
Improving calibration factors is not our primary target, instead improving the invariant masses (of intermediate particles), which are used as training variables in FEI will impact efficiency and purity

TABLE VI: Contents of the DECAY file concerning the $B^+ \rightarrow \bar{D}^0 \pi^+ \pi^+ \pi^-$ final state and corresponding measurements in PDG [in %].

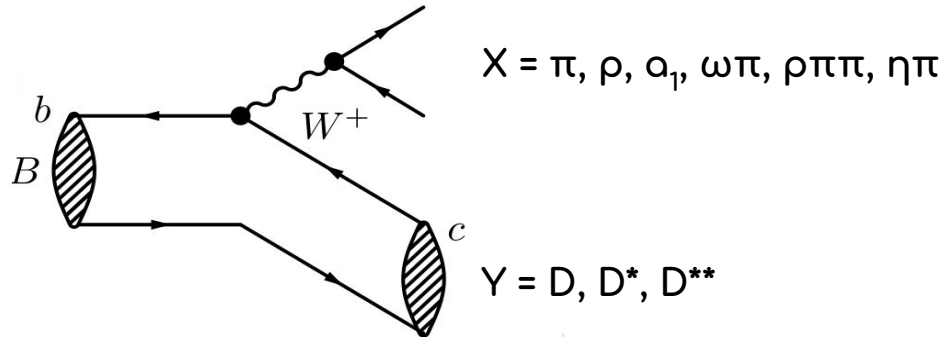
Decay	Belle	Belle II	Marker	Ref
$B^+ \rightarrow \bar{D}^0 \pi^- \pi^+ \pi^+$	0.46	0.51	■	[8]
$B^+ \rightarrow \bar{D}^0 \rho(770)^0 \pi^+; \rho(770)^0 \rightarrow \pi^+ \pi^-$	0.39	0.42	★	[8]
$B^+ \rightarrow \bar{D}^0 a_1(1260)^+; a_1(1260)^+ \rightarrow \rho(770)^0 \pi^+; \rho(770)^0 \rightarrow \pi^+ \pi^-$	0.13	0.14	★	[8]
$B^+ \rightarrow \bar{D}^0 a_1(1260)^+; a_1(1260)^+ \rightarrow f_0(500) \pi^+; f_0(500) \rightarrow \pi^+ \pi^-$	0.05	0.05	★	[8]
$B^+ \rightarrow \bar{D}_1(2420)^0 \pi^+; \bar{D}_1(2420)^0 \rightarrow D^*(2010)^- \pi^+; D^*(2010)^- \rightarrow \bar{D}^0 \pi^-$	0.04	0.02		[10], [9]
$B^+ \rightarrow \bar{D}_1(2430)^0 \pi^+; \bar{D}_1(2430)^0 \rightarrow D^*(2010)^- \pi^+; D^*(2010)^- \rightarrow \bar{D}^0 \pi^-$	0.03	0.02		[10], [9]
$B^+ \rightarrow \bar{D}_2^*(2460)^0 \pi^+; \bar{D}_2^*(2460)^0 \rightarrow D^*(2010)^- \pi^+; D^*(2010)^- \rightarrow \bar{D}^0 \pi^-$	0.01	0.01		[10], [9]
$B^+ \rightarrow D^*(2010)^- \pi^+ \pi^+; D^*(2010)^- \rightarrow \bar{D}^0 \pi^-$	-	0.09	■	[10]
$B^+ \rightarrow \bar{D}^0 a_1(1260)^+; a_1(1260)^+ \rightarrow \pi^+ \pi^+ \pi^-$	-	0.07	★	[8]
$B^+ \rightarrow \bar{D}_1(2420)^0 \pi^+; \bar{D}_1(2420)^0 \rightarrow \bar{D}^0 \pi^- \pi^+$	-	0.02		[10], [9]
$B^+ \rightarrow \bar{D}^0 K^*(892)^+; K^*(892)^+ \rightarrow K^0 \pi^+; K^0 \rightarrow K_S^0; K_S^0 \rightarrow \pi^+ \pi^-$	-	0.01		
Rest of Exclusive	0.03	0.03		
Sum of Exclusive	1.12	1.38		
Sum of PYTHIA	0	0		
Total Sum	1.12	1.38		

[\[BELLE2-NOTE-PH-2022-002\]](#)

By restudying the CLEO and LHCb measurements for this mode, we realized that the NR and ρ components should be almost 0 and should be dominated by a_1^+



Model for $B \rightarrow D^{(*,**)} \eta\pi \rho\pi^0$ decays



Happens through 2 channels, one with spectator quarks (call Y) and one from the W (call X).

We want to modify the DECAY table to latest PDG/paper interpretations and this model to see the impact.

2 primary rules:

- $D^0 X : D^{*0} X : D^{**0} X \sim 1 : 1 : 1$
(based on observation from $D \pi^- : D^* \pi^- : D^{**} \pi^-$ and $D \rho^- : D^* \rho^-$)
- $Y \pi^- : Y \rho^- : Y a_1^- \sim 1 : 2.5 : 2.5$
(based on predictions and confirmed with $\tau \rightarrow h \nu$ decays)

Essentially validation, we do not want to fine-tune (except set 0 there is no signal*).

Additional information:

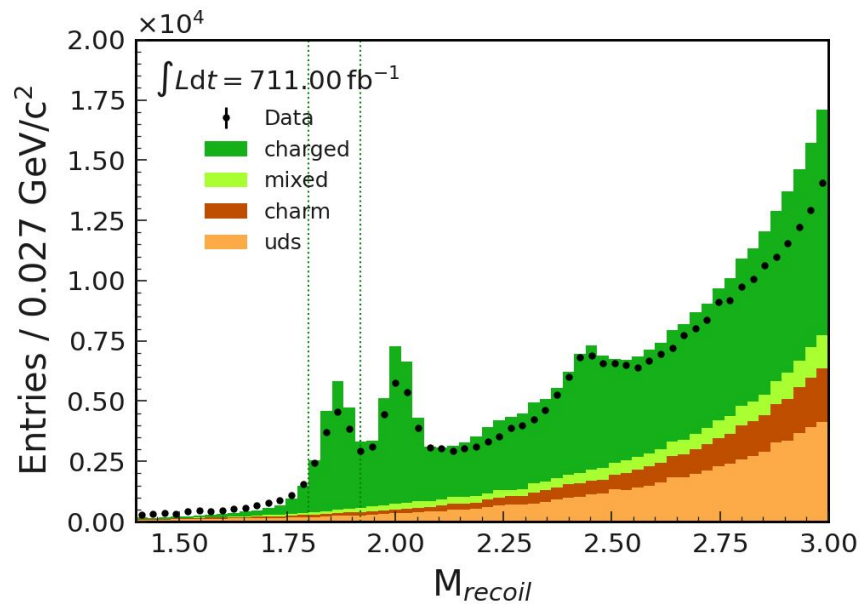
- $3\pi \pi^0$ is hard to model without some sort of ρ' resonance
 - For $\omega\pi$ we fix from measurements.
 - For $\rho\pi\pi$ and $\eta\pi$, we let PYTHIA generate it.
- Decays of D^{**} particles is synchronized with Belle II
- The fraction of 4 different D^{**} is fixed based on observations.

*See backup

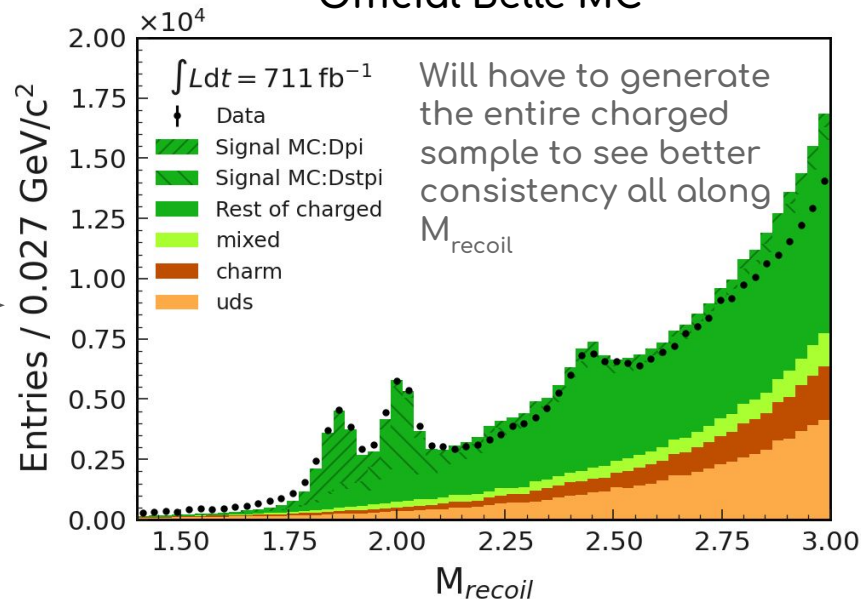
Validation by embedding signal MC

To quickly study the impact of the modified DECA.Y.DEC file, generated Signal MC of $B \rightarrow D^{(*)}\pi$ (other B decays updated) and replaced corresponding events in the generic Charged MC:

Official Belle MC



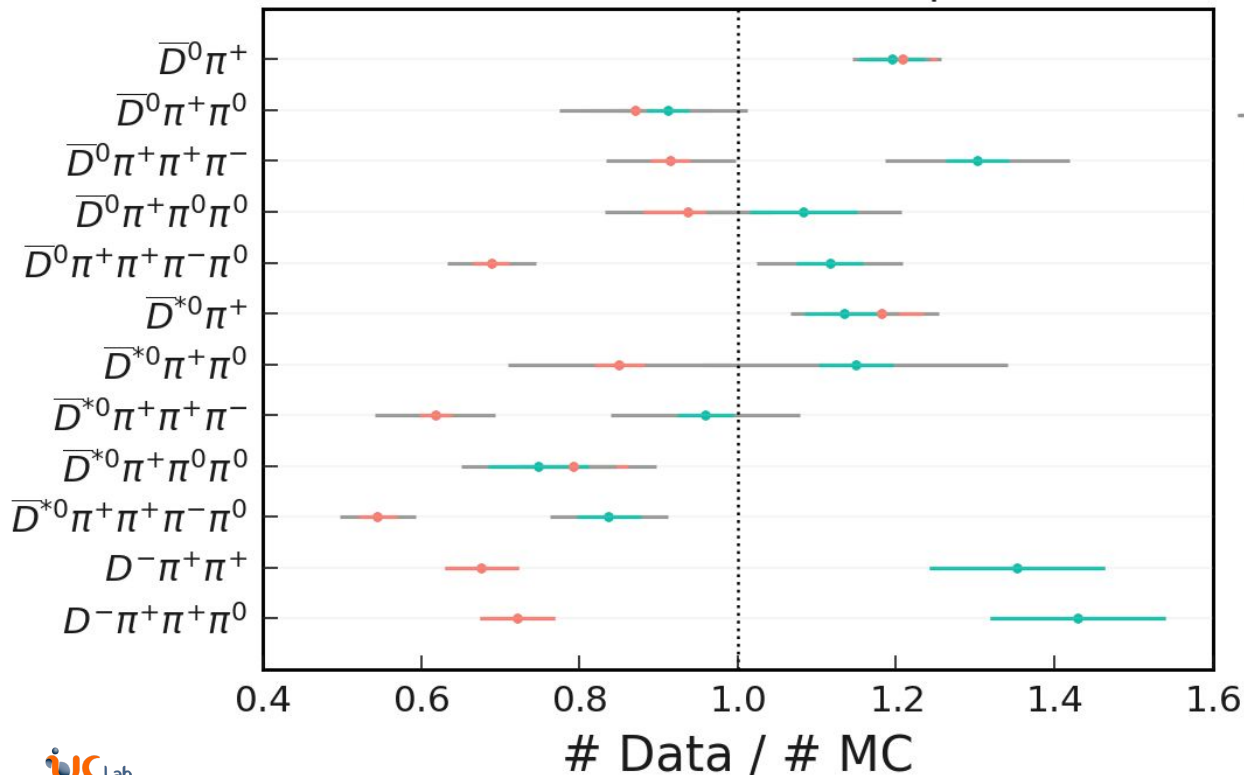
Modified Signal MC embedded in Official Belle MC



Updated calibration factors

per mode

3σ window around D^0 peak



$\int L dt = 711 \text{ fb}^{-1}$

- PDG uncertainty
- Official MC
- Proposed MC

Overall calibration factor:

$(82.6 \pm 0.9)\%$

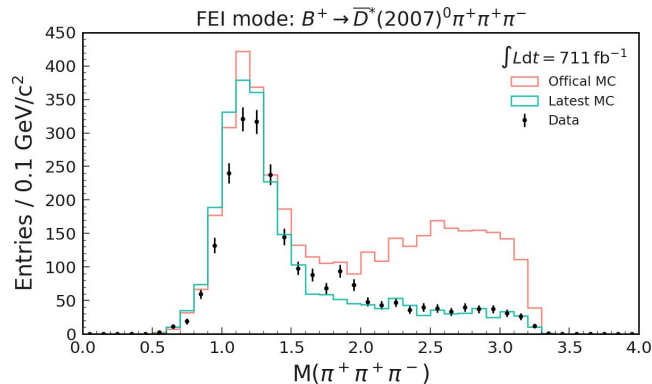
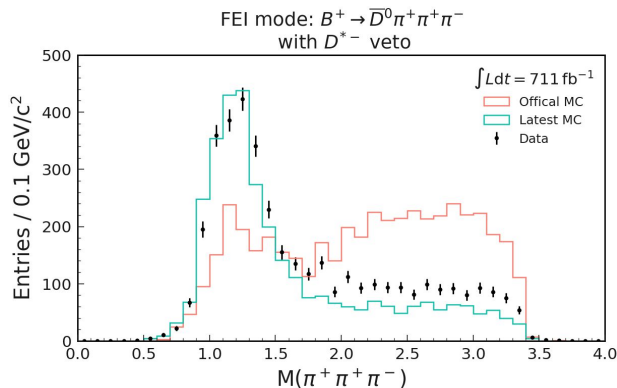


$(104.2 \pm 1.2)\%$

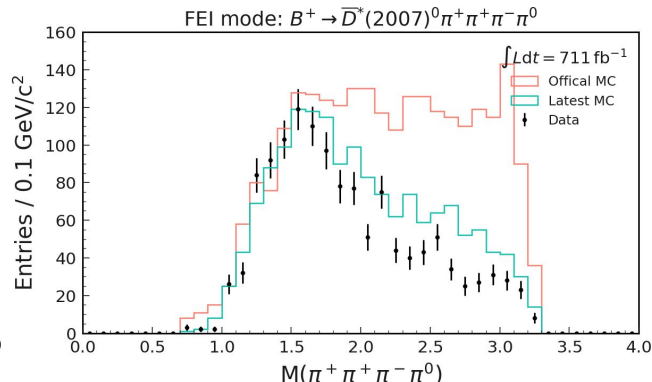
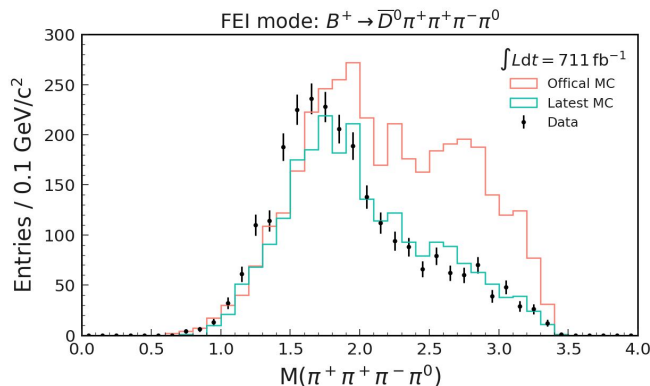
Decay description is improved!

The improvement is not limited to calibration factors, but more importantly in the invariant masses (of intermediate particles), which are used as training variables in FEI

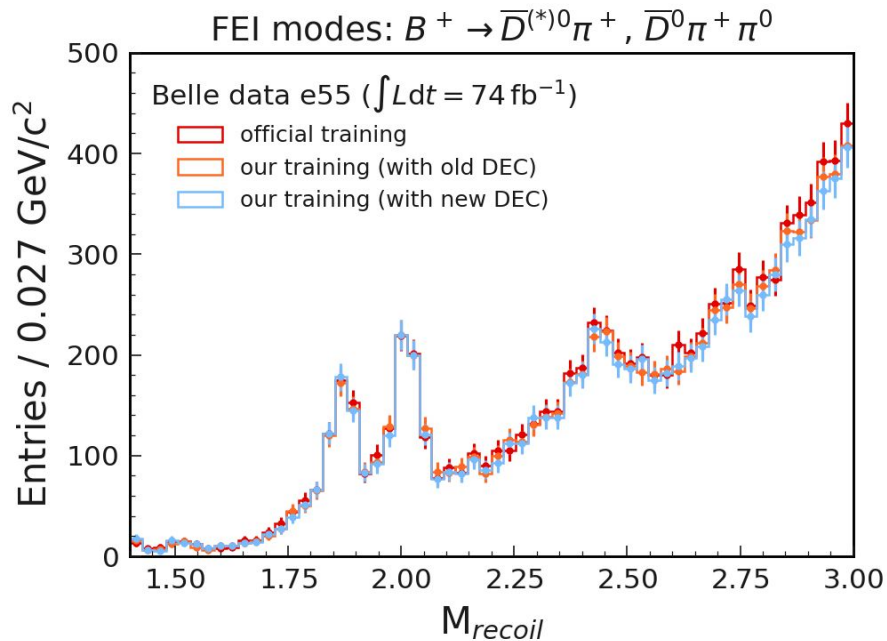
$3\pi^\pm$ case:



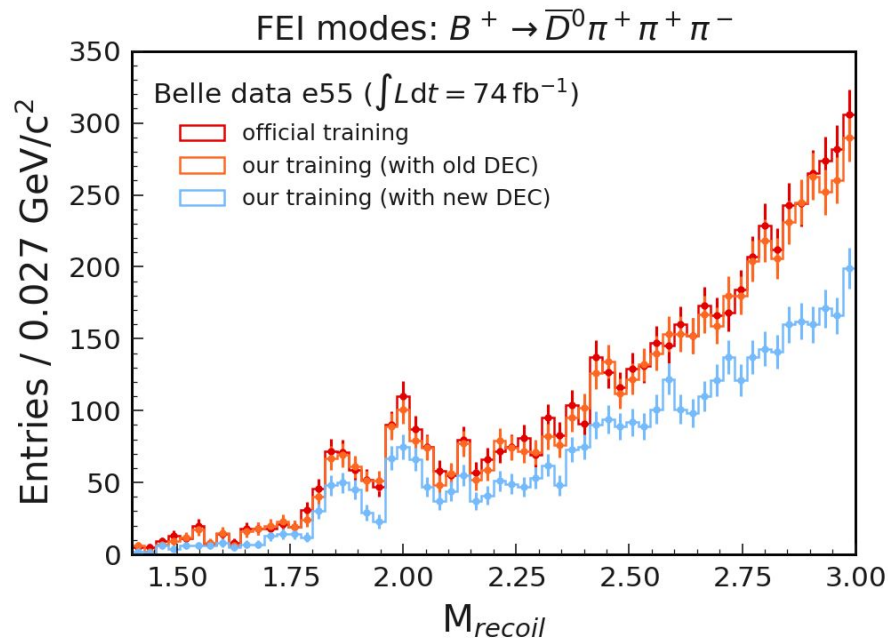
$3\pi^\pm \pi^0$ case:



Retraining FEI: Validation



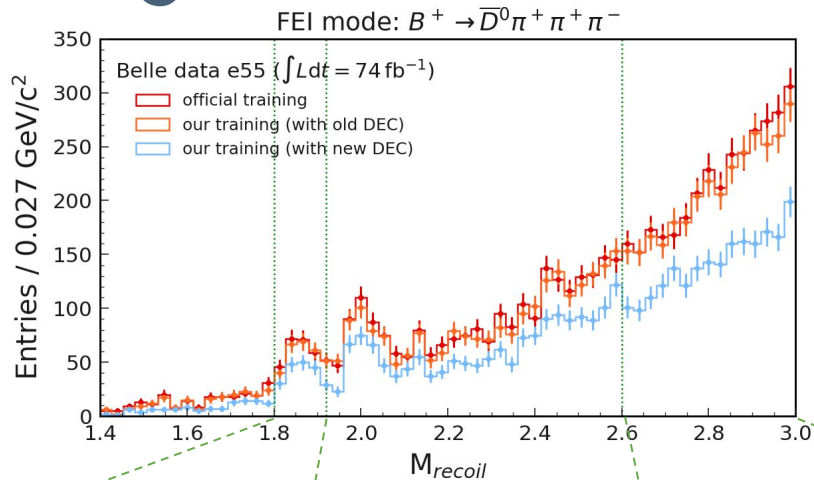
Nothing changes in the FEI modes where we did not change anything.



There is a significant background reduction in FEI modes where MC model is improved.

Our training has some issues while reconstructing modes with π^0 , under investigation... (see backup) 14

Retraining FEI: Effective cuts

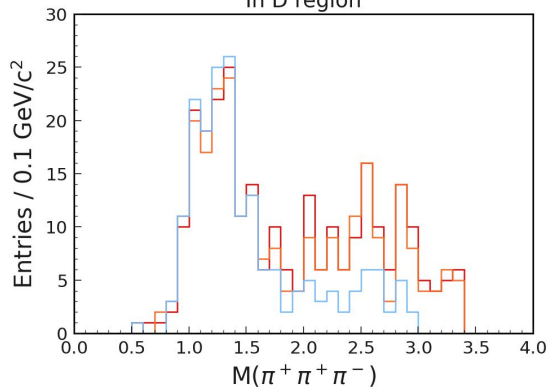


The new training is effectively applying a a_1^+ cut!

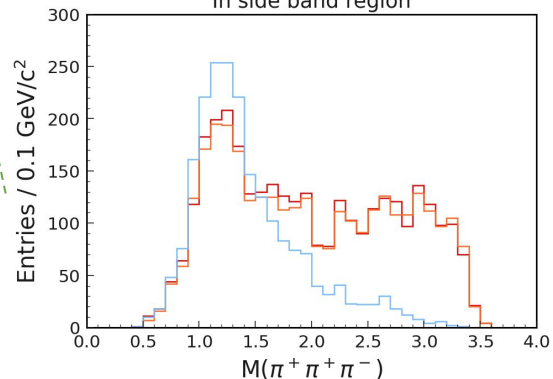
Can we apply this cut manually?
Is that enough?

Can we have a fully cut-based B-tagging? i.e., no training?

In D region

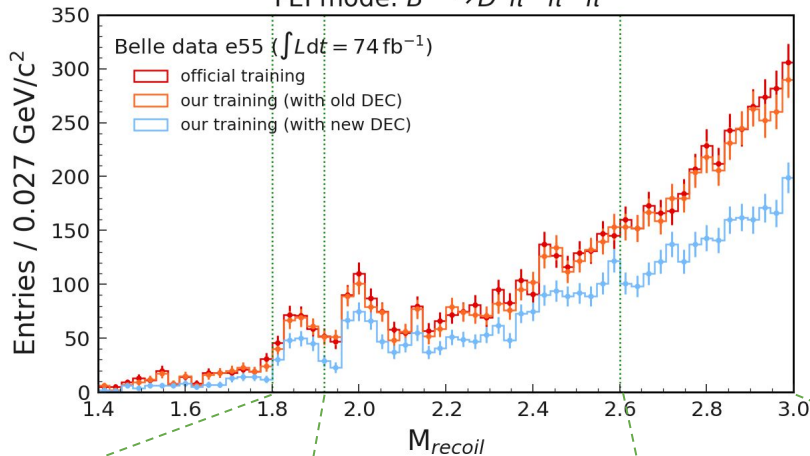


In side band region

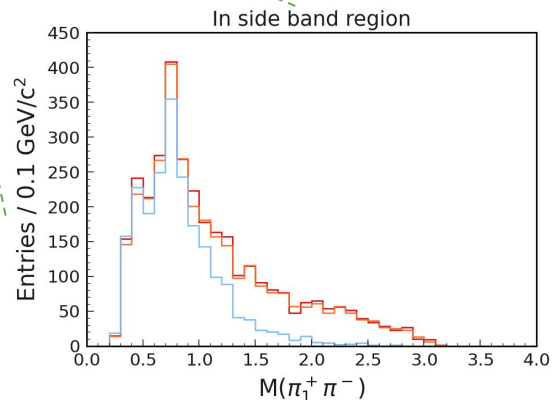
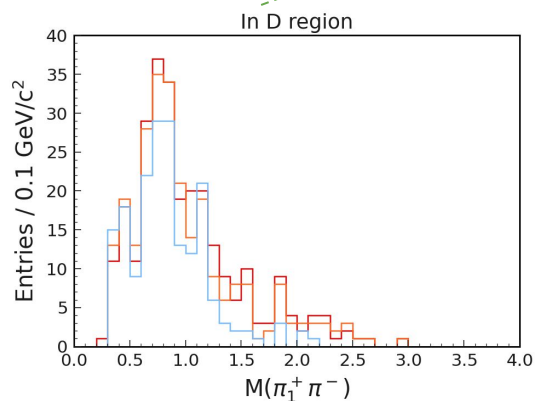


Retraining FEI: Effective cuts

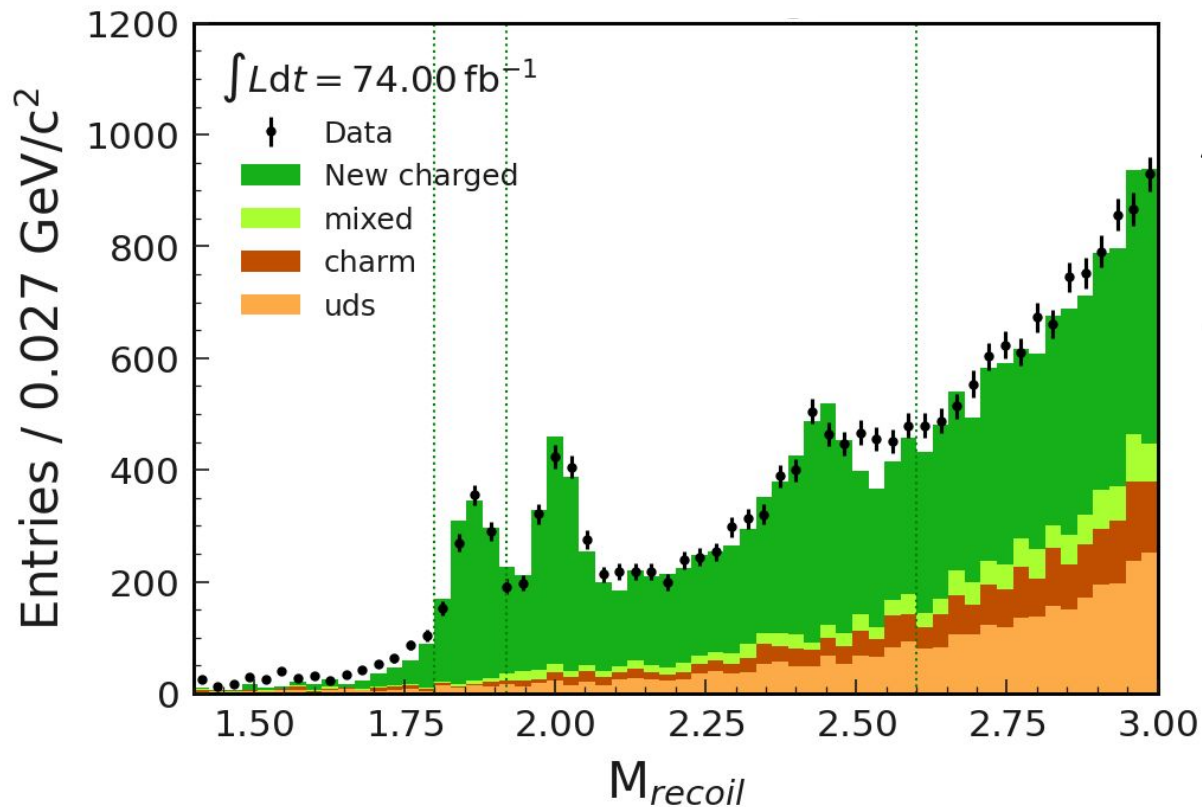
FEI mode: $B^+ \rightarrow \bar{D}^0 \pi^+ \pi^+ \pi^-$



$M(3\pi)$ is the dimension we usually look at, but the changed kinematics is visible in other dimensions like $M(2\pi)$ also.



Retraining FEI: Data-MC agreement



After reconstructing all MC and data with the training based on new DEC, the Data - MC agreement improves too! (even at higher M_{recoil} !)

$B^+ \rightarrow D\pi$ selection procedure

These cuts could be further optimized, but seem good enough for preliminary studies.

We start by reconstructing a FEI-Hadronic B with cuts:

- $M_{bc} > 5.27 \text{ GeV}/c^2$
- $|\Delta E| < 0.05 \text{ GeV}$
- FEI Signal Probability > 0.01

Select a π with:

- $|d0| < 1$ and $|z0| < 3$
- $L_{K/\pi} < 0.9$ and $\mu\text{-id} < 0.9$ and $e\text{-id} < 0.9$

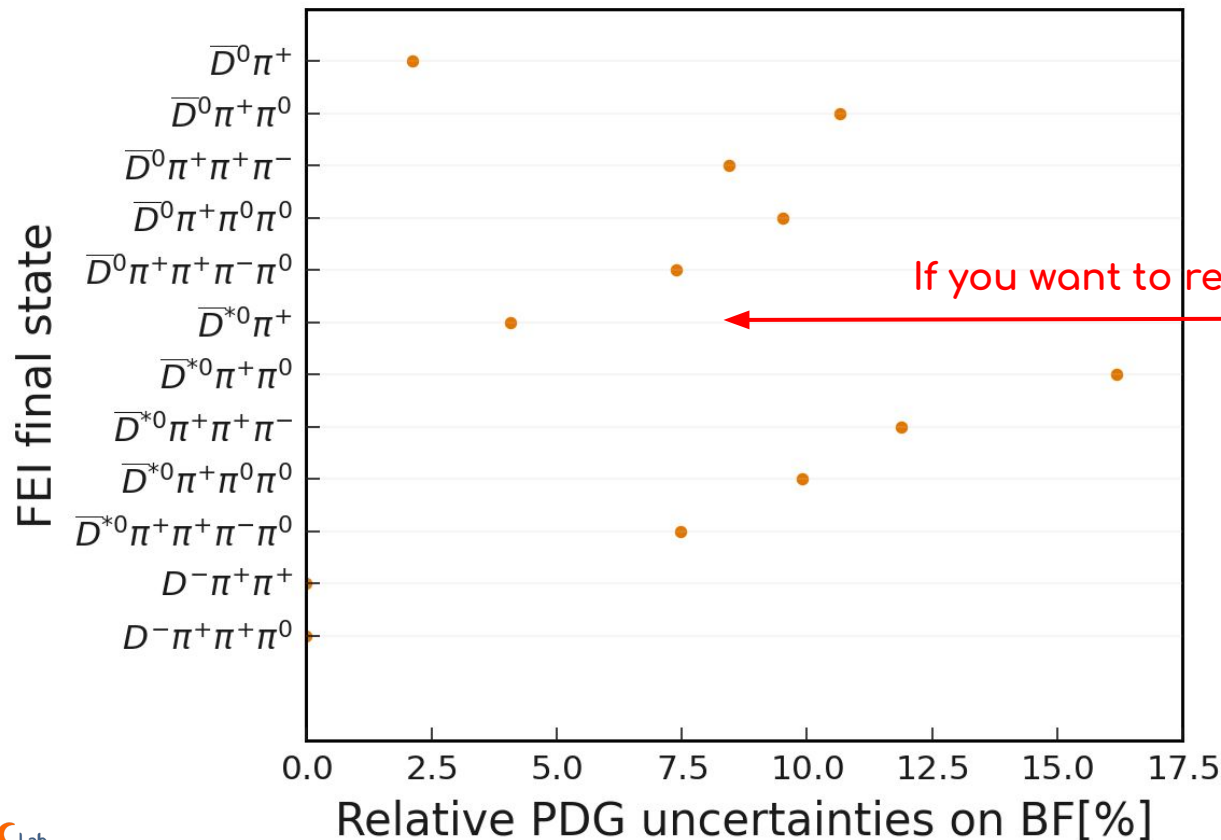
Simple continuum suppression:

- Event sphericity > 0.2
- B_{tag} 's $\cos\text{TBTO} < 0.9$

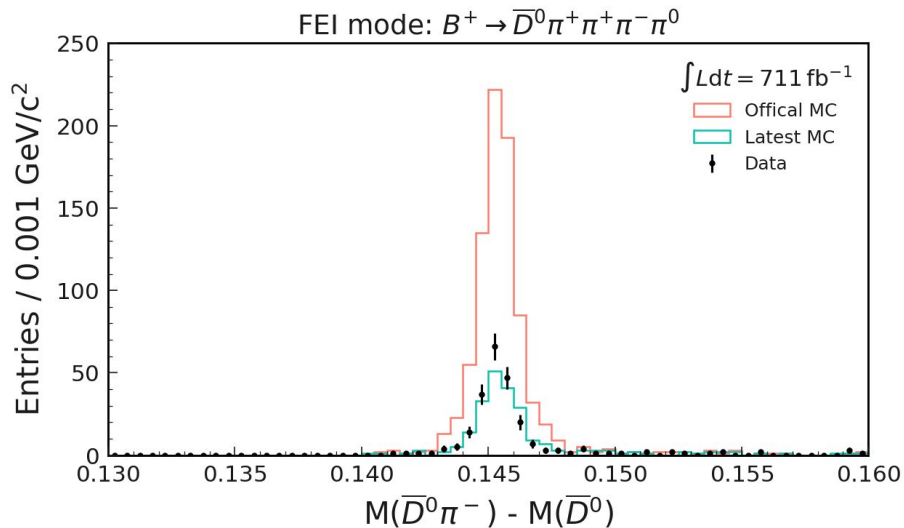
After all this, if there are multiple candidates, we select the one with highest FEI signal probability and highest π momentum in CMS

The code is present [\[here\]](#)

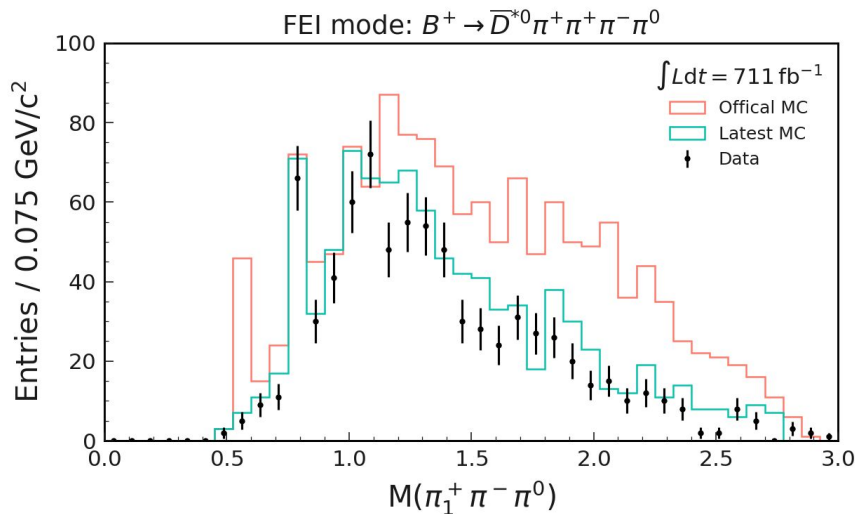
Relative PDG uncertainties



Changes in DEC not based on measurements: 1/2

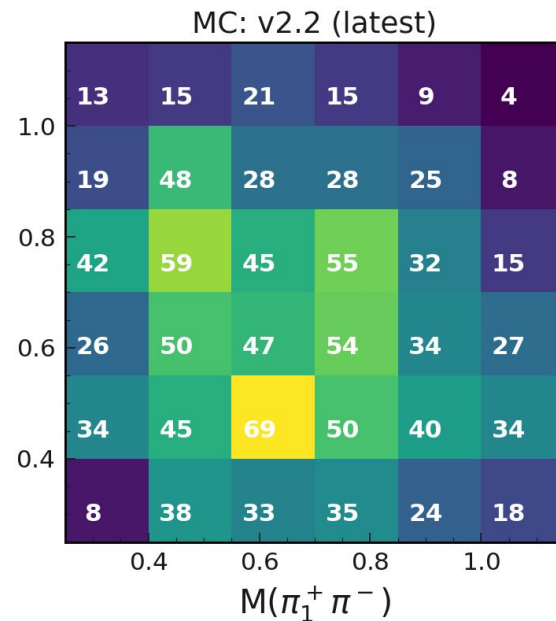
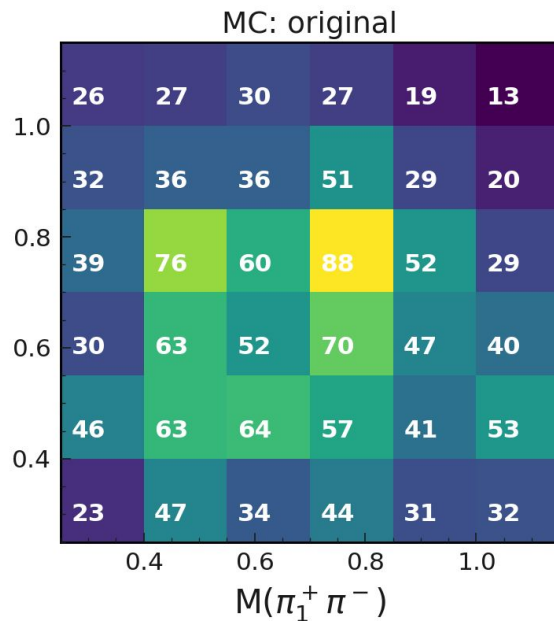
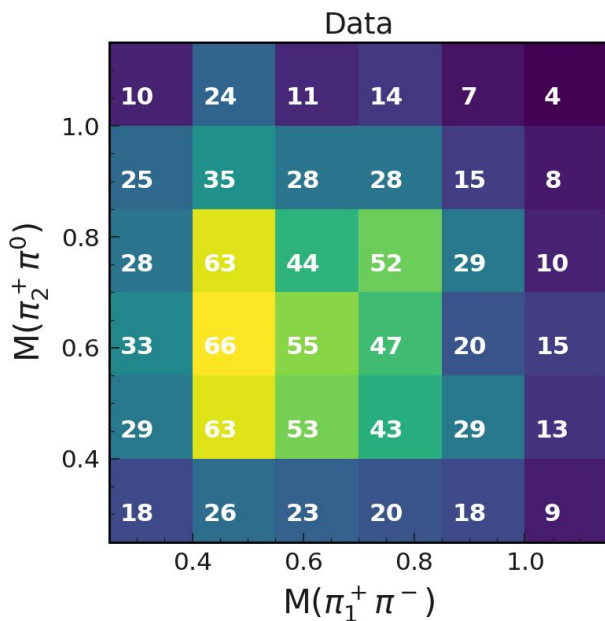


ARGUS measured it to be $(1.5 \pm 0.7)\%$
But we see that the contribution coming
from D^{**} is enough



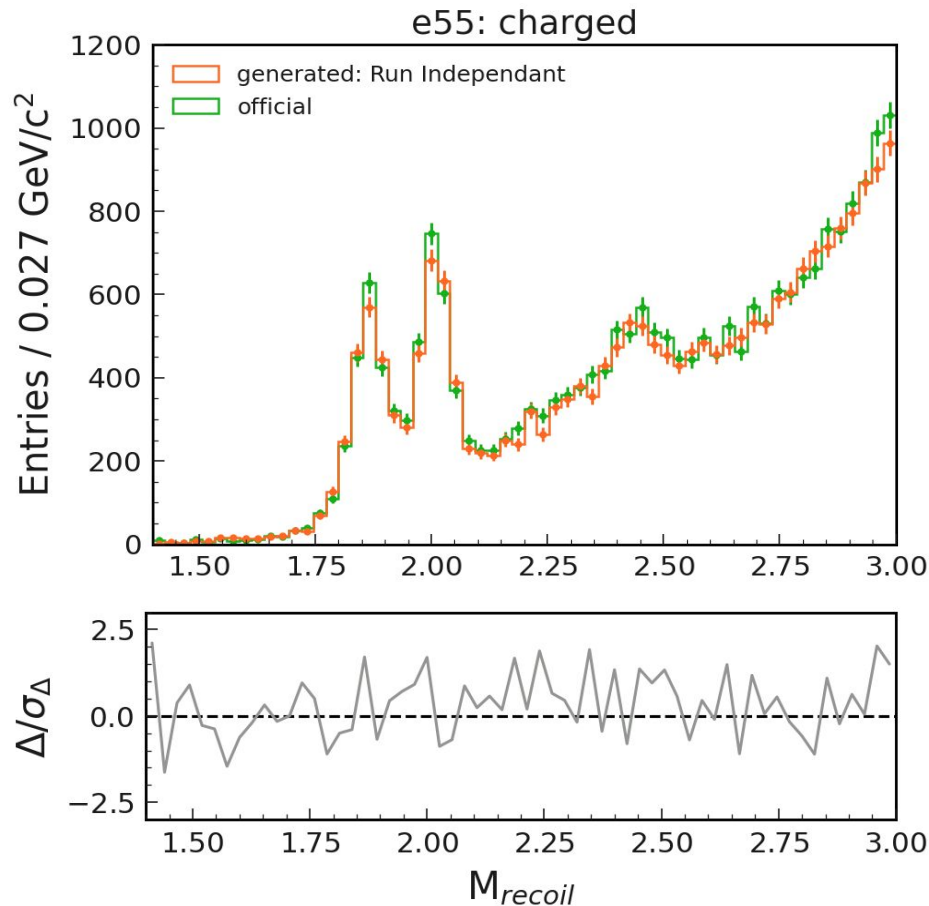
No measurement, but overestimated by
PYTHIA.

Changes in DEC not based on measurements: 2/2



Regenerating run-independent* samples

*still exp-dependent BG



Run-Independent sample of 10% seems good enough for comparison?

Regenerating run-independent* samples

*still exp-dependent BG

With new DEC file:

