

$$B \rightarrow X_c \ell \bar{v}_\ell$$

Cocktail please!

– or known & unknowns
for $X_c \notin \{D, D^*\}$

Thanks to Markus, Mirco, Dean, Michele and Zoltan for useful feedback!

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$$X_c \notin \{D, D^*\}$$



(Quark-model)
Resonances

$D^{**}(1P)$
 $D^{**}(1D)$
...

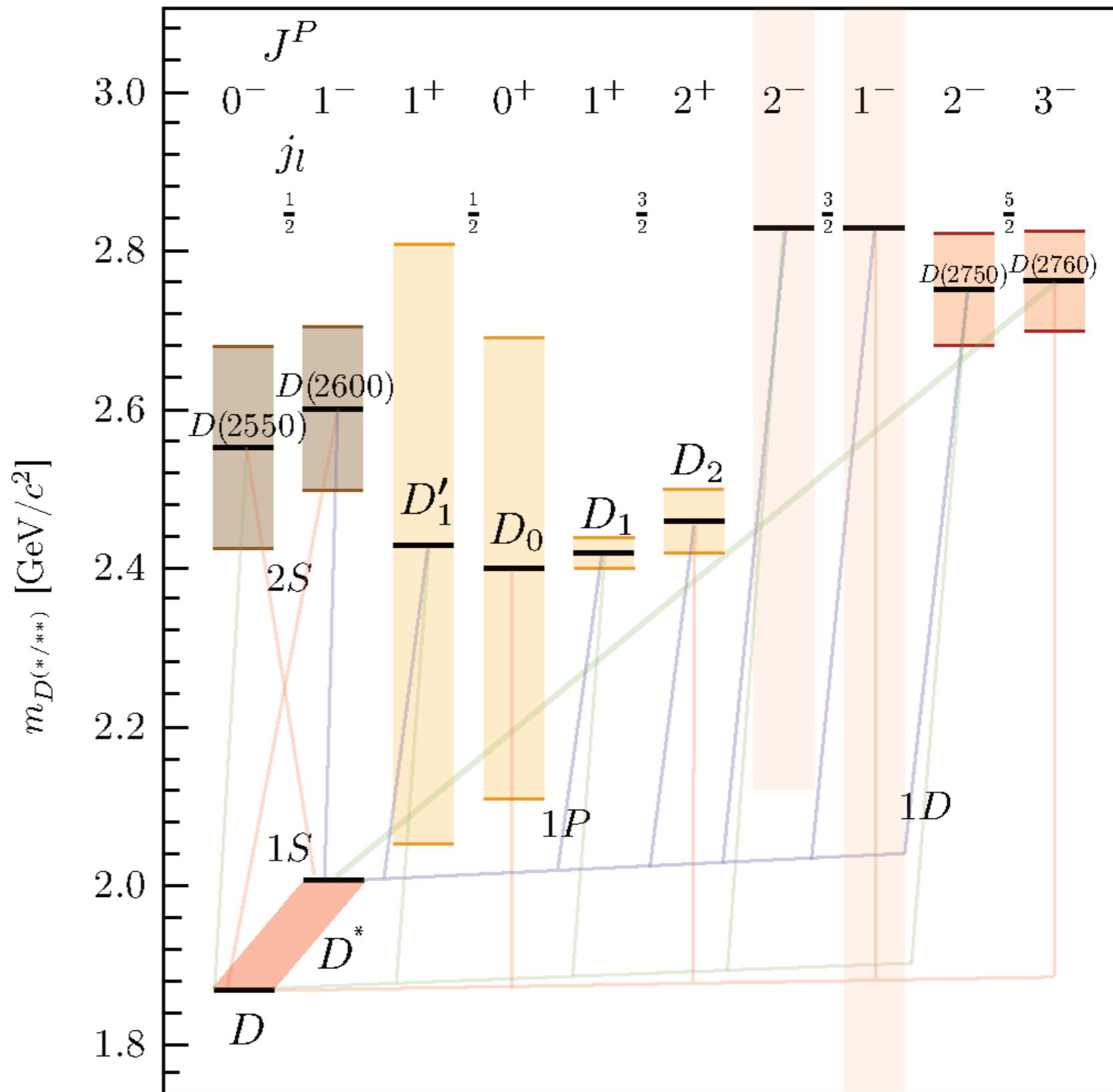
Non-resonant
contributions

$$D^{(*)}\pi, D^{(*)}\eta$$

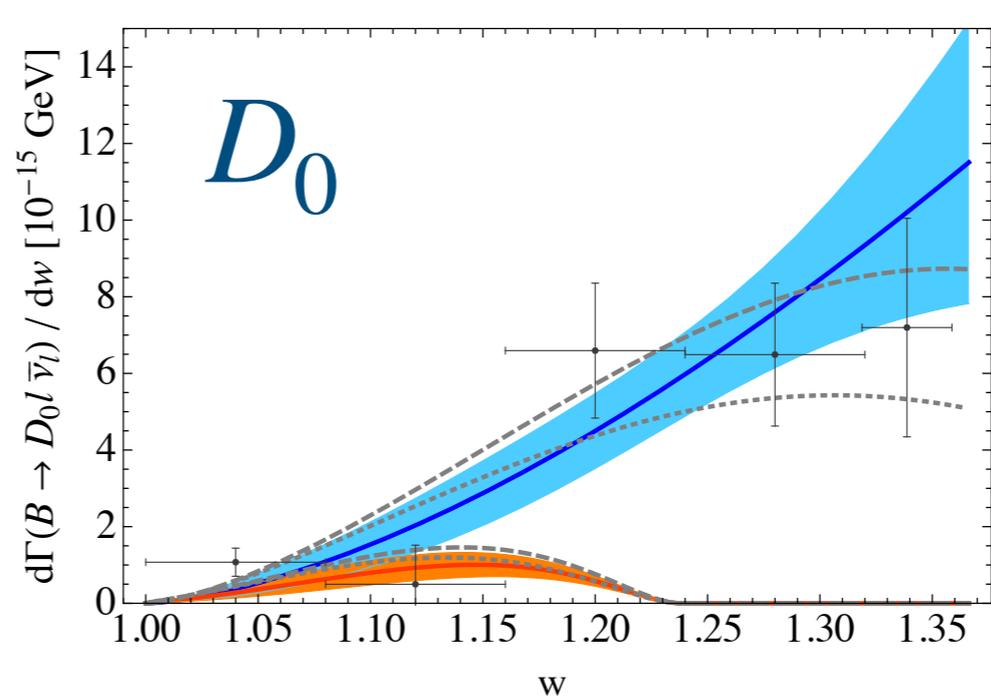
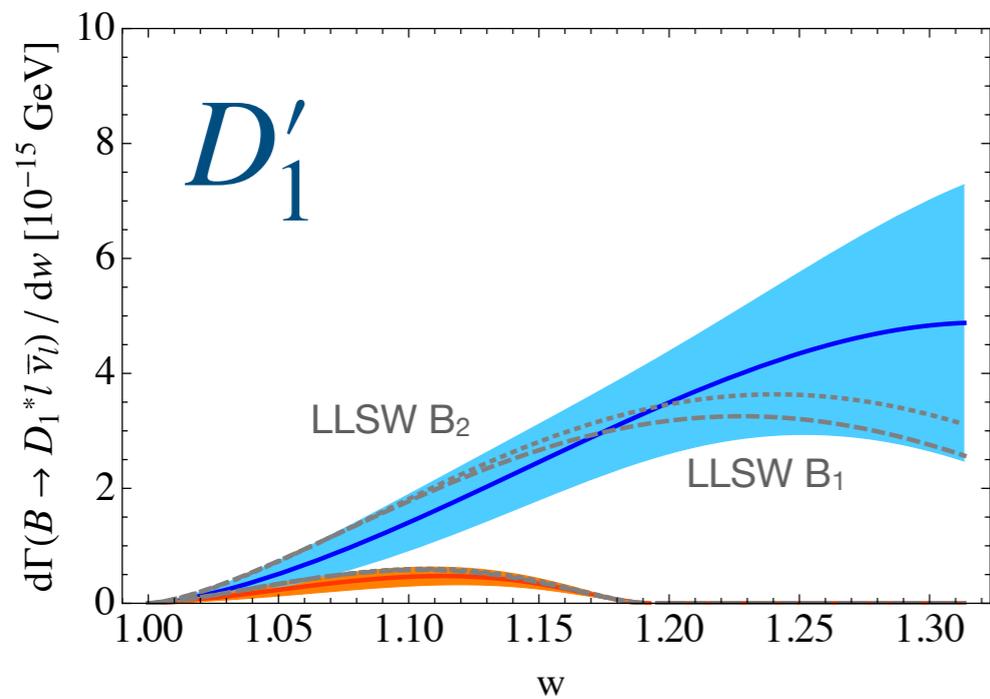
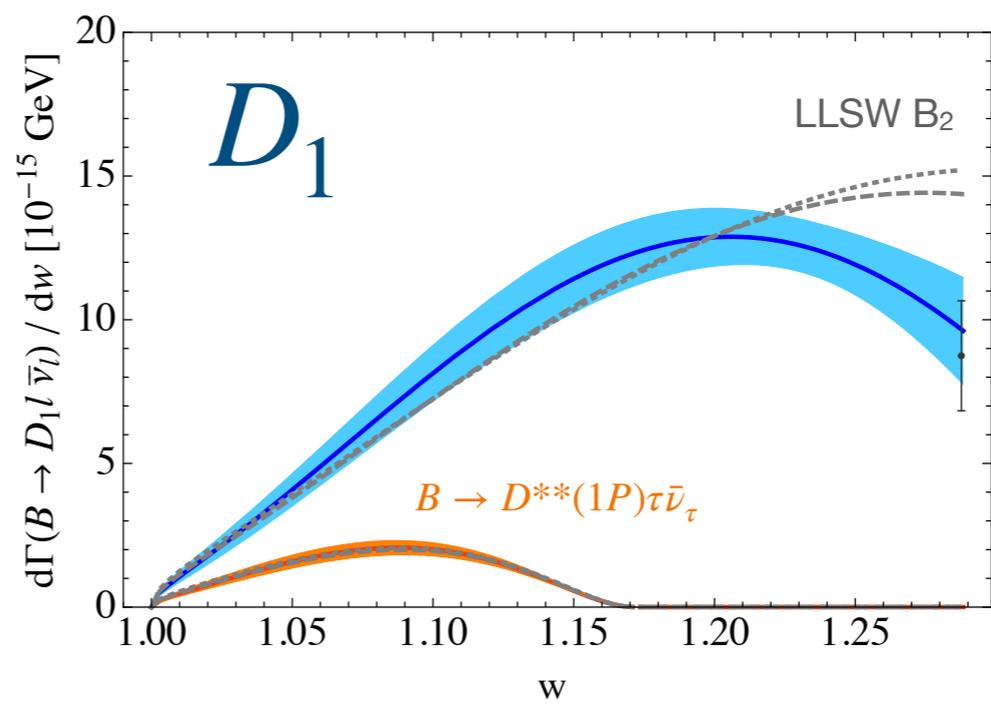
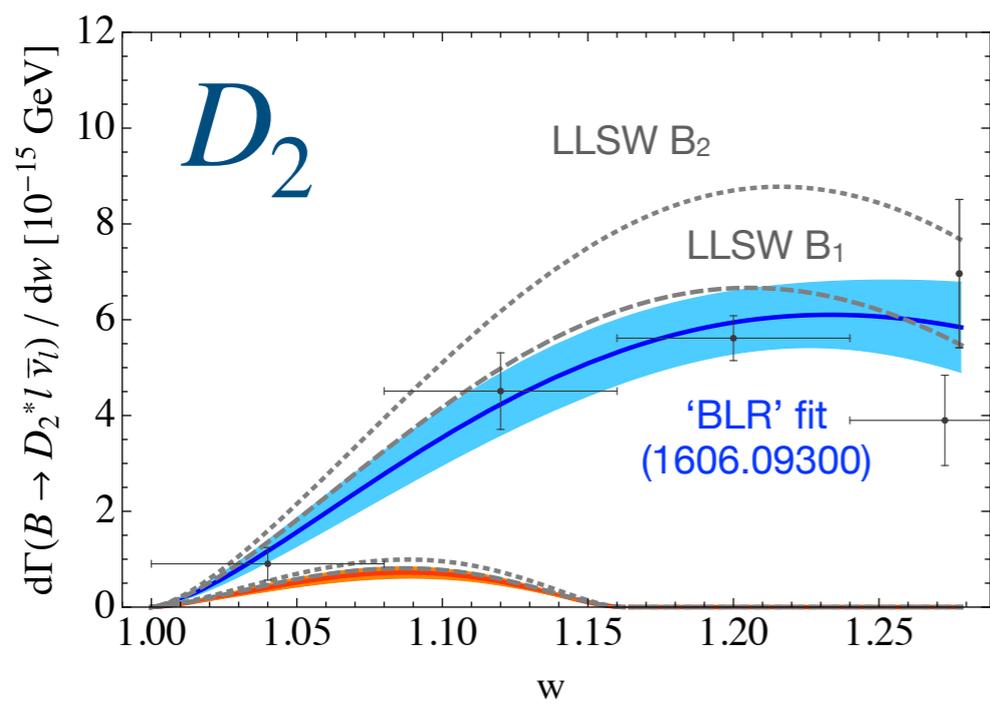
$$D^{(*)}\pi\pi, \dots$$

More exotic contributions ?

X

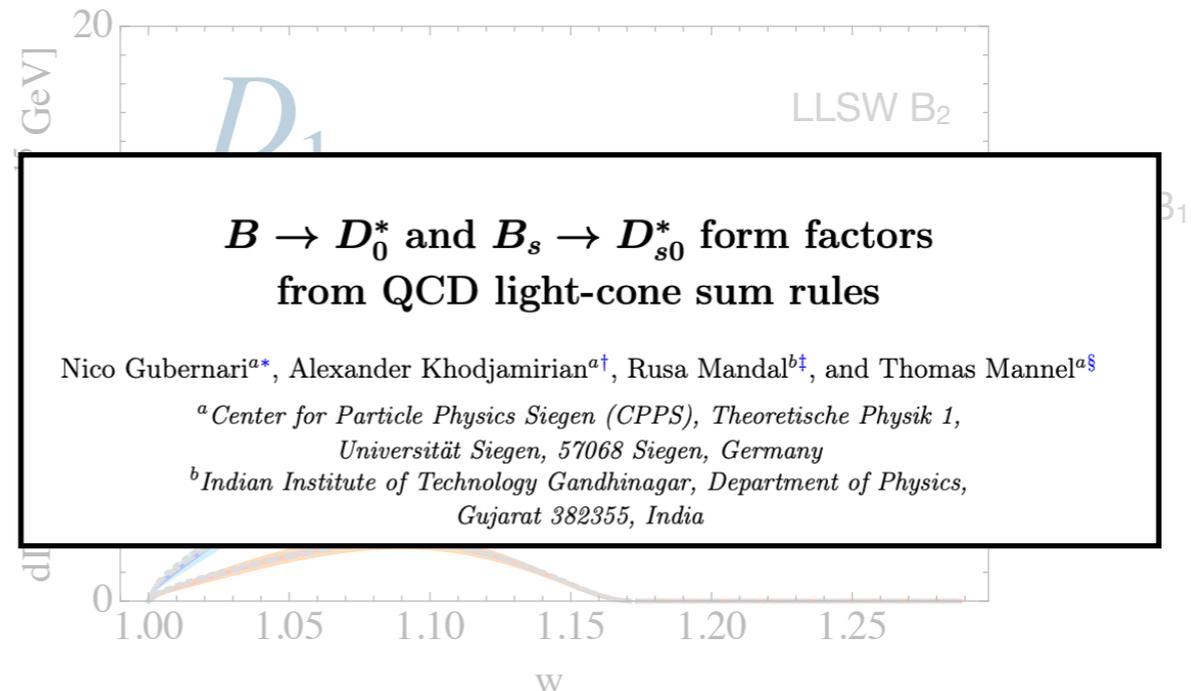
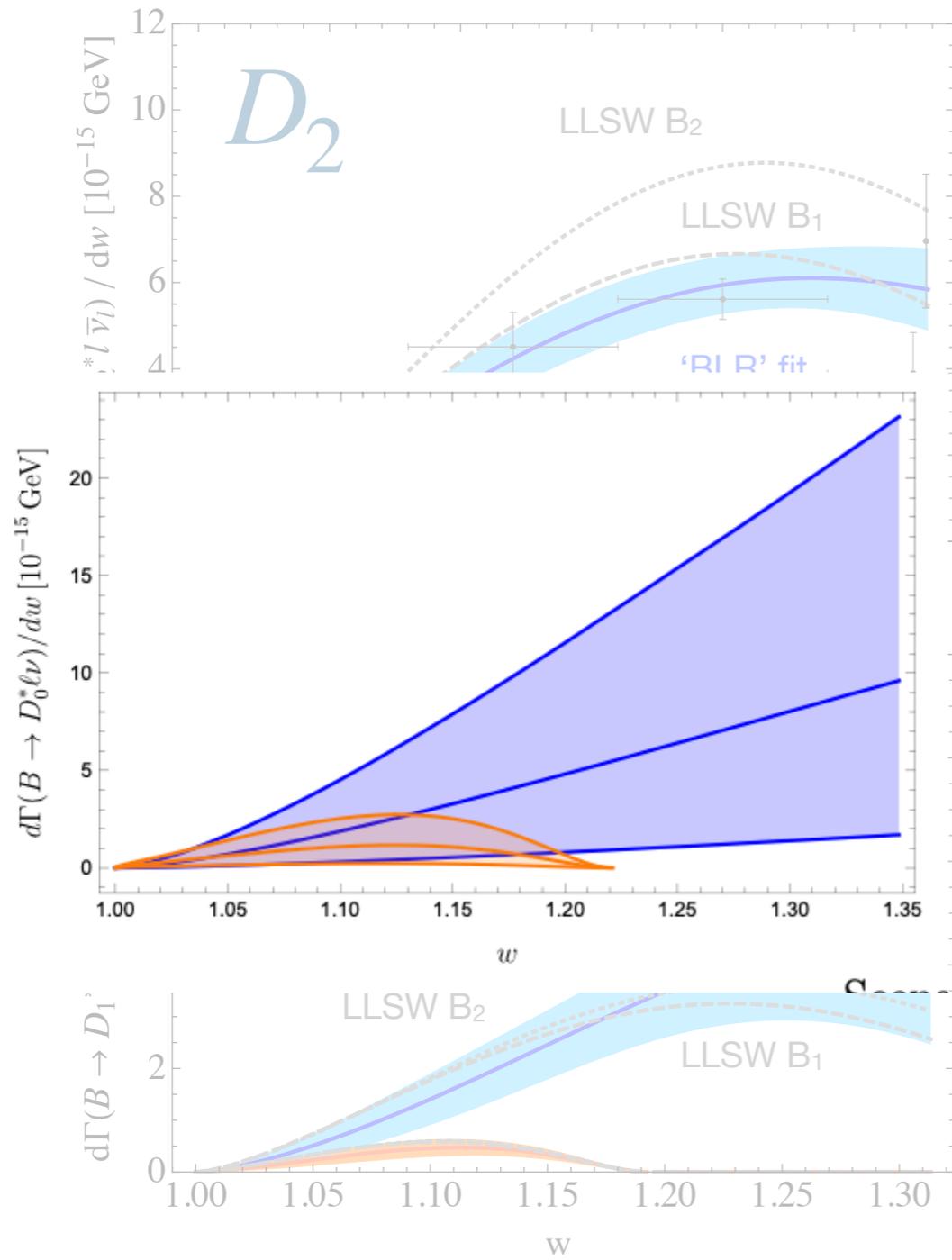


Form Factors



$$w \sim q^2$$

Also ISGW2 and LCSR predictions



$B \rightarrow D_0^*$ and $B_s \rightarrow D_{s0}^*$ form factors from QCD light-cone sum rules

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$B \rightarrow D_1(2420)$ and $B \rightarrow D_1'(2430)$ form factors from QCD light-cone sum rules

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$$w \sim q^2$$

Decay	$\mathcal{B}(B^+)$	$\mathcal{B}(B^0)$
$B \rightarrow D \ell^+ \nu_\ell$	$(2.4 \pm 0.1) \times 10^{-2}$	$(2.2 \pm 0.1) \times 10^{-2}$
$B \rightarrow D^* \ell^+ \nu_\ell$	$(5.5 \pm 0.1) \times 10^{-2}$	$(5.1 \pm 0.1) \times 10^{-2}$
$B \rightarrow D_1 \ell^+ \nu_\ell$	$(6.6 \pm 0.1) \times 10^{-3}$	$(6.2 \pm 0.1) \times 10^{-3}$
$B \rightarrow D_2^* \ell^+ \nu_\ell$	$(2.9 \pm 0.3) \times 10^{-3}$	$(2.7 \pm 0.3) \times 10^{-3}$
$B \rightarrow D_0^* \ell^+ \nu_\ell$	$(4.2 \pm 0.8) \times 10^{-3}$	$(3.9 \pm 0.7) \times 10^{-3}$
$B \rightarrow D_1' \ell^+ \nu_\ell$	$(4.2 \pm 0.9) \times 10^{-3}$	$(3.9 \pm 0.8) \times 10^{-3}$
$B \rightarrow D \pi \pi \ell^+ \nu_\ell$	$(0.6 \pm 0.9) \times 10^{-3}$	$(0.6 \pm 0.9) \times 10^{-3}$
$B \rightarrow D^* \pi \pi \ell^+ \nu_\ell$	$(2.2 \pm 1.0) \times 10^{-3}$	$(2.0 \pm 1.0) \times 10^{-3}$
$B \rightarrow D \eta \ell^+ \nu_\ell$	$(4.0 \pm 4.0) \times 10^{-3}$	$(4.0 \pm 4.0) \times 10^{-3}$
$B \rightarrow D^* \eta \ell^+ \nu_\ell$	$(4.0 \pm 4.0) \times 10^{-3}$	$(4.0 \pm 4.0) \times 10^{-3}$
$B \rightarrow X_c \ell \nu_\ell$	$(10.8 \pm 0.4) \times 10^{-2}$	$(10.1 \pm 0.4) \times 10^{-2}$

fairly well known,
some tensions in
isospin assumption

broad states based
mostly on
three measurements
(BaBar, Belle, DELPHI)

BaBar result

?

Additional Measurements: **resonant + non-resonant** $B \rightarrow D^{(*)} \pi \ell \bar{\nu}_\ell$

Note: $B \rightarrow D \pi \ell \bar{\nu}_\ell$ applies a selection to exclude $B \rightarrow D^*(\rightarrow D \pi_{\text{slow}}) \ell \bar{\nu}_\ell$

Is there space for non-resonant $B \rightarrow D^* \pi \ell \bar{\nu}_\ell$?

resonant $D^{**}(1P) \rightarrow D^{(*)}\pi$

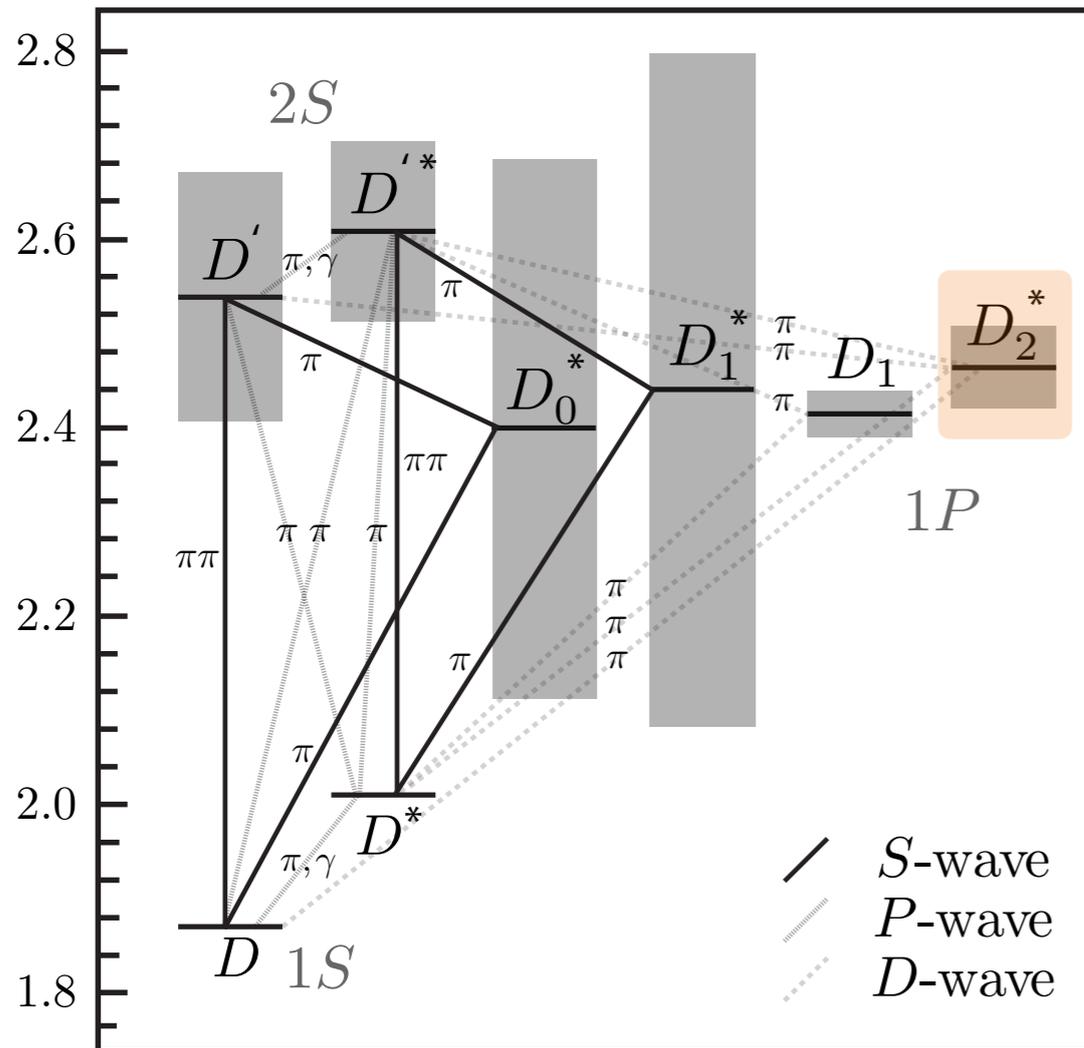
non-resonant and resonant $D^{(*)}\pi$

$$\sum \mathcal{B}(B \rightarrow D^{**}(\rightarrow D\pi)\ell\bar{\nu}_\ell) - \mathcal{B}(B \rightarrow D\pi\ell\bar{\nu}_\ell) = (0.03 \pm 0.09) \times 10^{-2}$$

$$\sum \mathcal{B}(B \rightarrow D^{**}(\rightarrow D^*\pi)\ell\bar{\nu}_\ell) - \mathcal{B}(B \rightarrow D^*\pi\ell\bar{\nu}_\ell) = (-0.11 \pm 0.11) \times 10^{-2}$$

not much space

But not that simple; there might be a sizable contribution we **overlooked**



HFLAV 2019 Averages for narrow D_2

Experiment	$\mathcal{B}(B^- \rightarrow D_2^0(D^{*+}\pi^-)\ell^-\bar{\nu}_\ell)[\%]$ (rescaled)	$\mathcal{B}(B^- \rightarrow D_2^0(D^{*+}\pi^-)\ell^-\bar{\nu}_\ell)[\%]$ (published)
CLEO [506]	$0.055 \pm 0.066_{\text{stat}} \pm 0.011_{\text{syst}}$	$0.059 \pm 0.066_{\text{stat}} \pm 0.011_{\text{syst}}$
D0 [507]	$0.086 \pm 0.018_{\text{stat}} \pm 0.020_{\text{syst}}$	$0.088 \pm 0.018_{\text{stat}} \pm 0.020_{\text{syst}}$
Belle tagged [508]	$0.190 \pm 0.060_{\text{stat}} \pm 0.025_{\text{syst}}$	$0.18 \pm 0.06_{\text{stat}} \pm 0.03_{\text{syst}}$
BABAR tagged [509]	$0.075 \pm 0.013_{\text{stat}} \pm 0.009_{\text{syst}}$	$0.078 \pm 0.013_{\text{stat}} \pm 0.010_{\text{syst}}$
BABAR untagged B^- [510]	$0.087 \pm 0.009_{\text{stat}} \pm 0.007_{\text{syst}}$	$0.087 \pm 0.013_{\text{stat}} \pm 0.007_{\text{syst}}$
BABAR untagged B^0 [510]	$0.065 \pm 0.010_{\text{stat}} \pm 0.004_{\text{syst}}$	$0.087 \pm 0.013_{\text{stat}} \pm 0.007_{\text{syst}}$
Average	$0.077 \pm 0.006 \pm 0.004$	$\chi^2/\text{dof} = 5.4/5$ (CL=36.7%)

Used branching fraction ratio for $D_2 \rightarrow D^*\pi$ over $D_2 \rightarrow D\pi$

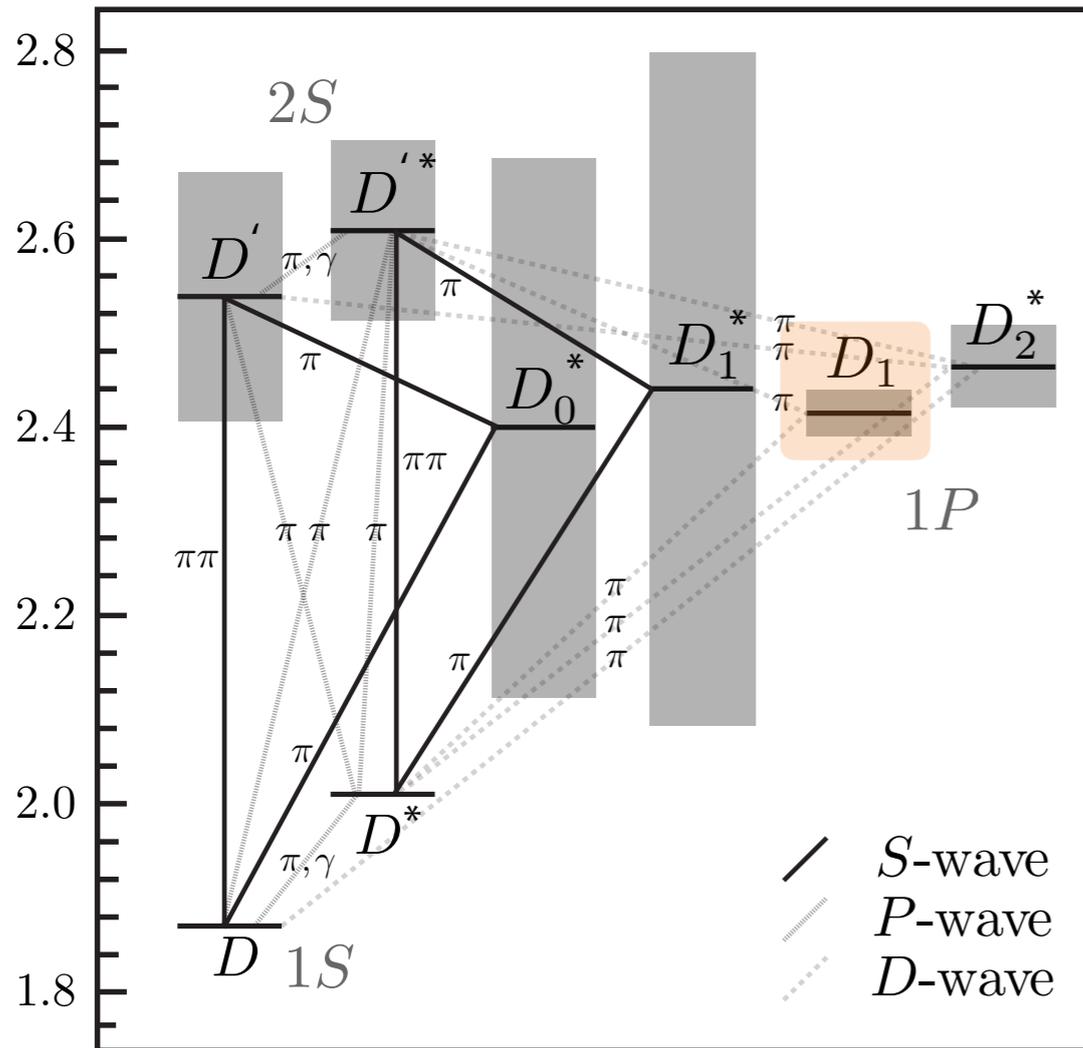
$$f_{D_2} = \frac{\mathcal{B}(D_2 \rightarrow D^-\pi^+)}{\mathcal{B}(D_2 \rightarrow D^{*-}\pi^+)} = 1.54 \pm 0.15$$

Isospin factor for conjugate 2-body pion modes:

$$f_\pi = \frac{\mathcal{B}(D^{**} \rightarrow D^{(*)-}\pi^+)}{\mathcal{B}(D^{**} \rightarrow D^{(*)}\pi)} = \frac{2}{3}$$

Full BF: $\overline{\mathcal{B}}(B^+ \rightarrow D_2\ell^+\bar{\nu}_\ell) = (0.29 \pm 0.03) \times 10^{-2}$

HFLAV 2019 Averages for narrow D_1



Experiment	$\mathcal{B}(B^- \rightarrow D_1^0(D^{*+}\pi^-)\ell^-\bar{\nu}_\ell)[\%]$ (rescaled)	$\mathcal{B}(B^- \rightarrow D_1^0(D^{*+}\pi^-)\ell^-\bar{\nu}_\ell)[\%]$ (published)
ALEPH [504]	$0.436 \pm 0.085_{\text{stat}} \pm 0.056_{\text{syst}}$	$0.47 \pm 0.10_{\text{stat}} \pm 0.07_{\text{syst}}$
OPAL [505]	$0.568 \pm 0.210_{\text{stat}} \pm 0.100_{\text{syst}}$	$0.70 \pm 0.21_{\text{stat}} \pm 0.10_{\text{syst}}$
CLEO [506]	$0.349 \pm 0.085_{\text{stat}} \pm 0.056_{\text{syst}}$	$0.373 \pm 0.085_{\text{stat}} \pm 0.057_{\text{syst}}$
D0 [507]	$0.214 \pm 0.018_{\text{stat}} \pm 0.035_{\text{syst}}$	$0.219 \pm 0.018_{\text{stat}} \pm 0.035_{\text{syst}}$
Belle Tagged B^- [508]	$0.430 \pm 0.070_{\text{stat}} \pm 0.059_{\text{syst}}$	$0.42 \pm 0.07_{\text{stat}} \pm 0.07_{\text{syst}}$
Belle Tagged B^0 [508]	$0.593 \pm 0.200_{\text{stat}} \pm 0.076_{\text{syst}}$	$0.42 \pm 0.07_{\text{stat}} \pm 0.07_{\text{syst}}$
BABAR Tagged [509]	$0.277 \pm 0.030_{\text{stat}} \pm 0.029_{\text{syst}}$	$0.29 \pm 0.03_{\text{stat}} \pm 0.03_{\text{syst}}$
BABAR Untagged B^- [510]	$0.293 \pm 0.017_{\text{stat}} \pm 0.016_{\text{syst}}$	$0.30 \pm 0.02_{\text{stat}} \pm 0.02_{\text{syst}}$
BABAR Untagged B^0 [510]	$0.282 \pm 0.026_{\text{stat}} \pm 0.023_{\text{syst}}$	$0.30 \pm 0.02_{\text{stat}} \pm 0.02_{\text{syst}}$
Average	$0.281 \pm 0.010 \pm 0.015$	$\chi^2/\text{dof} = 12.3/8$ (CL=13.8%)

Measurements only consider two-body final states, but there are also 3-body contributions:

$$f_{D_1} = \frac{\mathcal{B}(D_1 \rightarrow D^- \pi^+)}{\mathcal{B}(D_1 \rightarrow D^0 \pi^+ \pi^-)} = 2.32 \pm 0.54$$

Full BF: $\overline{\mathcal{B}}(B^+ \rightarrow D_1 \ell^+ \bar{\nu}_\ell) = (0.66 \pm 0.11) \times 10^{-2}$

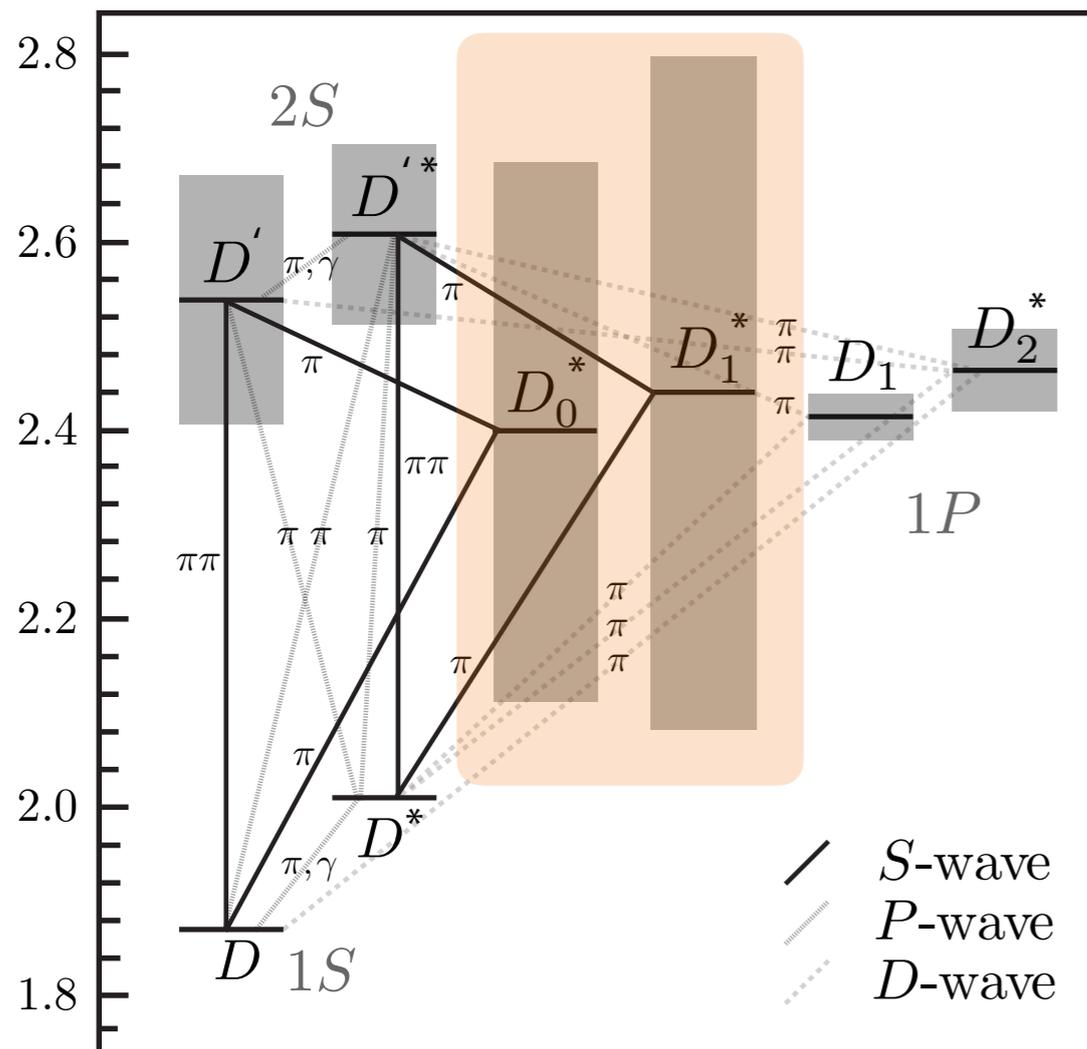
Isospin factor for conjugate 2-body pion modes:

$$f_\pi = \frac{\mathcal{B}(D^{**} \rightarrow D^{(*)-} \pi^+)}{\mathcal{B}(D^{**} \rightarrow D^{(*)} \pi)} = \frac{2}{3}$$

Isospin factor for conjugate 3-body pion modes:

$$f_{\pi\pi} = \frac{\mathcal{B}(D^{**} \rightarrow D^{(*)-} \pi^+ \pi^-)}{\mathcal{B}(D^{**} \rightarrow D^{(*)} \pi \pi)} = \frac{1}{2} \pm \frac{1}{6}$$

uncertainties cover assumptions for pure 3 body decay and the decay via intermediate states ($\rho \rightarrow \pi\pi, f_0 \rightarrow \pi\pi$)



Isospin factor for missing
2-body pion modes:

$$f_\pi = \frac{\mathcal{B}(D^{**} \rightarrow D^{(*)-}\pi^+)}{\mathcal{B}(D^{**} \rightarrow D^{(*)}\pi)} = \frac{2}{3}$$

HFLAV 2019 Averages for broad states

Experiment	$\mathcal{B}(B^- \rightarrow D_1^{\prime 0}(D^{*+}\pi^-)\ell^-\bar{\nu}_\ell)[\%]$ (rescaled)	$\mathcal{B}(B^- \rightarrow D_1^{\prime 0}(D^{*+}\pi^-)\ell^-\bar{\nu}_\ell)[\%]$ (published)
DELPHI [511]	$0.73 \pm 0.17_{\text{stat}} \pm 0.18_{\text{syst}}$	$0.83 \pm 0.17_{\text{stat}} \pm 0.18_{\text{syst}}$
Belle [508]	$-0.03 \pm 0.06_{\text{stat}} \pm 0.07_{\text{syst}}$	$-0.03 \pm 0.06_{\text{stat}} \pm 0.07_{\text{syst}}$
BABAR [509]	$0.26 \pm 0.04_{\text{stat}} \pm 0.04_{\text{syst}}$	$0.27 \pm 0.04_{\text{stat}} \pm 0.05_{\text{syst}}$
Average	$0.19 \pm 0.03 \pm 0.04$	$\chi^2/\text{dof} = 11.9/2$ (CL=0.003%)

Extremely poor agreement for D_1'

Average w/o Belle

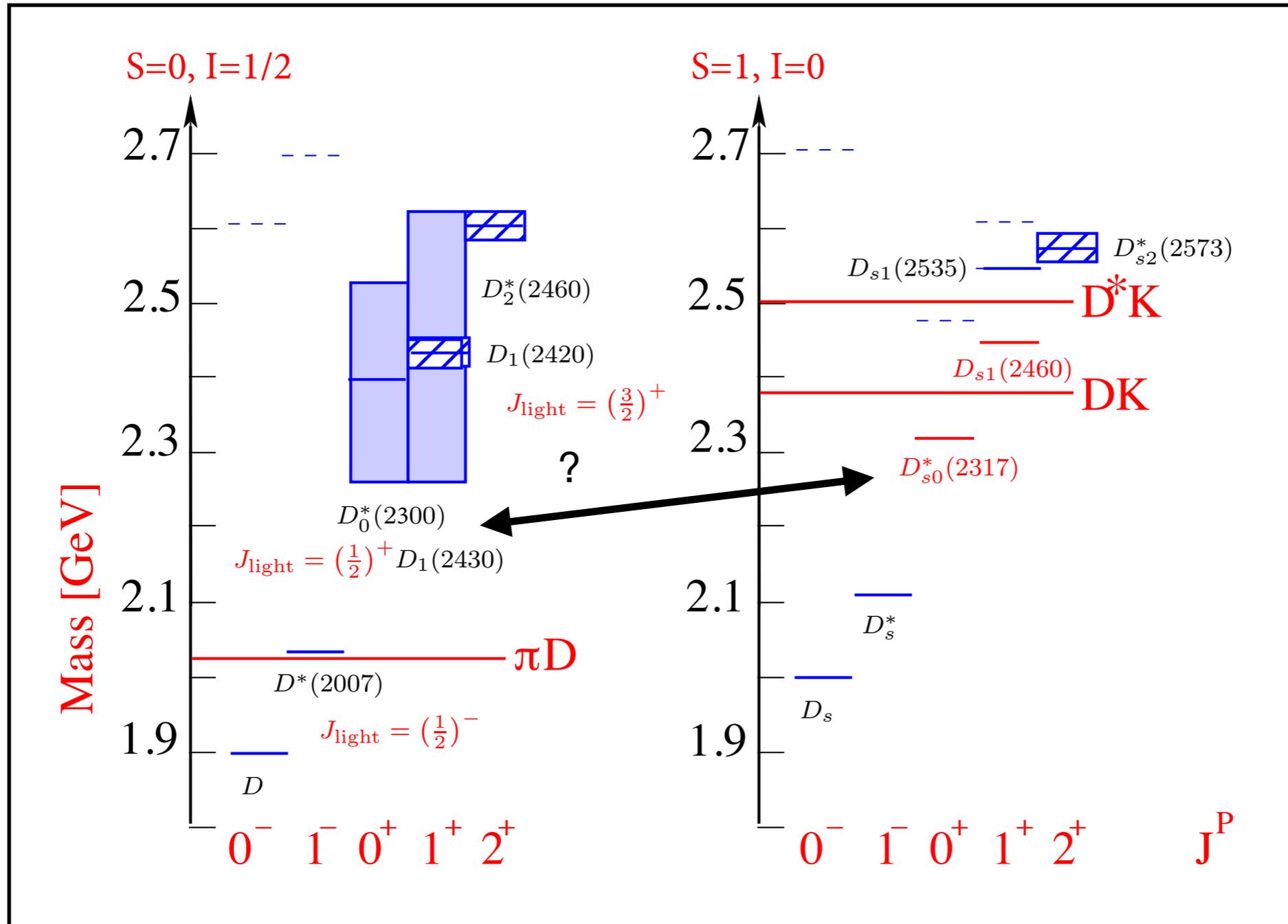
$$\overline{\mathcal{B}}(B^+ \rightarrow D_1'(\rightarrow D^{*-}\pi^+)\ell^+\bar{\nu}_\ell) = (0.28 \pm 0.06) \times 10^{-2}$$

Experiment	$\mathcal{B}(B^- \rightarrow D_0^{*0}(D^+\pi^-)\ell^-\bar{\nu}_\ell)[\%]$ (rescaled)	$\mathcal{B}(B^- \rightarrow D_0^{*0}(D^+\pi^-)\ell^-\bar{\nu}_\ell)[\%]$ (published)
Belle Tagged B^- [508]	$0.25 \pm 0.04_{\text{stat}} \pm 0.06_{\text{syst}}$	$0.24 \pm 0.04_{\text{stat}} \pm 0.06_{\text{syst}}$
Belle Tagged B^0 [508]	$0.22 \pm 0.08_{\text{stat}} \pm 0.06_{\text{syst}}$	$0.24 \pm 0.04_{\text{stat}} \pm 0.06_{\text{syst}}$
BABAR Tagged [509]	$0.32 \pm 0.04_{\text{stat}} \pm 0.05_{\text{syst}}$	$0.26 \pm 0.05_{\text{stat}} \pm 0.04_{\text{syst}}$
Average	$0.28 \pm 0.03 \pm 0.04$	$\chi^2/\text{dof} = 0.82/2$ (CL=66.4%)

Full BF:

$$\overline{\mathcal{B}}(B^+ \rightarrow D_1'\ell^+\bar{\nu}_\ell) = (0.42 \pm 0.09) \times 10^{-2}$$

$$\overline{\mathcal{B}}(B^+ \rightarrow D_0\ell^+\bar{\nu}_\ell) = (0.42 \pm 0.08) \times 10^{-2}$$

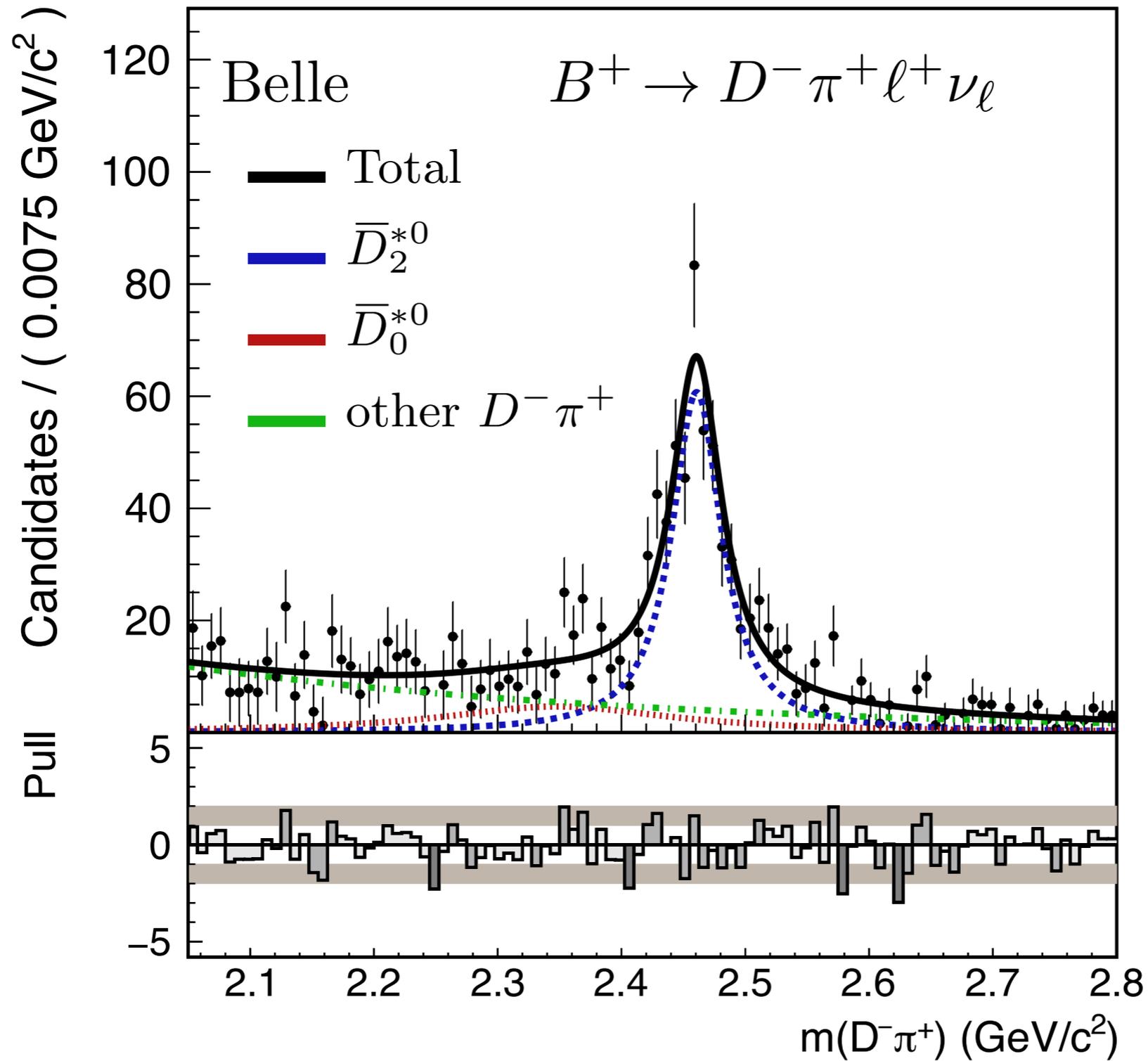


Looking into the strange sector :

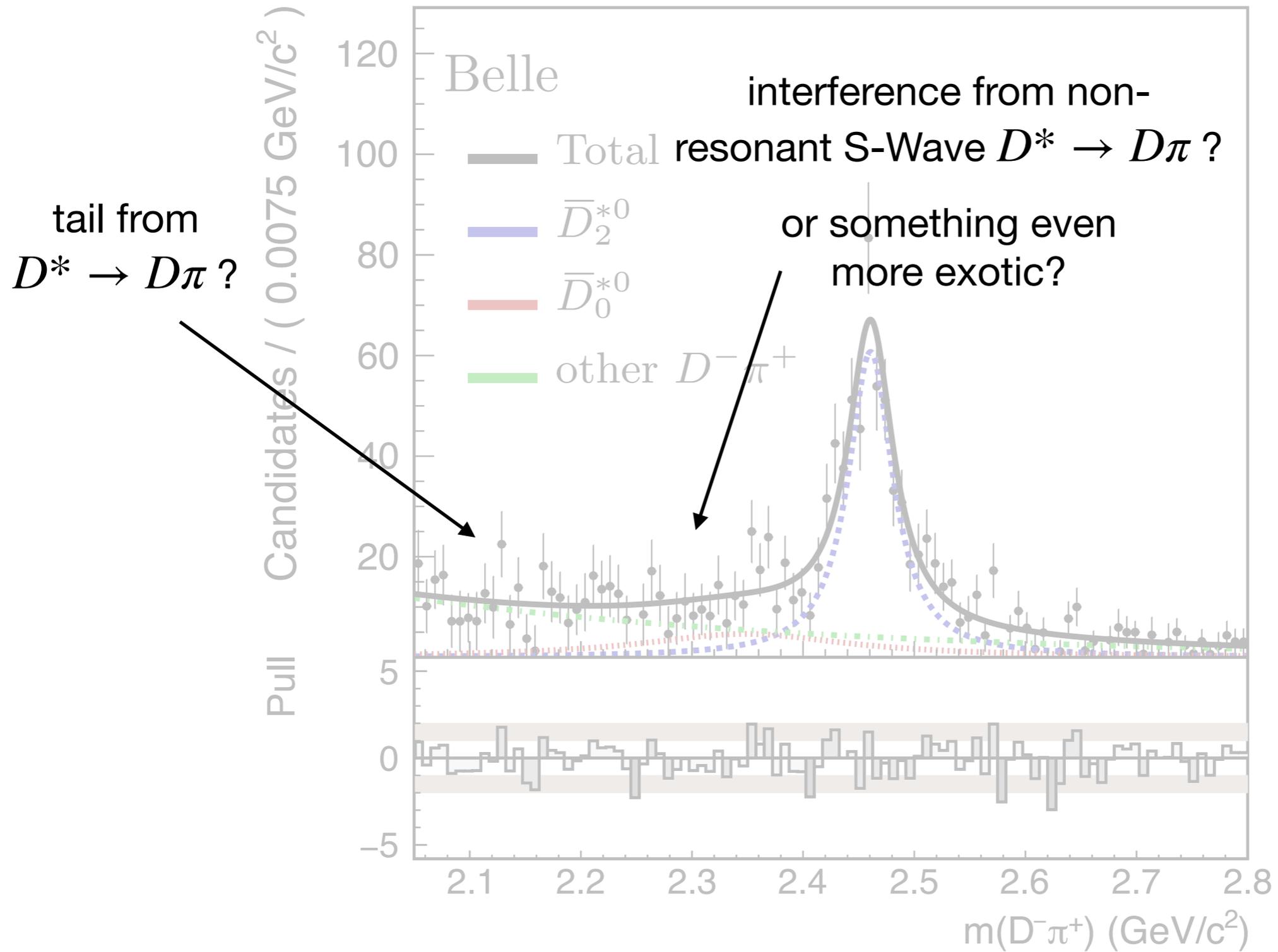
$$\begin{aligned} M(D_0^*) &\simeq M(D_{s0}^*) \\ M(D_1) &\simeq M(D_{s1}) \end{aligned}$$

Why so light?

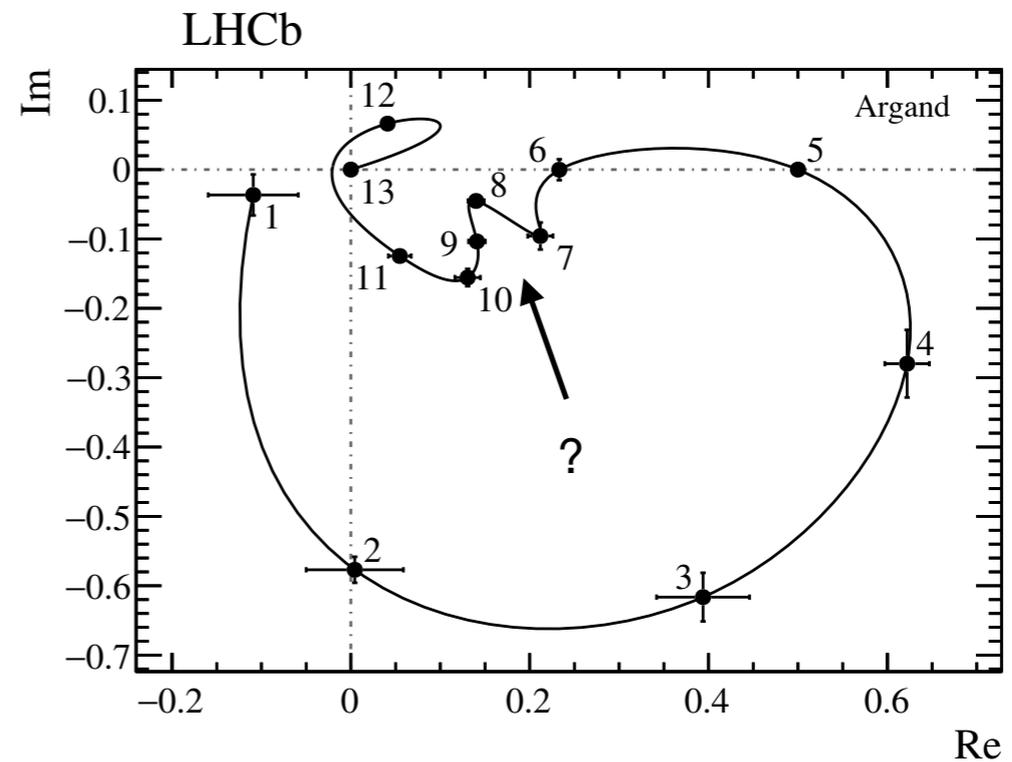
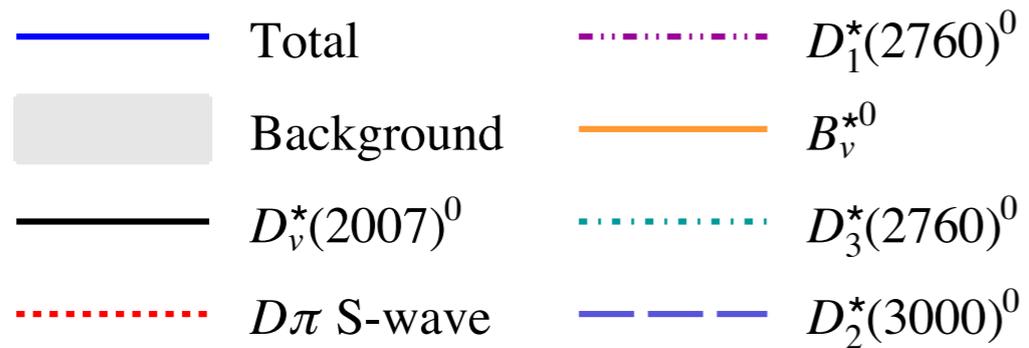
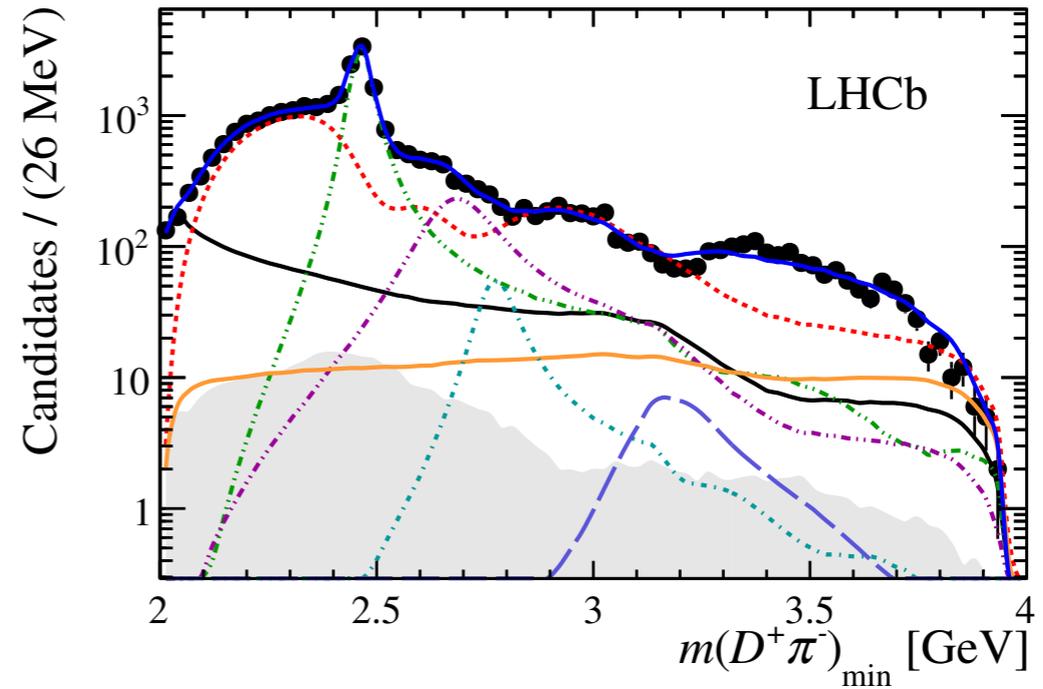
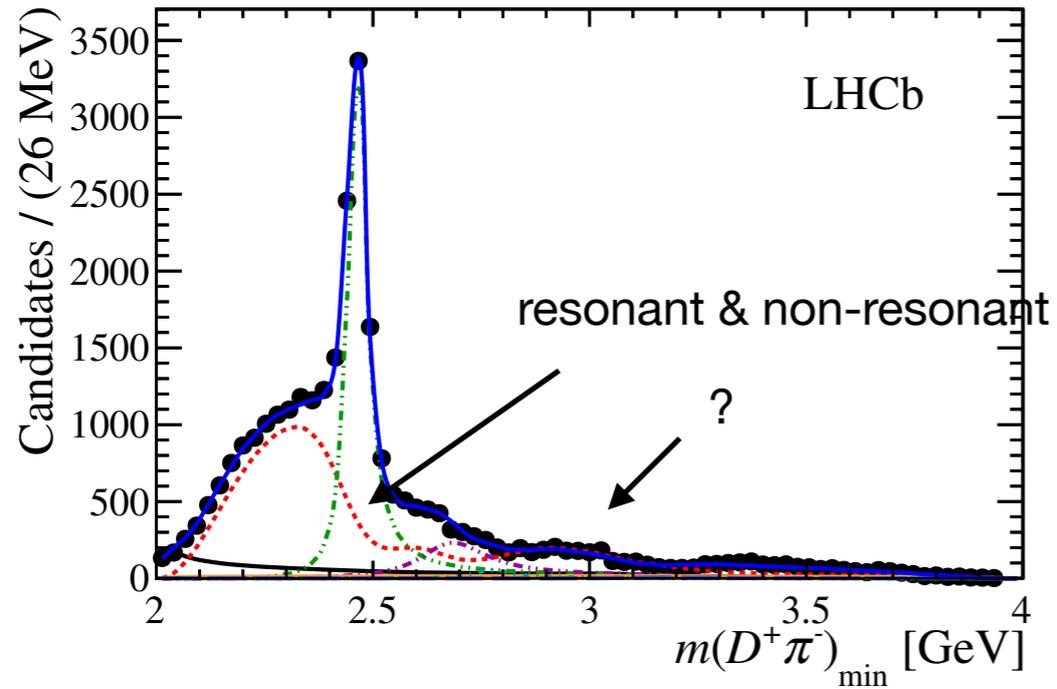
The Mass spectrum is intricate



The Mass spectrum is intricate



Mass spectrum in $B^+ \rightarrow D^- \pi^+ \pi^+$



PDG masses & widths:

$$D_0^*(2300) : M = 2300 \pm 19 \text{ MeV}, \Gamma = 274 \pm 40 \text{ MeV}, I(J^P) = \frac{1}{2}(0^+),$$

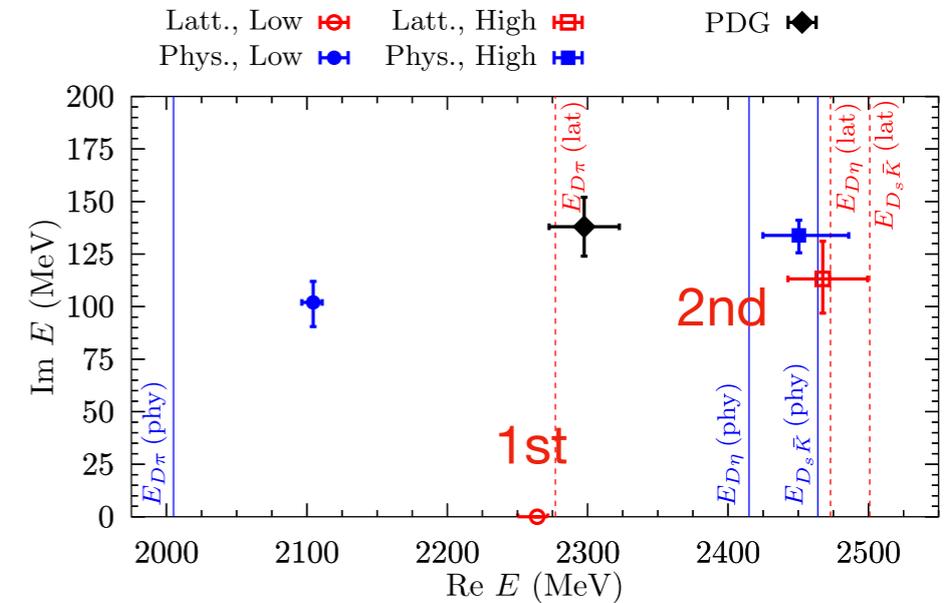
$$D_{s0}^*(2317) : M = 2318.0 \pm 0.7 \text{ MeV}, \Gamma < 3.8 \text{ MeV}, I(J^P) = 0(0^+).$$

1) **Lattice studies** with non-physical pion masses imply the existence of a second pole:

U χ PT prediction

extrapolated to physical pion masses

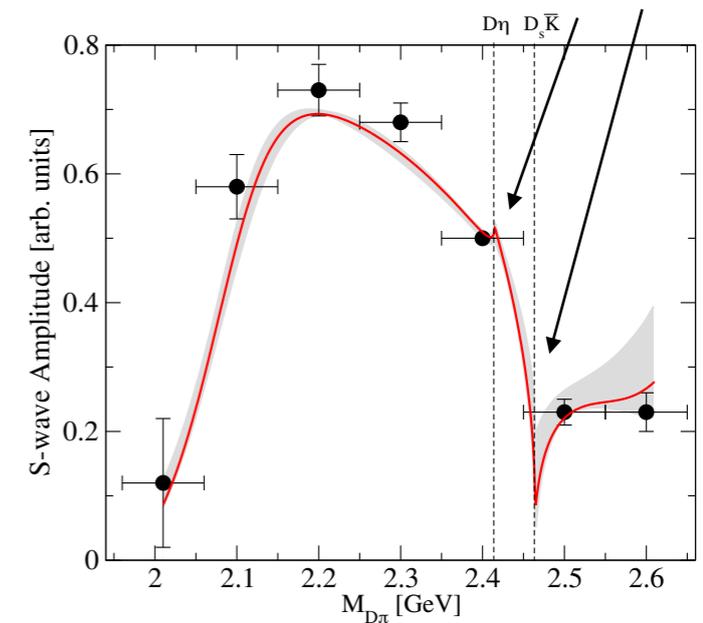
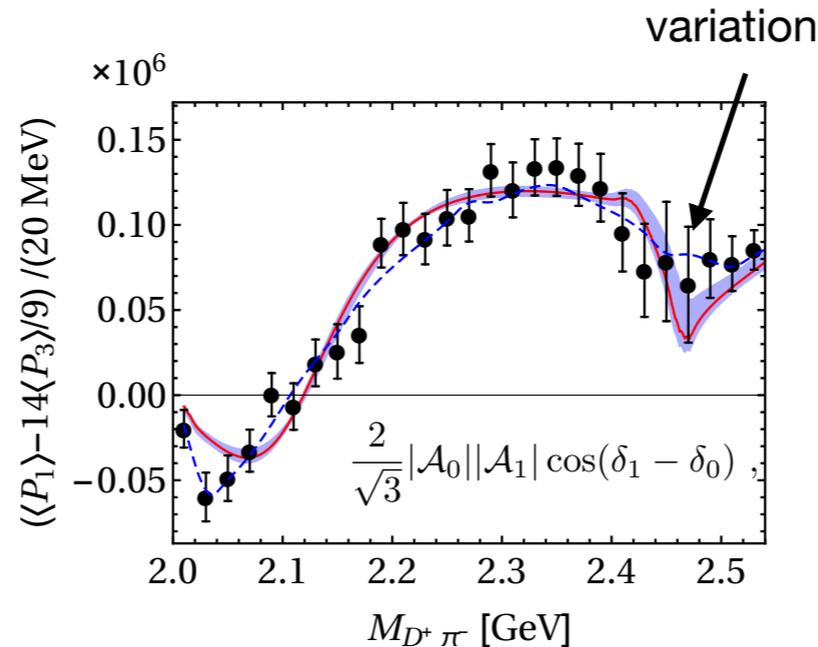
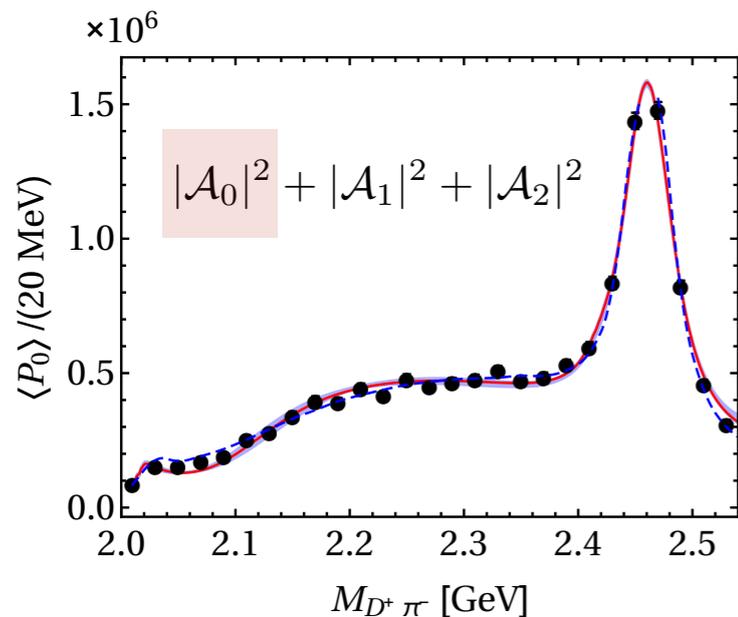
M (MeV)	$\Gamma/2$ (MeV)
2105_{-8}^{+6}	102_{-12}^{+10}
2451_{-26}^{+36}	134_{-8}^{+7}



2) Fit to Legendre moments of LHCb analysis with two poles:

two cusps indicating the presence of $D\eta, D_s K$

$$A(B^- \rightarrow D^+ \pi^- \pi^-) = \sum_{L=0}^2 \sqrt{2L+1} \mathcal{A}_L(s) P_L(z),$$

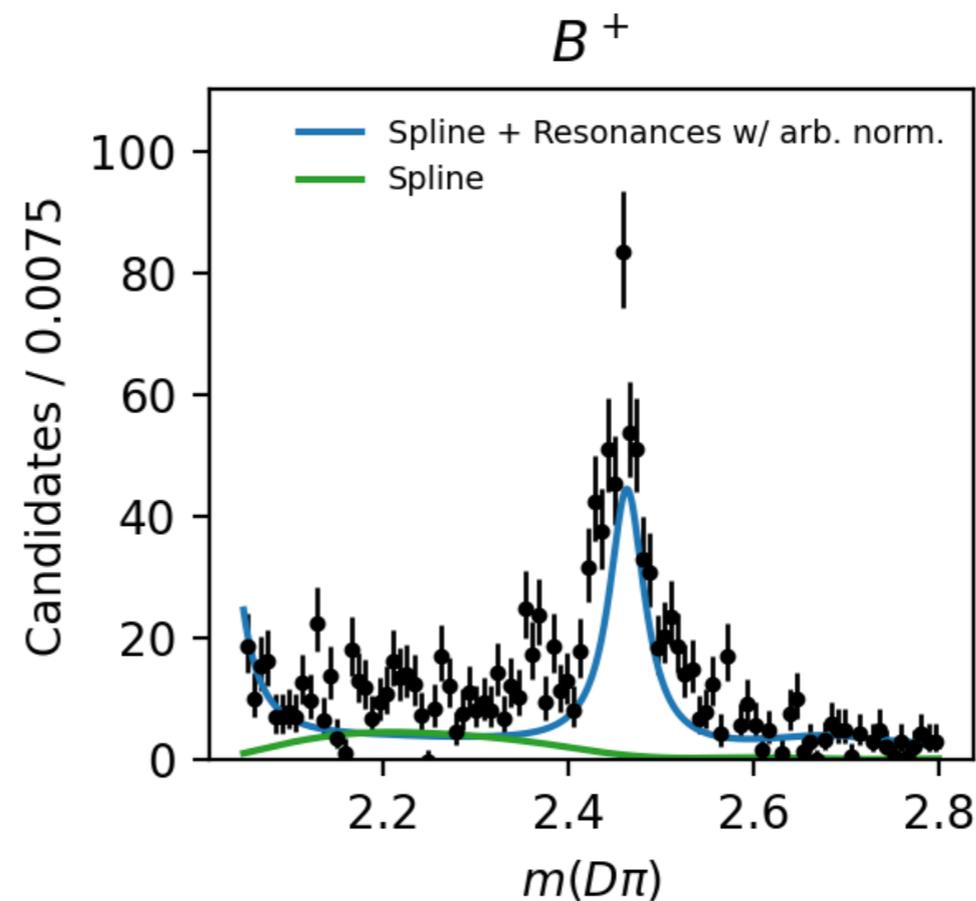
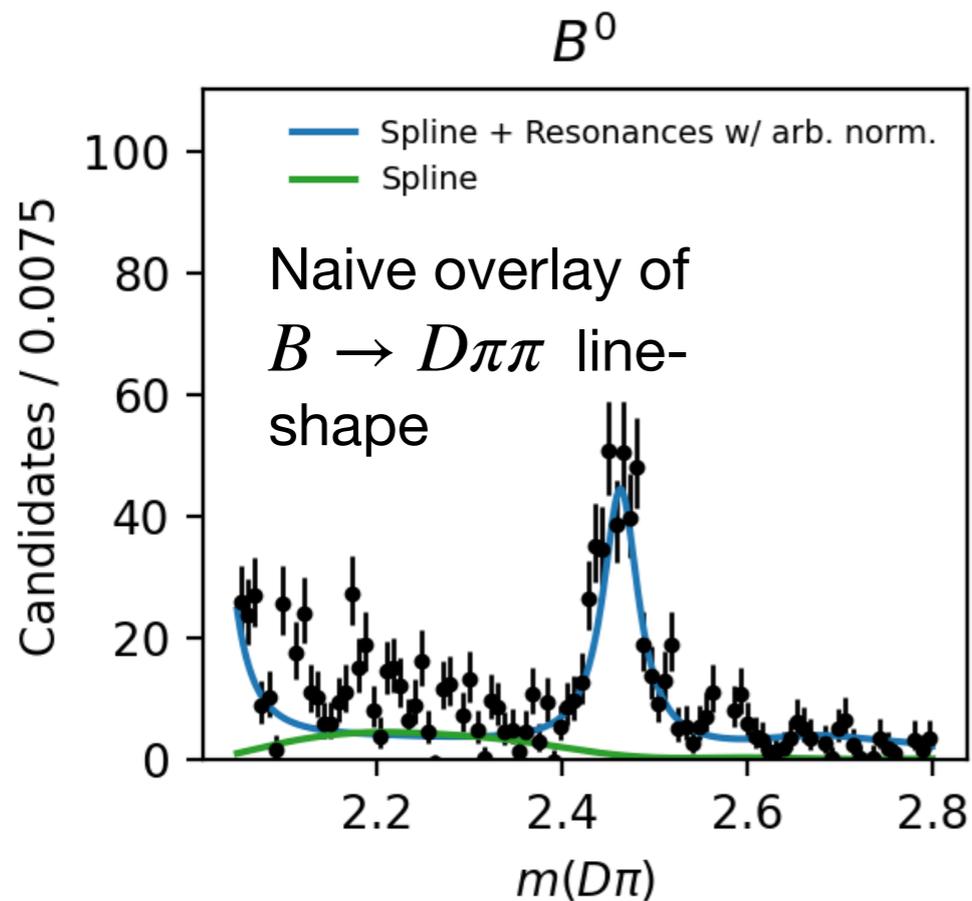
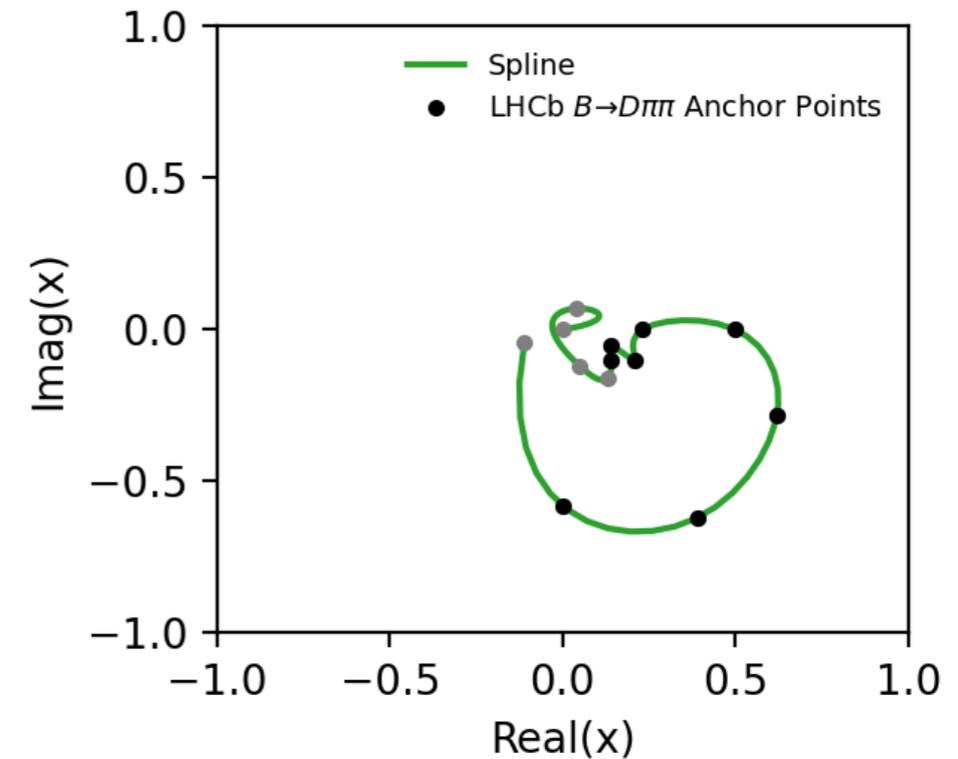


Overlaying the line-shape

Can the information from hadronic decays be combined with the semileptonic spectrum?

A priori yes, but need to account for **different production amplitudes** of $(D\pi)_{S\text{-Wave}}$, D^* and D_2

The poles and widths though are identical between both types of decays



First (!) Fit to the semileptonic spectrum

First fit to the semileptonic spectrum:

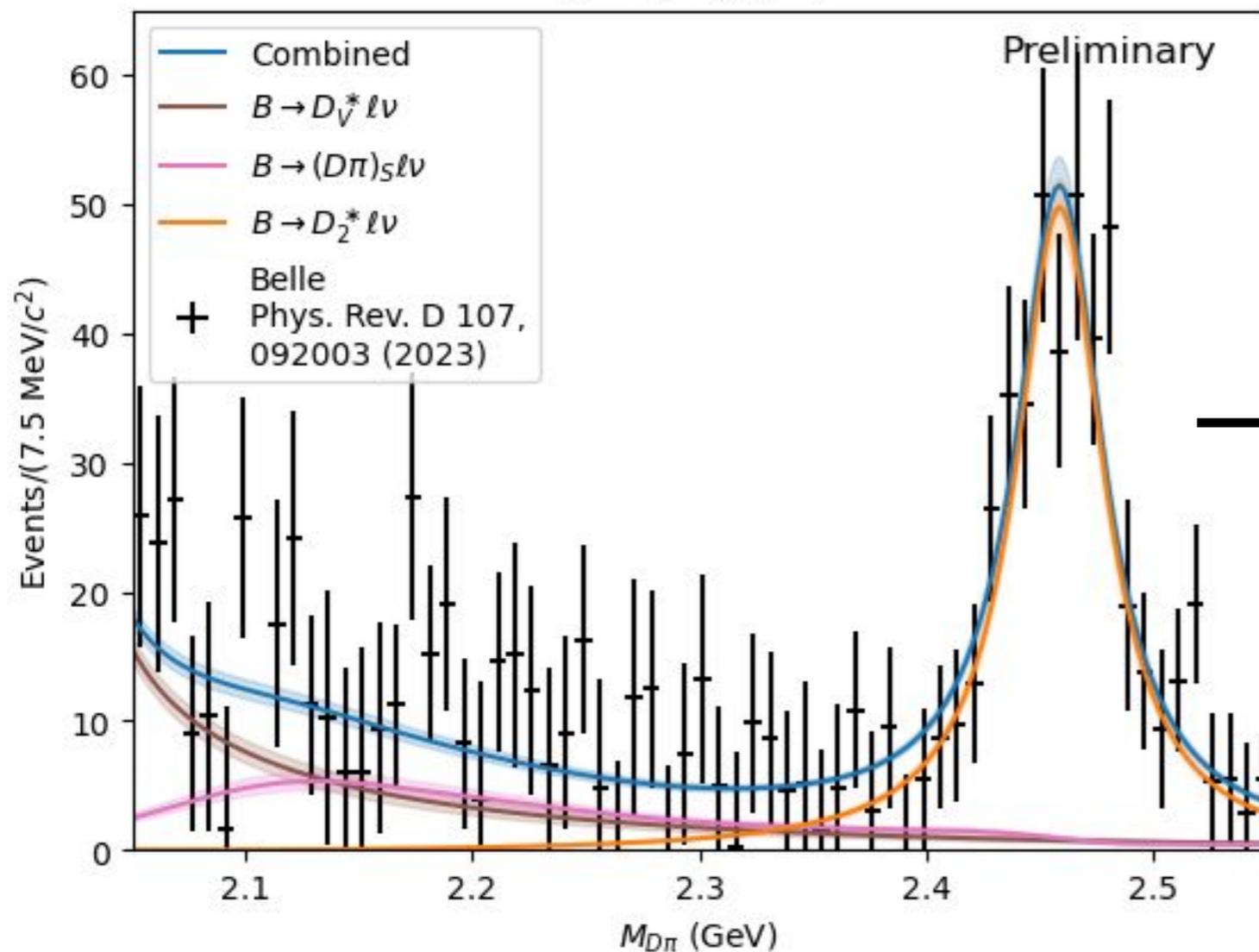
A model-independent description of $B \rightarrow D\pi l\nu$ decays

Erik J. Gustafson,¹ Florian Herren,^{1,2} Ruth S. Van de Water,¹ Raynette van Tonder,³ and Michael L. Wagman¹

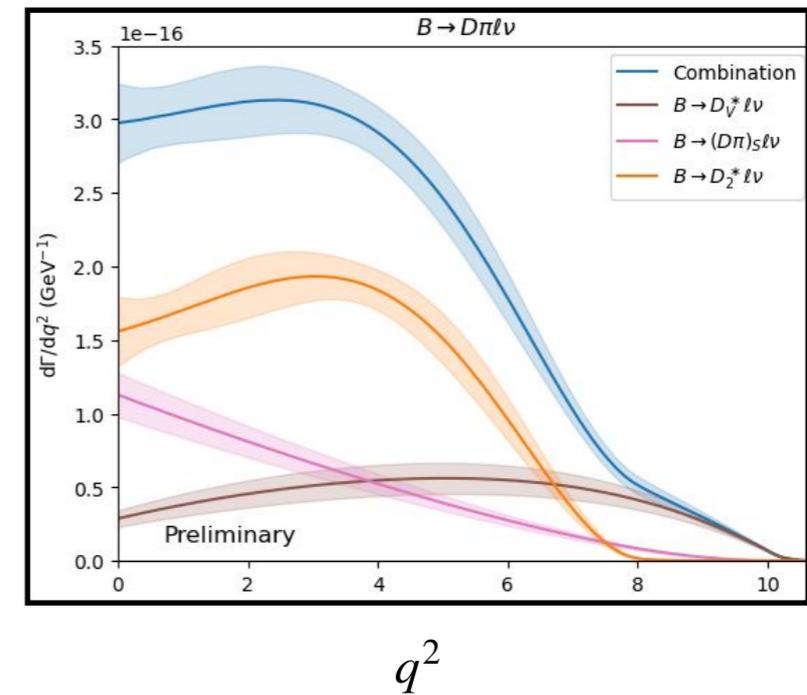
Fit which includes tail of D^* , $(D\pi)_{P\text{-Wave}}$ and D_2 : no interference in $m_{D\pi}$ due to different spins

Using Omnes matrix for coupled-channel $D\text{-}\pi$ scattering

$$B^0 \rightarrow \bar{D}^0 \pi^- l^+ \nu$$



Combining with FNAL/MILC D^* LQCD pred. & fit to experimental D_2 spectrum:



Our entire parametrization of FFs relies on the **narrow-width approximation**

$$\langle B | H_\mu | D^{**} \rangle = \text{on-shell to on-shell transition}$$

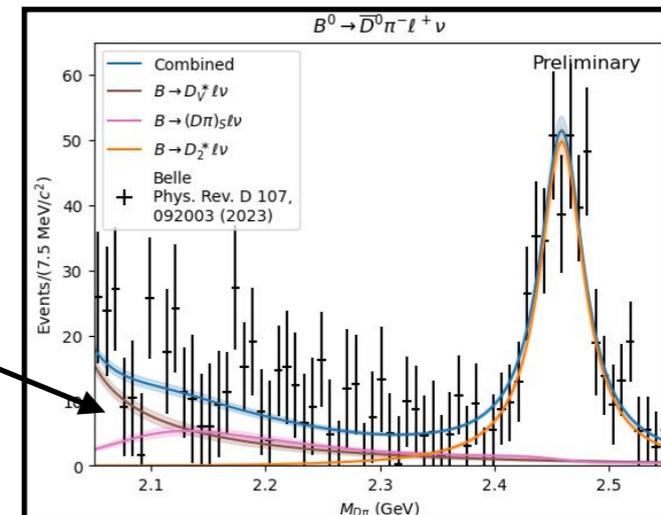
Non-trivial deviation from rate expressions : my understanding is

i.e. expression of rate near mode of BW \neq expression of rate in off-shell region

Can we derive from **first principles** expressions for this?

Important for broad states, but **also** for $D^* \rightarrow D\pi$

(Tail often parametrized as an exponential of a polynomial, float then slope and curvature)



Narrow-width approximation

- Ignore longitudinal contributions
- Product of transverse amplitudes times lineshape (Breit-Wigner, Blatt-Weisskopf etc)

Cannot formally describe off-shell contributions, but critical for moderately/very broad states!

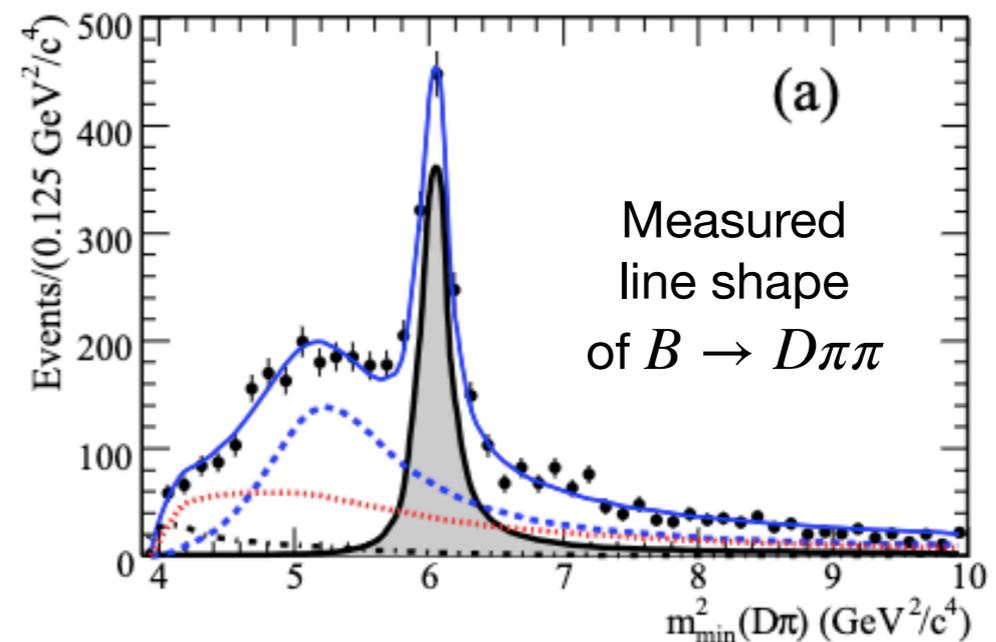
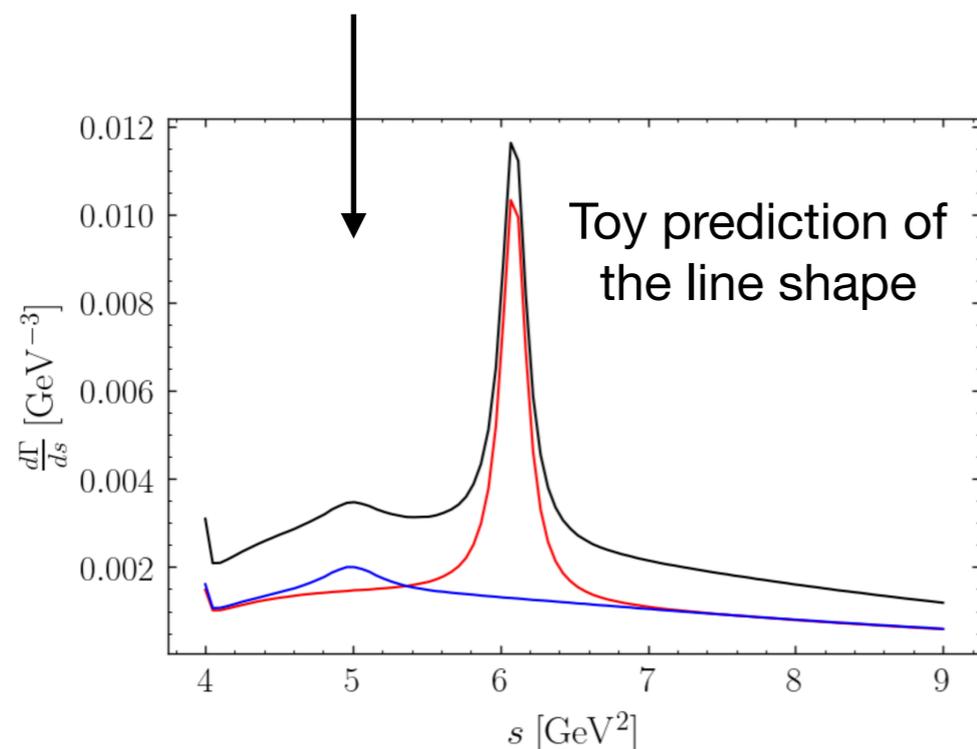
(Incorporating longitudinal piece inconsistent with form factor construction and/or effective theories (like HQET))

Partonic Approach (Robinson *et al*, soon on arXiv)

- On-shell recursion techniques for full amplitude
- 'holomorphic' generalization of FFs/HQET

Hadronic Approach (Papucci *et al*, in progress)

- On-shell recursion techniques for Isgur-Wise functions
- Adaptation of Heavy Hadron Chiral Perturbation Theory



<https://arxiv.org/abs/0901.1291>

$$B \rightarrow D^{(*)} \pi \pi \ell \bar{\nu}_\ell$$

Note: does not include Frank's measurement yet
(need to average it with BaBar) # 19

Channel	$R_{\pi^+\pi^-}^{(*)} \times 10^3$	$\mathcal{B} \times 10^5$
$D^0 \pi^+ \pi^- \ell^- \bar{\nu}$	$71 \pm 13 \pm 8$	$161 \pm 30 \pm 18 \pm 8$
$D^+ \pi^+ \pi^- \ell^- \bar{\nu}$	$58 \pm 18 \pm 12$	$127 \pm 39 \pm 26 \pm 7$
$D^{*0} \pi^+ \pi^- \ell^- \bar{\nu}$	$14 \pm 7 \pm 4$	$80 \pm 40 \pm 23 \pm 3$
$D^{*+} \pi^+ \pi^- \ell^- \bar{\nu}$	$28 \pm 8 \pm 6$	$138 \pm 39 \pm 30 \pm 3$
$D \pi^+ \pi^- \ell^- \bar{\nu}$	$67 \pm 10 \pm 8$	$152 \pm 23 \pm 18 \pm 7$
$D^* \pi^+ \pi^- \ell^- \bar{\nu}$	$19 \pm 5 \pm 4$	$108 \pm 28 \pm 23 \pm 4$

How much of $B \rightarrow D \pi \pi \ell \bar{\nu}_\ell$ is saturated by $B \rightarrow D_1(\rightarrow D \pi \pi) \ell \bar{\nu}_\ell$?

$$\mathcal{B}(B^+ \rightarrow D \pi \pi \ell \bar{\nu}_\ell) - \overline{\mathcal{B}}(B^+ \rightarrow D_1(\rightarrow D \pi \pi) \ell \bar{\nu}_\ell) = 0.06 \pm 0.09$$

most of it

Isospin factor for conjugate
3-body pion modes:

$$f_{\pi\pi} = \frac{\mathcal{B}(D^{**} \rightarrow D^{(*)-} \pi^+ \pi^-)}{\mathcal{B}(D^{**} \rightarrow D^{(*)} \pi \pi)} = \frac{1}{2} \pm \frac{1}{6}$$

Here $B \rightarrow D \pi \pi \ell \bar{\nu}_\ell$ only contains the
non- D_1 contributions

Full BF:

$$\begin{aligned} \mathcal{B}(B^+ \rightarrow D \pi \pi \ell \bar{\nu}_\ell) &= (0.06 \pm 0.09) \times 10^{-2} \\ \mathcal{B}(B^+ \rightarrow D^* \pi \pi \ell \bar{\nu}_\ell) &= (0.22 \pm 0.10) \times 10^{-2} \end{aligned}$$

The 'Gap'

Decay	$\mathcal{B}(B^+)$	$\mathcal{B}(B^0)$
$B \rightarrow D \ell^+ \nu_\ell$	$(2.4 \pm 0.1) \times 10^{-2}$	$(2.2 \pm 0.1) \times 10^{-2}$
$B \rightarrow D^* \ell^+ \nu_\ell$	$(5.5 \pm 0.1) \times 10^{-2}$	$(5.1 \pm 0.1) \times 10^{-2}$
$B \rightarrow D_1 \ell^+ \nu_\ell$	$(6.6 \pm 0.1) \times 10^{-3}$	$(6.2 \pm 0.1) \times 10^{-3}$
$B \rightarrow D_2^* \ell^+ \nu_\ell$	$(2.9 \pm 0.3) \times 10^{-3}$	$(2.7 \pm 0.3) \times 10^{-3}$
$B \rightarrow D_0^* \ell^+ \nu_\ell$	$(4.2 \pm 0.8) \times 10^{-3}$	$(3.9 \pm 0.7) \times 10^{-3}$
$B \rightarrow D_1' \ell^+ \nu_\ell$	$(4.2 \pm 0.9) \times 10^{-3}$	$(3.9 \pm 0.8) \times 10^{-3}$
$B \rightarrow D\pi\pi \ell^+ \nu_\ell$	$(0.6 \pm 0.9) \times 10^{-3}$	$(0.6 \pm 0.9) \times 10^{-3}$
$B \rightarrow D^*\pi\pi \ell^+ \nu_\ell$	$(2.2 \pm 1.0) \times 10^{-3}$	$(2.0 \pm 1.0) \times 10^{-3}$
$B \rightarrow D\eta \ell^+ \nu_\ell$	$(4.0 \pm 4.0) \times 10^{-3}$	$(4.0 \pm 4.0) \times 10^{-3}$
$B \rightarrow D^*\eta \ell^+ \nu_\ell$	$(4.0 \pm 4.0) \times 10^{-3}$	$(4.0 \pm 4.0) \times 10^{-3}$
$B \rightarrow X_c \ell \nu_\ell$	$(10.8 \pm 0.4) \times 10^{-2}$	$(10.1 \pm 0.4) \times 10^{-2}$

Consistently counted
 $D_1 \rightarrow D\pi\pi$ contribution

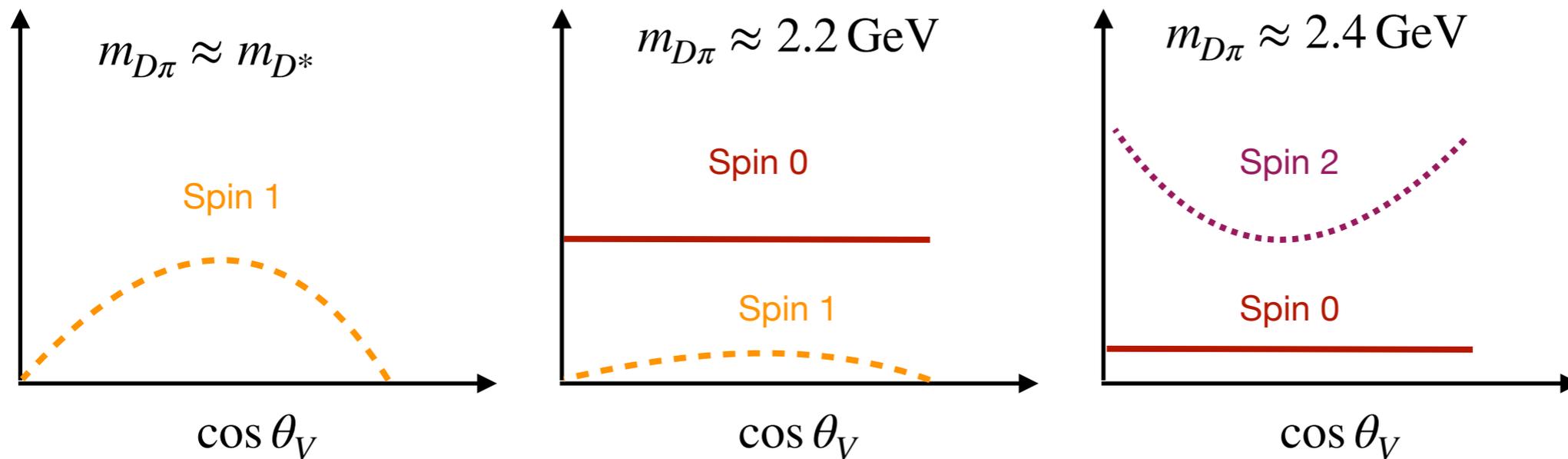
$$\mathcal{B}(B^+ \rightarrow X_c \ell \bar{\nu}_\ell) - \mathcal{B}(B^+ \rightarrow D^{(*)} \ell \bar{\nu}_\ell) - \mathcal{B}(B^+ \rightarrow D^{**}(1P) \ell \bar{\nu}_\ell) - \mathcal{B}(B^+ \rightarrow D^{(*)}\pi\pi \ell \bar{\nu}_\ell) = \mathcal{B}(B^+ \rightarrow X_c^{\text{Gap}} \ell \bar{\nu}_\ell)$$

$$\rightarrow \mathcal{B}(B^+ \rightarrow X_c^{\text{Gap}} \ell \bar{\nu}_\ell) = (0.8 \pm 0.5) \times 10^{-2}$$

Since we have no clue what populates this 'gap' a 100% error seems prudent, a possible candidate is $B \rightarrow D^{(*)}\eta \ell \bar{\nu}_\ell$

1. Measure the $B \rightarrow D^{(*)}\pi\ell\bar{\nu}_\ell$ $\{q^2, \cos\theta_V, m_{D^*\pi}\}$ spectra

Can disentangle different states :



Important to do this in 2D or 3D — maybe can be done **untagged**?

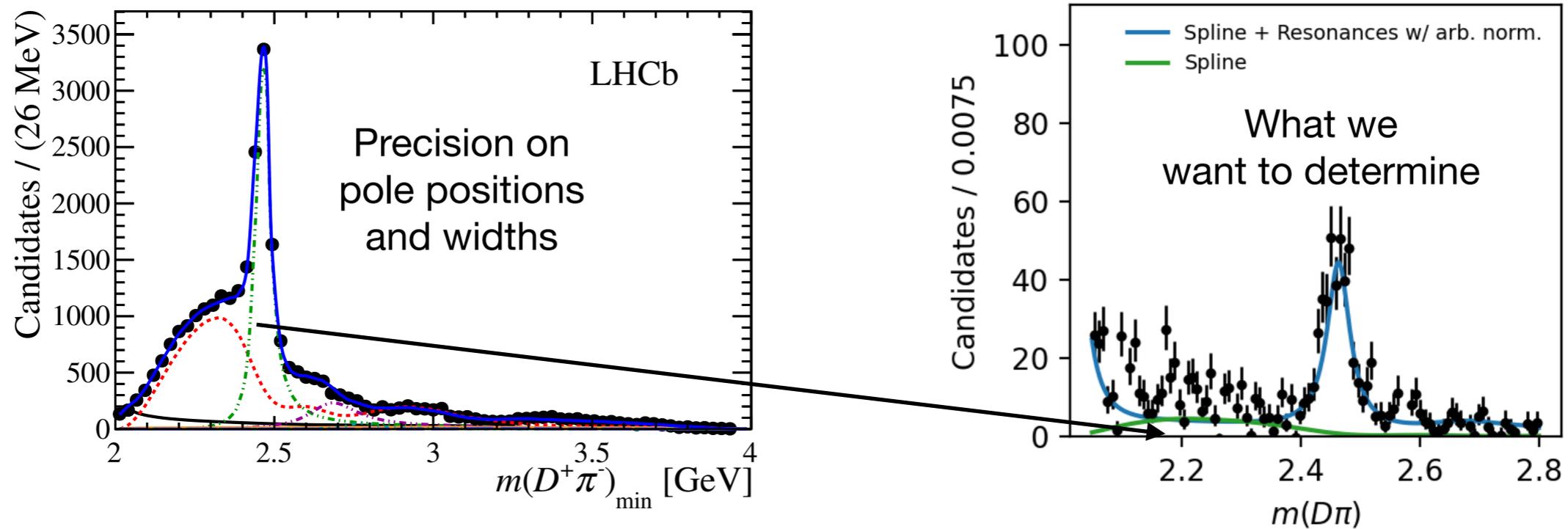
Can **extract** with this the q^2 **spectrum** for the different spin types, can e.g. study of size of break down of on-shell assumptions of $B \rightarrow D^*\ell\bar{\nu}_\ell$

The q^2 spectrum is **extremely important** to improve our knowledge on $B \rightarrow D^{**}(1P)\ell\bar{\nu}_\ell$

2. Combined analysis / interpretation of $B \rightarrow D^{(*)}\pi\ell\bar{\nu}_\ell$ and $B \rightarrow D^{(*)}\pi\pi$

Use **precision** of hadronic analyses and **combine** it with semileptonic modes to determine BFs.

Huge **difference** in sensitivity to **determine** poles and widths

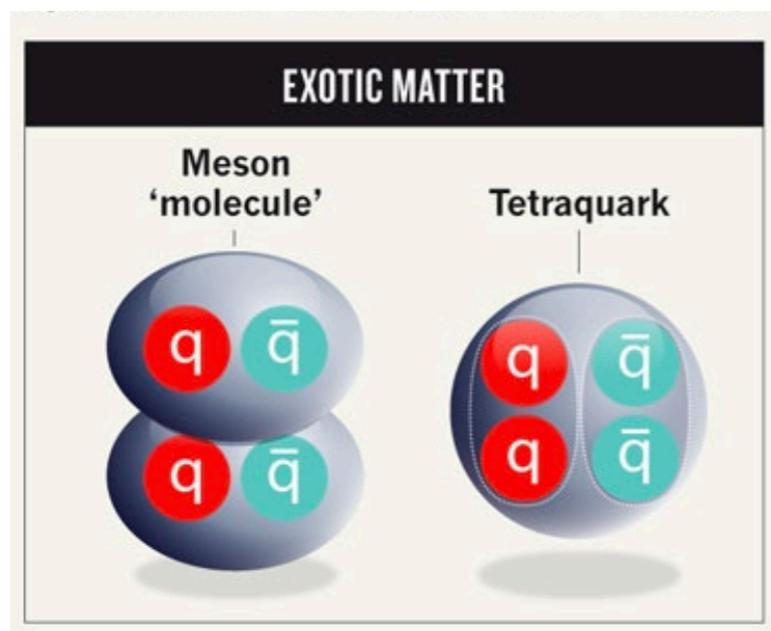


The second pole could be a molecule state, we should study it

2. Combined analysis / interpretation of $B \rightarrow D^{(*)}\pi\ell\bar{\nu}_\ell$ and $B \rightarrow D^{(*)}\pi\pi$

Use **precision** of hadronic analyses and **combine** it with semileptonic modes to determine BFs.

Huge **difference** in sensitivity to **determine** poles and widths

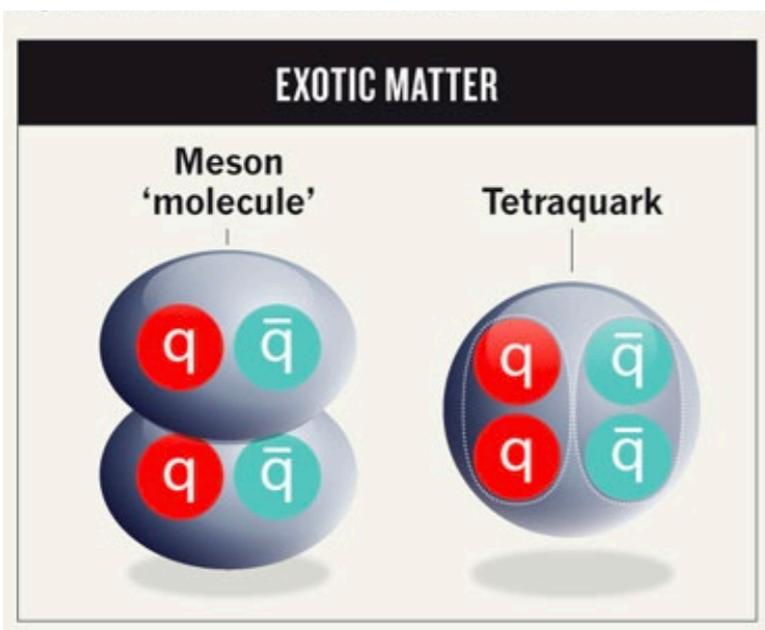


We should study the nature of the second pole. Is it indeed something exotic? Or another (expected) higher state?

3.a Searches for the missing and exotic $B \rightarrow X\ell\bar{\nu}_\ell$ contributions

Dedicated searches
for $B \rightarrow D^{(*)}\eta\ell\bar{\nu}_\ell$,
 $B \rightarrow D^{(*)}n\pi\ell\bar{\nu}_\ell$,
...

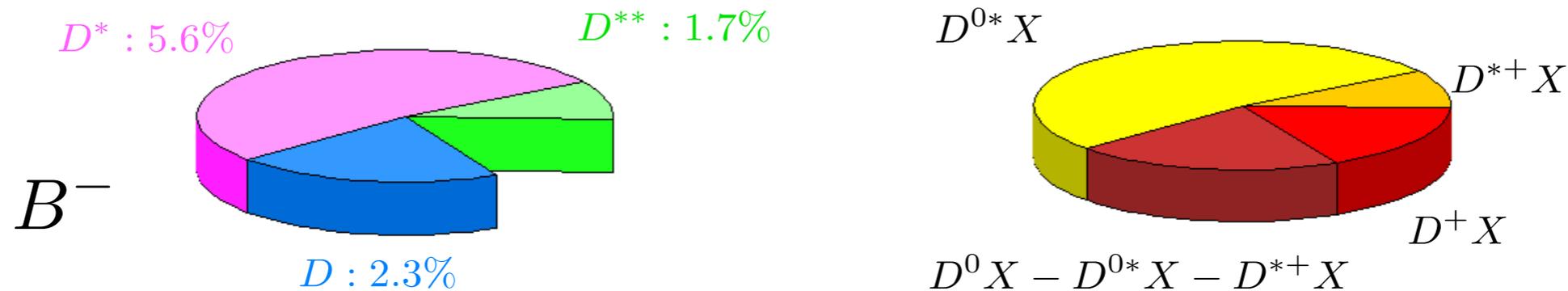
Decay	$\mathcal{B}(B^+)$	$\mathcal{B}(B^0)$
$B \rightarrow D\ell^+\nu_\ell$	$(2.4 \pm 0.1) \times 10^{-2}$	$(2.2 \pm 0.1) \times 10^{-2}$
$B \rightarrow D^*\ell^+\nu_\ell$	$(5.5 \pm 0.1) \times 10^{-2}$	$(5.1 \pm 0.1) \times 10^{-2}$
$B \rightarrow D_1\ell^+\nu_\ell$	$(6.6 \pm 0.1) \times 10^{-3}$	$(6.2 \pm 0.1) \times 10^{-3}$
$B \rightarrow D_2^*\ell^+\nu_\ell$	$(2.9 \pm 0.3) \times 10^{-3}$	$(2.7 \pm 0.3) \times 10^{-3}$
$B \rightarrow D_0^*\ell^+\nu_\ell$	$(4.2 \pm 0.8) \times 10^{-3}$	$(3.9 \pm 0.7) \times 10^{-3}$
$B \rightarrow D_1'\ell^+\nu_\ell$	$(4.2 \pm 0.9) \times 10^{-3}$	$(3.9 \pm 0.8) \times 10^{-3}$
$B \rightarrow D\pi\pi\ell^+\nu_\ell$	$(0.6 \pm 0.9) \times 10^{-3}$	$(0.6 \pm 0.9) \times 10^{-3}$
$B \rightarrow D^*\pi\pi\ell^+\nu_\ell$	$(2.2 \pm 1.0) \times 10^{-3}$	$(2.0 \pm 1.0) \times 10^{-3}$
$B \rightarrow D\eta\ell^+\nu_\ell$	$(4.0 \pm 4.0) \times 10^{-3}$	$(4.0 \pm 4.0) \times 10^{-3}$
$B \rightarrow D^*\eta\ell^+\nu_\ell$	$(4.0 \pm 4.0) \times 10^{-3}$	$(4.0 \pm 4.0) \times 10^{-3}$
$B \rightarrow X_c\ell\nu_\ell$	$(10.8 \pm 0.4) \times 10^{-2}$	$(10.1 \pm 0.4) \times 10^{-2}$



Are there (potentially) any other exotic states in what we have not measured?

3.b Find what is missing: $B \rightarrow D^{(*)} X \ell \bar{\nu}_\ell$ with hadronic tagging

Semi exclusive measurement from Belle (shown at ICHEP 2012)



$$B^0 : \mathcal{B}(D^0 X + D^+ X) / \mathcal{B}(X) = 1.027 \pm 0.018_{\text{stat.}} \pm 0.012_{\mathcal{B}(D)} \pm 0.040_{\text{sys}}$$

$$B^- : \mathcal{B}(D^0 X + D^+ X) / \mathcal{B}(X) = 1.010 \pm 0.015_{\text{stat.}} \pm 0.011_{\mathcal{B}(D)} \pm 0.040_{\text{sys}}$$

https://indico.cern.ch/event/181298/contributions/309049/attachments/243226/340452/20120705_semilep_belle_oswald.pdf

We can look into the ROE and study how many pions, kaons, etc. there are and determine individual resonant + non-resonant contributions

What we should stop doing

Many problematic points in past measurements of broad states

Disagreement maybe a sign of complicated pictures

We need to comment on this somehow in the next HFLAV / PDG update, not sure the averages hold much meaning

We have fairly good measurements of resonant + non-resonant $B \rightarrow D^{(*)}\pi\ell\bar{\nu}_\ell$ which are not affected

To determine the size of the ‘gap’ in our MC, we should just use

$$\mathcal{B}(B \rightarrow X_c^{\text{gap}}\ell\bar{\nu}_\ell) = \mathcal{B}(B \rightarrow X_c\ell\bar{\nu}_\ell) - \mathcal{B}(B \rightarrow D^{(*)}\ell\bar{\nu}_\ell) - \mathcal{B}(B \rightarrow D^{(*)}\pi\ell\bar{\nu}_\ell) - \mathcal{B}(B \rightarrow D^{(*)}\pi\pi\ell\bar{\nu}_\ell)$$

HFLAV 2019 Averages for broad states

Experiment	$\mathcal{B}(B^- \rightarrow D_1^{\prime 0}(D^{*+}\pi^-)\ell^-\bar{\nu}_\ell)[\%]$ (rescaled)	$\mathcal{B}(B^- \rightarrow D_1^{\prime 0}(D^{*+}\pi^-)\ell^-\bar{\nu}_\ell)[\%]$ (published)
DELPHI [511]	$0.73 \pm 0.17_{\text{stat}} \pm 0.18_{\text{syst}}$	$0.83 \pm 0.17_{\text{stat}} \pm 0.18_{\text{syst}}$
Belle [508]	$-0.03 \pm 0.06_{\text{stat}} \pm 0.07_{\text{syst}}$	$-0.03 \pm 0.06_{\text{stat}} \pm 0.07_{\text{syst}}$
BABAR [509]	$0.26 \pm 0.04_{\text{stat}} \pm 0.04_{\text{syst}}$	$0.27 \pm 0.04_{\text{stat}} \pm 0.05_{\text{syst}}$
Average	$0.19 \pm 0.03 \pm 0.04$	$\chi^2/\text{dof} = 11.9/2$ (CL=0.003%)

Extremely poor agreement for D_1'

Experiment	$\mathcal{B}(B^- \rightarrow D_0^{*0}(D^+\pi^-)\ell^-\bar{\nu}_\ell)[\%]$ (rescaled)	$\mathcal{B}(B^- \rightarrow D_0^{*0}(D^+\pi^-)\ell^-\bar{\nu}_\ell)[\%]$ (published)
Belle Tagged B^- [508]	$0.25 \pm 0.04_{\text{stat}} \pm 0.06_{\text{syst}}$	$0.24 \pm 0.04_{\text{stat}} \pm 0.06_{\text{syst}}$
Belle Tagged B^0 [508]	$0.22 \pm 0.08_{\text{stat}} \pm 0.06_{\text{syst}}$	$0.24 \pm 0.04_{\text{stat}} \pm 0.06_{\text{syst}}$
BABAR Tagged [509]	$0.32 \pm 0.04_{\text{stat}} \pm 0.05_{\text{syst}}$	$0.26 \pm 0.05_{\text{stat}} \pm 0.04_{\text{syst}}$
Average	$0.28 \pm 0.03 \pm 0.04$	$\chi^2/\text{dof} = 0.82/2$ (CL=66.4%)

New Belle result

	yield	\mathcal{S}	branching fraction [%]
$B^0 \rightarrow D_0^{*-} \ell^+ \nu_\ell$ with $D_0^{*-} \rightarrow \bar{D}^0 \pi^-$	-	-	<0.044 at 90% CL

We treated $B \rightarrow D^{**} \ell \bar{\nu}_\ell$ for a long time just as a background that needs to be studied

e.g. in $B \rightarrow D^{(*)} \tau \bar{\nu}_\tau$ often dedicated sidebands are studied or used;
not many make it into our final publication

There exist (internally in experiments) much more cross checks on how well our models are describing (or not describing) the data.

Going forward it would be useful if we become more conscious on providing useful information that document and help improve our understanding.